

May 10, 2024

# Quantum Simulation of Schrödinger Equation using Qiskit Framework

Project

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## Introduction

# Why Quantum simulation?

- Initial idea was proposed by Feynman in his 1982 paper "Simulating Physics with Computers"
- Idea of a non-Turing machine that would simulate quantum systems (i.e., quantum computer)
- Simulation of quantum many-body problems on a classical computer is a difficult task.
- Size of the Hilbert space of the system grows exponentially with the no. of particles.

- For example, consider simulating a system of  $n$  spin-1/2 particles
- Size of the Hilbert space grows as  $2^n$ 
  - For classical simulation, no. of bits required =  $2^n$
  - For quantum simulation, no. of qubits required =  $n$
- As noted by Feynman, growth in memory requirement is only linear on a quantum computer.
- Of course this is only true if we can design a quantum algorithm to extract information efficiently.
- In an actual quantum computer, we would need more than  $n$  qubits to account for noise and error correction.

## Quantum Simulation: Approach and Software used

# Software

- We have used IBM's Qiskit library for our project.
- Qiskit has many quantum computer simulators, two among them are:
  - Statevector simulator
  - Quantum Assembly Language Simulator (QASM)
- Statevector simulator assumes an ideal quantum computer with no noise, which gives an ideal output for a circuit.
- QASM, on the other hand, attempts to mimic a quantum computer by adding little noise to the result.
- QASM operates by running the circuit multiple times and storing the no. of times an outcome has occurred.

# Approach

- We have used both the statevector simulator and QASM simulator for our project.
- The result for a given quantum circuit can be then visualized by plotting a histogram.
- We have performed 6-qubit simulations for a single particle quantum system in different exactly solvable potentials, namely:
  - Free particle
  - Infinite square well potential
  - Quantum tunneling across Barrier potential
  - Simple Harmonic Oscillator potential
- In addition to this, we have performed simulation of an interacting two-particle system in an infinite square well potential.

# Quantum Simulation of Schrödinger Equation



# Quantum simulations

- The goal is to model single-particle quantum systems using quantum computing.
- There are two primary parts to the simulation process:
  - Initializing the Wavefunction: The first stage is to initialize the wavefunction, which is the quantum system's initial state.
  - The second phase involves creating and putting into practice a time-evolution algorithm that can precisely evolve the wavefunction over time.
  - The noncommutativity of the Hamiltonian operators a feature of quantum systems must be handled by this technique.
- By changing  $V$  we have simulated for free particle, particle in a box, Quantum tunneling, SHO and two particle potential.

# Schrödinger Equation

- The time dependent Schrödinger Equation is given by:

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

- Solving this differential equation yields:

$$|\psi(t)\rangle = e^{-iHt} |\psi_{\text{init}}\rangle$$

- Since the Hamiltonian encompasses the total energy of the system, it can be expressed as  $H = K + V$ , where  $K$  and  $V$  denote the operators for kinetic and potential energies, respectively. It's worth noting, however, that  $V$  and  $K$  do not commute, implying that  $e^{V+K} \neq e^V e^K$ , which complicates the simulation of the system.

## Discretization of Variables

- We must first discretize the continuous variables  $x$  and  $t$ .  
Consider the motion inside a finite region  $-d \leq x \leq d$ . Then, we can decompose this region into  $2^n$  intervals of length  $\Delta = 2d/2^n$  and represent these intervals by means of the Hilbert space of an  $n$ -qubit quantum register.
- The wave function is approximated as follows:

$$\tilde{\psi}(t) = \frac{1}{\aleph} \sum_{i=0}^{2^n-1} \psi(x_i, t) |i\rangle$$

where

$$x_i = -d + (i + \frac{1}{2})\Delta$$

and

$$\aleph = \sqrt{\sum_{i=0}^{2^n-1} |\psi(x_i, t)|^2}$$

# Trotter Decomposition

- The Schrödinger equation may be integrated by propagating the initial wave function for each time step  $\epsilon$  as follows:

$$\psi(x, t + \epsilon) = e^{\frac{-i}{\hbar}[H_0 + V(x)]\epsilon} \psi(x, t)$$

- The Lie-Trotter formula is given by:

$$\lim_{n \rightarrow \infty} \left( e^{A \frac{t}{n}} e^{B \frac{t}{n}} \right)^n = e^{(A+B)t}$$

- If the time step is small enough, we may write:

$$e^{\frac{-i}{\hbar}[H_0 + V(x)]\epsilon} = e^{-\frac{i}{\hbar} H_0 \epsilon} e^{-\frac{i}{\hbar} V(x) \epsilon}$$

# Final Algorithm

- We now consider the momentum  $k$ , variable conjugate to  $x$ , i.e.  $-i(d/dx) = F^{-1}kF$
- The first operator (kinetic energy part) may be written as:

$$e^{-\frac{i}{\hbar}H_0\epsilon} = F^{-1}e^{\frac{i}{\hbar}(\frac{\hbar^2k^2}{2m})\epsilon}F$$

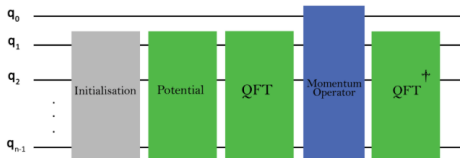
- Thus, the final wave function after time  $t=l\epsilon$  obtained from the initial wave function is given by applying the following unitary operator  $l$  times:

$$F^{-1}e^{\frac{i}{\hbar}(\frac{\hbar^2k^2}{2m})\epsilon}Fe^{-\frac{i}{\hbar}V(x)\epsilon}$$

# Final Algorithm

Below is the final simulation algorithm we arrive at after combining everything discussed above. The time evolution operator is given by:

$$|\psi(t)\rangle = F^{-1} \exp\left(-\frac{i}{\hbar} \frac{p^2}{2m} \Delta t\right) F \exp\left(-\frac{i}{\hbar} V(x) \Delta t\right) |\psi_{\text{init}}\rangle$$



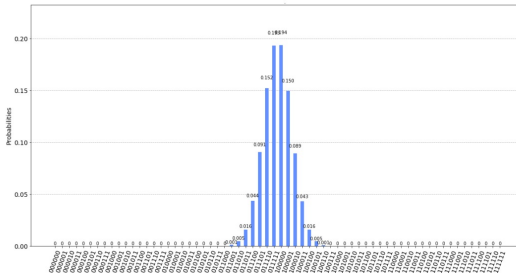
**Figura 1:** Quantum circuit representation of the time evolution operator for a single time step

## Results

# Free Particle

$V = 0$ , We start by initializing Gaussian wave packet

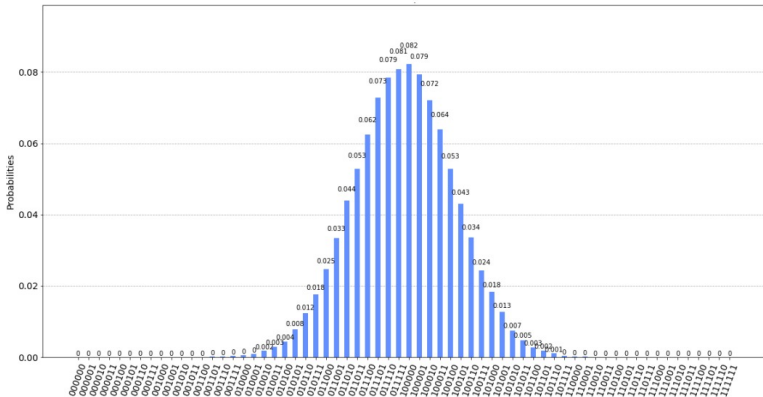
$$|\psi(x)\rangle = \frac{1}{\sigma\sqrt{2\pi}} e^{-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2},$$



Initialisation,  $t=0$  : Initial Wave packet

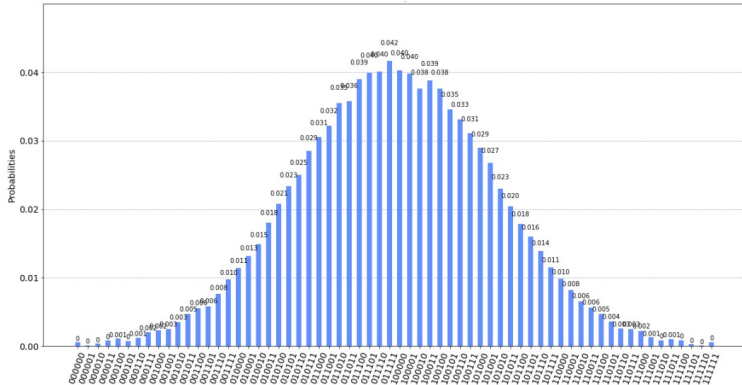


# Free Particle



After 1st iteration,  $t=1$  : Slightly Spread out

# Free Particle



After 8th iteration,  $t=8$  : Spread out increases with time as expected

# Particle in Infinite potential well

We need to solve,

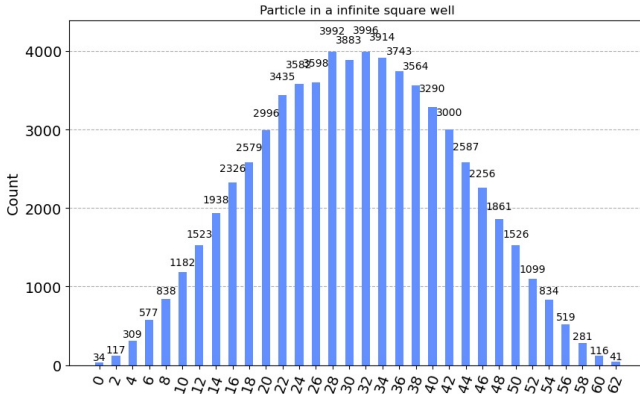
$$\left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)\right) \varphi_E(x) = E \varphi_E(x),$$

$$V(x) = \begin{cases} \infty, & x < 0 \\ 0, & 0 < x < L \\ \infty, & x > L. \end{cases}$$

Solution is

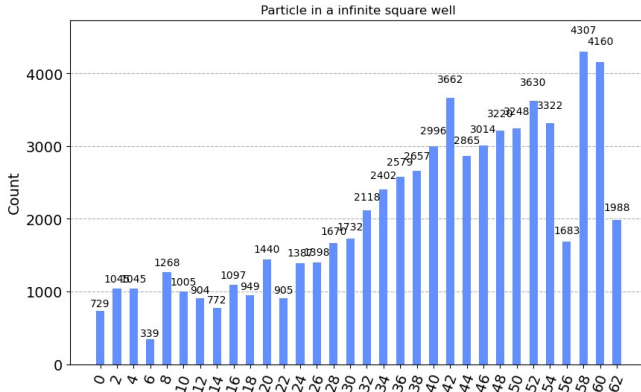
$$\varphi(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right), \quad n = 1, 2, 3, \dots$$

# Particle in Infinite potential well



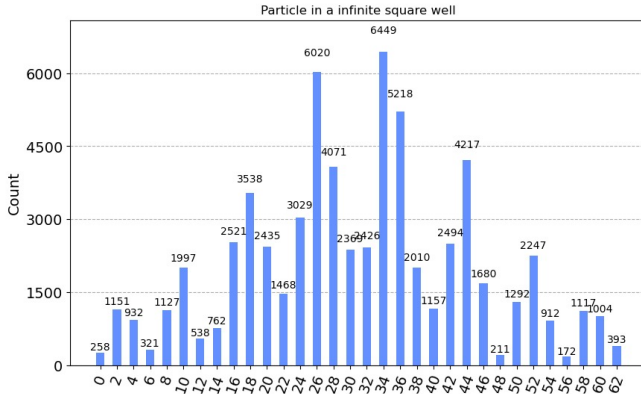
Initialisation,  $t=0$  : Initial Wave packet

# Particle in Infinite potential well



After 3rd iteration,  $t=3$  : Wave bounces of the boundary as expected

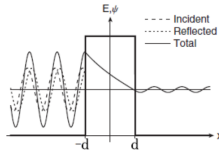
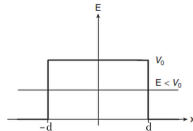
# Particle in Infinite potential well



After 7th iteration,  $t=7$

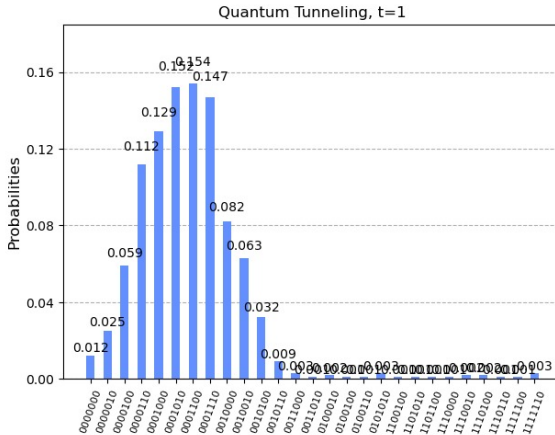
# Quantum Tunneling

A potential barrier with the incident particle energy below the potential barrier height.



An incident sinusoidal wavefunction quantum tunneling through a barrier. The wavefunction exponentially decays within the potential barrier.

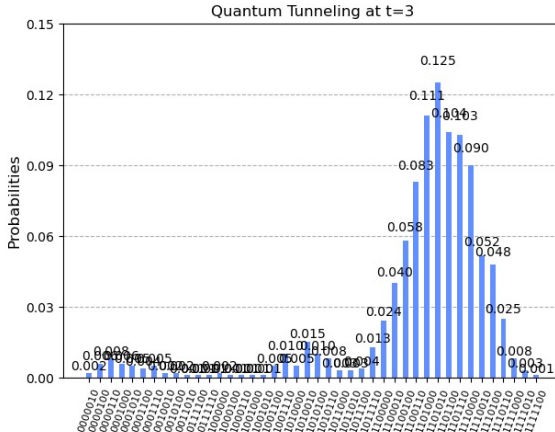
# Quantum Tunneling



Initialisation,  $t=1$  : Wave packet just before interacting with the potential barrier

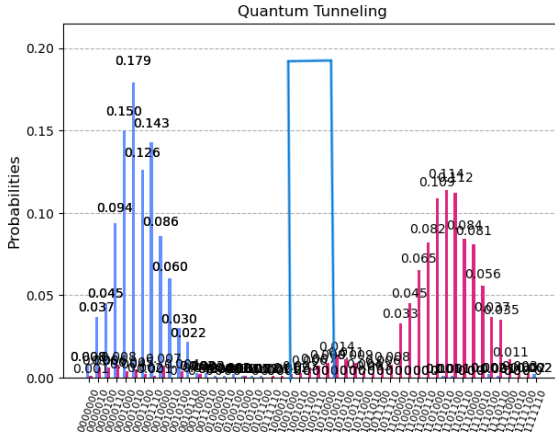


# Quantum Tunneling



After 3rd iteration,  $t=3$  : Tunneled wave

# Quantum Tunneling



initial state and 4th iteration: Better visualization of tunneling

# Simple Harmonic Potential

Hamiltonian:  $\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$

Time-independent Shrodinger Eqn:  $\hat{H}\psi_n(x) = E_n\psi_n(x)$

The wave functions of the harmonic oscillator are given by the Hermite polynomials multiplied by a Gaussian factor:

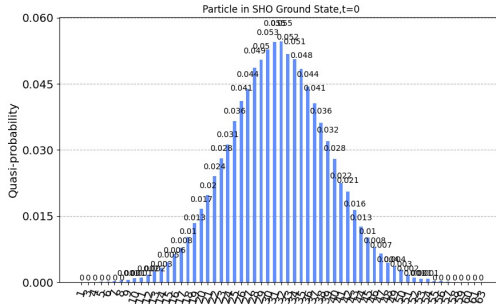
$$\psi_n(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n\left(\sqrt{\frac{m\omega}{\hbar}}x\right) e^{-\frac{m\omega x^2}{2\hbar}}$$

where  $H_n(x)$  are the Hermite polynomials, and  $\hbar$  is the reduced Planck constant.

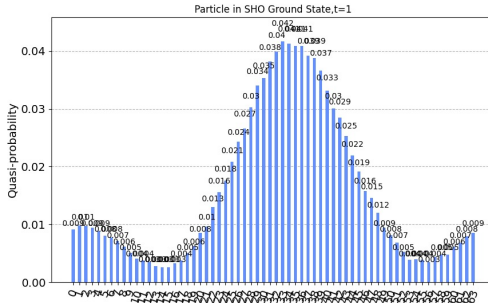
Energy eigenvalues:  $E_n = \left(n + \frac{1}{2}\right) \hbar\omega$ , where  $n = 0, 1, 2, \dots$

represents the quantum number associated with the energy levels.

Initialisation,  $t=0$  : Initial Wave packet

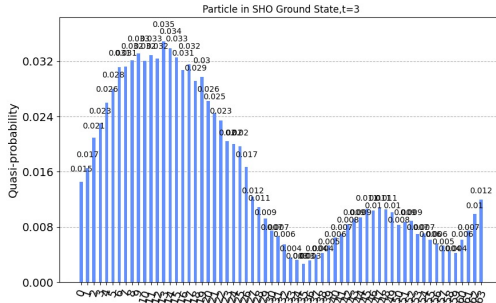


# Simple Harmonic Potential

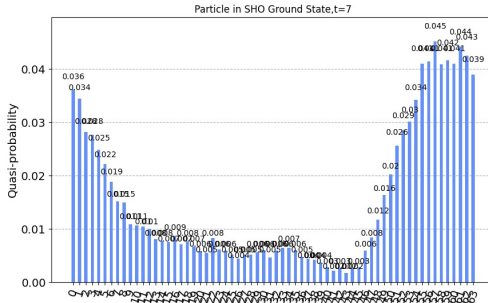


After 1st iteration,  $t=1$

# Simple Harmonic Potential



# Simple Harmonic Potential



After 7th iteration,  $t=7$

# Two particle Schrödinger Equation

- Hamiltonian of two interacting particle:

$$H = (H_{0A} + H_{0B} + V_A(x_A) + V_B(x_B) + V_{AB}(x_A - x_B)) \quad (1)$$

where  $x_A$  and  $x_B$  represent the position coordinates of the first and second particle, respectively.

- The operators  $H_{0i} = \frac{p_A^2}{2m_i} = -\frac{\hbar^2}{2m_i} \frac{\partial^2}{\partial x_i^2}$  is the the standard free Hamiltonians for the particles.
- Schrödinger equation:

$$i\hbar \frac{d}{dt} \psi_{AB}(x_A, x_B, t) = H \psi_{AB}(x_A, x_B, t), \quad (2)$$



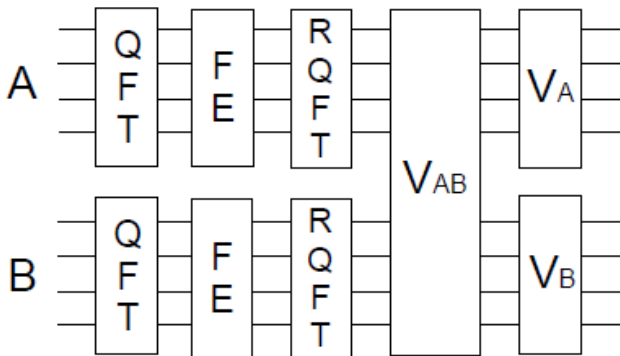
## Two Particle Potential

- $V_A(x_A)$  and  $V_B(x_B)$  are external potentials affecting particles A and B respectively.
- $V_{AB}$ : Interaction potential  
Depends on the relative distance  $x_A - x_B$ .
- For our case,  $V_{AB}(x_A - x_B) = v_0 \delta(x_A - x_B)$ ,  
 $v_0$  is a constant.
- System evolution from an initial state at time  $t_0$  to a state at  $t_1 = t_0 + \Delta t$  is given by:

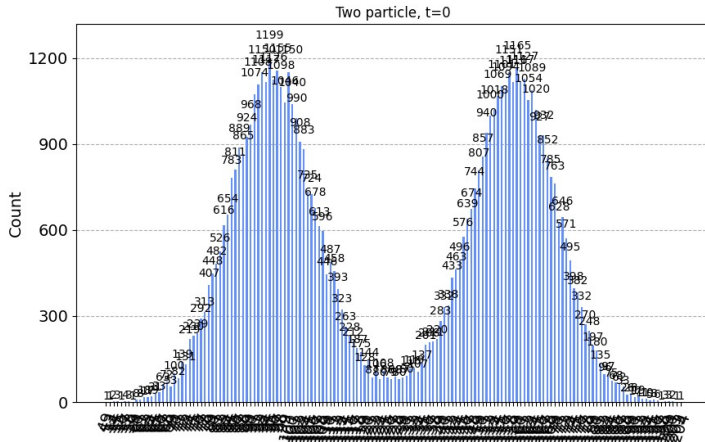
$$\psi_{AB}(x_A, x_B, t_1) = e^{(-\frac{i}{\hbar}(H)\Delta t)}\psi_{AB}(x_A, x_B, t_0),$$

(3)

## Two Particle Quantum circuit

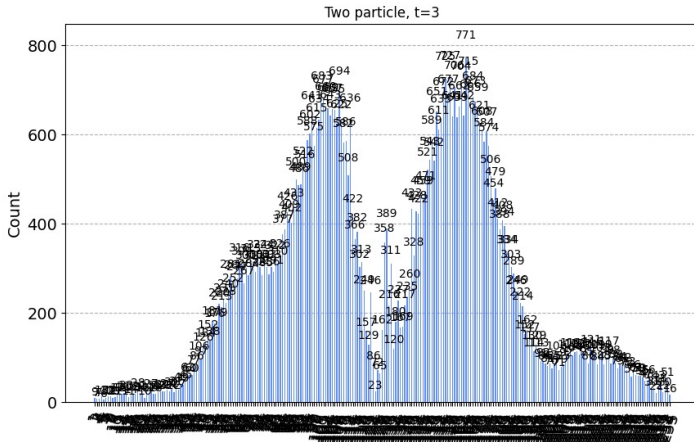


# Two Particle



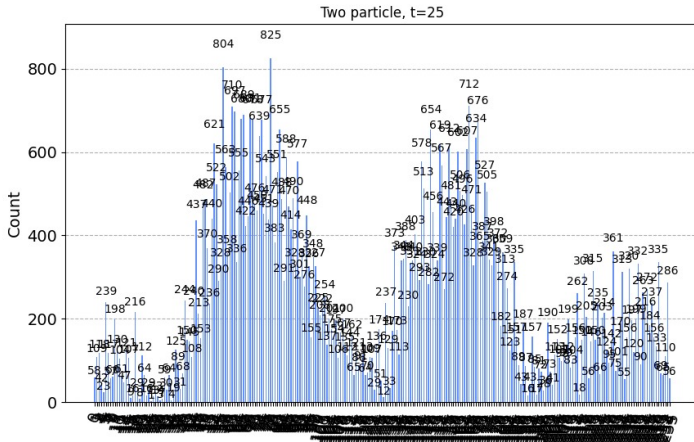
Initialisation of two particle wave packet,  $t=0$  : Initial Wave packet

# Two Particle

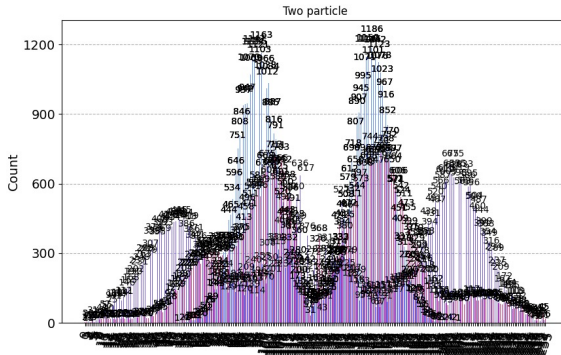


Two particle wave packet after 3rd iteration,  $t=3$

# Two Particle



Visualization of time evolution of two-particle wavefunction(blue:t=0,Pink:t=3,violet:t=7)

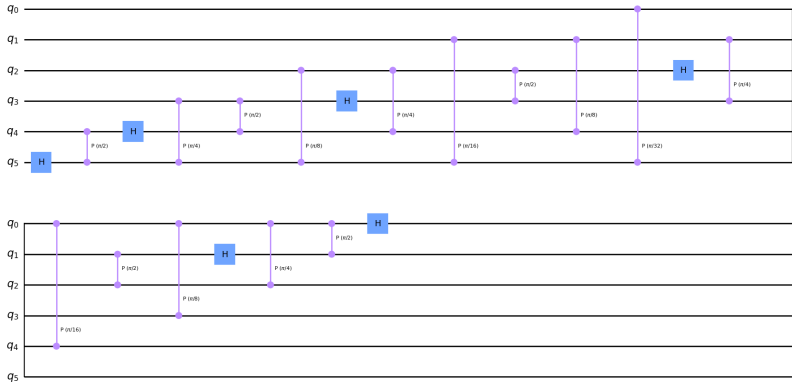


# References

- ① I. N. Michael Chuang, "Quantum computation and quantum information," 2005.
- ② M. Ostrowski, "Quantum simulation of two interacting schrödinger particles,"
- ③ S. Beneti Casati, "Principles of quantum computation and information," vol. 1,
- ④ M. Aboulela, "Quantum simulation of the Schrodinger equation using IBM's quantum computers,"

# Appendix

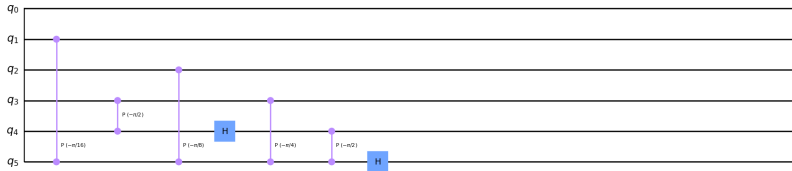
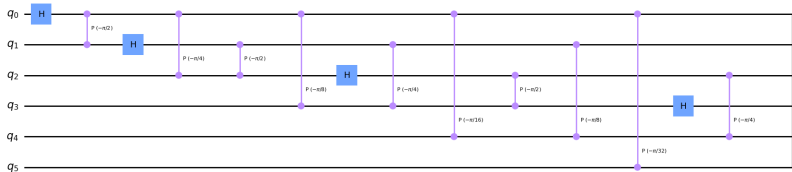
## 1. Quantum circuit for Quantum Fourier Transform for 6-qubit system





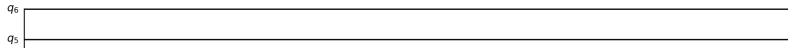
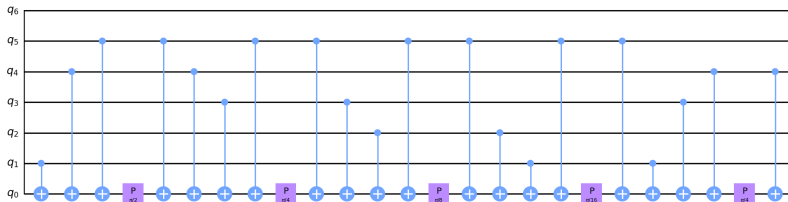
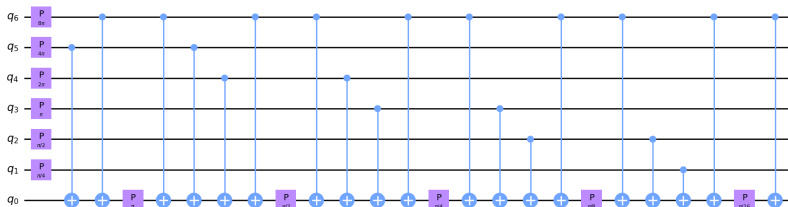
# Appendix

## 2. Quantum circuit for Inverse Quantum Fourier Transform for 6-qubit system

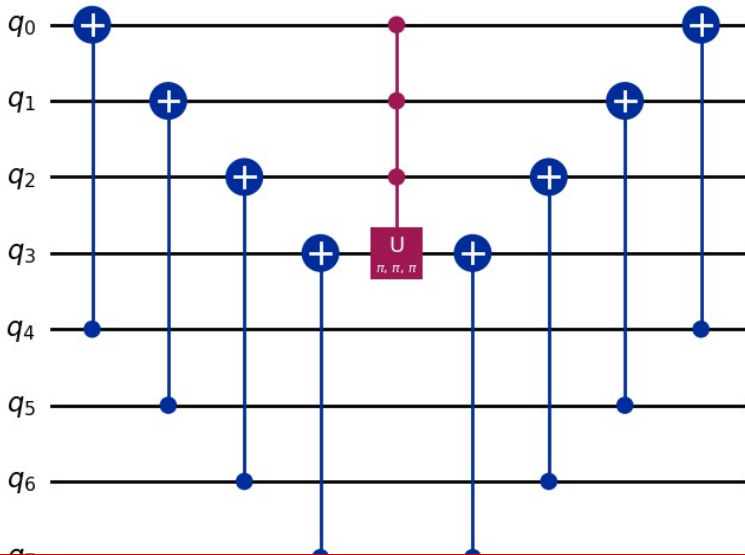


# Appendix

## 2. Quantum circuit for Kinetic Energy Operator for 6-qubit system



# Two Particle interaction potential



# Two Particle Reference

