Investigation of Homoclinic Orbits in the Circular Restricted Three Body Problem

Luke Bury & Don Kuettel April 27, 2018

Introduction

In 1885, a competition was held by Acta Mathematica in which participants were challenged to solve one of four outstanding math problems. Henri Poincaré, a prominent mathematician of the time, entered the contest and ultimately won. However, his submission contained a critical mistake that, when corrected, led to the discovery of homoclinic orbits. In mathematics, a homoclinic orbit is defined as a trajectory which joins a saddle equilibrium point to itself. More precisely, a homoclinic orbit is a union between the stable manifold, $W^s(p)$, and the unstable manifold, $W^u(p)$, of an equilibrium point. Figure 1 shows an example of a simple, two-dimensional homoclinic orbit about the saddle equilibrium point p. As the figure shows, as time approaches either negative or positive infinity, the homoclinic orbit will approach p.

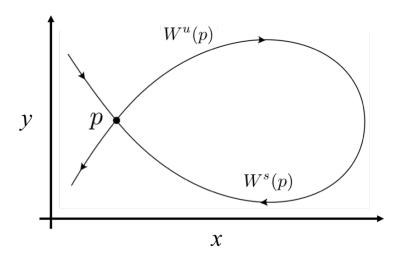


Figure 1: A 2-Dimensional homoclinic orbit.

Homoclinic orbits can reveal much about the chaotic behavior of dynamical systems, but locating these orbits requires a non-trivial, multi-step process. First, a periodic orbit about a chosen equilibrium point must be identified. From this periodic orbit, both stable and unstable manifolds must be constructed. Finally, intersections of the stable and unstable manifolds are found via Poincaré sections, and homoclinic orbits are thus identified. The following report briefly discusses the history of Poincare's discovery of homoclinic orbits and demonstrates a detailed procedure

for locating these orbits. Ultimately, homoclinic orbits in the Earth-Moon system are found and presented.

The Competition

As documented by Andersson and Barrow-Green [1, 2], in 1885, Acta Mathematica announced a mathematics competition to the world. This competition, sponsored by King Oscar II of Sweden, encouraged interested parties to make an attempt at solving one of four unanswered problems of the time. Henri Poincaré, a prominant mathematician (who was largely favored to win the competition) attempted the first problem, which essentially asked for a solution to the perplexing n-body problem. Rather than sticking to the exact prompt, Poincaré decided to instead attempt a subset of the n-body problem known as the three-body problem - the first order of the n-body problem remaining unsolved. To further simplify his initial effort, Poincaré restricted the three-body system in a manner known today as the Circular Restricted Three-Body Problem (CR3BP), which will be discussed further in the next section.

Poincaré's submission to the competition was the culmination of several strands of his work from the previous decade, which included geometrical and analytical theory, integral invariants, and periodic solutions to dynamical systems. Poincaré applied these theories in an attempt to rigorously prove stability and find periodic solutions for the CR3BP. The competition judges immediately recognized the importance of Poincaré's submission and unanimously crowned him the victor. However, around the time that his work was first being printed, a discussion with Lars Edvard Phragmén led to the discovery of an error within Poincaré's submission that held significant ramifications. The error was rooted in Poincaré's failing "to take proper account of the exact geometric nature of a particular curve" [3]. In Theorem III of the paper's first, and flawed, edition, Poincaré claimed that a particular invariant curve was closed (Figure 2(a,b)). He failed to consider that the curve could be self-intersecting (Figure 2(c)). After realizing his mistake, Poincaré reworked his derivations to discover that the asymptotic surfaces are not closed, but instead, they intersect along infinitely many asymptotic trajectories. Ironically, this mistake cemented Poincaré's place among legendary mathematicians for his resulting discovery of doubly asymptotic, or homoclinic, points/orbits in the CR3PB and dynamical systems in general. These curves will resurface in section ???????as the Poincaré sections of intersecting stable and unstable manifolds, where exact intersections of the curves correspond to homoclinic orbits. First, however, the CR3BP and periodic orbits must be well understood.

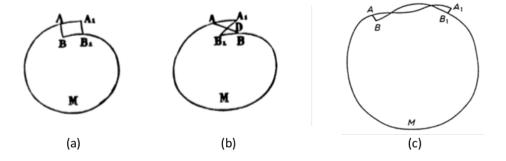


Figure 2: (a) Diagrams of incorrectly closed invariant curves from the first edition of Poincaré's memoir. (b)Invariant curve with self intersection from Poincaré's corrected work [3]

Circular Restricted Three Body Problem and Applications

- Description of CR3BP
- Application to find POs... manifolds... homoclinic orbits

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Before studying periodic orbits, their manifolds, and the process of locating homoclinic orbits, it is necessary to provide the reader a solid foundation of the Circular Restricted Three Body Problem (CR3BP) - the primary mathematical toolbox from which the aforementioned trajectories are created and analyzed. In the CR3BP (Figure 3), the origin of the system is set at the barycenter of the two main bodies in the system (e.g., the Earth & Moon), and the frame rotates so these bodies remain stationary on the x-axis. The bodies are assumed to move in perfectly circular orbits and act as point masses from a gravitational perspective. The restricted problem is then to ascertain the motion of the third body whose mass is considered negligible. The system is typically normalized so that the masses of the two bodies sum to 1 ($m_1 = \mu$, $m_2 = 1 - \mu$), the distance between the bodies is 1, and the gravitational constant G is equal to 1. Under these conditions, the equations of motion for the CR3BP are shown in Equations 1-3. Figure 4 demonstrates the difference between a trajectory about a single body with no external forces and a trajectory where a third body is included. Clearly, the addition of a third body has the capacity to significantly perturb an orbit.

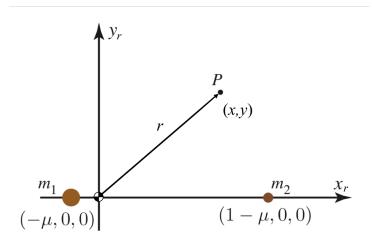


Figure 3: Layout of Circular Restricted Three Body Problem [4]

$$\ddot{x} = 2\dot{y} + x - (1 - \mu) \left(\frac{x + \mu}{R_1^3}\right) - \mu \left(\frac{x - 1 + \mu}{R_2^3}\right) \tag{1}$$

$$\ddot{y} = -2\dot{x} + y\left(-\frac{1-\mu}{R_1^3} - \frac{\mu}{R_2^3} + 1\right) \tag{2}$$

$$\ddot{z} = z \left(-\frac{1-\mu}{R_1^3} - \frac{\mu}{R_2^3} \right). \tag{3}$$

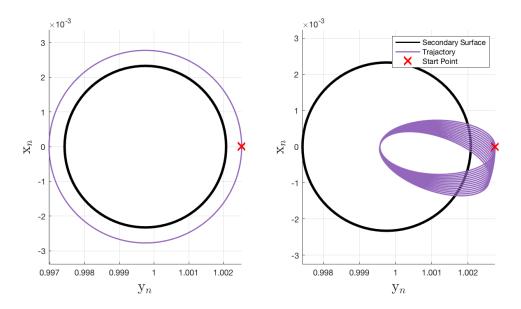


Figure 4: Demonstration of third-body-effects in orbital dynamics. On the left, a particle is propagated in a circular orbit about a secondary body with basic, two-body dynamics. On the right, the particle is propagated from the same initial conditions, but the CR3BP equations of motion (Equations 1-3) are implemented with a mass ratio, μ , of 2.5×10^{-4} .

Periodic/Homoclinic Orbits

- Eq points and stability
- zero velocity curves
- POs and algorithm to find them
- manifolds and the algorithms
- PO stability ... eigenvalues of monodromy matrix; Poincaré sections and their intersections ... homoclinic orbits
- Plotting and discussion of homoclinic orbits (piecewise)

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Conclusion

• Homoclinic orbits are cool/interesting, but for spacecraft applications they aren't very useful; Heteroclinic orbits are useful fa show

References

- [1] K. G. Andersson. Poincaré's discovery of homoclinic points. Archive for History of Exact Sciences, 48(2):133–147, 1994.
- [2] J. Barrow-Green. Oscar ii's prize competition and the error in poincaré's memoir on the three body problem. Archive for History of Exact Sciences, 48(2):107–131, 1994.
- [3] J. Barrow-Green. Poincaré and the Three Body Problem. American Mathematical Society, 1997.
- [4] Lo M. Marsden J. Ross S. Sang Koon, W. Dynamical Systems, The Three-Body Problem, and Space Mission Design. Marsden Books, 2011.