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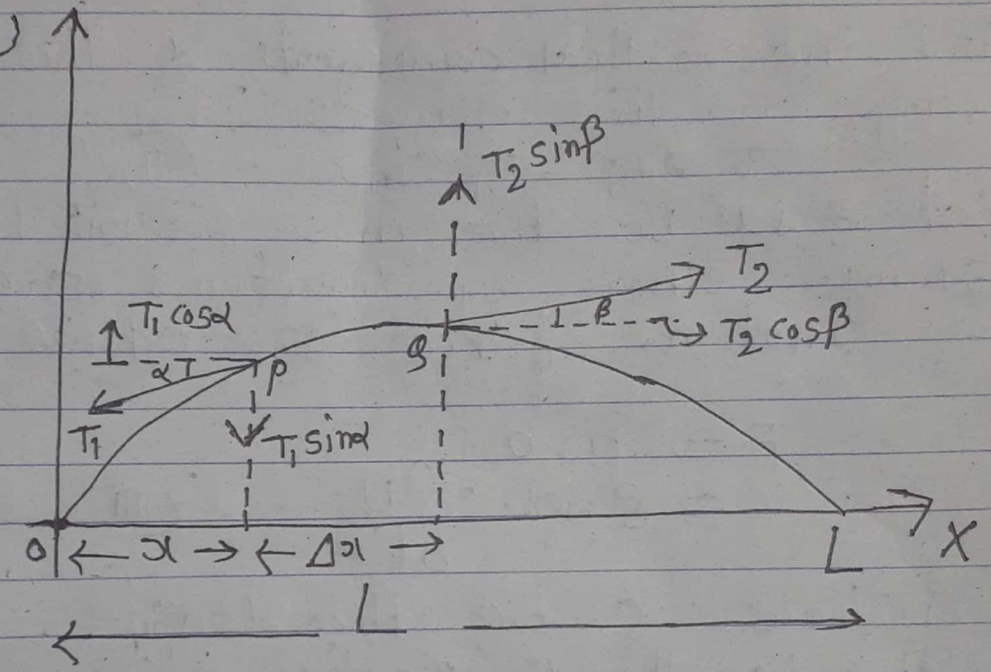
Wave equation (vibration of string)

Derivation of one dimensional wave equation



Consider a tightly stretched elastic string of length $OL = L$ and fix it at the ends O and L. When the string vibrates, each point of the string makes small vibration on U-axis which is shown in figure.

Let $P(x, u)$ and $Q(x+\Delta x, u+\Delta u)$ be two neighbouring points on the string. Let T_1 and T_2 be the tension at the points P and Q where the tensions makes angles α and β with X-axis. Since the point of the string moves vertically, there is no motion in the horizontal direction. Thus the horizontal components of the tension must be equal i.e constant

$$\therefore \quad T_1 \cos\alpha = T_2 \cos\beta = T \text{ (say)} \quad \ldots (1)$$

Also the vertical components of the force acting on this elements is

$$T_2 \sin\beta - T_1 \sin\alpha \quad \ldots (11)$$

Let $\rho \Delta x$ be the mass per unit length of string. Then by Newton's second law of motion, force is equal to

$$F = ma$$
$$= \rho \Delta x \cdot \frac{\partial^2 u}{\partial t^2} \quad \ldots (111)$$

Where ρ be the density, Δx is the length of portion string and $\frac{\partial^2 u}{\partial t^2}$ is acceleration

From eqn (11) and (111)

$$T_2 \sin\beta - T_1 \sin\alpha = \rho \Delta x \cdot \frac{\partial^2 u}{\partial t^2} \quad \ldots (1V)$$

Now dividing eqn (1V) by (11), we get

$$\frac{T_2 \sin\beta}{T_2 \cos\beta} - \frac{T_1 \sin\alpha}{T_1 \cos\alpha} = \frac{\rho \Delta x \cdot \frac{\partial^2 u}{\partial t^2}}{T}$$

$$\tan\beta - \tan\alpha = \frac{\rho \Delta x}{T} \cdot \frac{\partial^2 u}{\partial t^2} \quad \ldots (V)$$

We know that $\tan\alpha$ and $\tan\beta$ are the slopes of the string at x and $x+\Delta x$ respectively

Where $\tan\alpha = \left.\frac{\partial U}{\partial x}\right|_{at\ x}$

$$\tan\beta = \left.\frac{\partial U}{\partial x}\right|_{at\ x+\Delta x}$$

Eqn (V) becomes, we get

$$\left[\left.\frac{\partial U}{\partial x}\right|_{at\ x+\Delta x} - \left.\frac{\partial U}{\partial x}\right|_{at\ x}\right] = \frac{\rho\,\Delta x}{T}\cdot\frac{\partial^2 U}{\partial t^2}$$

or, $$\frac{1}{\Delta x}\left[\left.\frac{\partial U}{\partial x}\right|_{at\ x+\Delta x} - \left.\frac{\partial U}{\partial x}\right|_{at\ x}\right] = \frac{\rho}{T}\cdot\frac{\partial^2 U}{\partial t^2}$$

Taking limit $\Delta x \to 0$ on both sides, we get

or $$\lim_{\Delta x\to 0}\frac{1}{\Delta x}\left[\left.\frac{\partial U}{\partial x}\right|_{x+\Delta x} - \left.\frac{\partial U}{\partial x}\right|_{x}\right] = \lim_{\Delta x\to 0}\frac{\rho}{T}\frac{\partial^2 U}{\partial t^2}$$

or $$\lim_{\Delta x\to 0}\left[\frac{\left.\frac{\partial U}{\partial x}\right|_{x+\Delta x} - \left.\frac{\partial U}{\partial x}\right|_{x}}{\Delta x}\right] = \frac{\rho}{T}\cdot\frac{\partial^2 U}{\partial t^2}$$

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or $$\frac{\partial^2 u}{\partial x^2} = \frac{\rho}{T} \cdot \frac{\partial^2 u}{\partial t^2}$$

or, $$\frac{\rho}{T} \cdot \frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2}$$

or $$\frac{\partial^2 u}{\partial t^2} = \frac{T}{\rho} \frac{\partial^2 u}{\partial x^2}$$

or, $$\boxed{\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}}$$

Where $c^2 = \frac{T}{\rho}$

This is the required one dimensional wave equation.

Solution of one dimensional wave equation under certain initial and Boundary conditions

Solution:- We have one dimensional wave eqn is

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} \quad \cdots (1)$$

Where $u(x, t)$ is the deflection of the string with boundary condition

$u(0, t) = 0$ and $u(L, t) = 0$ for all t .... (2)

then its initial condition is

$u(x, 0) = f(x)$ ... (3) initial deflection

$\left.\frac{\partial u}{\partial t}\right|_{t=0} = g(x)$ ... (4) initial velocity

Suppose $u(x, t) = F(x) . G(t)$ ... (5)

be the Solution of eqn (1)

Differentiating eqn (5) w.r. to x and t

$$\frac{\partial^2 u}{\partial t^2} = F . \ddot{G} \quad \text{and} \quad \frac{\partial^2 u}{\partial x^2} = F'' G$$

Where

dots denotes derivative w.r. to t

prime denotes derivative w.r. to x

putting these values in eqn (1)

we get

$$F\ddot{G} = c^2 F'' G$$

$$\Rightarrow \frac{F''}{F} = \frac{\ddot{G}}{G \cdot c^2}$$

$$\Rightarrow \frac{F''}{F} = \frac{\ddot{G}}{G \cdot c^2} = K \text{ (say)}$$

which gives

$$F'' - FK = 0 \quad \cdots (6)$$

$$\ddot{G} - c^2 K G = 0 \quad \cdots (7)$$

These are second order ordinary differential equation. We have to determine F and G from eqn (6) and (7) under boundary conditions.

Also we have

$$U(x, t) = F(x) \cdot G(t)$$

i.e $U(0, t) = F(0) \cdot G(t) = 0$

$$U(L, t) = F(L) \cdot G(t) = 0$$

for all t

For solving eqn (6) we have to generate boundary condition (b.c.)

If $G = 0$, then $U = 0$ which is no meaning or interest. So $G \neq 0$

Then we get

$$F(0) = 0 = F(L) \quad \cdots (8)$$

case (i) Let $K = 0$ then eqn (6) becomes

$$F'' = 0$$

integrating

$$F' = a$$

again integrating

$$F = ax + b \quad \cdots (9)$$

Where a and b are constants. We have to find a and b under boundary condition

We have

$$F(0) = 0 \text{ and } F(L) = 0$$

$\therefore \quad F(x) = ax + b$

or $F(0) = a \cdot 0 + b$

$0 = b$

$\therefore \quad b = 0$

Again

$F(L) = aL + b$

$0 = aL + 0$

$\therefore a = 0$

We get $a = b = 0$

putting these values of a and b in eqn (9)

$$F = 0 \cdot x + 0 \Rightarrow F = 0$$

Which is no interest.

case (ii) Suppose that $K > 0$ ie K is positive

Let $K = \lambda^2$

Then eqn (6) becomes

$$F'' - \lambda^2 F = 0$$

its auxillary eqn is

$$m^2 - \lambda^2 = 0$$

$$m = \pm \lambda$$

its general solution is

$$F(x) = A \cdot e^{\lambda x} + B \cdot e^{-\lambda x} \quad \cdots (10)$$

Where A and B are constants to be determine under boundary condition.

i.e $F(0) = A \cdot e^0 + B \cdot e^0$

$$0 = A + B$$

$$A + B = 0 \quad \cdots (i)$$

Again $F(L) = A \cdot e^{\lambda L} + B \cdot e^{-\lambda L}$

$$0 = A \cdot e^{\lambda L} + B \cdot e^{-\lambda L}$$

$$A \cdot e^{\lambda L} + B \cdot e^{-\lambda L} = 0 \quad \cdots (ii)$$

Solving (i) and (ii) we get

$A = B = 0$.

$\therefore$ Eqn (10) becomes

$F = 0 \Rightarrow U = 0$

This is also no interest.

case (iii) Let $K < 0$

i.e K is negative

Suppose $K = -p^2$

Eqn (6) becomes

$$F'' + p^2 F = 0$$

A. E. is

$$m^2 + p^2 = 0$$

$$m = \pm pi$$

its general solution is

$$F(x) = e^{px}(A\cos qx + B\sin qx)$$ where p = real part, q = imaginary part

$$F(x) = e^{0\cdot x}(A\cos px + B\sin px)$$

$$F(x) = A\cos px + B\sin px \quad \text{--- (11)}$$

using boundary condition

$$F(0) = A\cos 0 + B\sin 0$$

$$0 = A\cdot 1 + 0$$

$$\therefore A = 0$$

Also

$$F(L) = A\cos pL + B\sin pL$$

$$0 = 0\cdot\cos pL + B\sin pL$$

$$\therefore B\sin pL = 0$$

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Let $B \neq 0$. So

$$\sin pL = 0$$

$$\sin pL = \sin n\pi$$

$$\cdot \quad pL = n\pi$$

$$p = \frac{n\pi}{L}$$ where n is an integer

We have $B \neq 0$, setting $B = 1$

Eqn (11) becomes

$$F(x) = 0 \cdot \cos px + 1 \cdot \sin \frac{n\pi}{L} \cdot x$$

$$F(x) = \sin \frac{n\pi}{L} \cdot x$$

We obtain infinitely many soln

$$F(x) = F_n(x)$$

i.e $F_n(x) = \sin \frac{n\pi}{L} \cdot x \quad \cdots (12)$

for $n = 1, 2, 3, \ldots$

Similarly for solution of eqn (7) i.e

$$\ddot{G} - c^2 K G = 0$$

We have $K = -p^2 = -\left(\frac{n\pi}{L}\right)^2$

Then Eqn (7) becomes

$$\ddot{G} + c^2\left(\frac{n\pi}{L}\right)^2 G = 0$$

or $$\ddot{G} + \left(\frac{cn\pi}{L}\right)^2 G = 0$$

$$\ddot{G} + \lambda_n^2 G = 0 \quad \text{where} \quad \lambda_n = \frac{cn\pi}{L}$$

its auxillary eqn is

$$m^2 + \lambda_n^2 = 0$$

$$\therefore m = \pm \lambda_n i \qquad \text{here } p = 0,\ q = \lambda_n$$

its general solution

$$G(t) = e^{pt}\left(B_n \cos\lambda_n t + B_n^* \sin\lambda_n t\right)$$

$$G(t) = e^0\left(B_n \cos\lambda_n t + B_n^* \sin\lambda_n t\right)$$

$$G(t) = B_n \cos\lambda_n t + B_n^* \sin\lambda_n t$$

i.e $$G_n(t) = B_n \cos\lambda_n t + B_n^* \sin\lambda_n t$$

Therefore required solution of eqn (1) is

$$U_n(x,t) = F_n(x) \cdot G_n(t)$$

$$= \sin\frac{n\pi}{L}x \left(B_n \cos\lambda_n t + B_n^* \sin\lambda_n t\right)$$

By fundamental theorem the solution of given wave eqn is the infinitely many u_n. Thus

$$u(x,t) = \sum_{n=1}^{\infty} u_n(x,t)$$

$$\Rightarrow u(x,t) = \sum_{n=1}^{\infty} \left(B_n \cos \lambda_n t + B_n^* \sin \lambda_n t\right) \sin \frac{n\pi}{L} \cdot x \quad \text{--(13)}$$

From eqn (3) [initial condition]

$$u(x,0) = f(x)$$

Then eqn (13) becomes

$$u(x,0) = \sum_{n=1}^{\infty} \left(B_n \cos 0 + B_n^* \sin 0\right) \cdot \sin \frac{n\pi}{L} x$$

$$\text{i.e} \quad f(x) = \sum_{n=1}^{\infty} B_n \sin \frac{n\pi}{L} x$$

Which is fourier sine series and its co-efficient B_n is given by

$$B_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi}{L} x \, dx$$

Also, we have

$$\left.\frac{\partial u}{\partial t}\right|_{t=0} = g(x)$$

Eqn (13) becomes

$$\Rightarrow \left[\sum_{n=1}^{\infty}\left(-B_n \lambda_n \sin \lambda_n t + B_n^* \lambda_n \cos \lambda_n t\right) \sin \frac{n\pi}{L} x\right]_{t=0} = g(x)$$

put $t = 0$

$$\Rightarrow \sum_{n=1}^{\infty} B_n^* \lambda_n \sin \frac{n\pi}{L} x = g(x)$$

or $$g(x) = \sum_{n=1}^{\infty} B_n^* \lambda_n \sin \frac{n\pi}{L} x$$

Which is also fourier sine series of $g(x)$. its co-efficient is given by

$$B_n^* \lambda_n = \frac{2}{L} \int_0^L g(x) \sin \frac{n\pi}{L} x \, dx$$

$$B_n^* = \frac{2}{\lambda_n L} \int_0^L g(x) \sin \frac{n\pi}{L} x \, dx$$

$$= \frac{2}{\frac{cn\pi}{L} \times L} \int_0^L g(x) \sin \frac{n\pi}{L} x \, dx$$

where $\lambda_n = \frac{cn\pi}{L}$

$$= \frac{2}{cn\pi} \int_0^L g(x) \sin \frac{n\pi}{L} x \, dx$$

Therefore we get required Solution of one dimensional wave eqn is

$$U(x,t) = \sum_{n=1}^{\infty}\left(B_n \cos\lambda_n t + B_n^* \sin\lambda_n t\right)\sin\frac{n\pi}{L}x$$

Where

$$B_n = \frac{2}{L}\int_0^L f(x)\sin\frac{n\pi}{L}x\,dx$$

$$B_n^* = \frac{2}{cn\pi}\int_0^L g(x)\sin\frac{n\pi}{L}x\,dx$$

Where $f(x)$ and $g(x)$ are initial deflection and initial velocity of the string respectively.