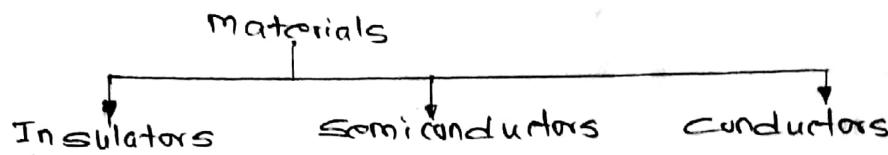


Chapter 1:

Semiconductor Diode

Based on the electrical conductivity, all the materials in nature are classified as insulators, semiconductors and conductors.



Insulators

An insulator is a material that are characterized by poor electrical conductivity. For eg; Glass, quartz, rubber, porcelain, bakelite, mica etc. The resistivity level of an insulator is of the order of 10^{10} to $10^{12} \Omega\text{-m}$.

The energy band structure of an insulator is as shown below:

Band structure of a material defines the band of energy level that an electron can occupy. Valence band is the range of electron energy where the electron remain bonded to the atom & do not contribute to the electric current.

Conduction band is the range of electron energies higher than valence band where e⁻s are free to accelerate under the influence of external voltage source resulting flow of charge.

The energy band between the valence band & conduction band is called Forbidden band gap or Forbidden energy gap. It is the energy required by an e⁻ to move from Valence band to Conduction band.

$$1\text{eV} = 1.6 \times 10^{-19} \text{J}$$

There is a large Forbidden band gap $\geq 5\text{eV}$. Because of this large gap, there are very few e⁻s in the Conduction Band & hence conductivity of insulator is poor. Even an increase in temperature or applied electric field is insufficient to transfer e⁻s from VB to CB.

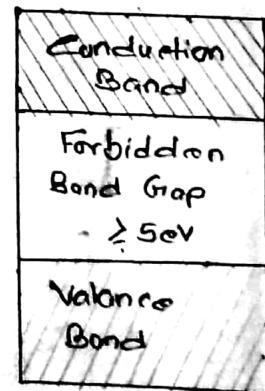


Fig: Energy Band structure of Insulator

Conductors:

Conductor or conducting materials are good conductors of electricity and characterized by a large conductivity & small resistivity. For eg; Copper, Aluminium, Silver etc. The resistivity of conductors is in the range of $10^{-8} \Omega\text{-m}$.

In a conductor, there is no forbidden energy gap between valence band & conduction band i.e. they overlap each other. Electrons randomly move through the solid. So electrons in the conductors are called free es. Therefore at a room temperature when electric field is applied, an electric current flows through a conductor. The resistance of conductor increases with increase in temp.

Semiconductor

A semiconductor is a material that has its conductivity somewhere between the insulator & conductor. The resistivity of semiconductor material lies in the range of 10^{-5} to $10^5 \Omega\text{-m}$.

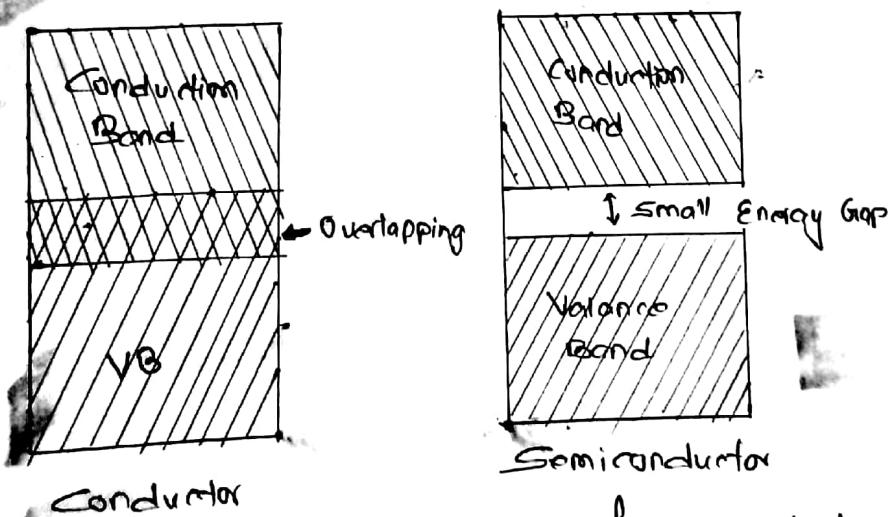


Fig: Energy Band diagrams for conductor and Semiconductors.

For eg; Silicon, Germanium etc. The Energy Gap for Si, Ge & GaAs is 1.1 eV, 0.785 eV & 1.42 eV respectively at absolute zero (i.e. zero Kelvin or -273.16°C). And at this temperature, the valence band does not have sufficient energy to move from valence band to conduction band.

Thus semiconductor acts as insulator at absolute zero. However at normal room temperature (25°C), semiconductors can easily jump to conduction band due to low energy gap between valence band and conduction band. Thus semiconductors are capable of conducting small current even at normal room temperature. The resistance of semiconductor decreases with an increase in temperature.

Types of Semiconductor

The resistance of semiconductor decreases with an increase of temperature i.e. temp coefficient of resistance of semiconductor is negative. They behave like an insulator at very low temperature but act as conductor at high temperatures.

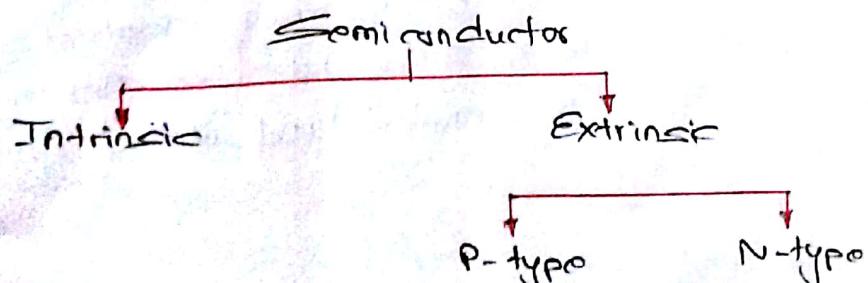
The mostly used semiconductors are Germanium (Ge) & Silicon (Si). Selenium (Se) & Tellurium (Te) are also a semiconductor.

Properties of Semiconductor (SiC)

1. The resistivity of SiC depends on illumination and decreases in bright surroundings.
2. The resistivity of SiC depends on the magnitude of electric field.
3. SiC are non-linear elements.
4. The conductivity of SiC changes considerably when small amount of impurities added to it. So that its conductivity can be controlled.

There are 2 types of SiC

- i) Intrinsic (pure) SiC
- ii) Extrinsic (impure) SiC.



Intrinsic Semiconductor

An intrinsic SiC is the purest form of SiC. The impurity content is less than one part in 100 million parts of the intrinsic SiC. There are many SiC materials such as Germanium, Silicon, Selenium, Tellurium, Gallium Arsenide etc & among these Ge & Si are widely used in a SiC devices.

Production of holes & free electrons

Let us consider a structure of Silicon atom. Silicon atom no. is 14 & has 4 valence electrons. These 4 e⁻s are shared by 4 neighbouring atoms in the crystal structure by means of covalent bond.

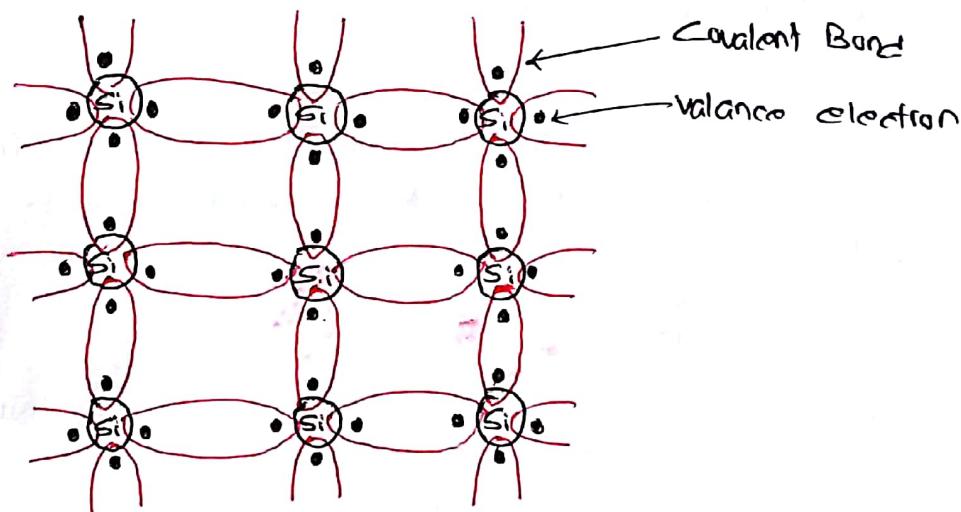


Fig: Crystal structure of Si at absolute zero (-273°C)

Above fig shows the crystal structure of Si at absolute zero temperature. Thus pure Si acts as poor conductor (due to lack of free e⁻s) at a low temperature or at absolute zero temperature.

At room temp., some of the covalent bonds break up. As a result, e⁻s are released and move to the conduction band, are called free e⁻s. The vacancy created in broken bond is called a Hole. Since hole is created after a removal of e⁻ from the covalent bond, it is considered as a positive charge.

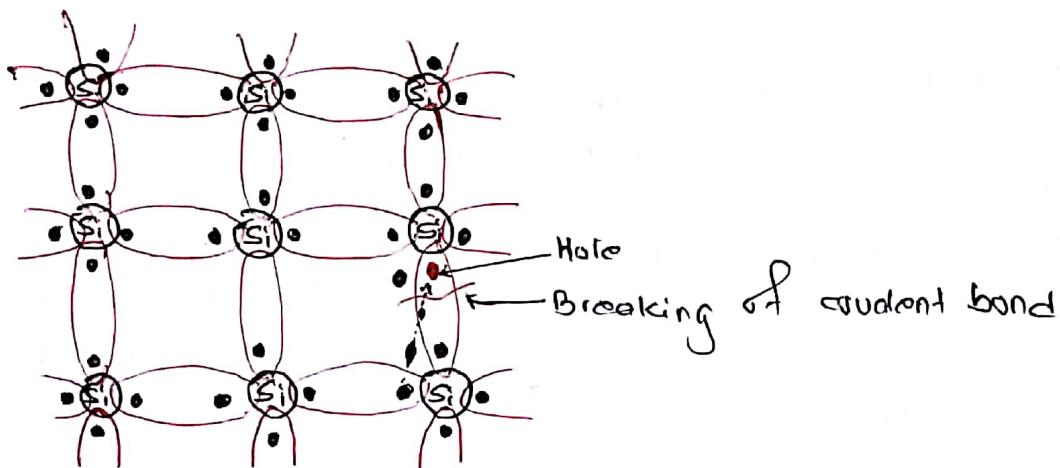


Fig: production of hole & free e^-

Electrons & holes are created in pairs. When a free electron approaches the hole, it gets attracted and falls into the hole. This process of merging of free e^- & a hole is called recombination. During this process, both the free e^- s & holes disappear. However, an energy is released and causes another covalent bond to breakdown & generate a new electron, hole pair.

Conduction in Intrinsic Semiconductors

When an e^- s are released during the breakdown of covalent bonds, they move randomly through a crystal lattice. When an external electric field is applied to the pure intrinsic SiC, the free e^- s in the conduction band move towards the +ve terminal while holes move towards the -ve terminal of power supply (battery). As a result, an electric current flows. This current is very small.

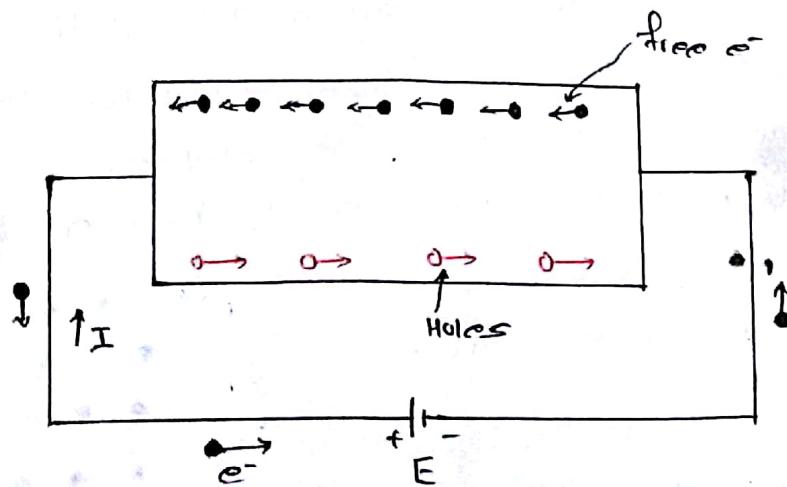


Fig: Conduction of current in Intrinsic SiC

Conduction in Semiconductor

SIC consists of 2 types of charge carriers: positive charge carriers called holes & negative charge carriers called electrons.

The movement of these charge carriers give rise to the flow of current in SiC. Thus we can say that SiC passes two types of current.

- i) Electron current
- ii) Hole current

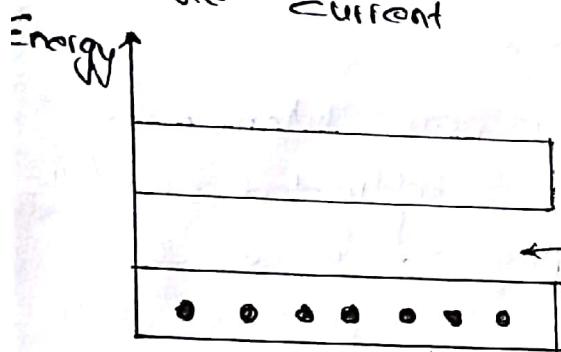
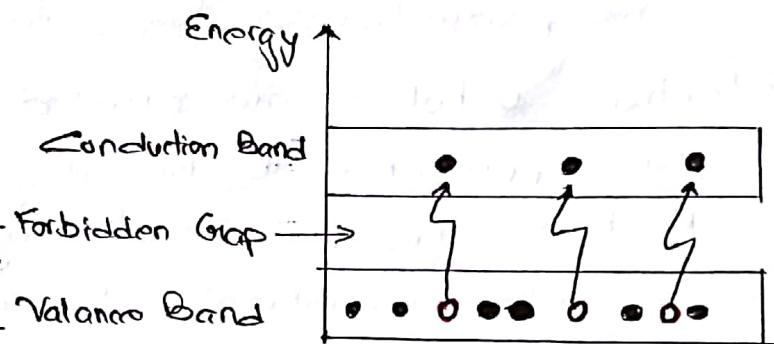


Fig: When no energy is applied.



Fig! When energy is applied

v) Electron Current

When an energy is applied to a SiC, the valence e^- at valence band can gain sufficient energy and jump to the conduction band and become free e^- .

The movement of these free e^- s gives rise to a electron current (I_e).

\Rightarrow Hole Current

When a valence e⁻ gain sufficient energy & jumped to conduction band, they leave a hole at the valence bond. Further valence electrons can move through this hole. Thus relative motion of valence e⁻s at valence band through holes give rise to hole current (I_h)

Thus, total current in a SIC is the sum of e⁻ current & hole current.

$$\text{i.e. } \mathcal{I} = \mathcal{I}_e + \mathcal{I}_h$$

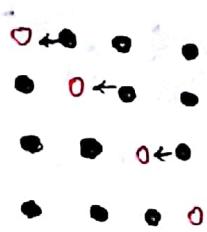


Fig: movement of hole

Extrinsic Semiconductor

An intrinsic SiC has very limited applications/significance as it conduct very small amount of current at room temperature. However, the electrical conductivity of intrinsic SiC can be increased significantly by adding small amount of impurity to it during a crystallization process. This process of adding an impurity to a pure SiC is called Doping and the resultant SiC formed is called extrinsic SiC. The amount of impurity added is 1 part in 10^6 atoms.

There are 2 types of Extrinsic SiC

1) N-type SiC

2) P-type SiC

1) P-type Semiconductor

When a small amount of trivalent impurity such as aluminium, Gallium, Indium, Boron is added to a pure SiC, then P-type SiC is formed.

The crystal structure of P-type SiC is as shown below:

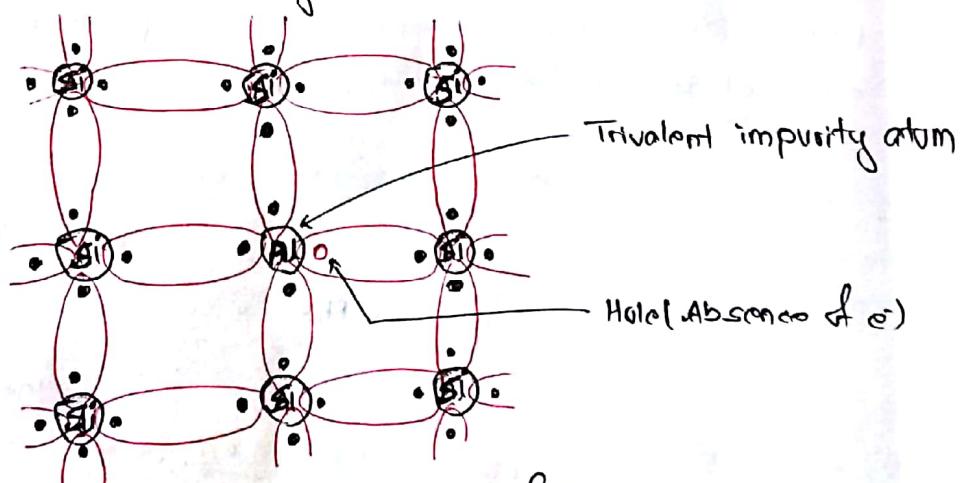


Fig: Crystal Structure of P-type SiC

The 3 valence e^- s of the trivalent impurity forms 3 covalent bonds with the neighboring atoms and a vacancy exists in the fourth bond giving rise to the holes, which can accept an electron. So, trivalent impurities are also called acceptor impurities or acceptor agents.

In P-type SiC, holes are majority carriers & e^- s are minority carriers.

The energy band diagram of P-type SiC is

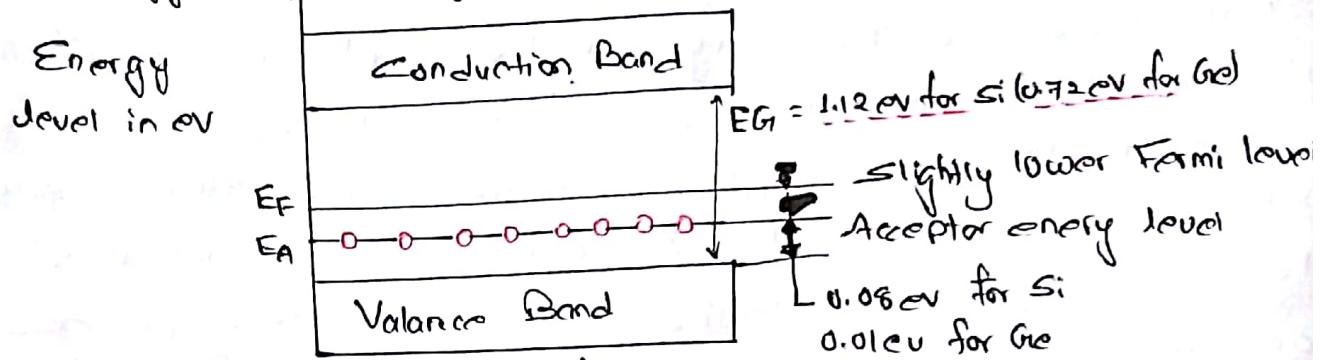


Fig: Energy band diagram of P-type SiC

A very small amount of energy (0.08eV in case of Silicon & 0.01eV in case of Germanium) is required for an e^- to leave the valence bond and occupy the acceptor energy level, holes are created in the valence bond by these electrons.

The holes so created constitute the larger number of carriers in the SiC materials.

3) N-Type Semiconductor

When a small amount of pentavalent impurity is added, such as Arsenic, Antimony, Bismuth, Phosphorous, to a pure SiC during crystallization process, the resulting crystal is called N-type SiC.

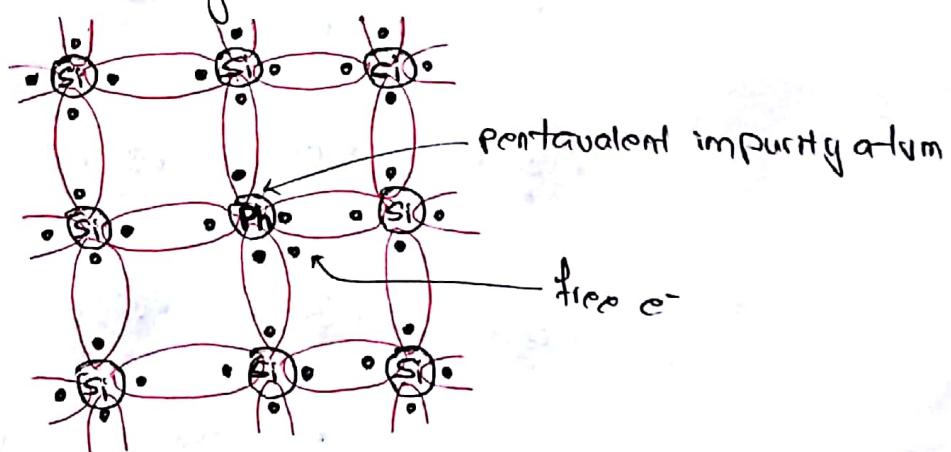


Fig: Crystal structure of N-type SiC

For each impurity added, one free e^- is generated, which then moves to conduction band. As the pentavalent impurity donates electron to crystal lattice, it is also called Donor impurity or donor agent. In N-type SiC, e^- s are majority charge carriers.

* Fermi Level is the highest energy state occupied by e^- in a material at absolute zero temp.

The energy band diagram for N-type Si is,

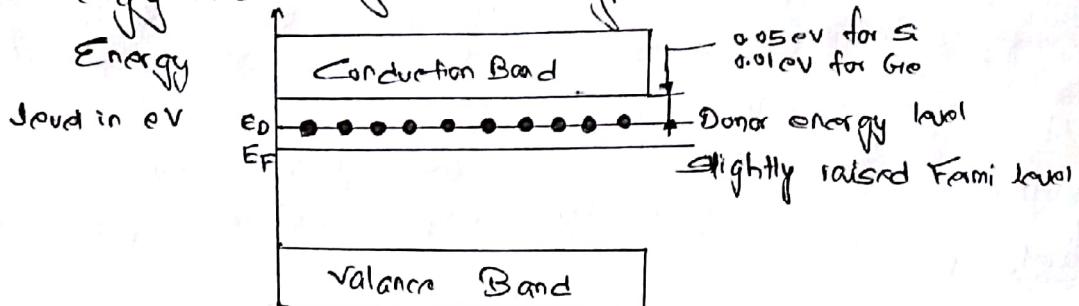


Fig: Energy band diagram for N-type Si.

Here, the energy required to detach the fifth e^- from the atom is only 0.05eV for Silicon & 0.01eV for Germanium. This energy is \approx small that at room temperature practically all such e^- s become free.

Conductivity of Semiconductor

In a pure SiC, the no. of holes is equal to the no. of electrons.

Thermal agitation continue to produce new electron-hole pairs and e^- -hole pairs disappear because of recombination. With each electron hole pair created, two charge carrying particles are formed. One is -ve which is free e^- with mobility μ_e . The other is +ve, i.e. hole with mobility μ_h .

[The e^- & hole move in opposite direction in an electric field E , but since they are of opposite sign, the current due to both is in the same direction.] Hence total current density J within the intrinsic SiC is given by,

$$J = J_e + J_h \quad J = J_n + J_p$$

$$= qn\mu_e E + qp\mu_h E$$

$$J = (n\mu_e + p\mu_h) qE = \sigma E$$

where, $n = \text{no. of } e^- \text{ per unit volume}$.

$p = \text{no. of holes } \dots \dots$

$q = \text{charge of } e^- \text{ or holes in Coulombs}$

$E = \text{applied electric field strength, V/m}$

Hence, σ is conductivity of SiC & is given by.

$$\begin{aligned} * J &= \text{electric current per unit area of cross section} = I/A \\ &= \frac{Ne}{AT} = \frac{Net}{ATL} \quad (\because \frac{dq}{dt} = \frac{Ne}{T}) \\ &= \frac{N \cdot L}{AT} \cdot \rho = nve \quad (n = N/L \leftarrow \text{length} \\ &= n \sigma E = \sigma qE \quad v = \frac{L}{t} \leftarrow \text{drift velocity} \\ &\quad (\because v = uE) \end{aligned}$$

The resistivity of SiC is reciprocal of conductivity.

$$\rho = \frac{1}{\sigma}$$

It is evident from the above eqn that, current density within a SiC is directly proportional to applied electric field E .

For pure SiC, $n = p = n_i$

where n_i = intrinsic concentration.

The value of n_i is given by,

$$n_i^2 = A T^3 \exp(-E_{G0}/kT)$$

where, A = constant & is independent of temp
 $= 2.88 \times 10^{43}$

k = Boltzmann constant in eV per K

E_{G0} = Energy gap in eV at absolute zero
temp. E_G

T = temp

Thus,

$$\begin{aligned} J &= (n_i u_e + p u_h) q E \\ &= (n_i u_e + n_i u_h) q E \\ &= n_i (u_e + u_h) q E \end{aligned}$$

Hence conductivity in intrinsic SiC vs,

$$\sigma_i = n_i (u_e + u_h) q$$

Intrinsic conductivity increases at the rate of 5% per °C for Ge & 7% per °C for Si

Conductivity in Extrinsic SiC

The conductivity of intrinsic SiC is given by

$$\sigma_i = n_i (u_e + u_h) q = (n u_e + p u_h) q$$

For N-type, $n \gg p$

$$\therefore \sigma = n u_e q$$

For P-type, $p \gg n$

Theory of PN Junction

When a block of P-type semiconductor is joined with a block of N-type SiC, then a junction is formed in between them, which is called PN junction.

A ^h	A ^h	A ^h	A ^h
A ^h	A ^h	A ^h	A ^h
A ^h	A ^h	A ^h	A ^h

P-type SiC

D ^e	D ^e	D ^e	D ^e
D ^e	D ^e	D ^e	D ^e
D ^e	D ^e	D ^e	D ^e

N-type SiC

A ^h	A ^h	A ^h	A ^h	D ^e	D ^e	D ^e	D ^e
A ^h	A ^h	A ^h	A ^h	D ^e	D ^e	D ^e	D ^e
A ^h	A ^h	A ^h	A ^h	D ^e	D ^e	D ^e	D ^e

↑
PN Junction

Fig: Formation of PN junction

Hence, A = Acceptor atom (Trivalent)

D = Donor atom (Pentaivalent)

e = electron

h = hole.

During the formation of PN junctions, the e⁻s from N-type SiC move towards the P-type SiC & combine with the holes. The donor atoms of N-type ~~loses~~ e⁻ so it becomes ~~the~~ charged ion (D⁺) whereas the Acceptor atom of P-type accepts e⁻s, so it becomes -ve charged ion (A⁻).

A ^h	A ^h	A ⁻	A ⁻	D ⁺	D ⁺	D ⁺	D ⁺
A ^h	A ^h	A ⁻	A ⁻	D ⁺	D ⁺	D ⁺	D ⁺
A ^h	A ^h	A ⁻	A ⁻	D ⁺	D ⁺	D ⁺	D ⁺

A ^h	A ^h	-	-	±	+	D ^e	D ^e
A ^h	A ^h	-	-	+	+	D ^e	D ^e
A ^h	A ^h	-	-	+	+	D ^e	D ^e

Depletion region
(Barrier region)

These +ve & -ve charged ions form a layer on both sides of PN junction which is called depletion region or barrier region. This depletion region prevents the further movements of e⁻s from N-type to P-type S/C.

The +ve & -ve charge ions of depletion region (also called space charge) offers certain potential difference or voltage, which is called barrier potential or depletion potential or junction potential.

The value of barrier potential for Silicon is 0.7V and is 0.3V for Germanium.

PN Junction Diode

When a block of P-type & N-type S/C are joined, a PN junction is formed. If a metallic wires are connected, to each side of P-type & N-type S/C, the resulting structure is called PN junction diode.

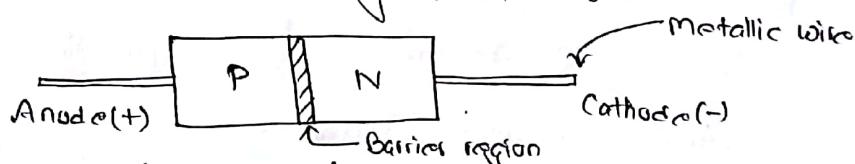


Fig: Structure of PN Junction Diode



Fig: Circuit symbol of PN junction diode.

Diode Biasing (Operation of PN Junction Diode)

The methods of connecting a voltage source to a diode is called diode biasing.

A PN junction diode can be biased in two ways:

- 1) Forward Biasing
- 2) Reverse Biasing

1) Forward Biasing

When a +ve terminal of a source is connected to positive terminal of a diode (P-type) & -ve terminal of a source is connected to a -ve terminal of a diode (N-type S/C), then the diode is said to be forward biased.

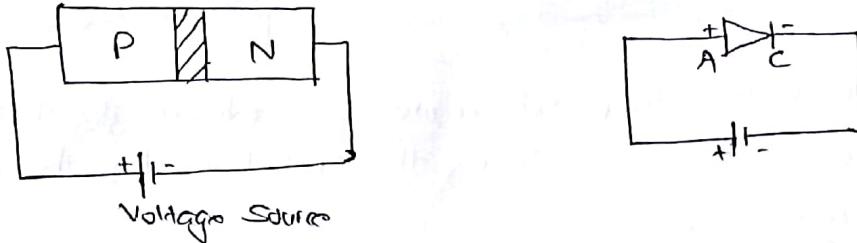


Fig: Forward Biasing

In forward biasing, the +ve terminal of source repels the holes whereas -ve terminal of source repels the electrons due to which there is flow of electron & hole across the junction. Thus, diode conduct current in forward biasing.

In forward bias, diode conducts current only if the supply voltage is greater than the barrier potential of a diode (0.07V for Si & 0.02V for Ge).

The depletion region becomes narrower when a diode is forward biased.

2) Reverse Biasing

When a +ve terminal of a source is connected to -ve terminal of a diode N-type Si & -ve terminal of a source is connected to +ve terminal of a diode (P-type Si), then the diode is said to be Reverse Biased.



Fig: Reverse biasing of PN junction diode

In reverse biasing, the +ve terminal of a source attracts the e⁻s from N-type Si whereas the -ve terminal of a source attracts the holes from P-type Si, away from the junction. Due to the movement of e⁻s & holes the depletion region widens & diode doesn't conduct current.

When a diode is reverse biased, the minority charge carriers gives rise to small current, which is negligible.

When reverse bias voltage is greater than a breakdown voltage, the diode break-down & acts as a short circuit so it conducts high current.

Voltage Current (VI) Characteristics of PN Junction Diode.

If 'V' is the voltage (applied) across a diode & 'I' be the current flowing through it, then the relationship between V & I is given by,

$$I = I_s (e^{\frac{V}{nV_T}} - 1) \quad (1)$$

where, I_s = reverse saturation current

V_T = thermal voltage = 0.026V at room temp.

n = constant ($1 \leq n \leq 2$)

When eqn (1) is plotted in a graph, then VI characteristics of diode is obtained as shown below:

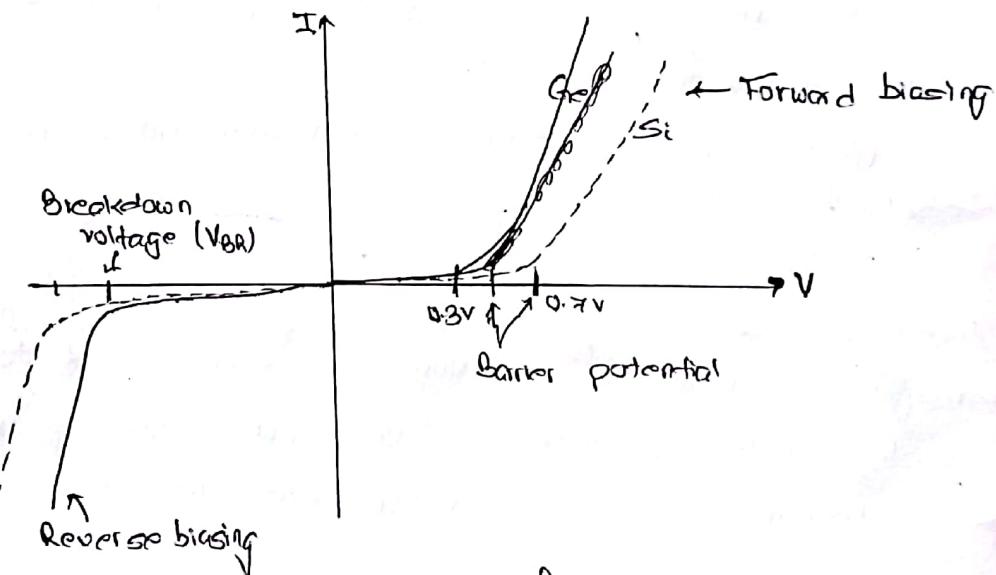


Fig: VI characteristics of a diode

When a diode is forward biased, small amount of current flows through it until the supply voltage is greater than barrier potential.

When the supply voltage exceeds the barrier potential, then the current rises sharply as shown in above figure.

Also, when the diode is reverse biased, negligible amount of current flows through it, but if the supply voltage exceeds the breakdown voltage (V_{BR}), the diode is broken down (destroyed) & it acts as short circuit due to which the current increases sharply as shown in above fig.

Effect of temperature in VI characteristics

The VI characteristics of PN junction diode is given by,

$$I = I_s (e^{\frac{V}{nV_T}} - 1) \quad (1)$$

In above eqn(1), reverse saturation current I_s & thermal voltage V_T are temperature dependant factors.

When temp increases V_T also increases which cause I to decrease. Similarly on increase in temp I_s also increases i.e. I_s doubles for every 10°C rise in temp, due to which I increases. Since the effect of increase of I_s is greater than the effect of increase of V_T , the current through diode increases with the increase in temperature.

The effect of temp in VI characteristics of diode can be shown as below:

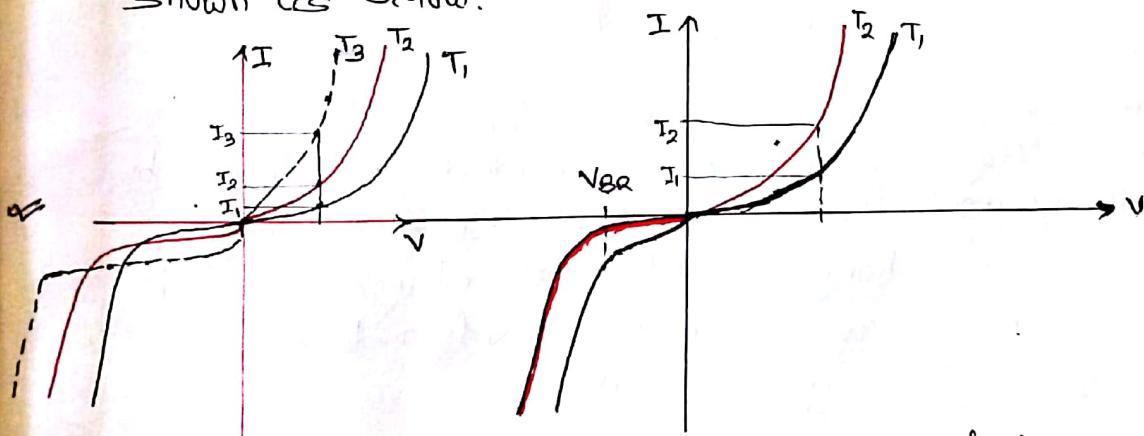


Fig: Effect of temp on VI characteristics of diode

The VI characteristics of diode shift towards left as temp increased. From figure(graph), it is observed that as $T_3 > T_2 > T_1$ thus $I_{T_3} > I_{T_2} > I_{T_1}$. i.e. as current temperature is increased, current increases sharply.

The reverse saturation current of diode is 5nA at 25°C . Find reverse saturation current at 125°C .

Soln: Given, $I_{S1} = 5\text{nA}$ at $T_1 = 25^\circ\text{C}$

$$I_{S2} = ? \quad \text{at } T_2 = 125^\circ\text{C}$$

As we know that reverse saturation current I_s doubles for every 10°C rise in temp.

$$\text{i.e. } I_{S2} = 2^n I_{S1}$$

$$\text{where, } n = \frac{T_2 - T_1}{10} = \frac{125 - 25}{10} = 10$$

$$\text{Now, } I_{S2} = 2^{10} \times I_{S1}$$

$$\therefore I_{S2} = 2^{10} \times 5 \text{nA} = 5120 \text{nA} = (1024 \times 5 \text{nA})$$

Diode Resistance

Diode shows two types of resistive effects:

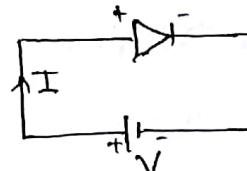
i) Static resistance (DC Resistance)

ii) Dynamic resistance (AC Resistance)

iii) Static Resistance (DC Resistance)

If 'V' be the dc voltage applied to a diode and 'I' be the current flowing through it, then the ratio of V and I is called static resistance of a diode.

$$\text{i.e. } R_{DC} = \frac{V}{I}$$

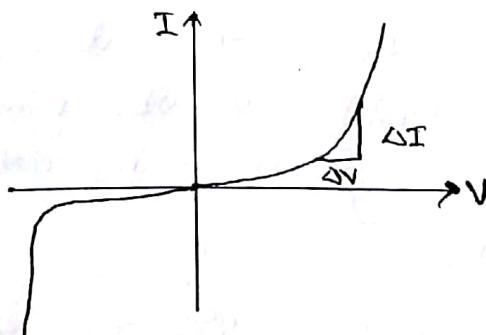


iv) Dynamic Resistance (AC Resistance)

When the voltage applied across a diode is changed by ΔV , then the current flowing through it will also change by ΔI . The ratio of change in voltage to change in current is called dynamic or AC resistance of a diode.

$$r_{AC} = \frac{\Delta V}{\Delta I} = \frac{dV}{dI} \quad (1)$$

The VI-characteristics of diode is,



$$\text{Here, slope} = \frac{\Delta I}{\Delta V} = \frac{dI}{dV}$$

$$\Rightarrow \frac{dV}{dI} = \frac{1}{\text{slope}}$$

$$\text{Thus, } r_{AC} = \frac{dV}{dI} = \frac{1}{\text{slope}} \quad (2)$$

The AC resistance of diode can also be defined as the reciprocal of slope of its VI characteristics.

$$\text{We know that, } I = I_s (e^{\frac{V}{nV_T}} - 1) = I_s e^{\frac{V}{nV_T}} \quad \text{--- (3)}$$

on Differentiating w.r.t. V,

$$\frac{dI}{dV} = I_s * \frac{1}{nV_T} * e^{\frac{V}{nV_T}}$$

$$\boxed{\frac{de^{3m}}{dm} = 3 \cdot 3^m}$$

$$\frac{dI}{dV} = \frac{I_s}{nV_T} e^{\frac{V}{nV_T}} = \frac{I_s e^{\frac{V}{nV_T}}}{nV_T} = \frac{I}{nV_T}$$

$$\text{or, } \frac{dV}{dI} = \frac{nV_T}{I}$$

$$\therefore r_{AC} = \frac{dV}{dI} = \frac{nV_T}{I}$$

Diode Capacitance / Junction Capacitance

Diode shows 2 types of capacitive effects:

The depletion region across the junction behaves as an insulating material or dielectric constant of the capacitor and the P & N sides act as parallel plates.

Thus, the junction exhibits the property of capacitance & is called junction capacitance or diode capacitance.

Diode shows 2 types of capacitive effects:

1) Transition capacitance (Space charge capacitance)

2) Diffusion capacitance

1) Transition Capacitance (C_T)

When a diode is reverse biased, the depletion region increases. The depletion region consists of a layer of positive & negative charged ions. These ions behave as the +ve & -ve plate of a parallel plate capacitor due to which capacitive effect is produced in a diode, known as transition capacitance.

Thus transition capacitance is the capacitive effect shown by diode when it is reverse biased.

The transition capacitance of a diode is given by,

$$C_T = \frac{\epsilon A}{w} \quad \text{--- (4)}$$

where, ϵ = Permittivity

A = Area of cross section (of a junction)

w = width of depletion region.

Also, Depletion width(w) is expressed by,

$$w^2 = \frac{2\epsilon V_B}{e} \left[\frac{1}{N_A} + \frac{1}{N_D} \right] \quad \text{--- (2)}$$

where, N_A = No. of Acceptor atoms

N_D = No. of Donor atoms

ϵ = Permittivity

e = charge of e^-

V_B = [junction voltage = $V_0 - V_R$] = applied reverse voltage.
(V_0 = contact potential and V_R is negative voltage applied to diode)

From eqn(1) & (2)

$$C_T = \frac{\epsilon A}{\sqrt{\frac{2\epsilon V_B}{e} \left[\frac{1}{N_A} + \frac{1}{N_D} \right]}} = \frac{\epsilon A}{\sqrt{\frac{2\epsilon V_B (N_A + N_D)}{e^2 N_A N_D}}}$$

$$\text{or, } C_T = \frac{\epsilon A * \sqrt{e N_A N_D}}{2\epsilon V_B (N_A + N_D)}$$

$$\Rightarrow C_T \propto \frac{1}{\sqrt{V_B}} \quad \text{--- (3)}$$

Thus, transition capacitance is inversely proportional to the square root of applied reverse voltage, V_B .

2) Diffusion Capacitance (C_D)

When a diode is forward biased, there is a movement of electron & holes through the PN junction. This movement of charged carriers give rise to a capacitive effect in a diode, which is called diffusion capacitance.

In forward biased diode, there is effect of transition capacitance, but the effect is suppressed by the capacitive effect due to movement of charge carriers. [Thus forward biased diode possess diffusion capacitance only.]

The Diffusion capacitance of diode is due to the movement of charge carriers across the junction.

$$\text{i.e. } C_D = \frac{dQ}{dV} \quad \text{--- (1)}$$

Also, $Q = Id = IT$, $\tau = T$, mean life time of carrier, constant

So, from eqn (1),

$$C_D = \frac{d(IT)}{dV} = T \frac{dI}{dV} \quad \text{--- (2)}$$

The characteristic eqn of Diode is,

$$I = I_s (e^{\frac{V}{\eta V_T}} - 1) = I_s e^{\frac{V}{\eta V_T}} \quad (\text{neglecting } I_s)$$

On differentiating w.r.t. V ,

$$\frac{dI}{dV} = \frac{I_s e^{\frac{V}{\eta V_T}}}{\eta V_T} = \frac{I}{\eta V_T}$$

Then from eqn (2),

$$C_D = T \times \frac{I}{\eta V_T} = \frac{T I}{\eta V_T}$$

$$\therefore \text{Diffusion Capacitance, } C_D = \frac{I}{\eta V_T} T$$

where, T = mean life time of carrier

I = forward current

η = constant $= 1 \leq \eta \leq 2$

V_T = thermal voltage.

Since, T , η and V_T are constant. Thus Diffusion capacitance is directly proportional to forward current I .

#Differences between Transition & Diffusion Capacitance

Transition Capacitance

1. The capacitive effect shown by diode when it is reverse biased is called transition capacitance.
2. When diode is reverse biased, the depletion region widens. The depletion region consists of +ve & -ve charge ions, behaves as parallel plate capacitor, Due to which capacitive effect arises in diode & is a transition capacitance.

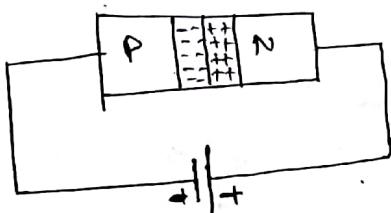


Fig: Reverse biased diode

3. The transition capacitance of diode can be expressed as;

$$C_T = \frac{\epsilon A}{W}$$

where, ϵ = permittivity

A = Area of cross-section
of depletion region

w = width of depletion
region.

4. The transition capacitance is inversely proportional to the reverse voltage.

Diffusion Capacitance

2. The capacitive effect shown by diode when it is forward biased is called diffusion capacitance.
2. When a diode is forward biased, the two & no charge carriers move through the PN junction. This movement of charge carrier gives an capacitive effect, known as diffusion capacitance.

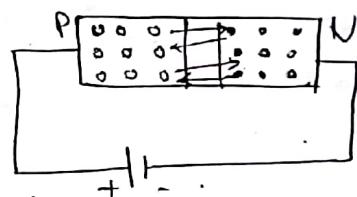


Fig: Forward biased diode

3. The diffusion capacitance of a diode is expressed as;

$$C_D = \frac{I}{\eta V_T} T$$

where T = mean lifetime of carriers

I = forward current of diode

η = constant ($1 \leq \eta \leq 2$)

V_T = Thermal voltage.

4. Diffusion capacitance is directly proportional to the forward current of diode.

Diode Switching Time (Reverse Recovery Time of Diode)

When a diode is forward biased it conducts current (switch ON) & when it is reverse biased it doesn't conduct current (switch OFF).

When a forward biased diode is suddenly reverse biased, the diode doesn't change its state from conduction (switch ON) to non-conduction (switch OFF).

The diode takes certain time to change its state from conduction (forward biased) to non-conduction

(reverse biased), which is called reverse recovery time (diode switching time) of a diode.

When the forward biased diode is reverse biased, the current simply reverses & stay measurable for certain period of time, called storage time (t_s).

After the storage phase has passed, the current gradually decreases & becomes zero to take diode into non-conduction state (reverse bias). This period of time is called transition time (t_t).

The reverse recovery time (t_{rr}) is the sum of storage time (t_s) & transition time (t_t).

$$\text{i.e. } t_{rr} = t_s + t_t$$

Junction Breakdown / Diode Breakdown

In reverse biasing, a very small current flows due to minority charge carriers, which is known as leakage saturation current (I_s).

If reverse voltage is further increased, a current rises sharply at a certain reverse-bias voltage. This mechanism of producing a large reverse current is called breakdown.

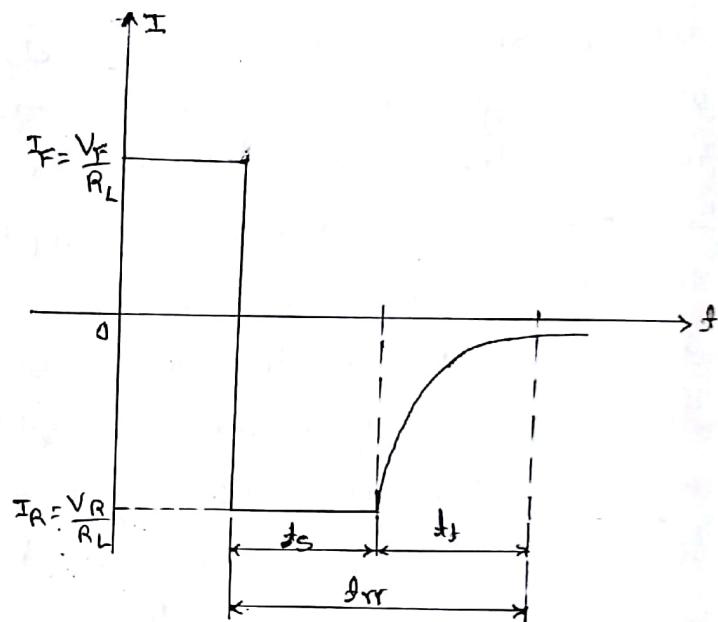


Fig: Reverse recovery time of diode

There are two types of breakdown mechanisms:

- 1) Zener breakdown
- 2) Avalanche breakdown

1) Avalanche Breakdown

When a diode is reverse biased, the free e^- s present on it can gain (high) sufficient energy from supply voltage. These highly energetic e^- s move with greater speed & can break-down (rupture) a covalent bond to generate free e^- (i.e. a e^- can rupture a covalent bond & generate two free e^- s). These newly generated free e^- s also gain energy & further generate free e^- s. by rupturing covalent bonds. The process continues & in a short period of time, large no. of free e^- s are generated in a diode. The flows of these e^- s give rise to large current. This phenomenon is called avalanche breakdown.

The avalanche breakdown occurs in a diode with higher breakdown voltage.

2) Zener Breakdown

When a diode is reverse biased, the supplied reverse voltage can generate an electric field. Due to which covalent bonds are ruptured & free e^- s are generated. If the diode has lower breakdown voltage, then the supply voltage can rupture (break-down) large no. of covalent bonds & generate large no. of free e^- s, due to which large no. of current flows through a diode. This phenomenon is called zener breakdown.

Zener breakdown occurs in zener diode with smaller breakdown voltage (zener voltage V_z).

→ The main differences between zener breakdown and avalanche breakdown are as belows:

Zener Breakdown

Avalanche Breakdown

1.

Zener Diode

Zener diode is a special purpose diode designed to be used as a voltage regulator. They are highly doped diode, thus have very small barrier region (less than 1mV). It can conduct current in reverse biasing. If it is forward biased it shows the property of forward biased PN Junction diode. It has sharp breakdown voltage, known as zener voltage (V_z).



Fig: Circuit symbol of zener diode.

The VI characteristics of zener diode is;

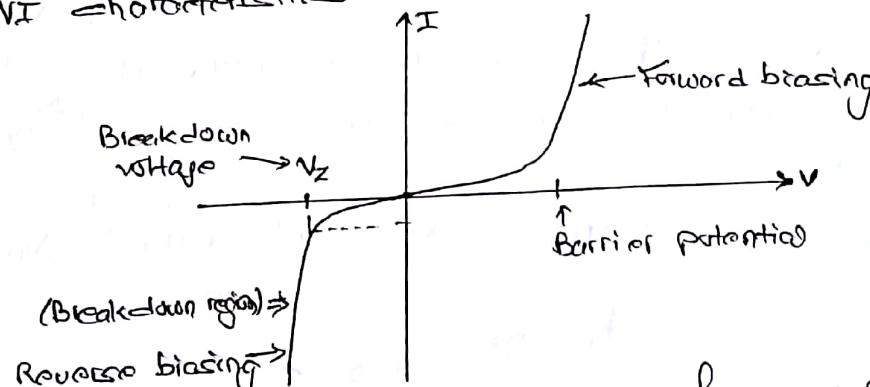


Fig: VI characteristics of zener Diode

Zener Diode as a Voltage Regulator

To use Zener diode as a voltage regulator, following condition must be satisfied,

- 1) Zener diode must be operated in reverse-bias & connected in parallel with load.
- 2) Input voltage must be greater than zener voltage ($V_{in} > V_z$)
i.e. it must be operated in a breakdown region.

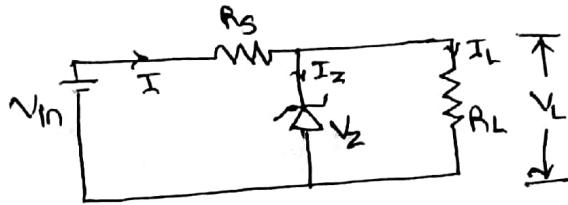


Fig: Zener diode as a voltage regulator

Here, V_{in} = input voltage

R_s = Source resistance

R_L = Load resistance

I = Source current

I_L = Load current

I_z = Zener current

V_z = Zener voltage

V_L = Load voltage

When zener diode regulates voltage across load, the load voltage will be equal to zener voltage.

$$\text{i.e. } V_L = V_z.$$

From above figure, $I = I_z + I_L$

$$\text{and, } V_L = I_L \cdot R_L$$

Case-I: When V_{in} changes

Let V_{in} is increased. As V_{in} increases, the current I also increases. If I increased, both I_L & I_z must be increased, but the zener diode conducts the increased current (i.e. only I_z increases) and maintain I_L constant.

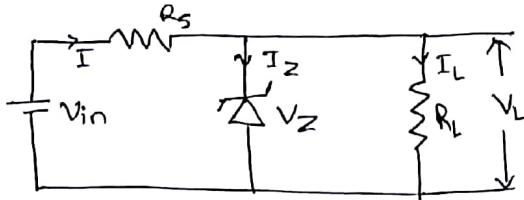
Thus, $V_L = I_L \cdot R_L$ becomes constant even if V_{in} increases.

Case-II: When R_L changes

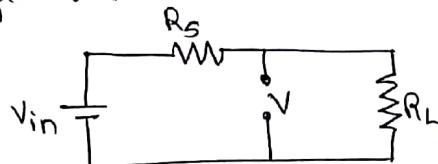
Suppose R_L increases. When R_L increases, the zener diode conducts more (i.e. I_z increases) & minimizes the value of I_L . So that even increase of R_L , as I_L decreases, $V_L = I_L \cdot R_L$ remains constant.

Zener Diode Numerical: Different cases

1) Case I: V_{in} = constant & R_L = constant



Here, we have to find whether the zener diode is 'ON' or 'OFF'. For that, remove zener diode from circuit & find open ckt voltage 'V'.



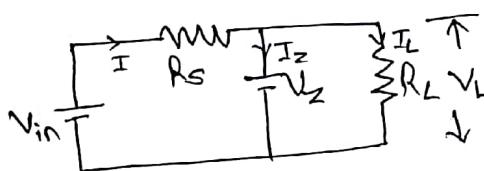
$$\text{Here, } V = \frac{R_L}{R_s + R_L} * V_{in}$$

a) If $V < V_z$, zener diode is OFF

b) If $V > V_z$, zener diode is ON

c) If zener diode is ON

Replace zener diode by voltage source with value V_z .



$$\text{Here, } V_L = V_z$$

$$I = \frac{V_{in} - V_z}{R_s}$$

$$I_L = \frac{V_L}{R_L}$$

$$I_z = I - I_L$$

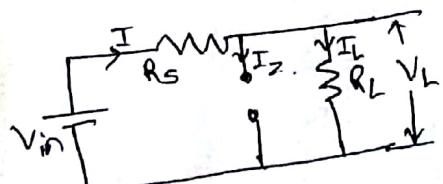
$$V_R = I \cdot R_s$$

$$P_z = V_z \cdot I_z$$

b) If zener diode is OFF

Replace zener diode by open circuit

$$\text{Here, } I_z = 0$$



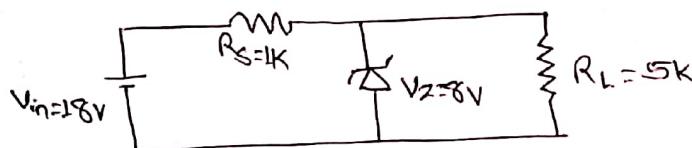
$$V_L = \frac{R_L}{R_s + R_L} * V_{in}$$

$$I = I_L = \frac{V_{in}}{R_s + R_L} \quad \text{or} \quad I_L = \frac{V_L}{R_L}$$

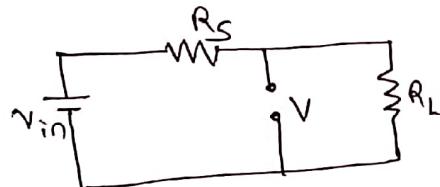
$$V_R = I \cdot R_s \quad \& \quad P_z = V_z \cdot I_z = 0$$

For the circuit shown below, find

- Load voltage (V_L)
- Zener current (I_z), (iii) Power across zener diode.



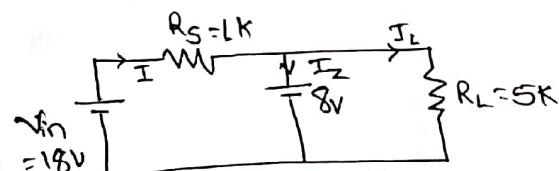
Sol: Removing zener diode & have to find open ckt voltage



Here, $V = \frac{R_L}{R_L + R_S} * V_{in} = \frac{5}{5+1} * 18 = 6V = 15V$

Since $V > V_Z$. So zener diode is ON.

Now, replacing zener diode by its (breakdown) voltage source with value V_Z .



From above circuit,

i) Load voltage (V_L) = $V_Z = 8V$

ii) Current flowing through zener diode (I_z)=?

We know that, $I_z = I - I_L$

where, $I = \frac{18-8}{R_S} = \frac{18-8}{1} = 10mA$

$$I_L = \frac{V_L}{R_L} = \frac{8}{5} = 1.6mA$$

$$\therefore I_z = I - I_L = 10 - 1.6 = 8.4mA$$

iii) Power across zener diode,

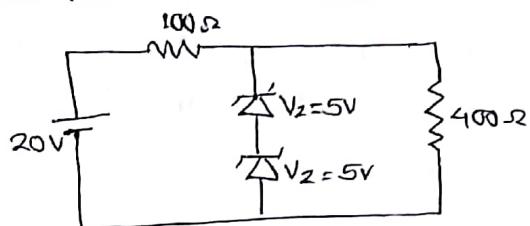
$$P_z = I_z * V_z$$

$$= 8.4mA * 8V$$

$$= 67.2mW$$

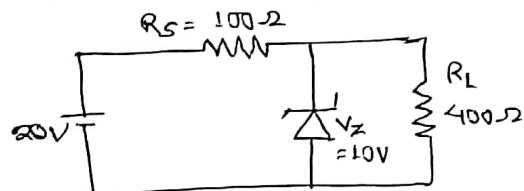
For the circuit shown below, find

- Load voltage (V_L).
- Current through Zener Diode (I_Z).
- Voltage across series resistance.
- Power dissipated across zener diode.

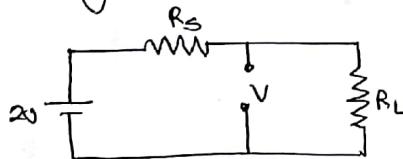


Sol:

The equivalent circuit is,



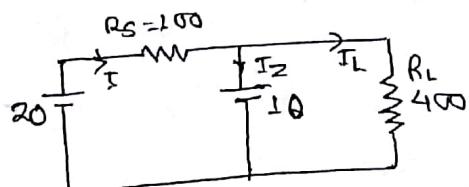
Calculating open circuit voltage



$$V = \frac{R_L}{R_S + R_L} \times 20 = \frac{400}{400 + 100} \times 20 = 16V$$

Here, $V > V_Z$. So Zener diode is ON.

Now,



Now,

i) Load voltage, $V_L = V_Z = 5 + 5 = 10V$

$$\text{ii) } I_Z = I - I_L = \frac{20 - 10}{100} - \frac{10}{40} = \frac{V_{in} - V_Z}{R_S} - \frac{V_L}{R_L}$$

$$= 0.1 - 0.025$$

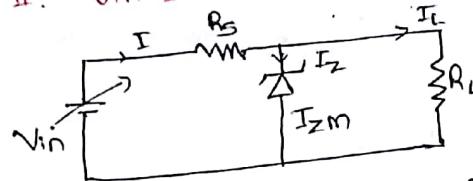
$$= 0.075A.$$

iii) Voltage across series resistance, $V_R = I R_S = 0.1 \times 100 = 10V$

iv) Power dissipated in zener diode,

$$P_Z = I_Z V_Z = 0.075 \times 10 = 0.75 \text{ Watt}$$

CASE II: V_{in} = Variable, R_L = constant



I_{zm} = maximum zener current.

Here, we have to find range of V_{in} [i.e. $V_{in(\min)}$ and $V_{in(\max)}$] required for safe turn 'ON' of zener diode. Zener must be ON. Thus,

$$V_L = V_z$$

For minimum value of V_{in} [$V_{in(\min)}$]

The open circuit voltage is given by,

$$V = \frac{R_L}{R_S + R_L} * V_{in}$$

For minimum condition, $V = V_z$

$$\therefore V_z = \frac{R_L}{R_S + R_L} * V_{in}$$

$$\text{or, } V_z (R_S + R_L) = R_L \cdot V_{in}$$

$$\Rightarrow V_{in(\min)} = \frac{V_z (R_S + R_L)}{R_L}$$

$$\text{Also, } I_{min} = \frac{V_{in(\min)} - V_z}{R_S}$$

For maximum value of V_{in} [$V_{in(\max)}$]

We know that, $I = \frac{V_{in} - V_z}{R_S}$

$$\text{or, } IR_S = V_{in} - V_z$$

$$\Rightarrow V_{in} = IR_S + V_z \quad \text{--- (1)}$$

Here, $V_{in} \rightarrow V_{in(\max)}$, when $I \rightarrow I_{max}$

$$V_{in(\max)} = I_{max} R_S - V_z$$

Again,

$$I = I_L + I_z$$

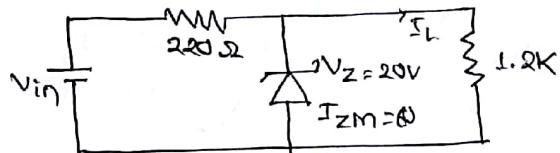
Here, $I \rightarrow I_{max}$ when $I_z \rightarrow I_{zm}$

$$\text{where, } I_L = V_z / R_L \Rightarrow I_{max} = I_L + I_{zm}$$

Now, from eqn(1),

$$\therefore V_{in(\max)} = (I_L + I_{zm}) R_S + V_z$$

Determine the range of values of V_{in} that maintain zener diode in ON-state. Also find the max power that can be dissipated in diode.



Sol:

$$V_{in(\max)} = (I_L + I_{Zm}) R_S + V_Z$$

$$\text{Hence, } V_Z = V_L = I_L R_L$$

$$\Rightarrow I_L = \frac{V_Z}{R_L} = \frac{20}{1.2 \times 10^3} = 16.67 \text{ mA}$$

$$\begin{aligned} \therefore V_{in(\max)} &= (16.67 + 60) \text{ mA} \times 220\Omega + 20 \\ &= 16.867 + 20 \\ &= 36.867 \text{ V} \end{aligned}$$

$$\text{Also, } V_{in(\min)} = \frac{R_S + R_L}{R_L} * V_Z = \frac{(1.2 \times 10^3 + 220)}{1.2 \times 10^3} \times 20$$

$$\therefore V_{in(\min)} = 23.67 \text{ V.}$$

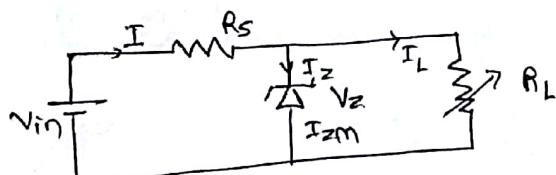
Thus, the range of values of V_{in} that maintain zener diode in ON-state = 23.67V to 36.867V

Again,

maximum power dissipated in zener diode is,

$$P_{max} = I_{Zm} * V_Z = 60 \times 20 = 1200 \text{ mW} = 1.2 \text{ Watt}$$

CASE III: V_{in} = constant, R_L = variable



Here, I_{Zm} = Maximum zener current

Here, we have to find range of R_L [$R_{L\min}$ & $R_{L\max}$] required for safe turn on of zener diode.

Since, Zener diode must be ON,

$$V_L = V_Z$$

For minimum value of R_L [$R_{L,\min}$]
 — The open circuit voltage is given by,

$$V = \frac{R_L}{R_L + R_S} \times V_{in}$$

For minimum condition, $V = V_Z$

$$\text{i.e. } V_Z \times (R_L + R_S) = R_L \times V_{in}$$

$$\therefore R_{L,\min} = \frac{V_Z \cdot R_S}{V_{in} - V_Z}$$

$$\& I_{L,\max} = \frac{V_L}{R_{L,\min}}$$

For maximum value of R_L [$R_{L,\max}$]

$$\text{As we know, } R_L = \frac{V_L}{I_L}$$

$$R_{L,\max} = \frac{V_L}{I_{L,\min}} \quad \text{--- (1)}$$

$$\text{Again, } I = I_L + I_Z$$

$$\Rightarrow I_L = I - I_Z$$

Here, $I_L \rightarrow I_{L,\min}$, when $I_Z \rightarrow I_{Zm}$

$$\therefore I_{L,\min} = I - I_{Zm} \quad \text{--- (2)}$$

From eqn (1), (2)

$$\therefore R_{L(\max)} = \frac{V_L}{I - I_{Zm}}$$

$$\text{where, } I = \frac{V_{in} - V_Z}{R_S}$$

A zener regulator has 12V zener voltage with variable load resistance. Calculate,

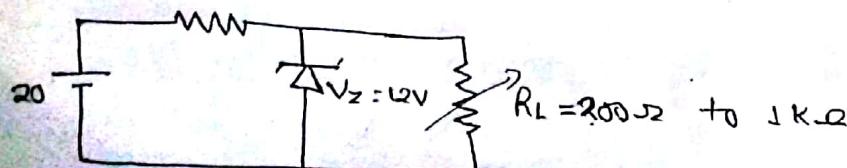
i) Maximum & minm load current.

ii) Maxm & minm power dissipation across load.

iii) Voltage drop across series resistance.

iv) minimum value of load resistance to ensure that zener diode $I = 0V$.

$$R_S = 100\Omega$$



Soln: Given, $V_{in} = 20V$

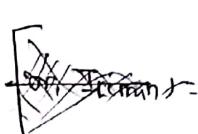
$$R_s = 100\Omega$$

$$V_z = 12V$$

$$R_{L(min)} = 200\Omega$$

$$R_{L(max)} = 1K = 1000\Omega$$

i) Maximum load current, $I_{L(max)} = \frac{V_L}{R_{L,min}} = \frac{V_z}{R_{L(min)}} = \frac{12}{200} = 60mA$

& $I_{L(min)} = \frac{V_L}{R_{L(max)}} = \frac{12}{1000} = 12mA$ 

ii) $P_{z(min)} = V_z \cdot I_{z(min)}$

$$P_{z(max)} = V_z \cdot I_{z(max)}$$

Here, $I = \frac{V_{in} - V_z}{R_s} = \frac{20 - 12}{100} = 80mA$

$$\therefore I_{z(min)} = I - I_{L(max)} = 80 - 60 = 20mA$$

$$\& I_{z(max)} = I - I_{L(min)} = 80 - 12 = 68mA$$

$$\therefore P_{z(min)} = 12 * 20 = 240mW$$

$$P_{z(max)} = 12 * 68 = 816mW$$

iii) Voltage drop across series resistance,

$$V_R = I R_s = 80mA * 100\Omega = 8V$$

$$(or, V_R = V_{in} - V_z = 20 - 12 = 8V)$$

iv) $R_{L(min)} = \frac{V_z \cdot R_s}{V_{in} - V_z} = \frac{12 * 0.1K}{20 - 12} = 0.15K\Omega$

For the given circuit, load current varies from 12 to 100mA. Find the value of series resistance R_s to maintain 7.2V across the load. Minimum zener diode current is 10mA.

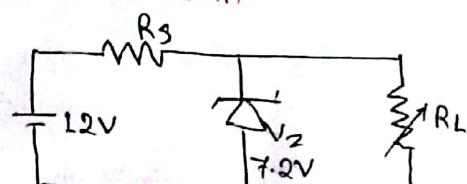
Soln: Given, $V_{in} = 12V$

$$V_z = 7.2V$$

$$I_{L(min)} = 12mA$$

$$I_{L(max)} = 100mA$$

$$I_{z(min)} = 10mA$$

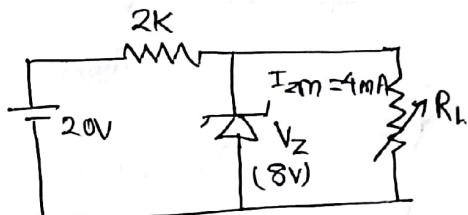


$$\text{Here, } I = I_{Z(\min)} + I_{L(\max)} = 10 + 100 = 110 \text{ mA}$$

$$\therefore \text{Series resistance, } R_s = \frac{V_{in} - V_Z}{I} = \frac{(12 - 7.2)V}{110 \times 10^{-3} A} = 43.68 \Omega.$$

If for the circuit given, find the maxm & minm value of load resistor R_L required for safe turn ON of zener diode. Also find the maximum & minimum power that can be dissipated across load.

Soln: Given, $V_{in} = 20V$
 $R_s = 2K$
 $V_Z = 8V$
 $I_{Zm} = 4\text{mA}$



$$\text{Now, } R_{L(\min)} = \frac{V_Z R_s}{V_{in} - V_Z} = \frac{8V \times 2 \times 10^3 \Omega}{20 - 8} = 1.833 K\Omega$$

$$\text{and, } R_{L(\max)} = \frac{V_L}{I - I_{Zm}} = \frac{V_Z}{I - I_{Zm}}$$

$$\text{where, } I = \frac{(20 - 8)V}{2 \times 10^3 \Omega} = 6 \text{ mA}$$

$$\therefore R_{L(\max)} = \frac{8V}{(6 - 4) \times 10^{-3} A} = 4 K\Omega$$

Also,

$$I_{L(\min)} = \frac{V_L}{R_{L(\max)}} = \frac{8V}{4 \times 10^3 \Omega} = 2 \text{ mA}$$

$$I_{L(\max)} = \frac{V}{R_{L(\min)}} = \frac{8V}{1.833 K} = 6 \text{ mA}$$

Thus, Power dissipated on load is,

$$P_L(\min) = V_L \cdot I_{L(\min)} = 8V \times 2 \times 10^{-3} A = 16 \text{ mW}$$

$$\& P_L(\max) = V_L \cdot I_{L(\max)} = 8V \times 6 \times 10^{-3} A = 48 \text{ mW.}$$

Schottky Diode

It is a metal semiconductor diode designed to be used for fast switching applications.

It consists of metal & sic region. The metals used are high conductive metals like silver, gold, platinum etc whereas sic used are normally N-type sic (because majority carrier is e^- & for conduction of maximum current).

It is a unipolar diode i.e. only e^- s conduct current. It is also called hot carrier diode, because the e^- must gain high energy to cross metal-sic junction to conduct current.

The metal-sic junction is called schottky junction.

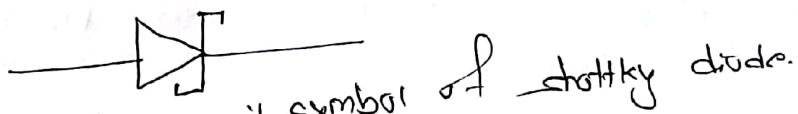
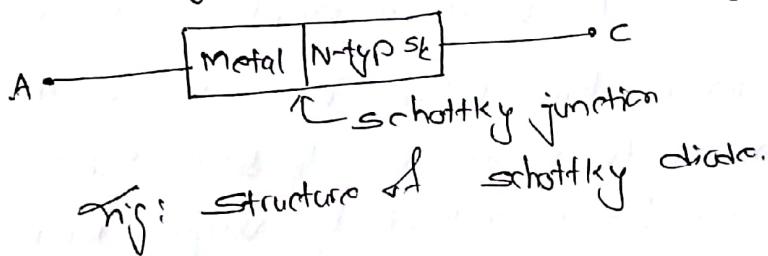
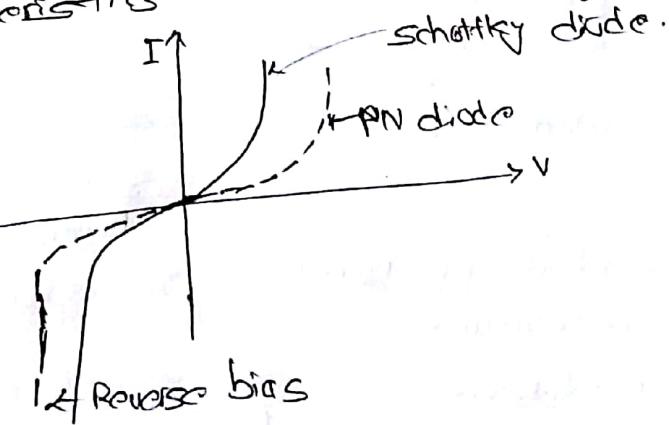


Fig: Circuit symbol of schottky diode.

The metal region of schottky diode is called anode or anode, whereas the sic region is called -ve terminal or cathode.

The VI characteristics of PN & schottky diode is,



Varactor Diode

The varactor diode is a variable capacitor diode. It is a reverse-biased pn junction diode whose transition capacitance varies with the applied reverse-bias voltage.

The varactor is also called varicap, voltcap, epicap, voltage variable capacitor, or tuning capacitor.

A reverse-biased PN junction diode possesses transition capacitance C_T which is given by

$$C_T = \frac{\epsilon A}{W}$$

where, ϵ = permittivity

A = Area of cross-section

W = width of depletion region

When the reverse-bias voltage 'V' is varied, the width of depletion region 'W' varies due to which C_T also varies. Thus, C_T can be controlled by varying reverse-voltage. This property of reverse-biased PN junction diode is called varactor.



Fig: Symbol of varactor

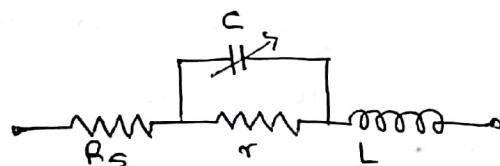


Fig: Equivalent circuit

Varactor diodes are used in:

- Tuning circuits
- FM modulators
- Adjustable bandpass filters
- Automatic frequency control devices.



Fig: Characteristic Curve

Tunnel Diode

A tunnel diode is a high conductivity two-terminal PN junction doped heavily about 1000 times greater than conventional junction diode. Because of heavy doping depletion width is reduced to an extremely small value at the order of 10^{-5} mm.

When a small +ve voltage is applied across tunnel diode, it can conduct significant amount of current due to a phenomenon called tunnelling phenomena.

Tunnel diode passes resistance (i.e. current decreases when voltage is increased).

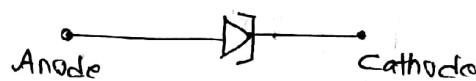


Fig: Symbol of Tunnel diode.

VI Characteristics of Tunnel Diode

When a small +ve voltage is applied across a tunnel diode, it starts to conduct current.

When the voltage is gradually increased the current also increases. When the voltage across tunnel diode

$V = V_p$ (peak voltage),

maximum current flows through it, called peak current I_p .

When the voltage across diode is increased beyond V_p (i.e. $V > V_p$), the current starts to decrease. The current becomes minimum called Valley current, I_v , when the voltage across diode is equal to valley voltage ($V = V_v$).

When the voltage is increased beyond V_v ($V > V_v$), the tunnel diode gains the property of normal PN junction diode & the current increases gradually like in PN junction diode.

Between peak point X & valley point Y, in VI characteristic the current decreases when voltage is increased. Thus a tunnel diode passes negative resistance in this region.

Tunnel diode passes negative resistance in this region.

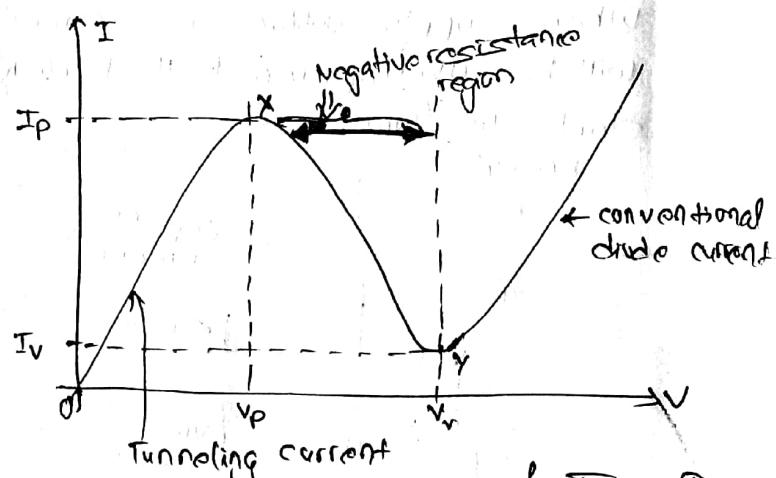


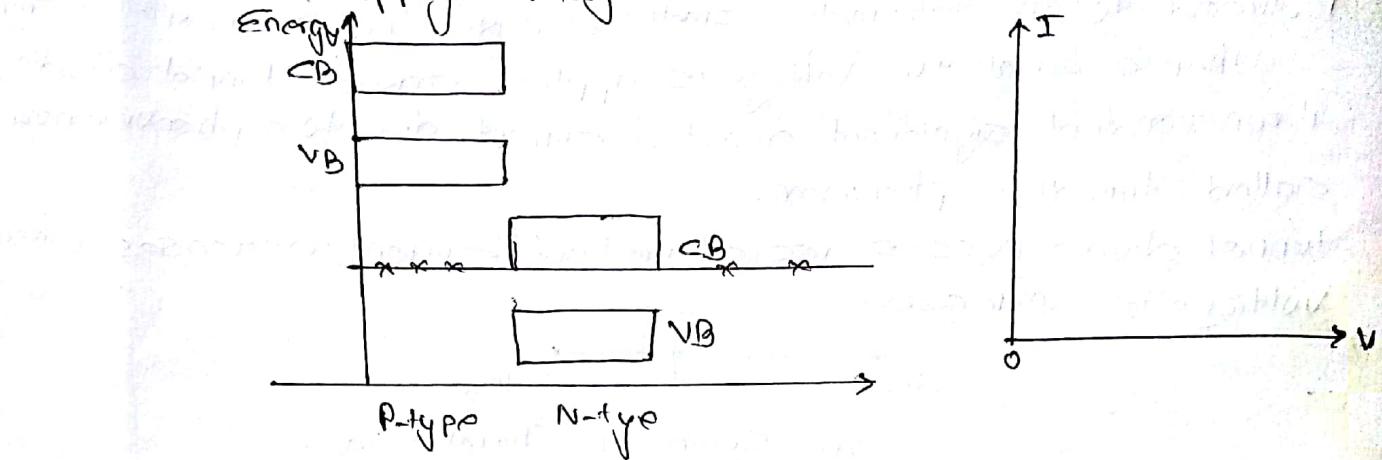
Fig: VI characteristics of Tunnel Diode.

Tunneling Phenomenon of Tunnel Diode

The tunneling phenomenon of tunnel diode can be explained by following steps:

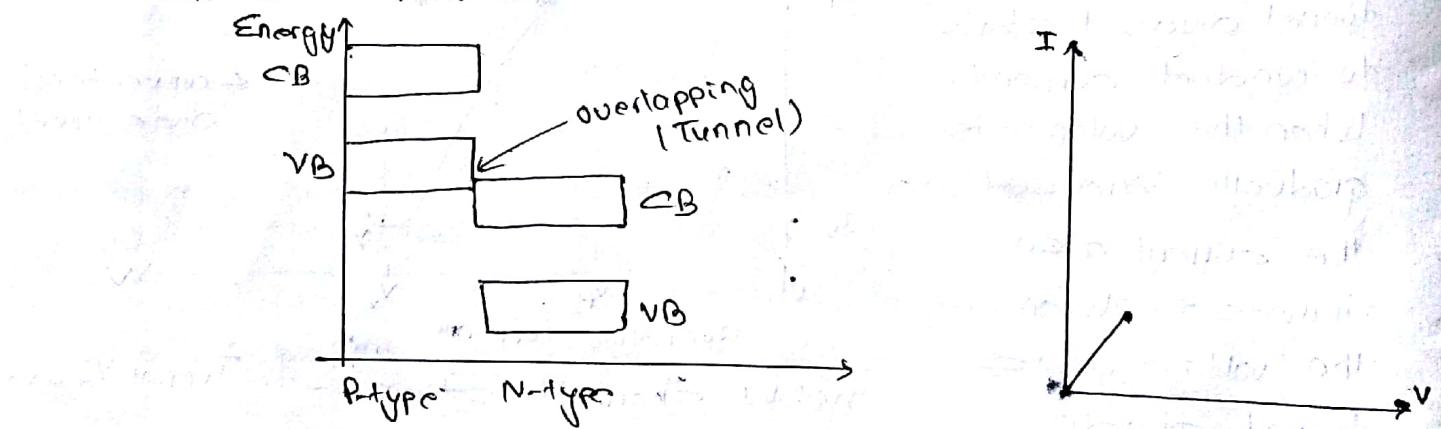
Step I: When $V=0$

When supply voltage is zero, there is flow of current.



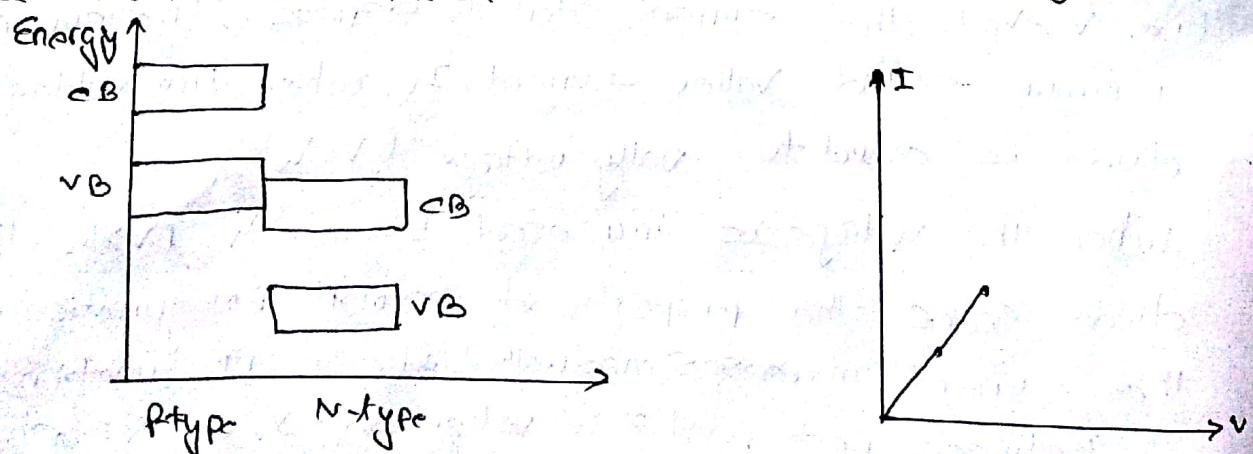
Step II: When $V = +ve$

When supply voltage is increased i.e. $V = +ve$, there will be overlapping both P- & N-type and current starts to flow.



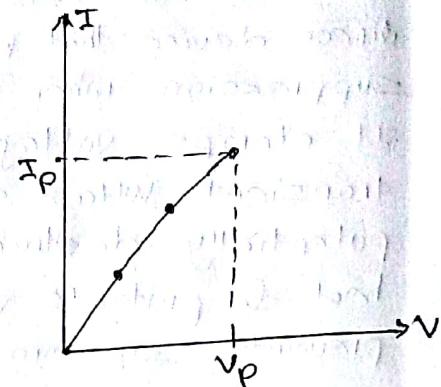
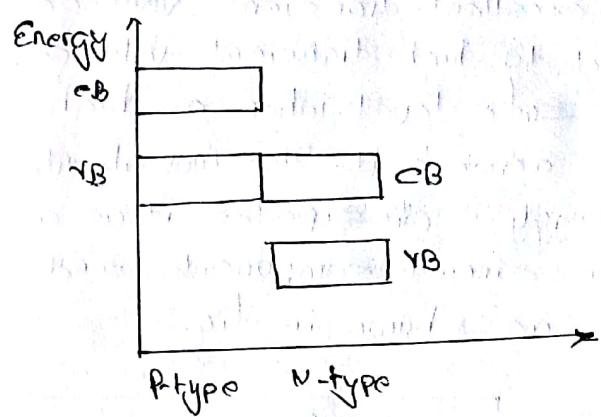
Step III: When V is increased

When V is further increases then overlapping increases & current also increases.



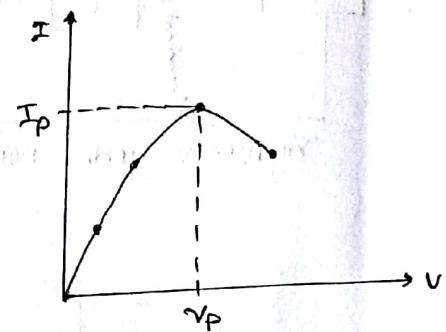
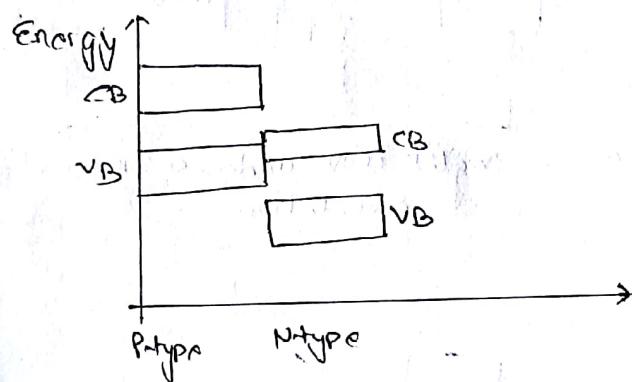
Step IV: When $V = V_p$

When $V = V_p$, at this condition current is maximum & overlapping also maximum.



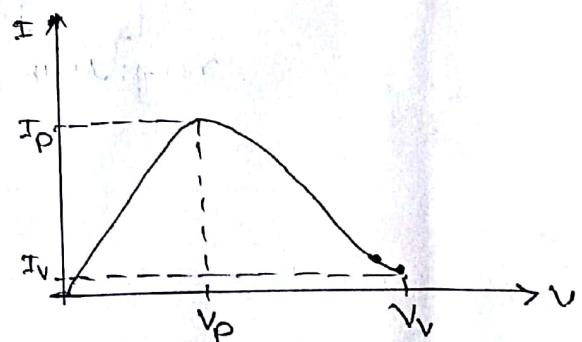
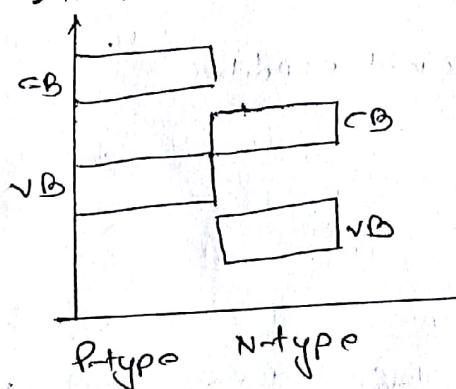
Step V: When $V > V_p$

when $V > V_p$, current starts to decreases

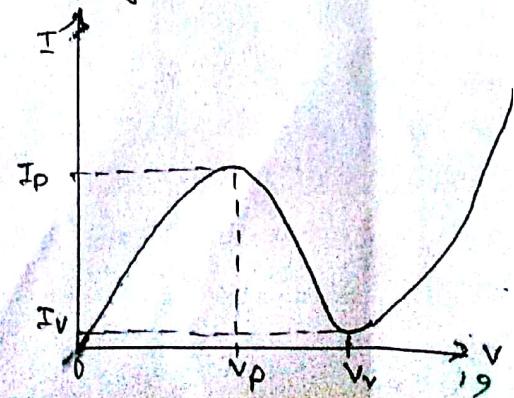
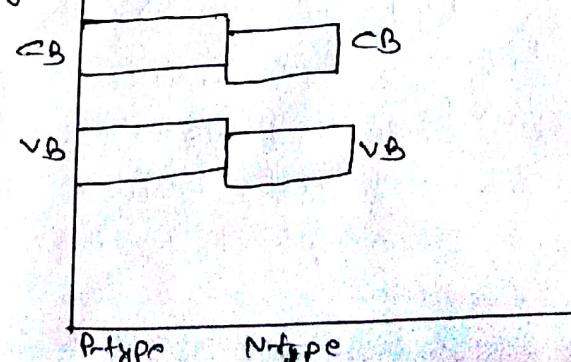


Step VI: When $V = V_v$

At this condition, current becomes minimum.

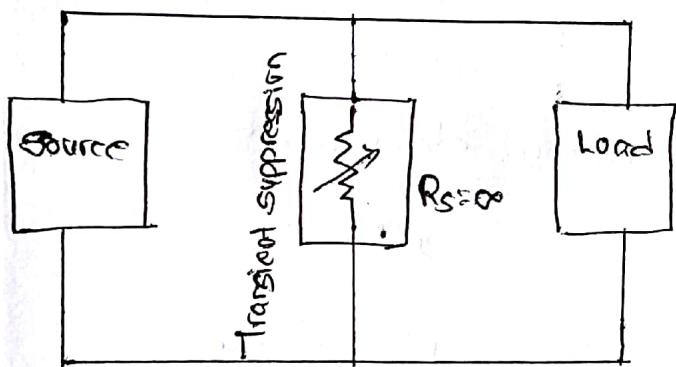


Step VII: When $V > V_v$
when further increasing voltage, it gain the property of PN junction diode & current increases gradually.

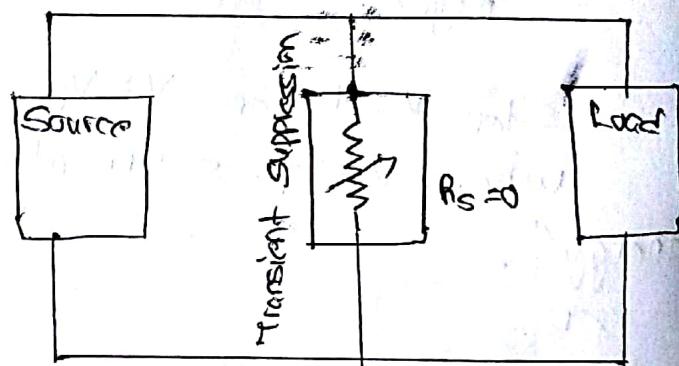


Metal Oxide Varistor

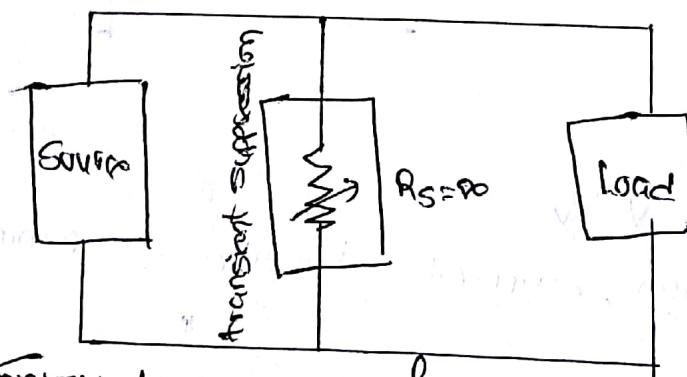
The metal oxide varistor, MOV, is a voltage dependent, non-linear device that provides excellent transient voltage suppression when exposed to high transient voltage. It clamps voltage to a safe level when a high transient voltage appears across it. The MOV absorbs potentially destructive energy & dissipates it as a heat & protects vulnerable circuit components and prevents system damage as shown in fig.



Fig(1): MOV under normal condition



Fig(2): MOV under a transient condition



Fig(3): At the end of transient condition