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<http://lukaseigentler.github.io>

# Pattern migration (or not?) of dryland vegetation stripes

MODIS, ICMS

15 September 2023

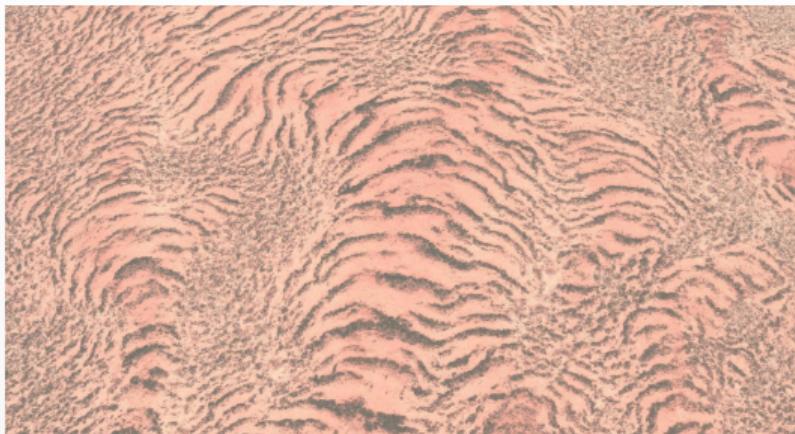
*Lukas Eigentler (Universität Bielefeld)*

*joint work with Jonathan A Sherratt (Heriot-Watt Univ.)*

# Vegetation patterns

Vegetation patterns are a classic example of a **self-organisation principle** in ecology.

Stripe pattern in Ethiopia<sup>1</sup>.



Gap pattern in Niger<sup>2</sup>.



- Plants increase water infiltration into the soil and thus induce a **positive feedback loop**.

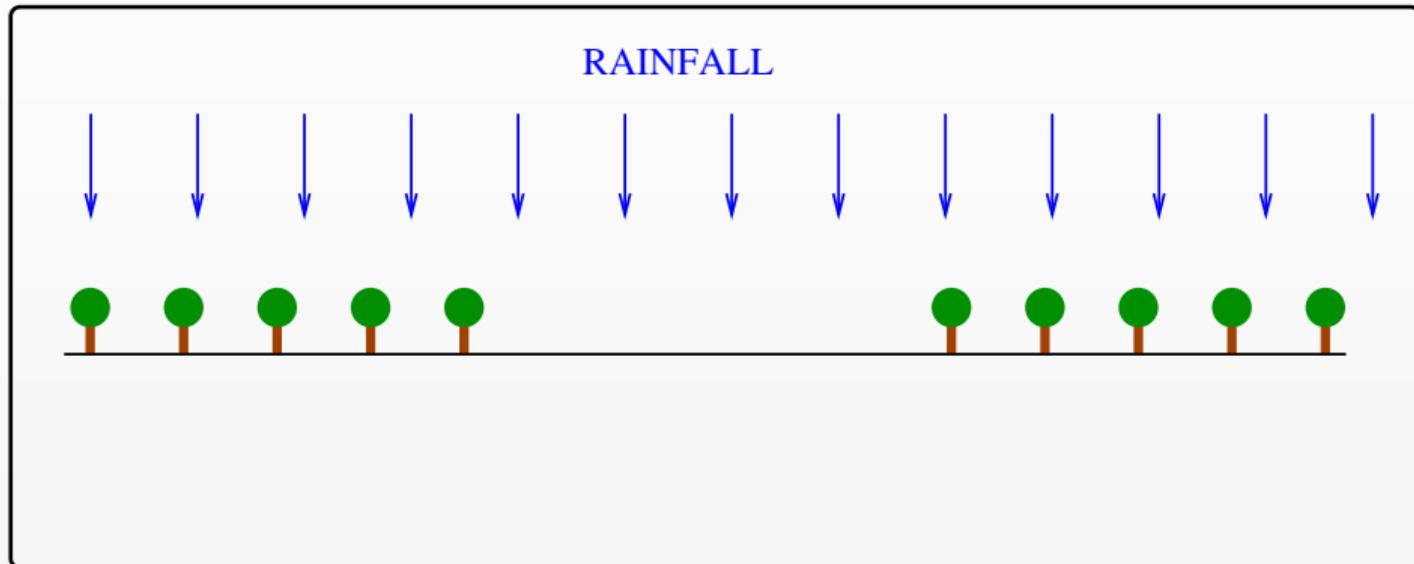
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<sup>1</sup>Source: Google Maps

<sup>2</sup>Source: Wikimedia Commons

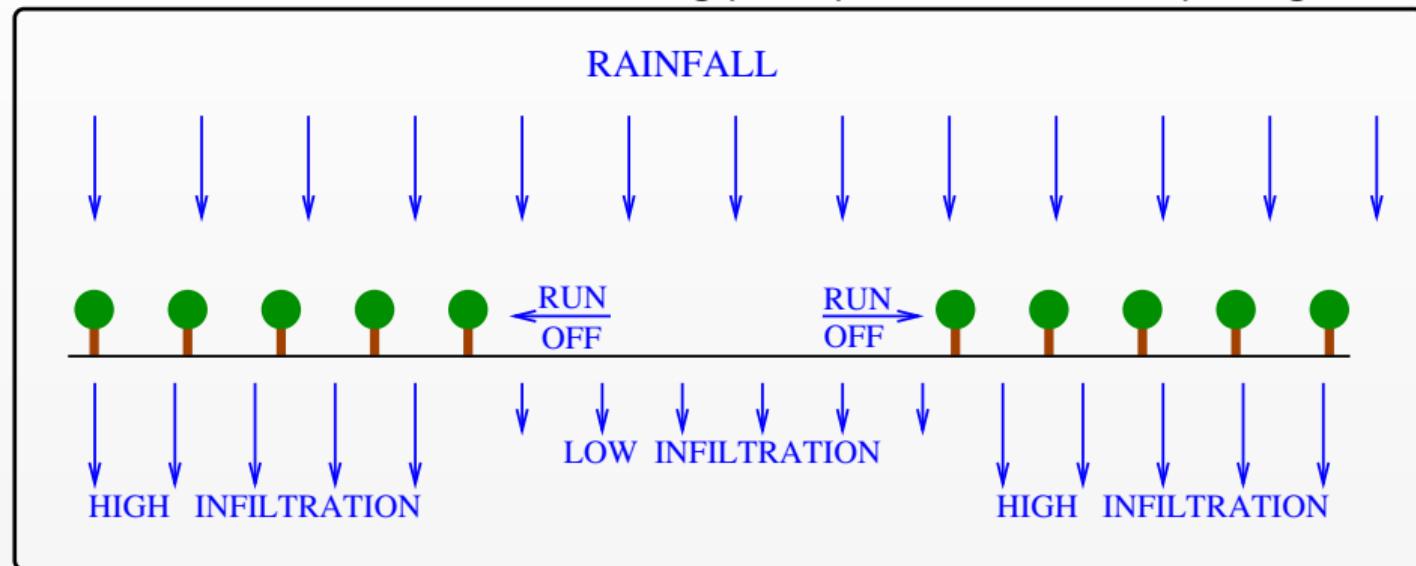
# Local facilitation in vegetation patterns

Positive feedback loop: Water infiltration into the soil depends on local plant density ⇒ redistribution of water towards existing plant patches ⇒ further plant growth.



# Local facilitation in vegetation patterns

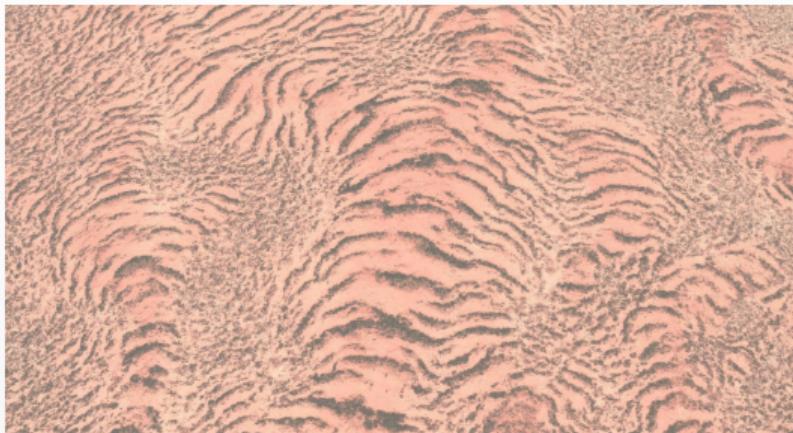
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# Vegetation patterns

Vegetation patterns are a classic example of a **self-organisation principle** in ecology.

Stripe pattern in Ethiopia<sup>3</sup>.



Gap pattern in Niger<sup>4</sup>.



- On sloped ground, stripes grow **parallel to the contours**.

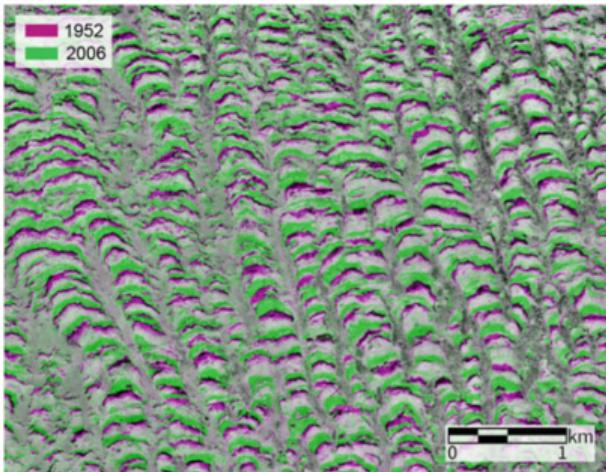
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<sup>3</sup>Source: Google Maps

<sup>4</sup>Source: Wikimedia Commons

# Vegetation patterns

Timeseries data.<sup>5</sup>



Uphill migration due to water gradient.<sup>6</sup>



- Contrasting field data: stripes either **move uphill** (< 1m per year) or are **stationary**<sup>7</sup>.
- **No reports of downhill movement.**

<sup>5</sup>Gandhi, P. et al.: *Dryland ecohydrology*. Springer International Publishing, 2019, pp. 469–509.

<sup>6</sup>Dunkerley, D.: *Desert* 23.2 (2018).

<sup>7</sup>Deblauwe, V. et al.: *Ecol. Monogr.* 82.1 (2012).

# Klausmeier model

One of the most basic phenomenological models is the **extended Klausmeier reaction-advection-diffusion model**.<sup>8</sup>

$$\begin{aligned}\frac{\partial u}{\partial t} &= \underbrace{u^2 w}_{\text{plant growth}} - \underbrace{Bu}_{\text{plant loss}} + \underbrace{\frac{\partial^2 u}{\partial x^2}}_{\text{plant dispersal}}, \\ \frac{\partial w}{\partial t} &= \underbrace{A}_{\text{rainfall}} - \underbrace{w}_{\text{evaporation}} - \underbrace{u^2 w}_{\text{water uptake by plants}} + \underbrace{\nu \frac{\partial w}{\partial x}}_{\text{water flow downhill}} + \underbrace{d \frac{\partial^2 w}{\partial x^2}}_{\text{water diffusion}}.\end{aligned}$$

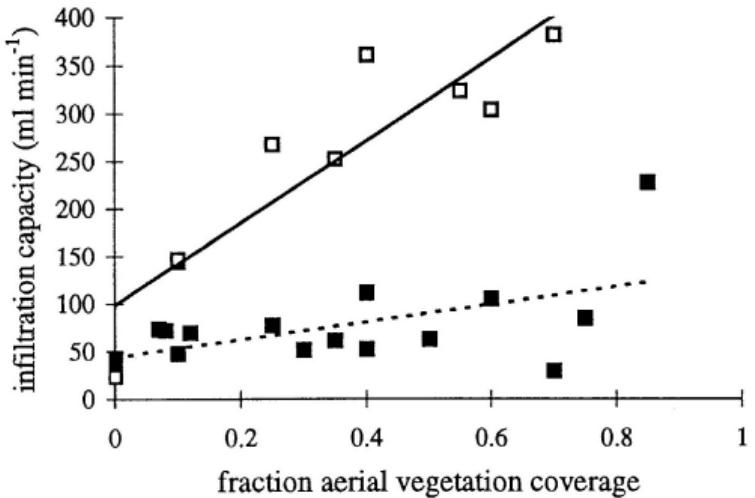
<sup>8</sup>Klausmeier, C. A.: *Science* 284.5421 (1999).

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# Water uptake



Infiltration capacity increases with plant density<sup>9</sup>

The nonlinearity in the water uptake and plant growth terms arises because plants increase the soil's water infiltration capacity.

⇒ Water uptake = Water density × plant density × infiltration rate.

<sup>9</sup>Rietkerk, M. et al.: *Plant Ecol.* 148.2 (2000)

## Research question

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- How can the contrasting field data on uphill movement be explained?

# Main research question

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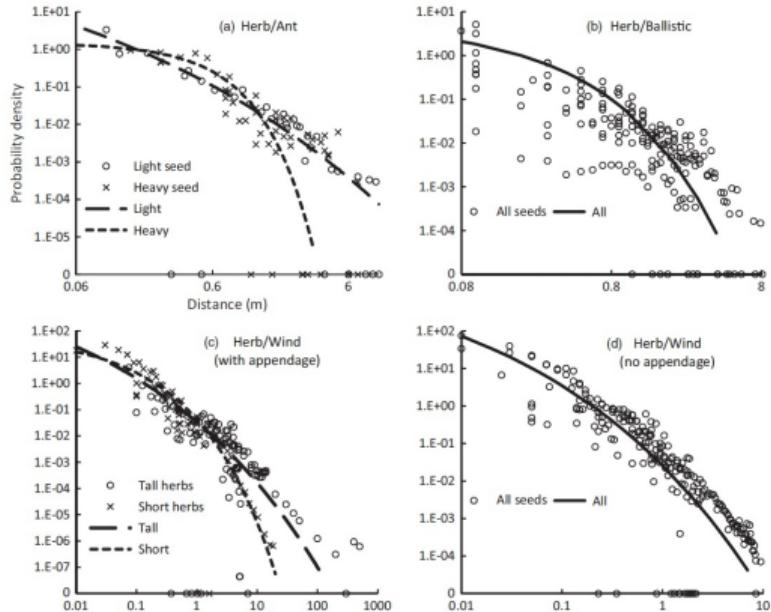
- How does nonlocal seed dispersal affect onset, existence and stability of patterns?  
⇒ How can the contrasting field data on uphill movement be explained?

# Local Model

The Klausmeier model models plant dispersal by a diffusion term, i.e. a local process.

$$\frac{\partial u}{\partial t} = \underbrace{u^2 w}_{\text{plant growth}} - \underbrace{Bu}_{\text{plant loss}} + \underbrace{\frac{\partial^2 u}{\partial x^2}}_{\text{local plant dispersal}},$$
$$\frac{\partial w}{\partial t} = \underbrace{A}_{\text{rainfall}} - \underbrace{w}_{\text{evaporation}} - \underbrace{u^2 w}_{\text{water uptake by plants}} + \underbrace{\nu \frac{\partial w}{\partial x}}_{\text{water flow downhill}} + \underbrace{d \frac{\partial^2 w}{\partial x^2}}_{\text{water diffusion}}.$$

# Nonlocal seed dispersal



Data of long range seed dispersal<sup>10</sup>

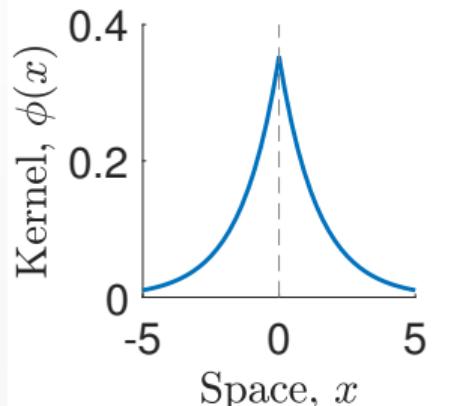
<sup>10</sup>Bullock, J. M. et al.: *J. Ecol.* 105.1 (2017)

More realistic: **Include effects of nonlocal processes**, such as dispersal by wind or large mammals.

# Nonlocal model

Diffusion is replaced by a **convolution of the plant density  $u$**  with a **dispersal kernel  $\phi$** . The scale parameter  $a$  controls the width of the kernel.

$$\begin{aligned}\frac{\partial u}{\partial t} &= \underbrace{u^2 w}_{\text{plant growth}} - \underbrace{Bu}_{\text{plant loss}} + \overbrace{C(\phi(\cdot; a) * u(\cdot, t) - u)}^{\text{nonlocal plant dispersal}}, \\ \frac{\partial w}{\partial t} &= \underbrace{A}_{\text{rainfall}} - \underbrace{w}_{\text{evaporation}} - \underbrace{u^2 w}_{\text{water uptake by plants}} + \underbrace{\nu \frac{\partial w}{\partial x}}_{\text{water flow downhill}} + \underbrace{d \frac{\partial^2 w}{\partial x^2}}_{\text{water diffusion}}.\end{aligned}$$



## Laplacian kernel

If  $\phi$  decays exponentially as  $|x| \rightarrow \infty$ , and  $C = 2/\sigma(a)^2$ , then the nonlocal model tends to the local model as  $\sigma(a) \rightarrow 0$ .

E.g. Laplace kernel

$$\phi(x) = \frac{a}{2} e^{-a|x|}, \quad a > 0, \quad x \in \mathbb{R}.$$

Useful because

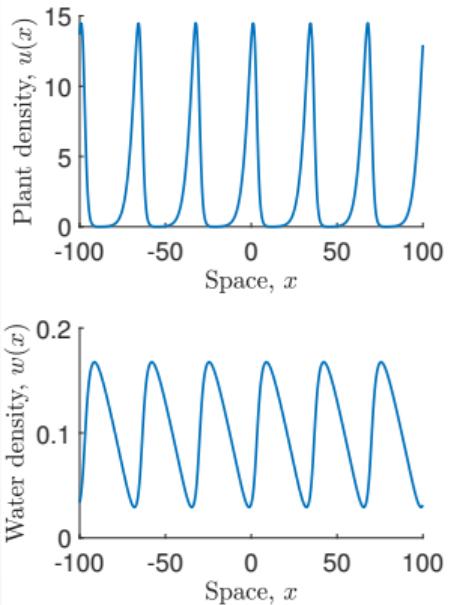
$$\hat{\phi}(k) = \frac{a^2}{a^2 + k^2}, \quad k \in \mathbb{R}.$$

and allows transformation into a local model. If  $v(x, t) = \phi(\cdot; a) * u(\cdot; t)$ , then

$$\frac{\partial^2 v}{\partial x^2}(x, t) = a^2(v(x, t) - u(x, t))$$

# Travelling waves

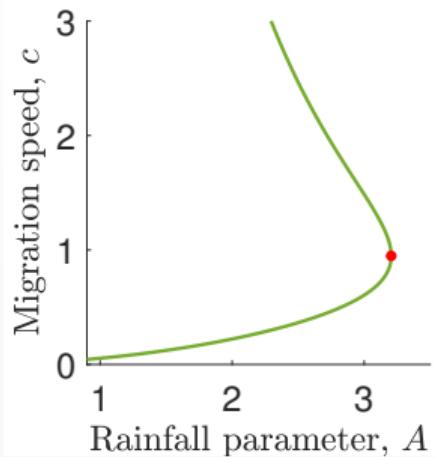
- Numerical simulations of the model on sloped terrain suggest uphill movement  $\Rightarrow$  Periodic travelling waves.
- Patterns correspond to **limit cycles** of the travelling wave integro-ODEs.



Numerical simulation.

# Travelling waves

- Numerical simulations of the model on sloped terrain suggest uphill movement  $\Rightarrow$  Periodic travelling waves.
- Patterns correspond to **limit cycles** of the travelling wave integro-ODEs.
- Numerical continuation shows that **patterns emanate from a Hopf bifurcation** and terminate at a homoclinic orbit.
- In the PDE model, pattern onset occurs at a threshold  $A = A_{\max}$ , the maximum rainfall level of the Hopf bifurcation loci in the travelling wave ODEs.



Location of the Hopf bifurcation in  $A$ - $c$  plane.

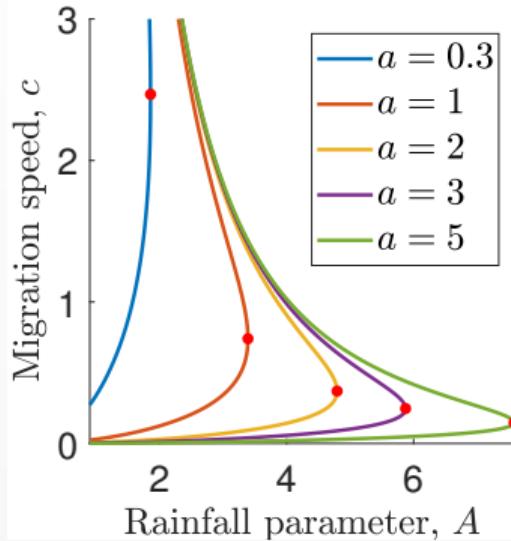
## Pattern onset

Using that  $\nu \gg 1$ ,

$$A_{\max} = \left( \frac{3C - B - 2\sqrt{2C}\sqrt{C - B}}{(B + C)^2} \right)^{\frac{1}{4}} a^{\frac{1}{2}} B^{\frac{5}{4}} \nu^{\frac{1}{2}},$$

to leading order in  $\nu$  as  $\nu \rightarrow \infty$ .

- Note that  $A_{\max} = O(\sqrt{\nu})$ .
- Decrease in  $a$  (i.e. increase in kernel width) causes decrease of  $A_{\max}$ .
- Increase in dispersal rate  $C$  causes decrease of  $A_{\max}$ .

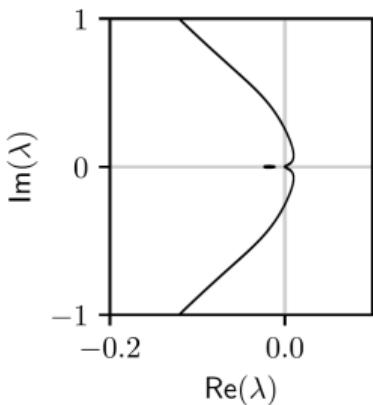


Locus of Hopf bifurcation for fixed  $C$  and varying  $a$ .<sup>11</sup>

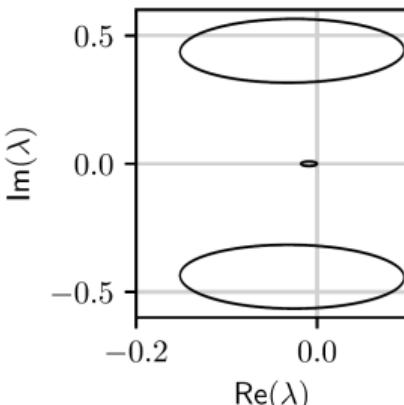
<sup>11</sup>EL and Sherratt, J. A.: *J. Math. Biol.* 77.3 (2018)

# Pattern stability

- The **essential spectrum** of a periodic travelling wave determines the behaviour of small perturbations.  $\Rightarrow$  Tool to determine pattern stability.
- Two different types stability boundaries: **Eckhaus-type** and **Hopf-type**.
- Essential spectra and stability boundaries are calculated using the numerical continuation method by Rademacher et al.<sup>12</sup>



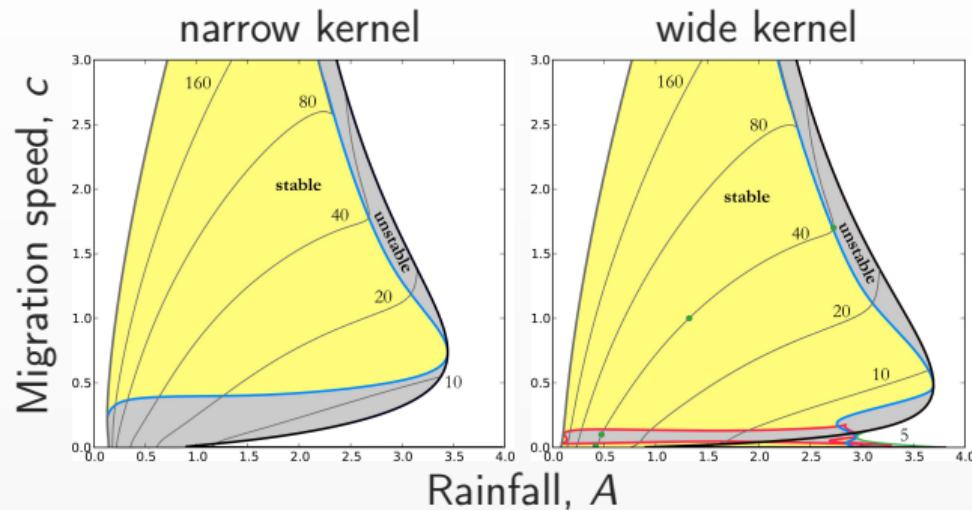
Eckhaus-type



Hopf-type

<sup>12</sup>Rademacher, J. D., Sandstede, B. and Scheel, A.: *Physica D* 229.2 (2007)

# Pattern existence and stability



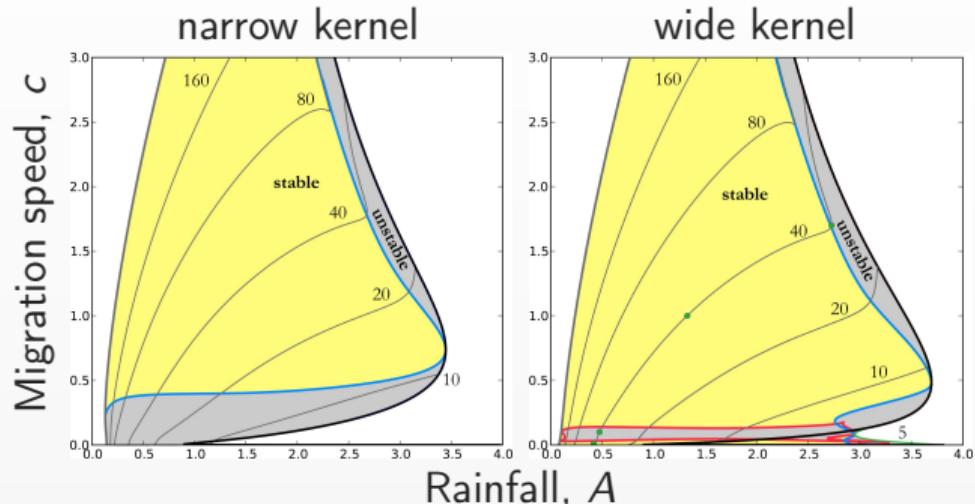
Stability of patterns in the  $A$ - $c$  plane.<sup>13</sup>

For wide kernels, the stability boundary towards the desert state changes from Eckhaus to Hopf-type. This yields

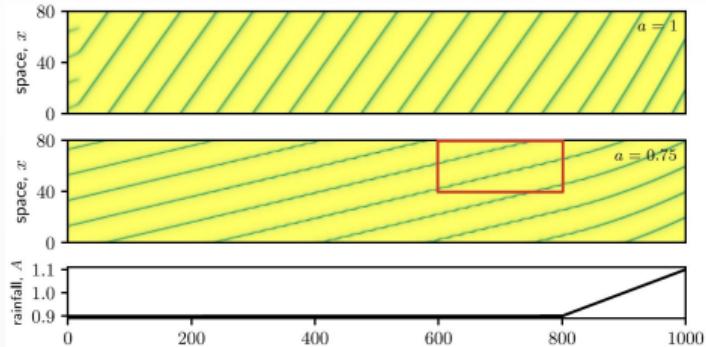
- increased resilience due to oscillating vegetation densities in peaks,

<sup>13</sup> Bennett, J. J. R. and Sherratt, J. A.: *J. Theor. Biol.* 481 (2018)

# Pattern existence and stability



Stability of patterns in the  $A$ - $c$  plane.<sup>14</sup>

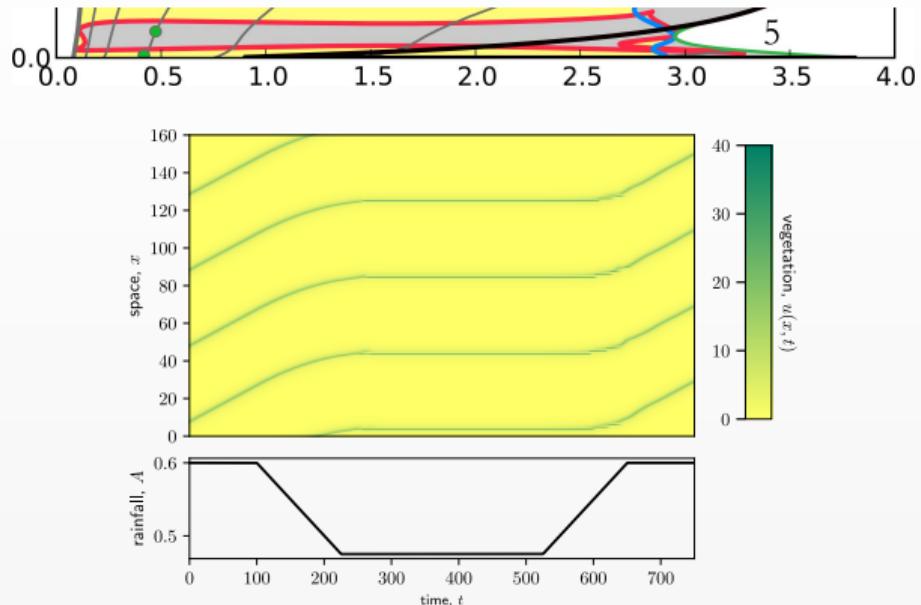


For wide kernels, the stability boundary towards the desert state changes from Eckhaus to Hopf-type. This yields

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# Pattern existence and stability



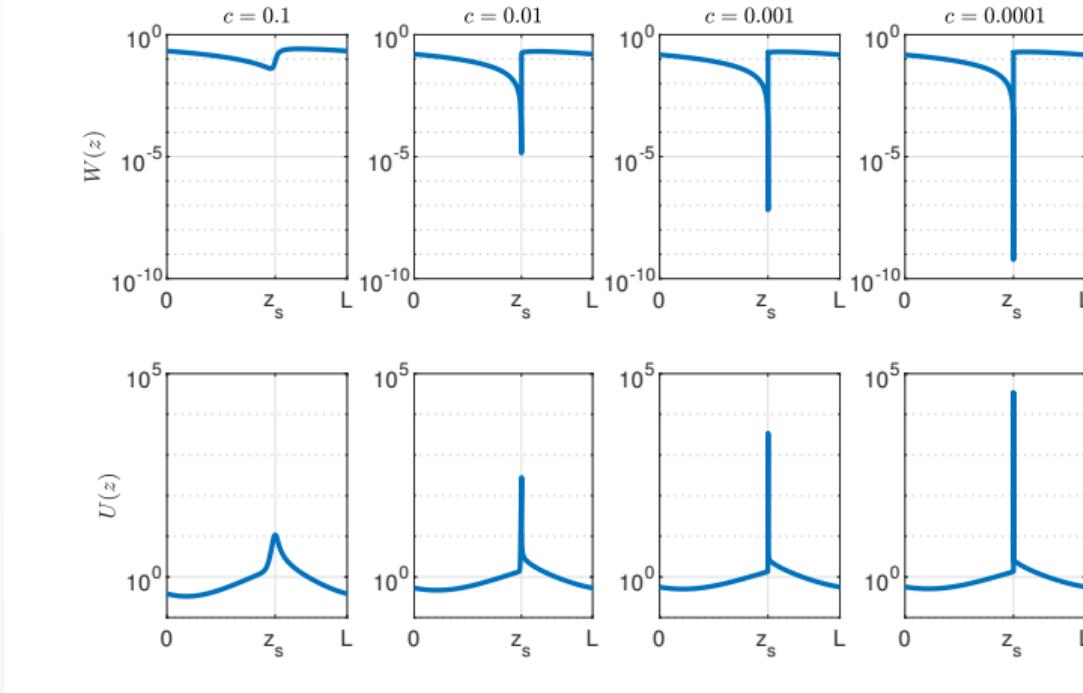
For wide kernels, the stability boundary towards the desert state changes from Eckhaus (sideband) to Hopf-type. This yields

- increased resilience due to oscillating vegetation densities in peaks,
- existence of stable patterns with small migration speed ( $c \ll 1$ ).

Existence of stable (almost) stationary patterns.<sup>15</sup>

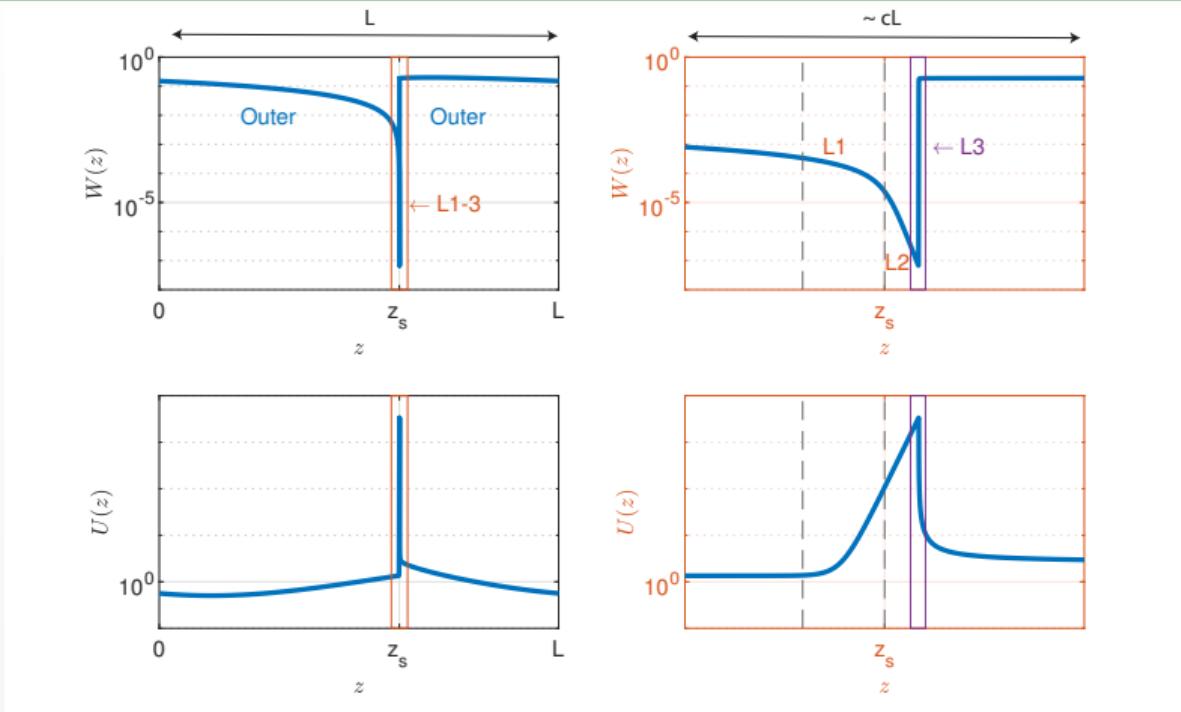
<sup>15</sup> Bennett, J. J. R. and Sherratt, J. A.: *J. Theor. Biol.* 481 (2018)

# Almost stationary spike patterns



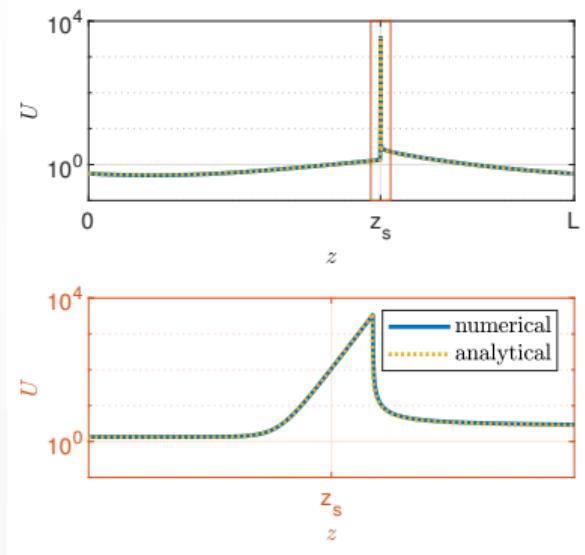
As  $c$  decreases, plant density develops a “spike”.

# Almost stationary spike patterns



Layered structure of spike solution

# Almost stationary spike patterns



Existence of almost stationary patterns is confirmed analytically using a singular perturbation theory approach, exploiting  $c \ll 1$ .

Analytical calculation of (almost) stationary patterns.<sup>16</sup>

<sup>16</sup>EL and Sherratt, J. A.: *J Math Biol* 86.15 (2023)

## Main conclusion

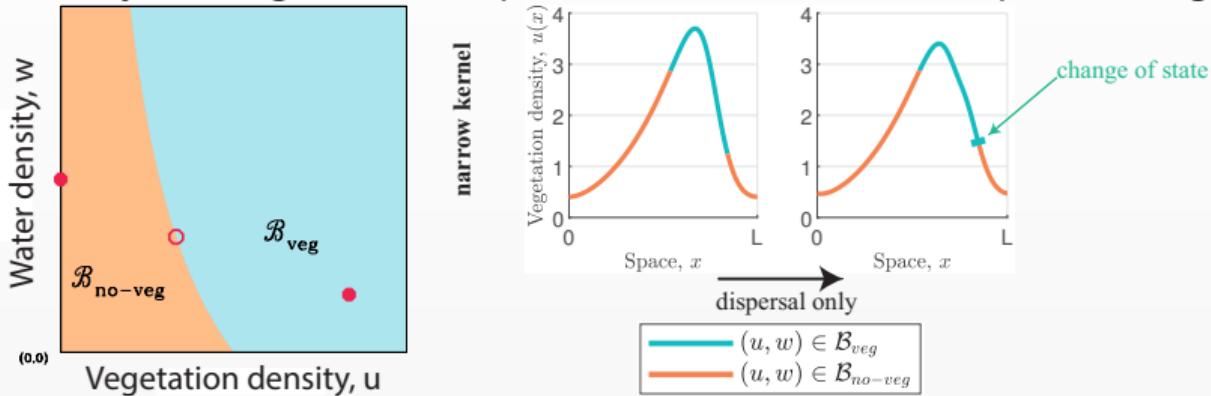
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- How can the contrasting field data on uphill movement be explained?

For long seed dispersal distances moving (uphill) and stationary patterns can occur for the same parameter values.

# Almost stationary patterns

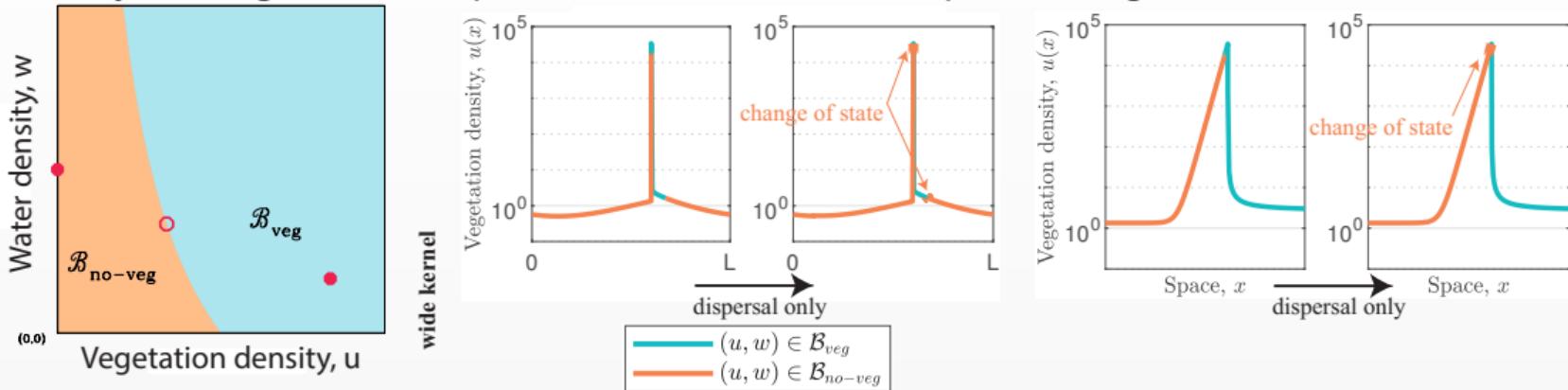
Q: Why do longer mean dispersal distances slow down pattern migration?



- narrow kernel: dispersal-induced plant increase at pattern edge causes transition from basin of attraction of desert state to vegetated state.

# Almost stationary patterns

Q: Why do longer mean dispersal distances slow down pattern migration?



- Narrow kernel: dispersal-induced plant increase at pattern edge causes transition from basin of attraction of desert state to vegetated state.
- Wide kernel: less dispersal to stripe edges  $\rightarrow$  insufficient to cause transition from basin of attraction of desert state to vegetated state.

## Other conclusions

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- Wider kernels and higher dispersal rates inhibit pattern onset.
- Stability analysis of periodic travelling waves provides ecological insights into pattern dynamics: Long-range seed dispersal increases the resilience of a pattern and stabilises (almost) stationary patterns.
- Numerical simulations (pattern onset) and space discretisation to avoid nonlocality (calculation of essential spectra) show no qualitative differences for other kernel functions.

## Future Work

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- Empirical tests of these hypotheses?
- How can a system reach a non-migration state?
- How resilient are non-migration states to environmental change?

# References

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- [1] Bennett, J. J. R. and Sherratt, J. A.: *J. Theor. Biol.* 481 (2018), pp. 151–161.
- [2] Eigentler, L. and Sherratt, J. A.: *J. Math. Biol.* 77 (2018), pp. 739–763.
- [3] Eigentler, L. and Sherratt, J. A.: *J. Math. Biol.* 86.15 (2023).
- [4] EL and Sherratt, J. A.: *J. Math. Biol.* 77.3 (2018), pp. 739–763.