

FLEXURAL WAVE LOCALIZATION IN STRUCTURED AND DISORDERED THIN PLATES

A DIRECT VALIDATION OF THE LANDSCAPE THEORY

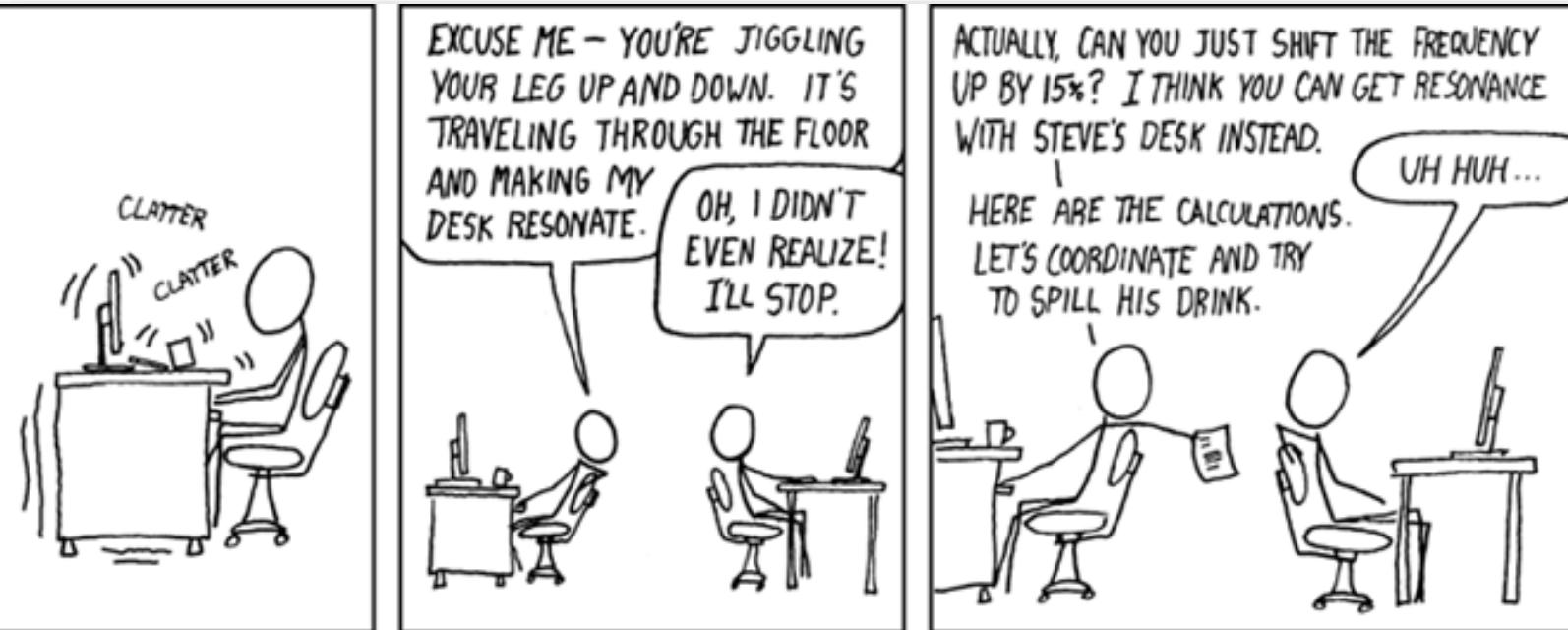
Patrick Sebbah

Department of Physics, Bar Ilan University, Ramat Gan, Israel

Formerly at

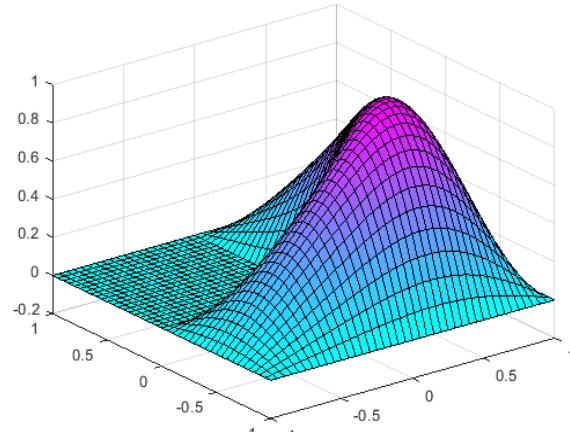
Institut Langevin, ESPCI ParisTech , CNRS, UMR 7587 Paris, France

THIS TALK ... IN A CARTOON



EIGENMODES AND EIGENVALUES

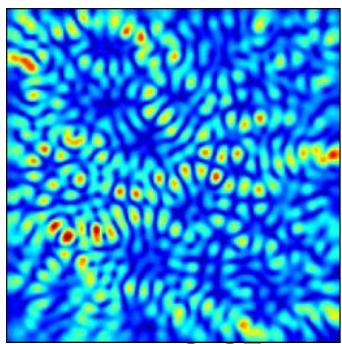
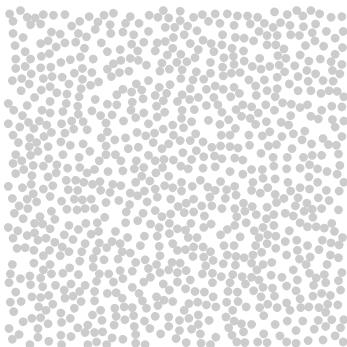
- Key ingredients to understand the physical properties of most vibratory systems



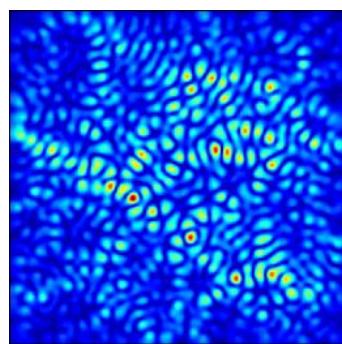
Eigenmode of L-Shaped Membrane

- Yet their prediction remains an open problem in many cases

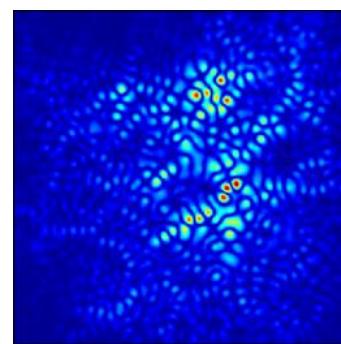
EXAMPLE 1 : DISORDERED SYSTEMS



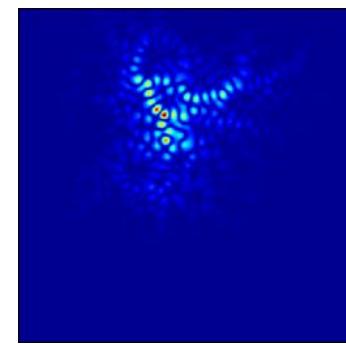
$\Delta n = 0.25$



$\Delta n = 0.50$



$\Delta n = 0.85$



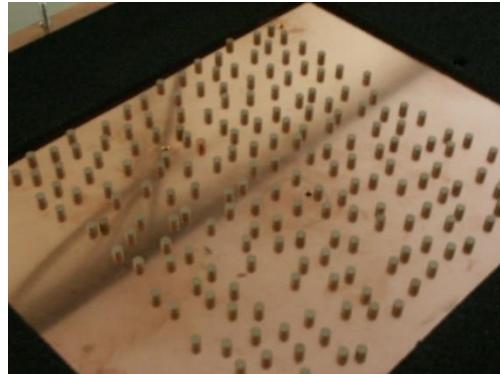
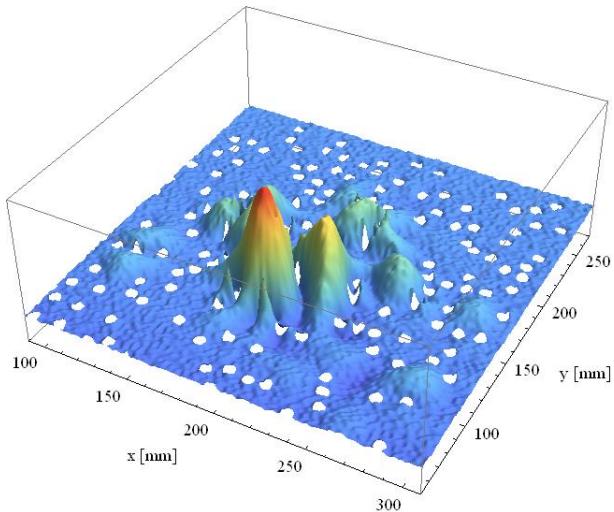
$\Delta n = 1.00$

Δn = index contrast between scatterers and matrix

DISORDERED-INDUCED LOCALIZATION

Anderson Localization

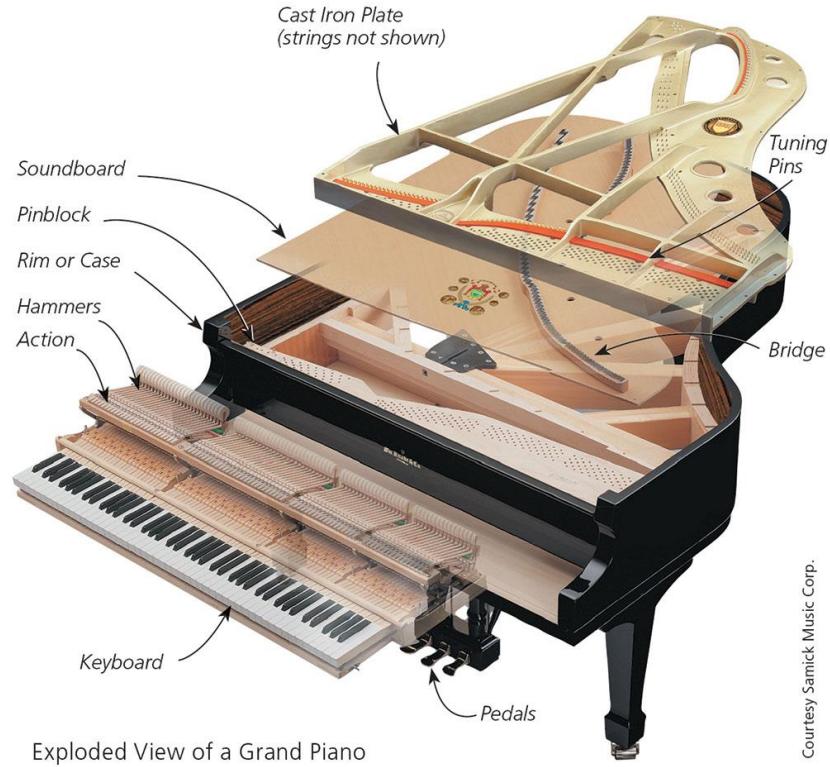
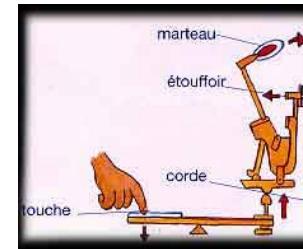
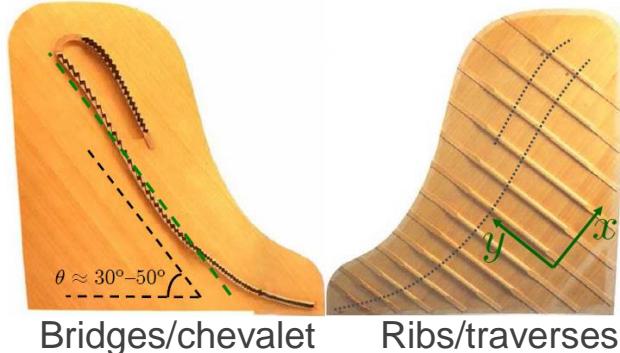
“Absence of diffusion
in certain random lattices”,
P.W. Anderson, *Phys. Rev* 109(1958).



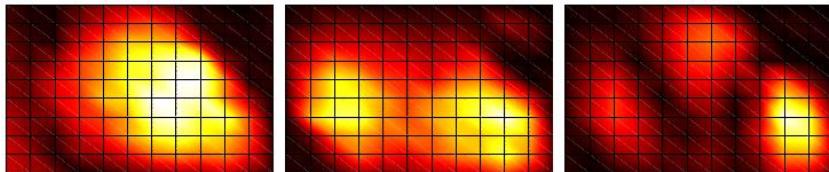
Laurent et al., PRL99, 253902 (2007)

EXAMPLE 2: ARCHITECTURE OF MUSIC INSTRUMENTS

The design of soundboard



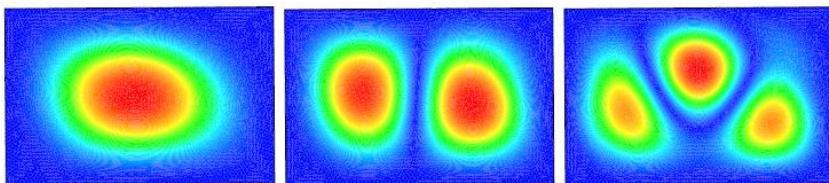
LOW AND HIGH FREQUENCY MODES



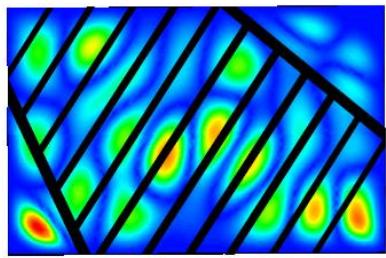
(a) (1,1)-mode

(b) (2,1)-mode

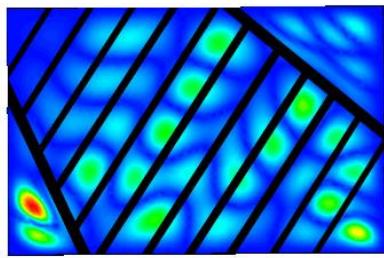
(c) (3,1)-mode



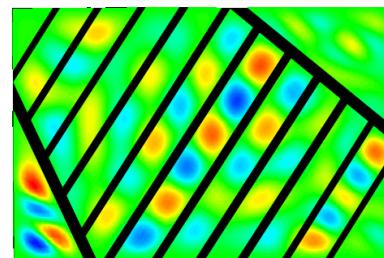
First 3 modal shapes of the upright piano soundboard :
measured and simulated



$f_n = 776 \text{ Hz}$



$f_n = 1089 \text{ Hz}$



$f_n = 1593 \text{ Hz}$

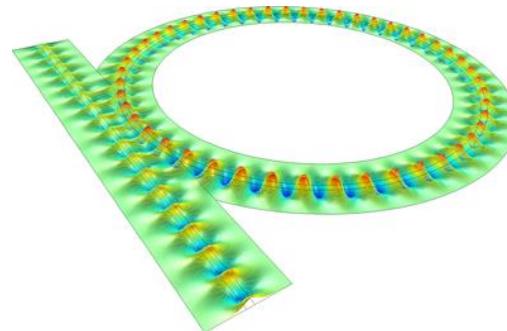
Intensity (and field) distribution of higher modes of the upright piano soundboard

OTHER STRING INSTRUMENTS



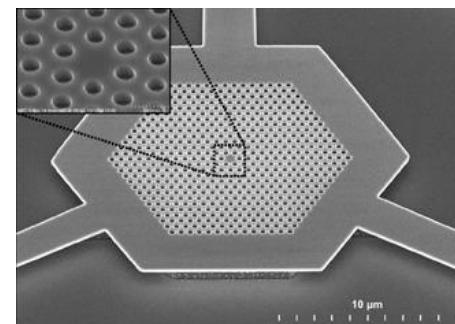
OTHER EXAMPLES

- Optical resonators



- Laser cavities

- microcavity lasers
- external cavity semiconductor laser
- dfb laser
- random lasers
- ...

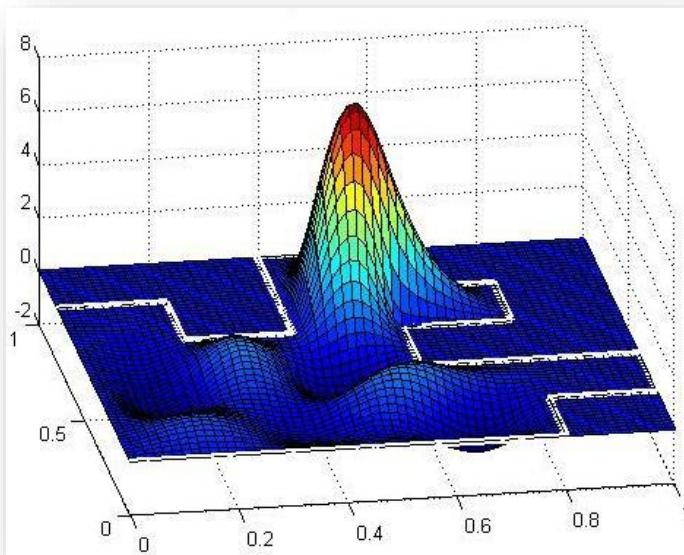


- Microwave resonators

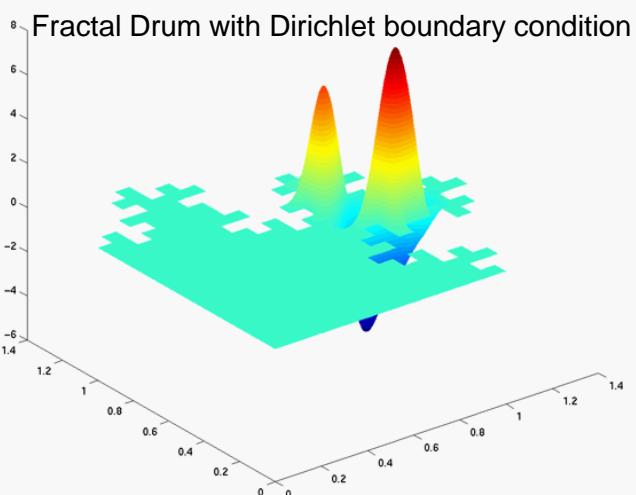
- ...

WAVE LOCALIZATION : A BROAD DEFINITION

A state is said to be “localized” if the vibration amplitude is large only in a restricted region of the total domain



WAVE LOCALIZATION IN COMPLEX STRUCTURES



Félix et al., J Sound Vib, 2007



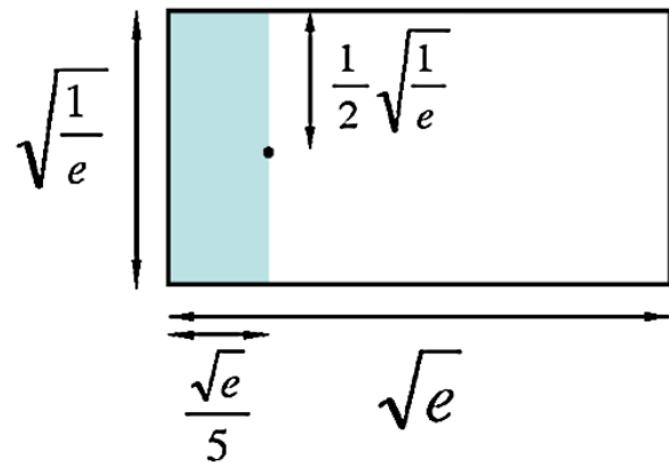
Heilman & Strichartz, Not Am Math Soc, 2010

"Localized eigenfunctions : Here you see them, there you don't"

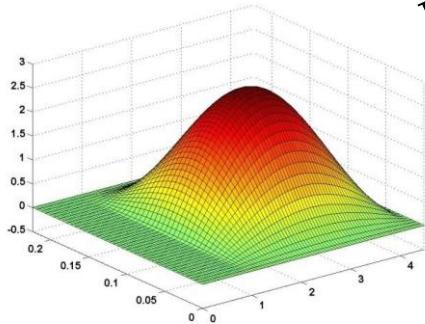
Applications:

- Structural engineering
- Noise abatement walls
- Clamped plate problem
- control of vibrations
- Musical instrument design

WAVE LOCALIZATION IN A CLAMPED PLATE

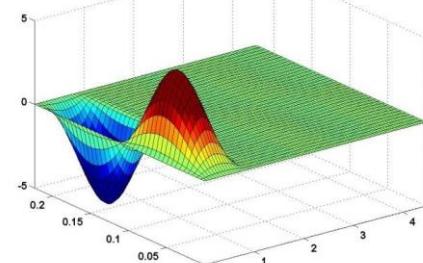


Mode 1

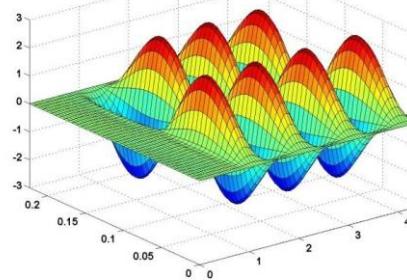


$\sqrt{e}=20$

Mode 40

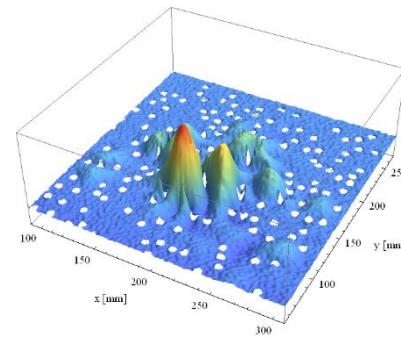


Mode 44



M.Filoche & Mayboroda, *PRL* 103, 254301 (2009)

QUESTIONS



- What is the effective potential that traps the waves ?
Are there “**invisible confining boundaries**”?
- Need for a **predictive tool** :
 - To retrieve the spatial and spectral properties of the localized modes, without resorting to a full eigenvalue problem
 - To predict the frequency dependence of the mode extension/distribution
- The **inverse problem**: Can we engineer plates with particular vibrational properties ?

THE LOCALIZATION LANDSCAPE

A new theoretical tool to retrieve crucial information on the spatial and frequency properties of these localized waves

THE LOCALIZATION LANDSCAPE

Elliptic operator L

$$L = H = -\Delta + V \rightarrow H\varphi = -i\hbar \frac{\partial \varphi}{\partial t}$$

$$L = -\Delta \quad \rightarrow -\Delta\varphi = -\frac{1}{c^2} \frac{\partial^2 \varphi}{\partial t^2}$$

$$L = -\Delta^2 \quad \rightarrow -\Delta^2\varphi = -\frac{1}{c^2} \frac{\partial^2 \varphi}{\partial t^2}$$

The properties of the system are entirely contained in the eigenfunctions and eigenvalues of the operator L

$$L\varphi_k = \lambda_k \varphi_k \quad \text{with } \varphi_k|_{\partial\Omega} = 0 \quad \lambda_k > 0$$

THE LOCALIZATION LANDSCAPE

Proposition :

for every $x \in \Omega$

$$|\varphi_k(x)| \leq \lambda_k u(x)$$

with normalization $\|\varphi_{k \max}\| = 1$

where u is the solution of the Dirichlet boundary problem:

$$Lu = 1 \text{ with } u|_{\partial\Omega} = 0$$

u is called the Landscape function the solution of

$$\forall x \in \Omega, \quad u(x) = \int_{\Omega} |G(x, y)| dy$$

u corresponds to the “static deformation” of the system under uniform load



THE LOCALIZATION LANDSCAPE

- The map of u is a landscape with many valleys and peaks.
- It partitions the domain into sub-regions.
- These sub-regions correspond to the localization regions of the different modes.

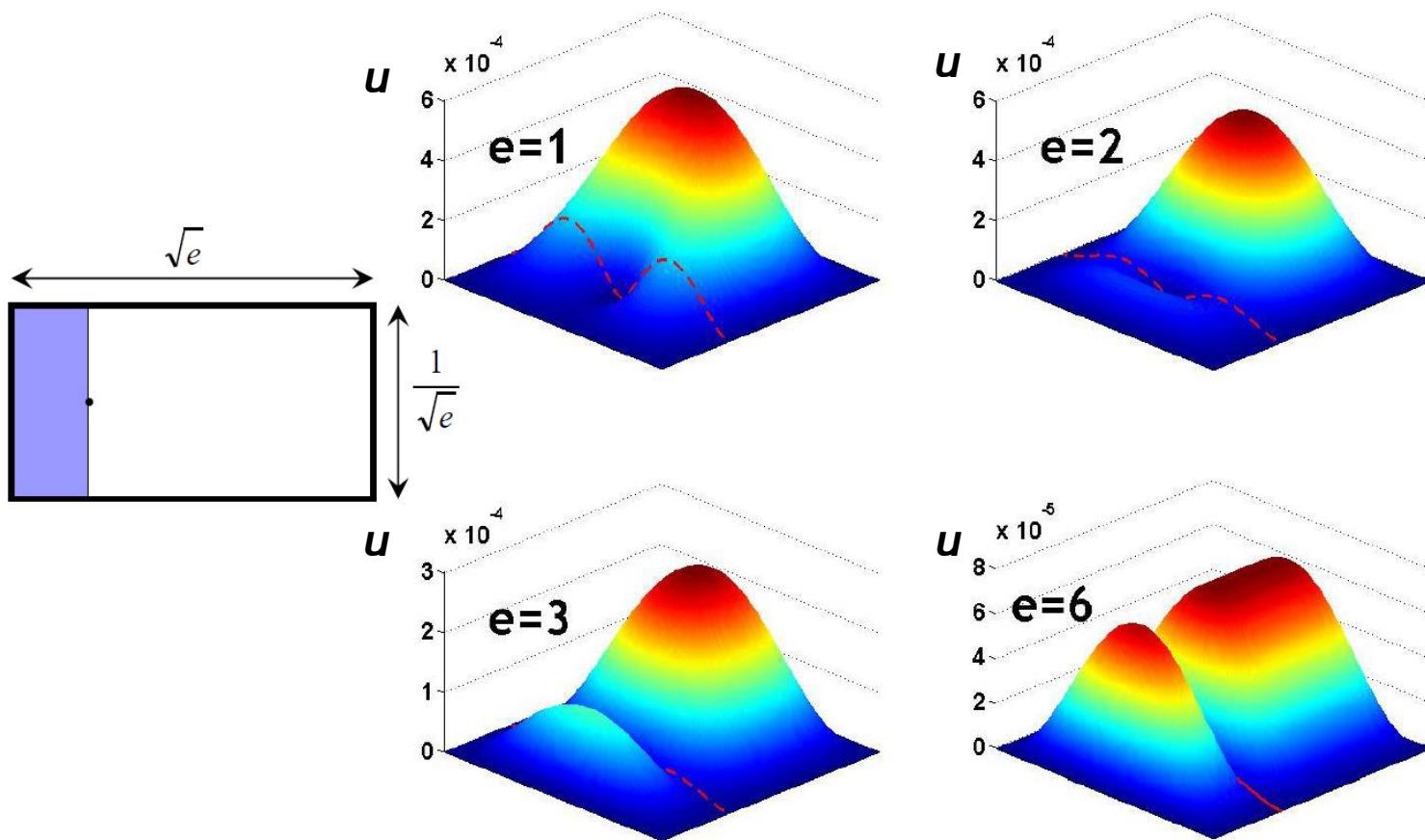
Indeed all eigenfunction φ satisfies :

$$\forall x \in \Omega, \quad |\varphi_k(x)| \leq \lambda_k u(x)$$

with normalization $\|\varphi_{k \max}\|=1$

The eigenmodes are “small” when u is “small”

THE LANDSCAPE IN THE CLAMPED PLATE



THE LOCALIZATION LANDSCAPE

- Describes the sub-regions of localization
- Gives, in a single measurement
 - the eigenfrequencies
 - The eigenmode profile
 - The frequency dependence of the mode spatial extension
- This is a static measurement, which does not require a precise knowledge of the microscopic geometry
- It can be used to address the inverse problem and design the system geometry corresponding that a desired spatial distribution of the modes

OUTLINE

- **Part 1: Elastic waves confined in a structured thin plate**
Direct demonstration of the predictive power of the localization landscape in a simple geometry
- **Part 2: Disordered-induced localization in thin plates**
Localized modes in
 - (a) randomly-pinned plates
 - (b) randomly-distributed blind holes in thin plates.

COLLABORATORS

- ⦿ **Marcel FILOCHE** : PMC, Ecole Polytechnique, CNRS
- ⦿ **Svitlana MAYBORODA** : School of Mathematics, University of Minnesota
- ⦿ **Michael ATLAN** : Institut Langevin, ESPCI-CNRS, Paris
- ⦿ **Post Docs** : Gautier LEFEBVRE (Langevin), Kun TANG (BIU), Eitam LUZ (BIU)
- ⦿ **PhD** : Olivier Xéridat (LPMC), Marc DUBOIS (Langevin)
- ⦿ **Undergraduates** : Alexane GONDEL (Ecole des Mines), Florian FEPPON, Aime LABBE, Camille GILLOT, Alix GARELLI, Maxence ERNOULT (Ecole Polytechnique)

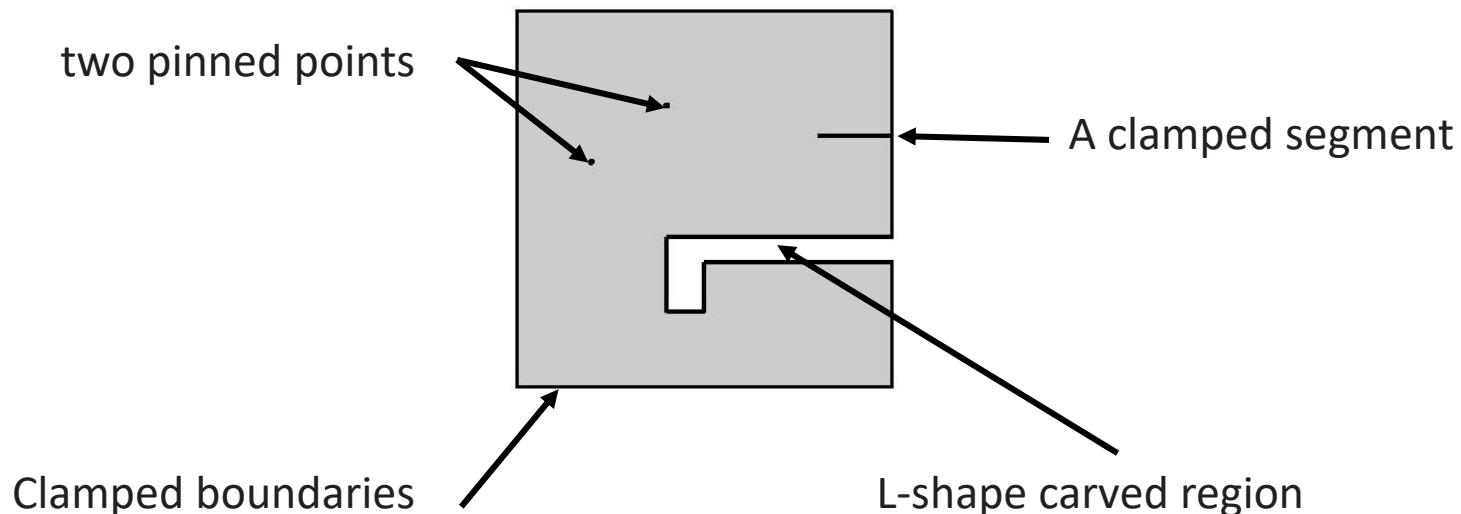
Supported by the *Agence Nationale de la Recherche (ANR PLATON 20012-2015)*, the *Labex WIFI (ANR-10-IDEX-0001-02 PSL*)*, *DGA*, *PICS-ALAMO (CNRS)*, *The Israel Science Foundation (ISF 1781/15, ISF 2074/15, ISF 2630/20)*, *The United States-Israel Binational Science Foundation (NSF/BSF 2015694)*.

PART 1: ELASTIC WAVES CONFINED IN A STRUCTURED THIN PLATE

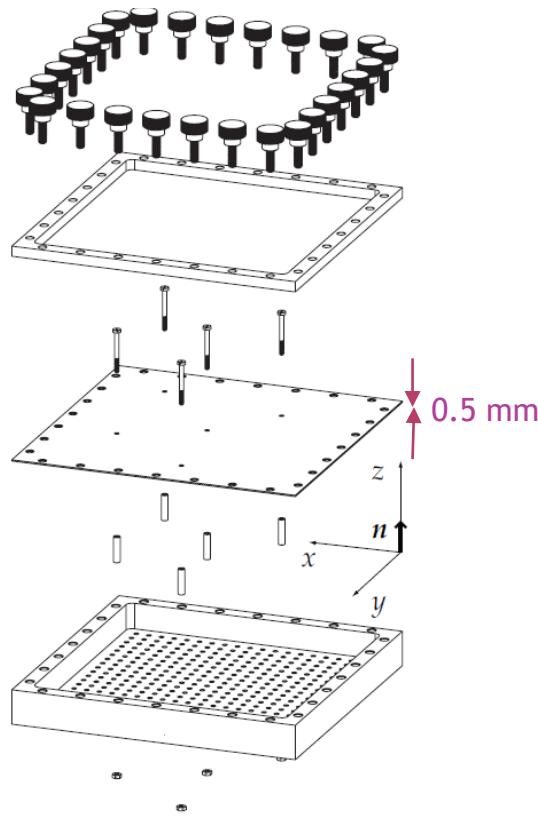
A direct demonstration of the predictive power of the localization landscape theory

G. Lefebvre *et al.*, Phys. Rev. Lett. (2017)

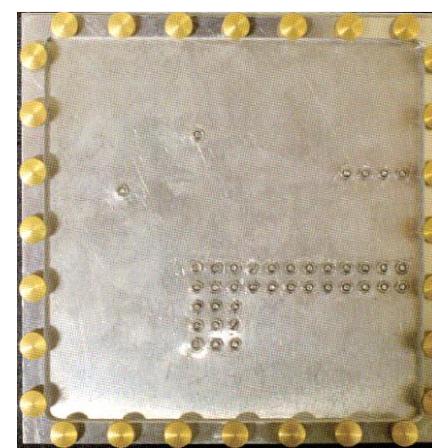
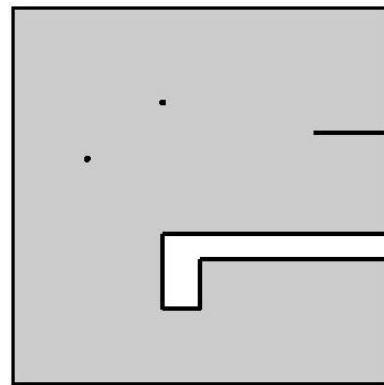
A STRUCTURED SYSTEM



EXPERIMENTAL REALISATION

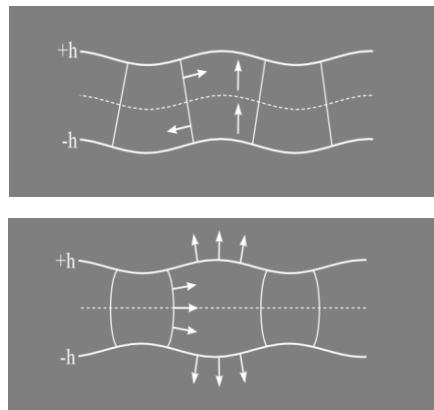
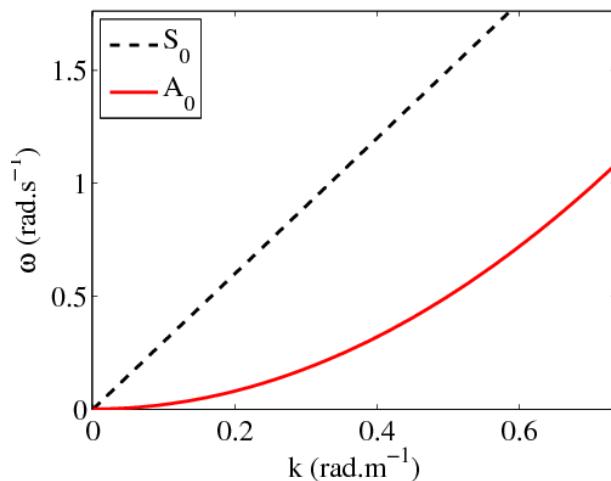


Exploded view



THIN PLATE EQUATION

Lamb waves



A_0 Vertical displacement

S_0

Low frequency approximation: Flexural waves

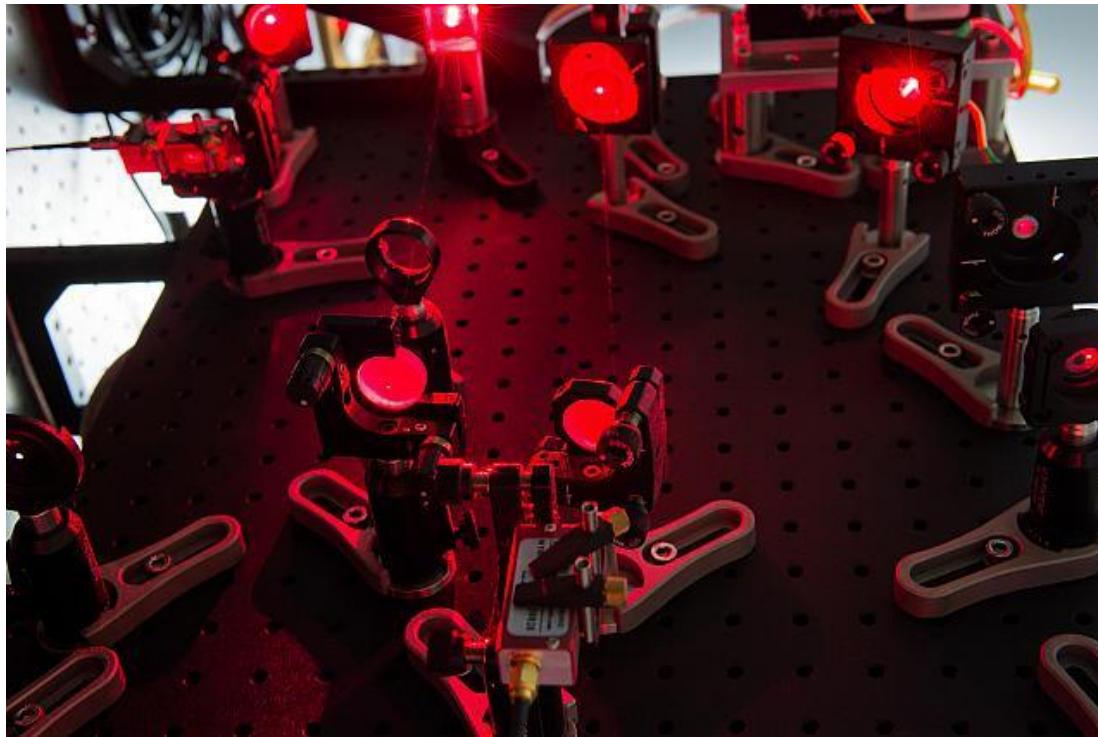
Kirchhoff-Love equation for vertical displacement u :

$$\Delta^2 u + \frac{\rho h}{D} \frac{\delta^2 u}{\delta t^2} = 0, \quad \omega^2 = \frac{D}{\rho h} k^4$$

$$\text{with } D = \frac{Eh^3}{12(1-\nu^2)}$$

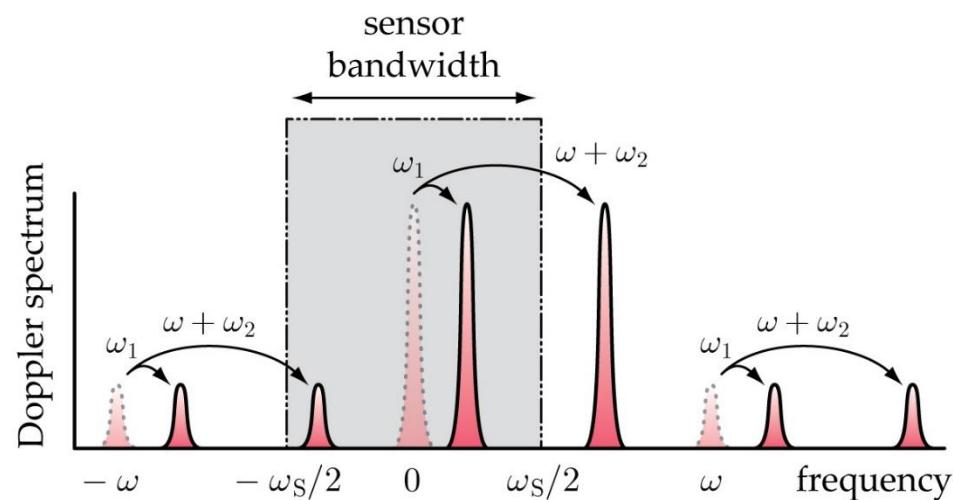
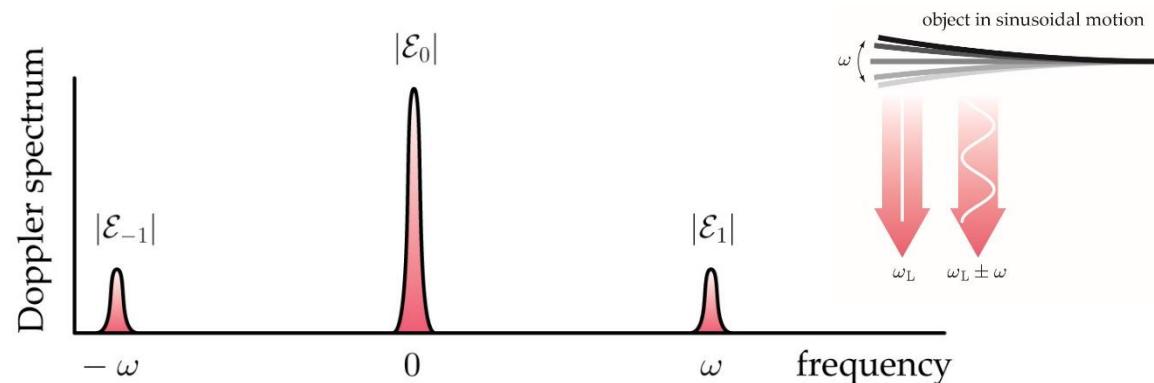
HETERODYNE HOLOGRAPHIC IMAGING

Wide field imaging of out of plane vibration vs. frequency

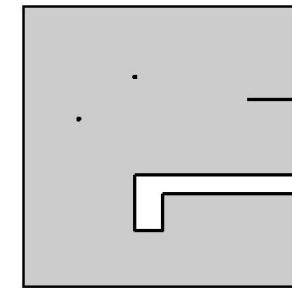


M. Atlan *et al.* Opt. Lett. 39, 2014

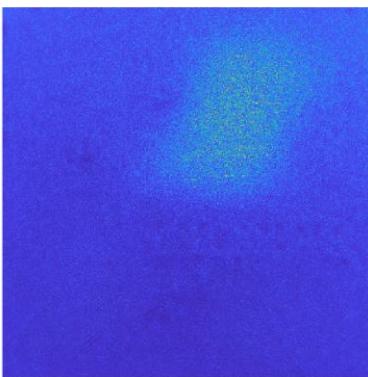
DOUBLE SIDEBAND MODULATION



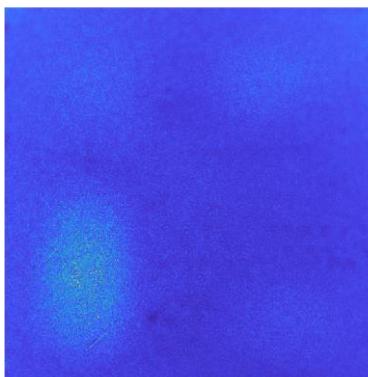
FIRST MODES: INTENSITY



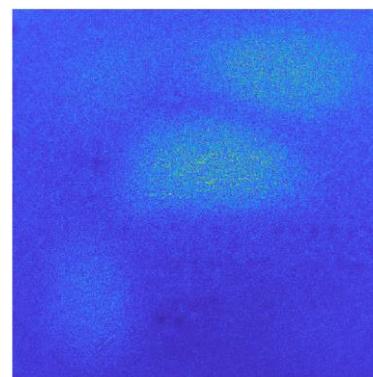
Mode 1 $f = 1790$ Hz



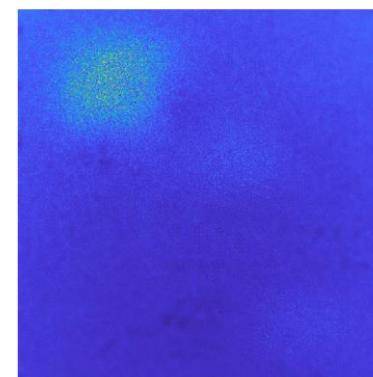
Mode 2 $f = 2160$ Hz



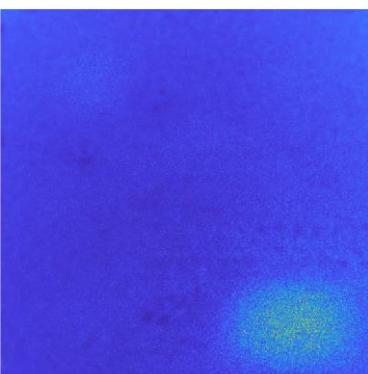
Mode 3 $f = 2430$ Hz



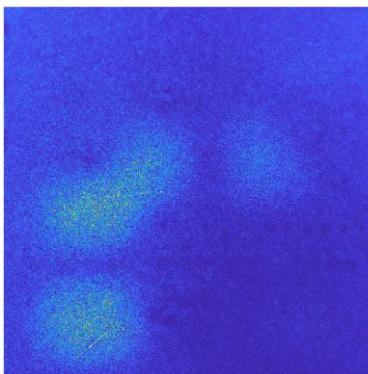
Mode 4 $f = 2580$ Hz



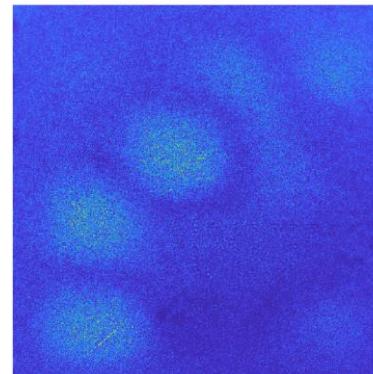
Mode 5 $f = 2820$ Hz



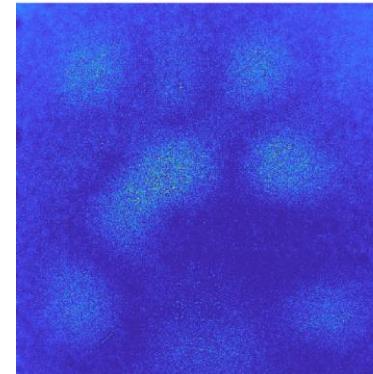
Mode 6 $f = 3040$ Hz



Mode 7 $f = 3410$ Hz



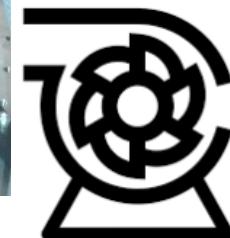
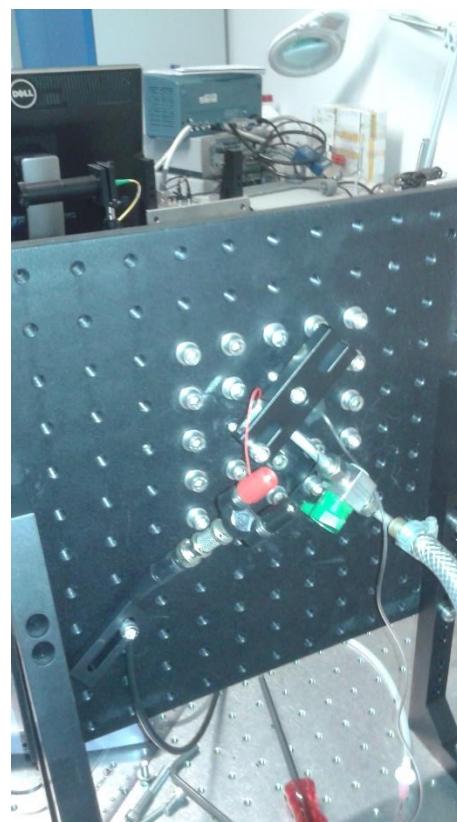
Mode 8 $f = 3850$ Hz



OBJECTIVE

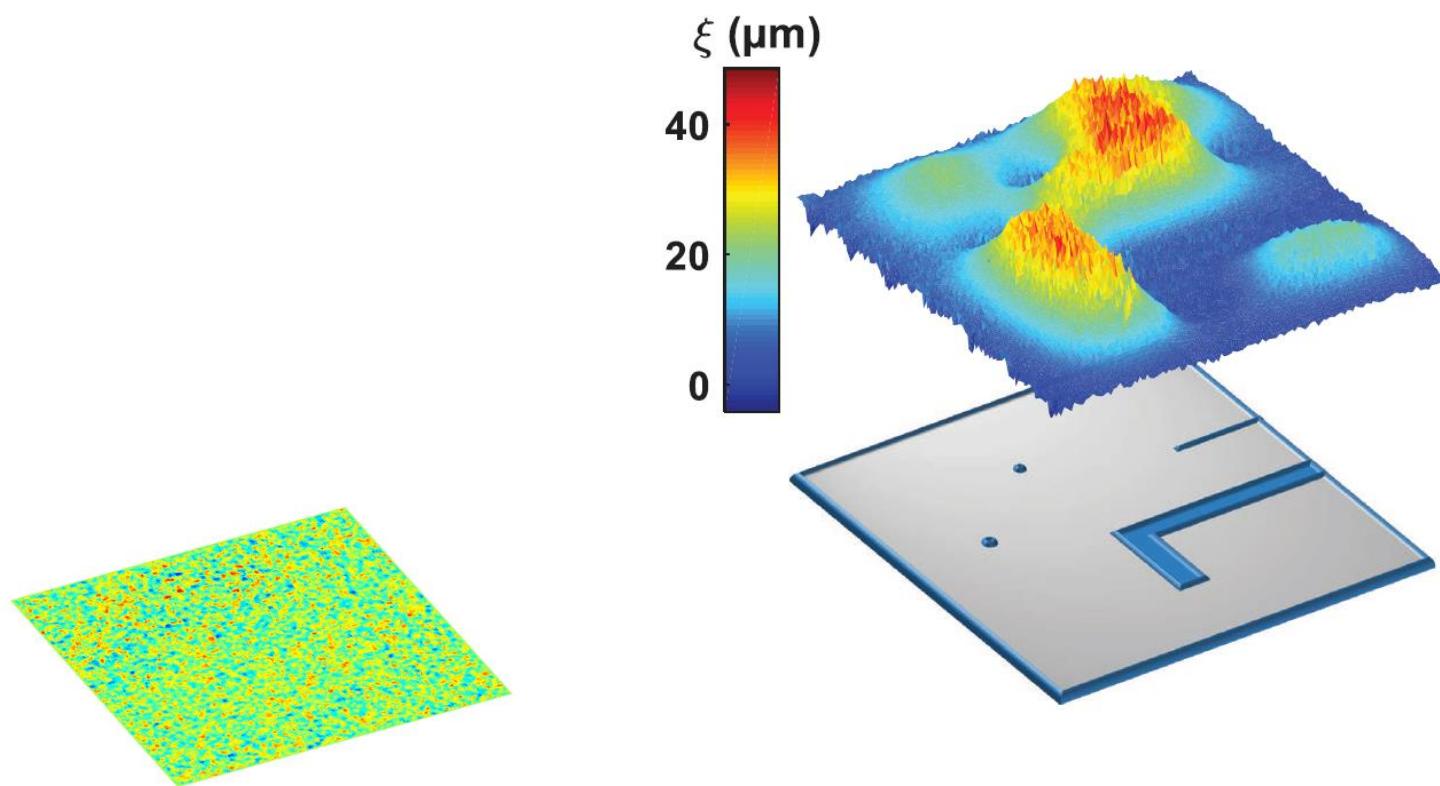
- Can we measure the *localization landscape* ?
- Does the theory apply to a real system ?

MEASURING THE LANDSCAPE



Measure of the static deformation using holographic interferometry as low pressure is progressively released

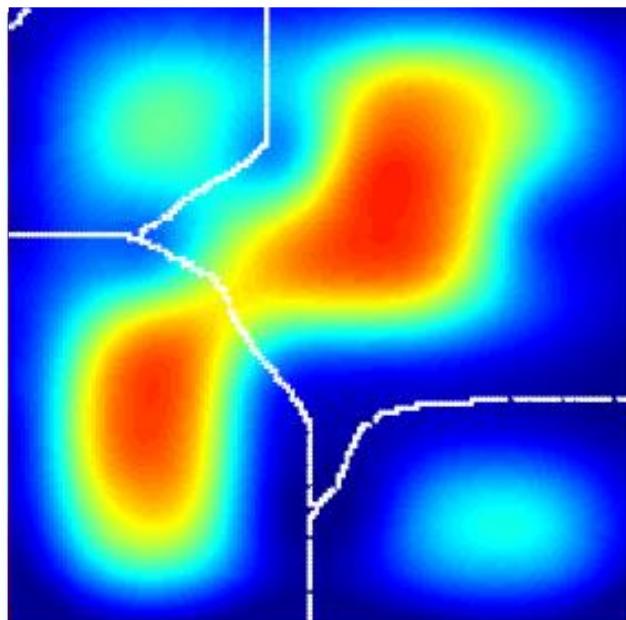
MEASURING THE LANDSCAPE



G. Lefebvre *et al.*, Phys. Rev. Lett. (2017)

1 - LOCALIZATION REGIONS

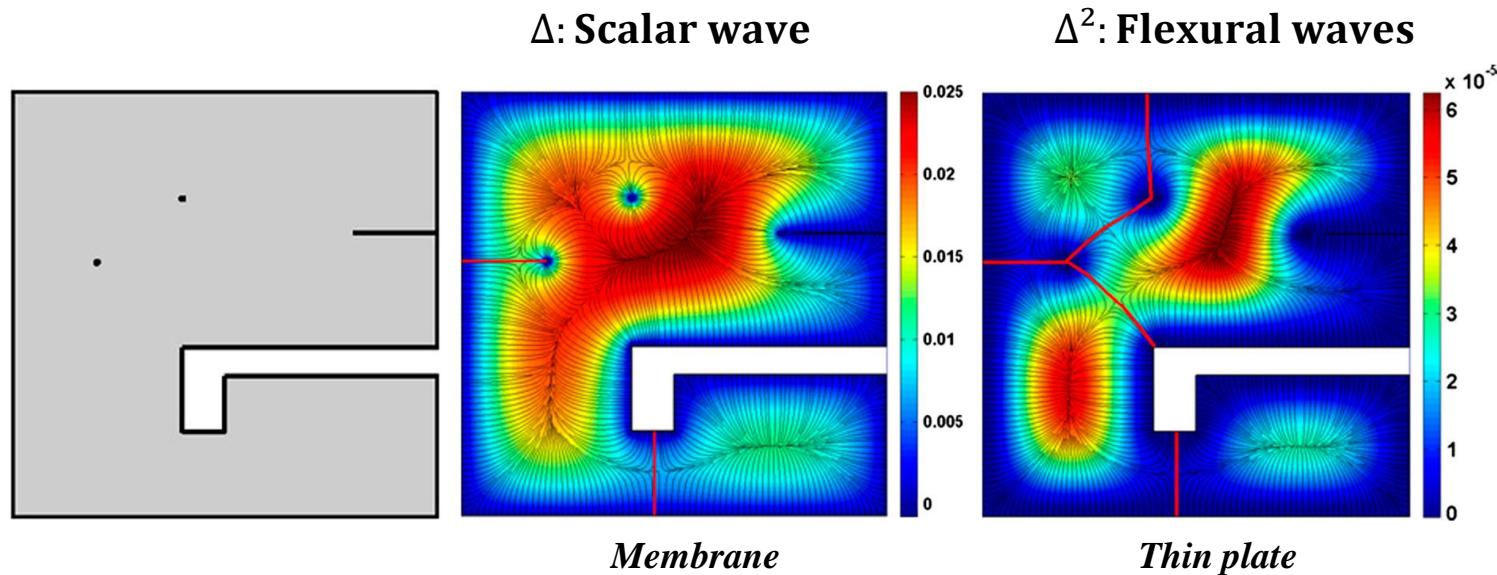
The map of u is a landscape with many valleys and peaks



The white lines are the watershed lines of
steepest ascent or steepest descent

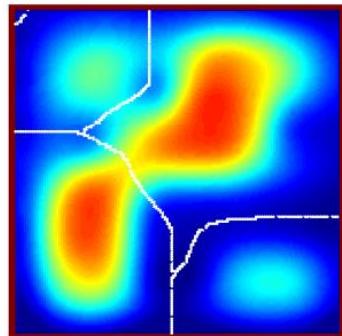
LANDSCAPES FOR DIFFERENT OPERATORS

The nature of the operator is embedded in the landscape function

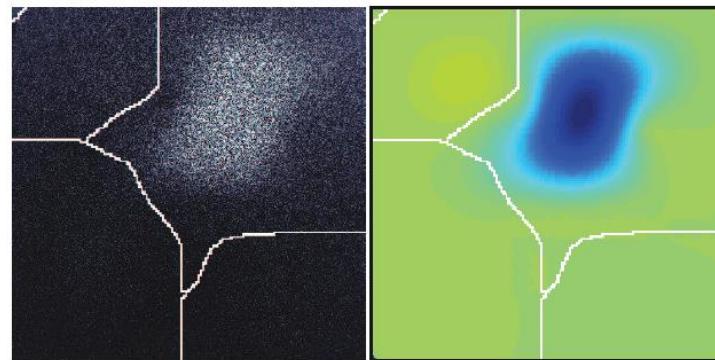


1 - LOCALIZATION REGIONS

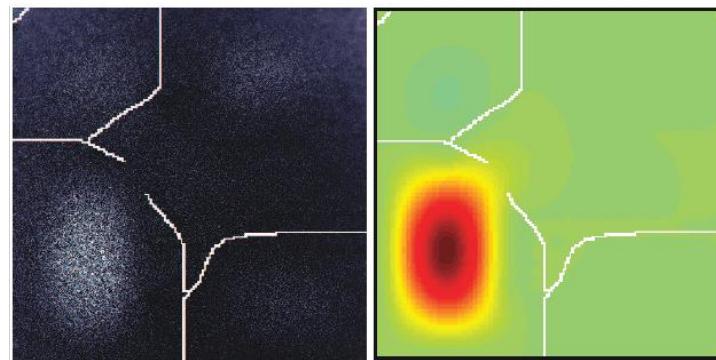
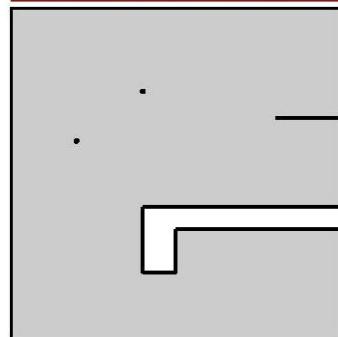
$u(x)$



Mode 1 $f = 1790$ Hz



Mode 2 $f = 2160$ Hz

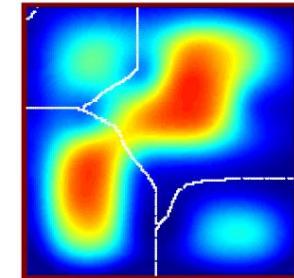


Experiment

Simulation

1 - LOCALIZATION REGIONS

These sub-regions correspond to the localization regions of the different modes.



All eigenfunctions φ satisfies :

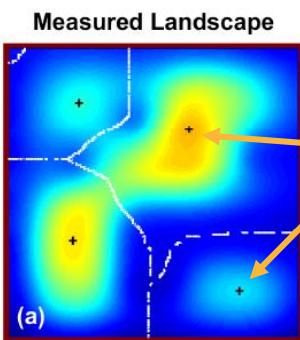
$$\forall x \in \Omega, \quad |\varphi_k(x)| \leq \lambda_k u(x)$$

with normalization $\|\varphi_{k \max}\| = 1$

The eigenmodes are “small” where u is “small”

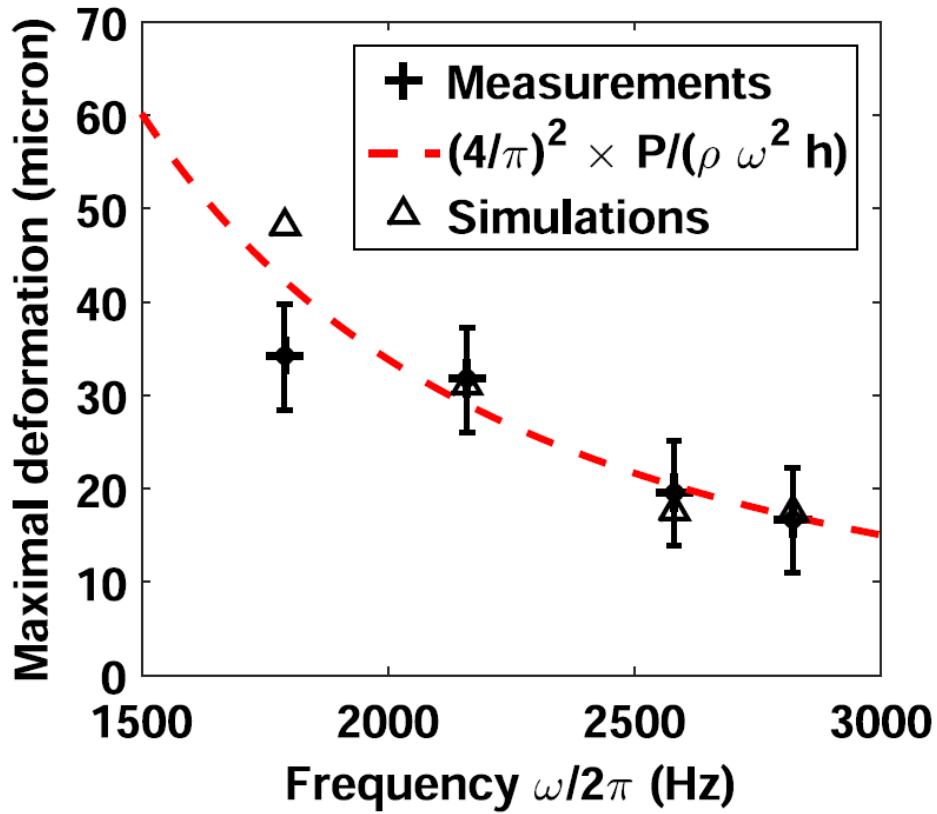
The “valleys” of $u(x)$ are the invisible barriers for waves

2 - EIGENVALUE PREDICTION



$\max_i(u) = \frac{1.33^2}{\omega_i^2}$ where ω_i is the frequency
of the fundamental mode of region i

2 - EIGENVALUE PREDICTION



$$\max_i(u) = \frac{1.33^2}{\omega_i^2}$$

3 - EFFECTIVE VALLEY NETWORK

$$|\varphi_k(x)| \leq \lambda_k u(x)$$

Because φ_k is normalized ($\|\varphi_{k \max}\|=1$), the constraint exists only if

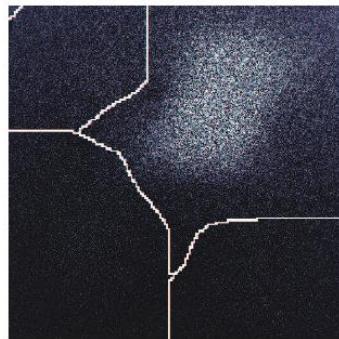
$$u(x) \leq \lambda_k^{-1}$$

which is referred to as the **effective valley network**

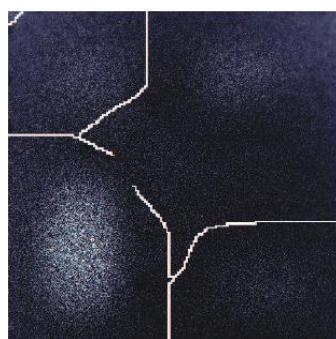
While the landscape u is **uniquely defined**,
the valley network **evolves** with the growth of λ .

4 - EXPERIMENTAL TEST OF THE VALLEY NETWORK OPENING

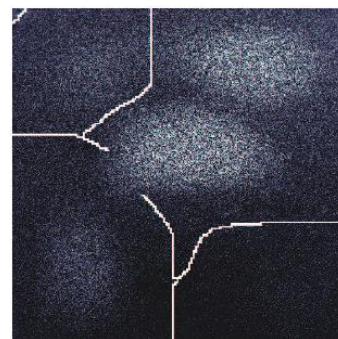
Mode 1



Mode 2



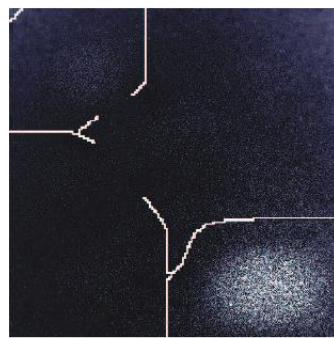
Mode 3



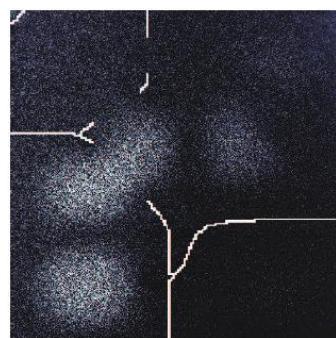
Mode 4



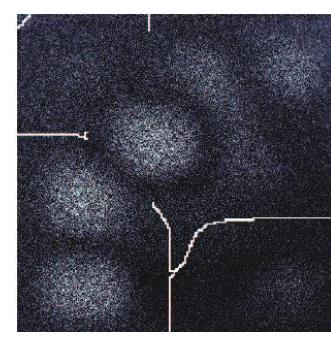
Mode 5



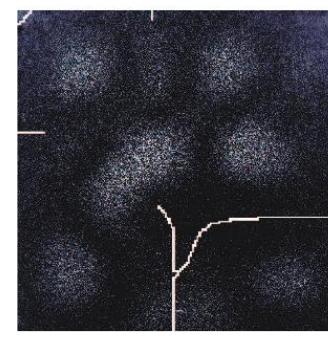
Mode 6



Mode 7



Mode 8

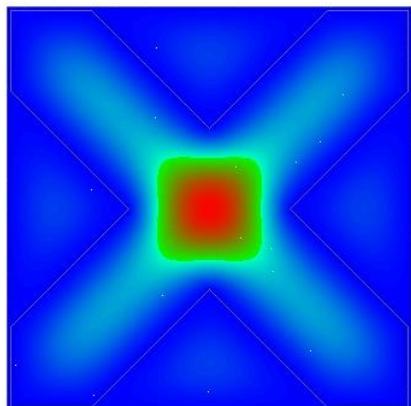


As frequency grows, the valleys “disappear” and the sub-regions begin to merge opening “free” passages for propagation.
At some point, **states are fully delocalized**.

SUMMARY

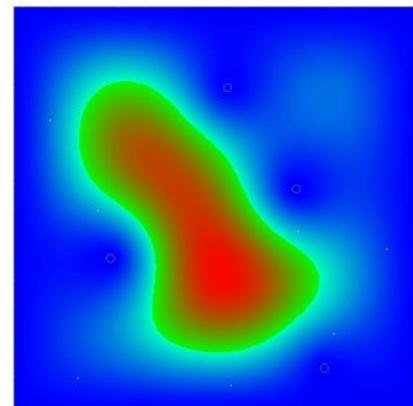
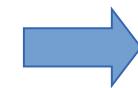
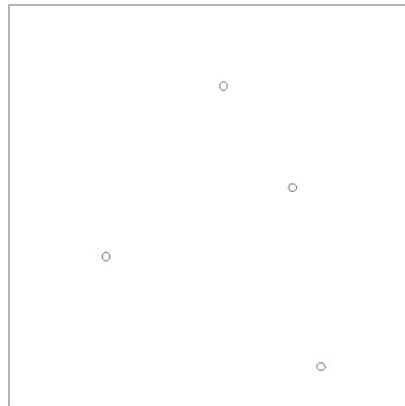
- The *localization landscape* describes the sub-regions of localization
- The local maxima of the *localization landscape* provide with the eigenvalues.
- The effective valley networks gives the frequencies at which gap open and modes extend in the system.

SOLVING THE INVERSE PROBLEM



u^*

Define an ideal
landscape function u^*



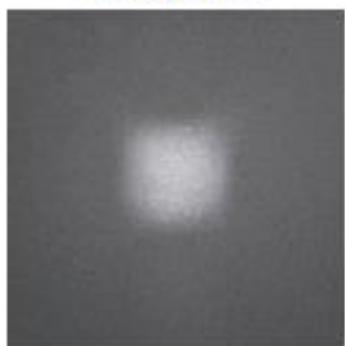
u



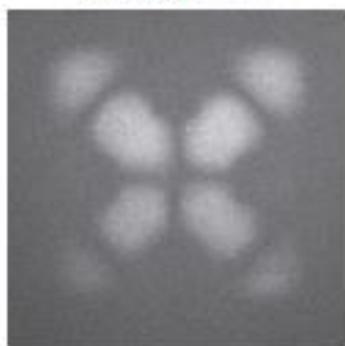
Minimize $||u - u^*||$

PLATE WITH 16 PINS

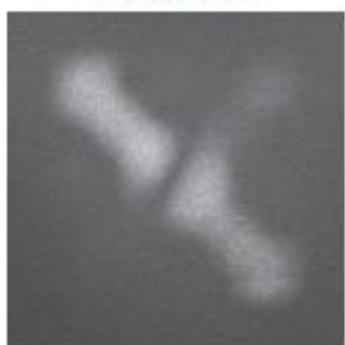
1650 Hz



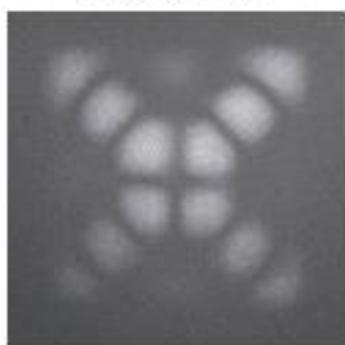
3730 Hz



2950 Hz



5540 Hz



CONCLUSION

- First experimental validation of the theory of the *localization landscape*
- The localization landscape has been directly measured.
- The essential information about localized states is retrieved without resorting to a full eigenvalue problem
 - one simple static measurement provides a rather good picture of the low frequency behavior of the system
 - The opening frequencies of the valleys are accurately predicted

PART 2: DISORDERED-INDUCED LOCALIZATION IN THIN PLATES

- I- Localized modes in randomly-pinned plates
- II- Localization by resonant scatterers

MOTIVATION

- Exploring Anderson localization in 2D disordered media
 - Observing localized states, mode coupling, necklace states, exceptional points
 - Transition from localized to extended modes
 - Dynamics of Localization
- Trapping elastic waves
 - Energy confinement and control
 - Understanding and controlling the acoustic behavior of complex structures
 - Application to noise reduction and wave control
- Predicting localized states of vibration :
Structural engineering
 - Application to instrument design
 - design of structures with predefined modal characteristics (Optical cavities, lasers)

OBJECTIVES

We explore two different routes to localize elastic waves in disordered thin plates

- 1 - Localization in a randomly pinned plate



- 2 - Localization by resonant scatterers



I - LOCALIZATION IN RANDOMLY-PINNED PLATE

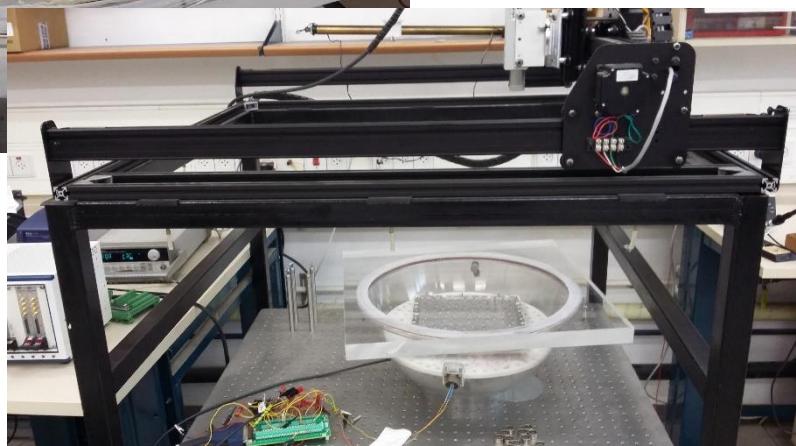
RANDOMLY PINNED PLATE



23cmx23cm
100 pins, 3mm-diam
Absorbing boundary
conditions

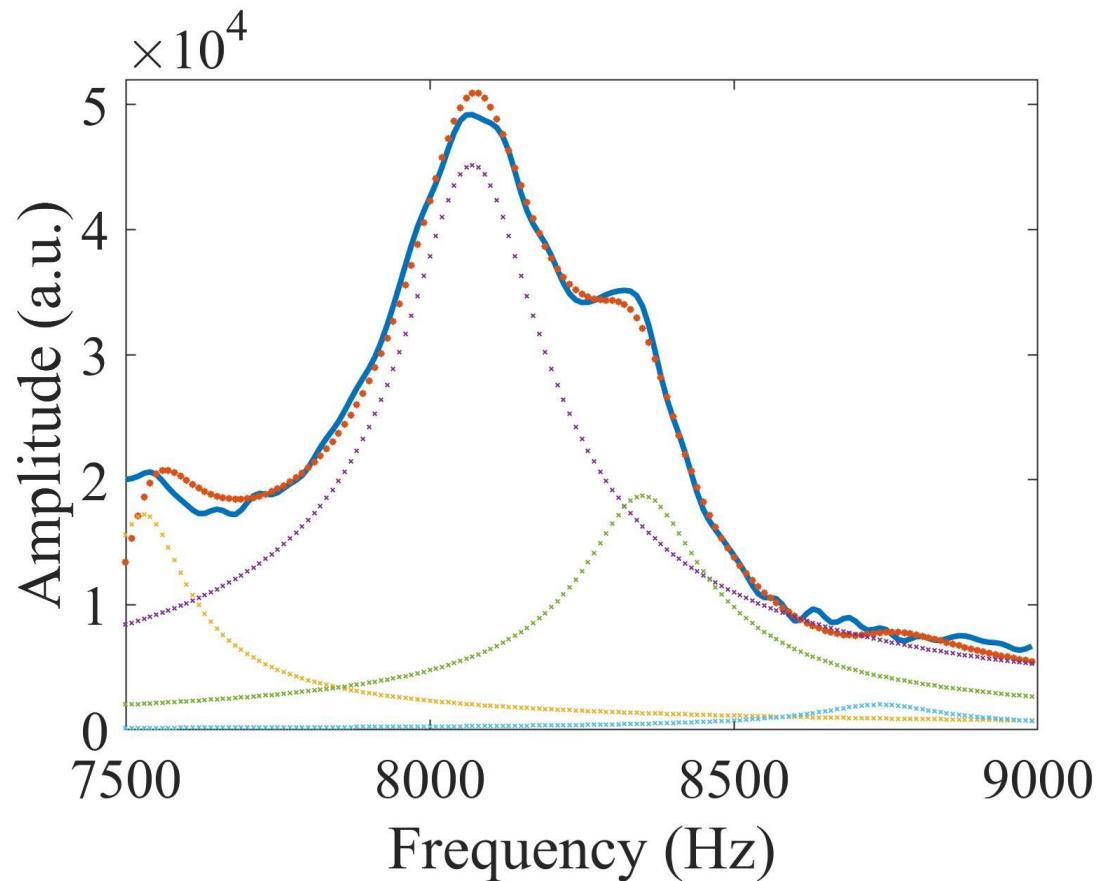


Vacuum chamber



2D-scan of a laser
vibrometer probe

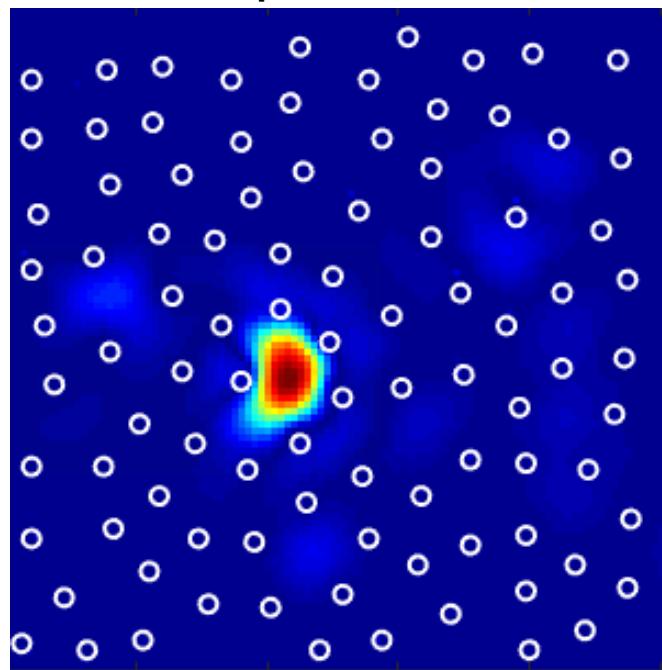
AMPLITUDE SPECTRUM



Spectral decomposition as
a sum of complex Lorentzian functions

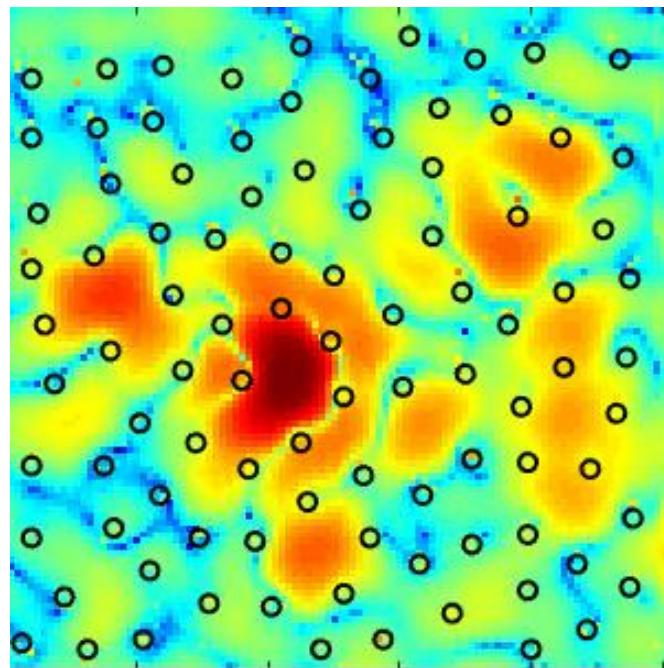
SPATIAL DISTRIBUTION

freq=7940 Hz



SPATIAL DISTRIBUTION

freq=7940 Hz



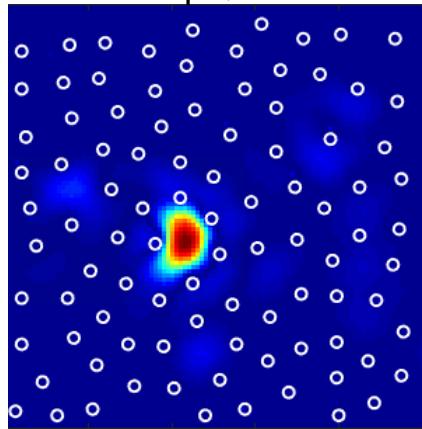
Logarithmic scale

SPATIAL DISTRIBUTION

Linear scale

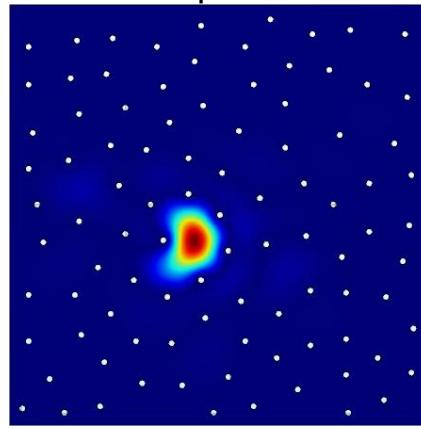
experiment

freq=7940 Hz

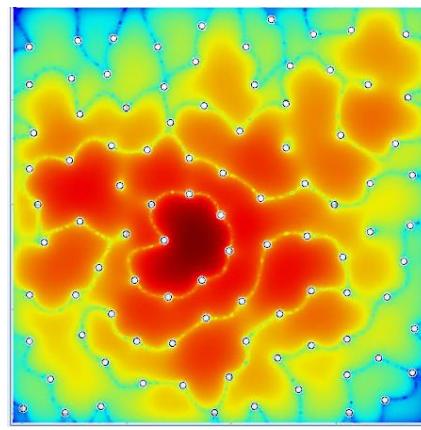
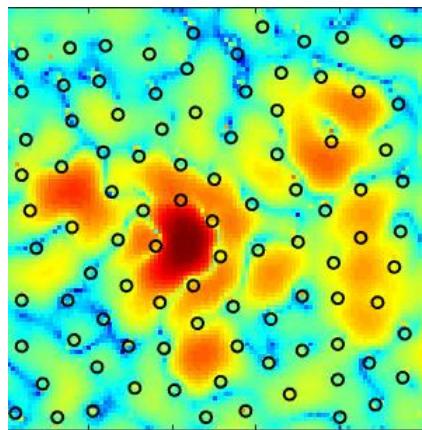


simulations

freq=7556 Hz



Log scale

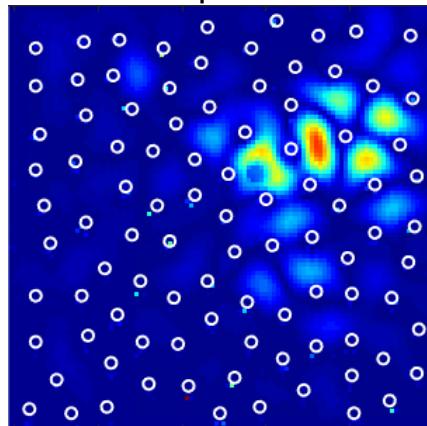


SPATIAL DISTRIBUTION

Linear scale

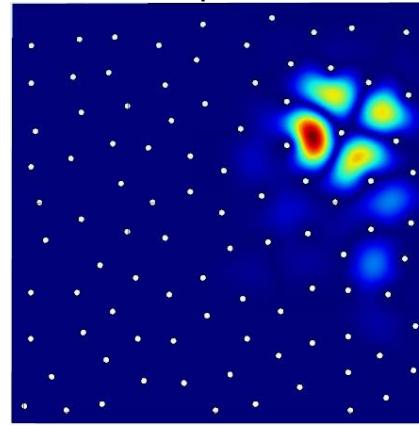
experiment

freq=10440 Hz

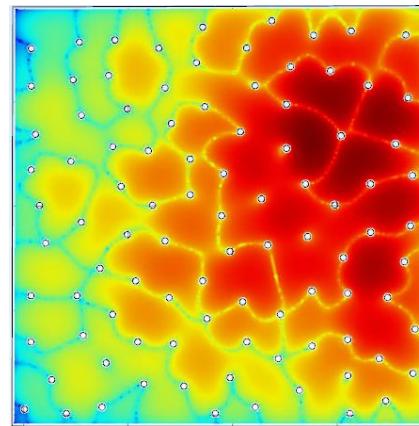
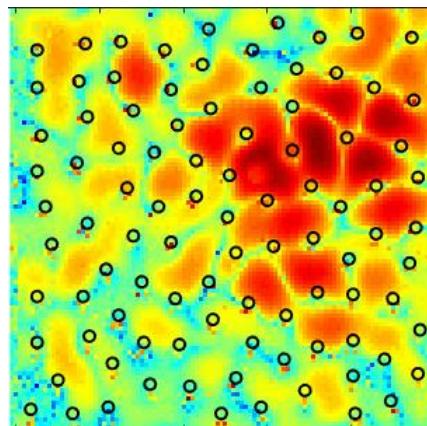


simulations

freq=9415 Hz

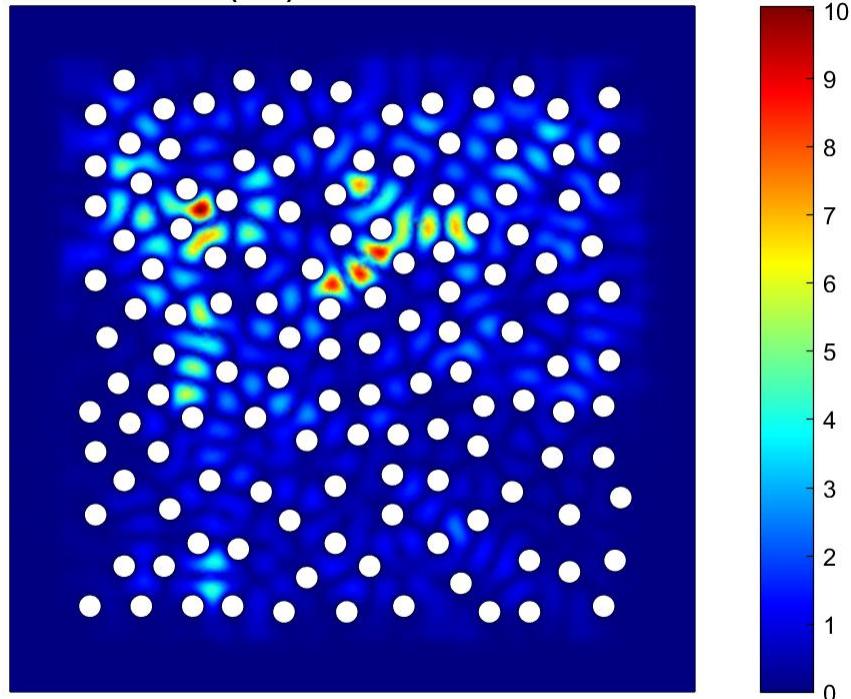


Log scale

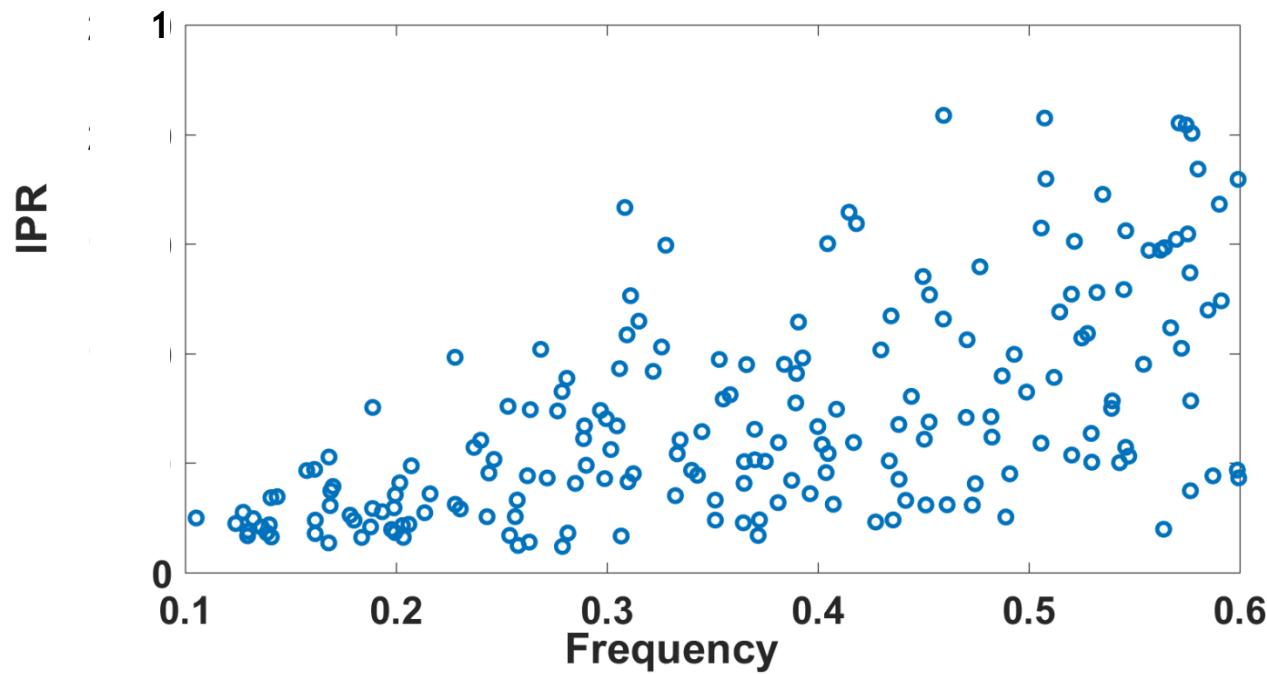


MODE EXTENSION

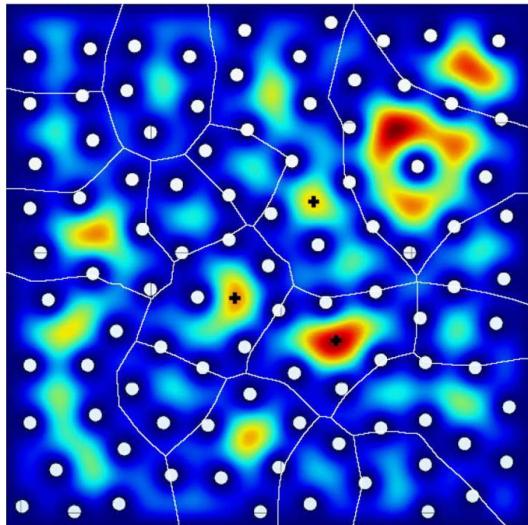
$\lambda(500)=0.60151+0.0018458i$



INVERSE PARTICIPATION RATIO



THE LOCALIZATION LANDSCAPE



Computed
landscape function u

The landscape function u is the solution of

$$Lu = \mathbf{1} \text{ with } u|_{\partial\Omega} = 0$$

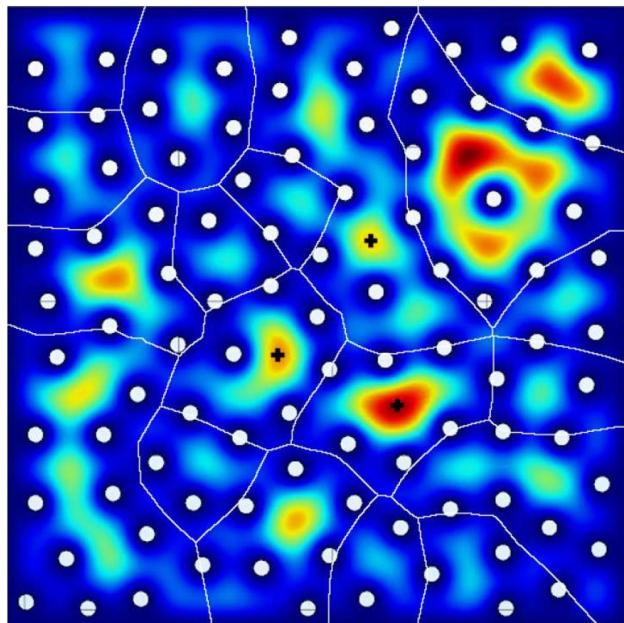
u corresponds to the “static deformation” of the system under uniform load

The regions delimited by the valleys correspond to the localization regions of the different modes.

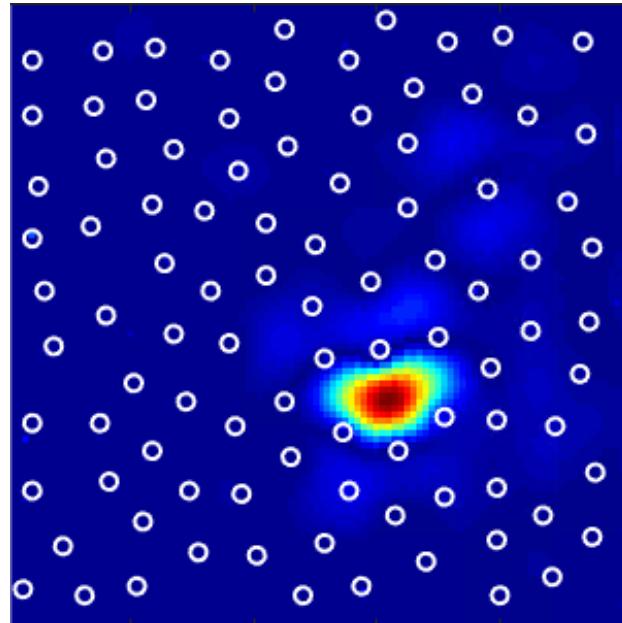
Local maxima give the inverse of the resonant frequencies

LANDSCAPE VS. MODES

landscape function

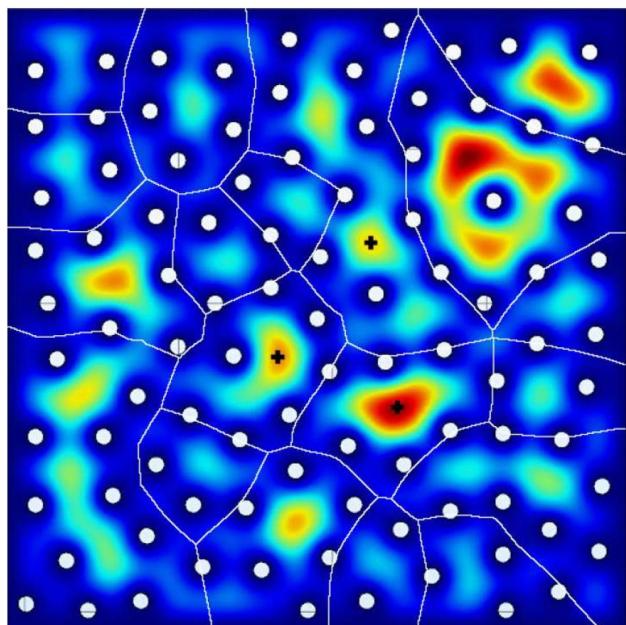


freq=6612 Hz

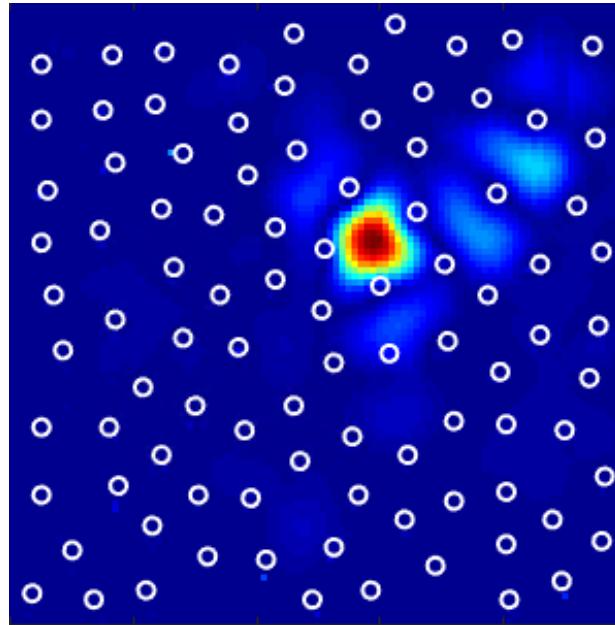


LANDSCAPE VS. MODES

landscape function

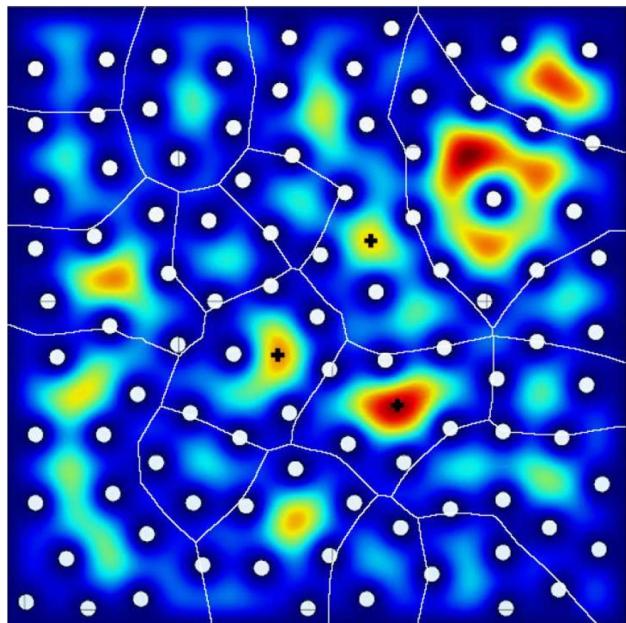


freq=7510 Hz

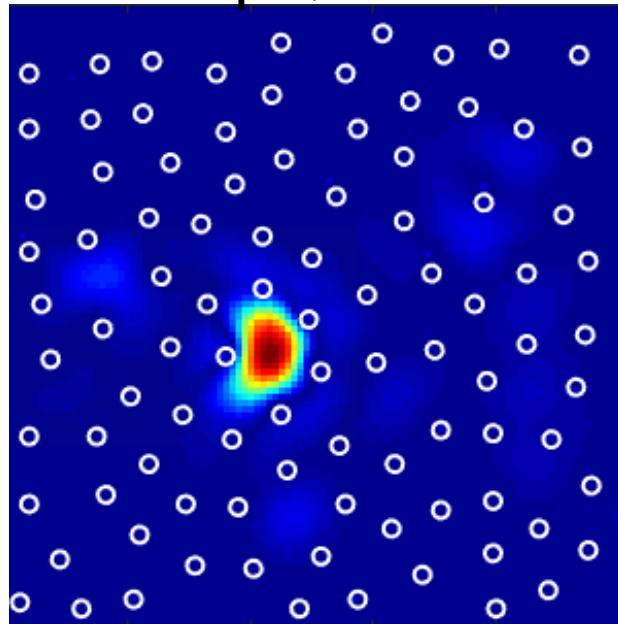


LANDSCAPE VS. MODES

landscape function

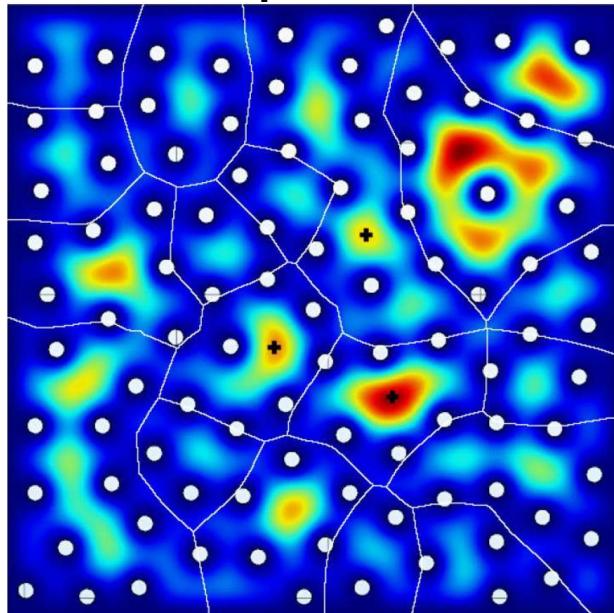


freq=7940 Hz

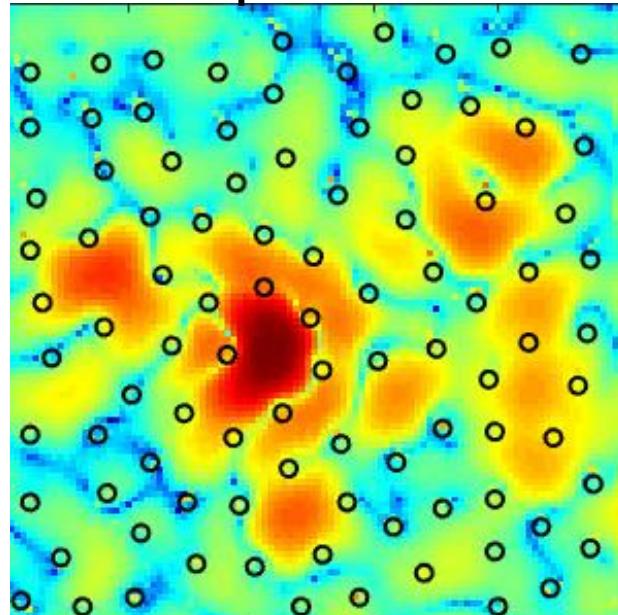


LANDSCAPE VS. MODES

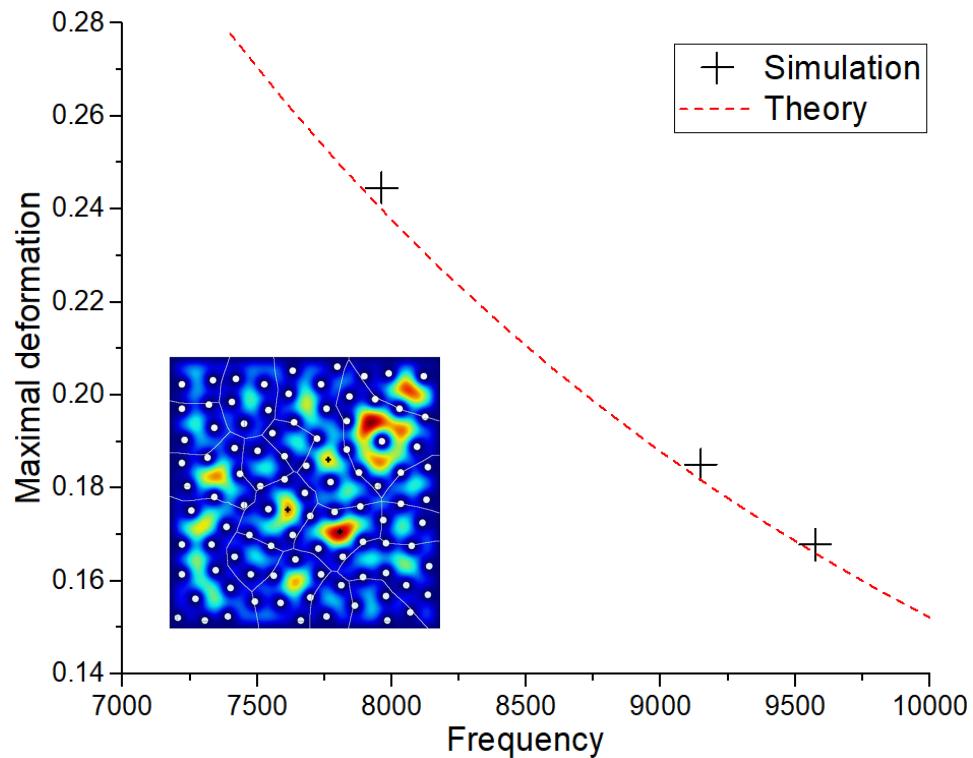
landscape function

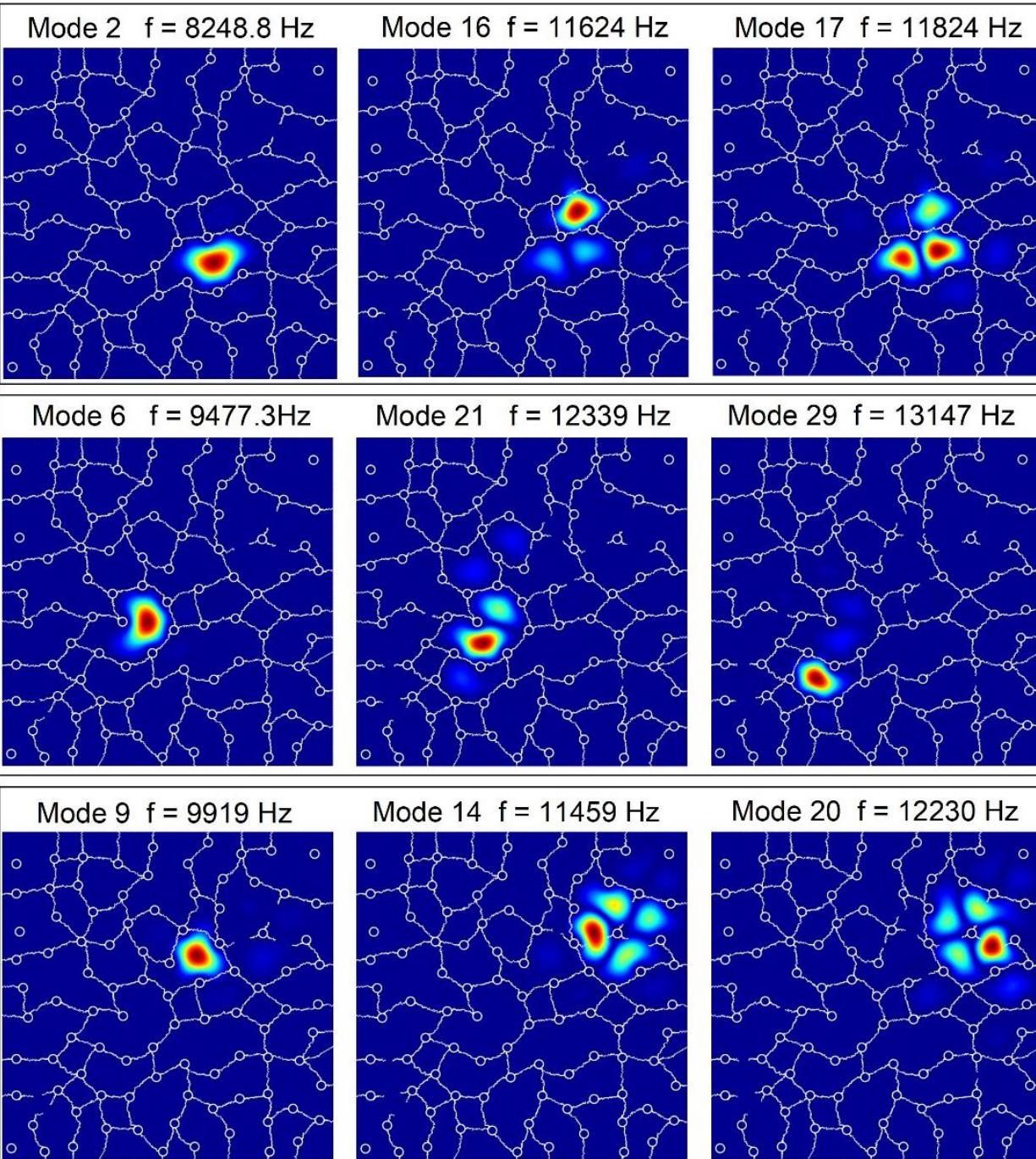


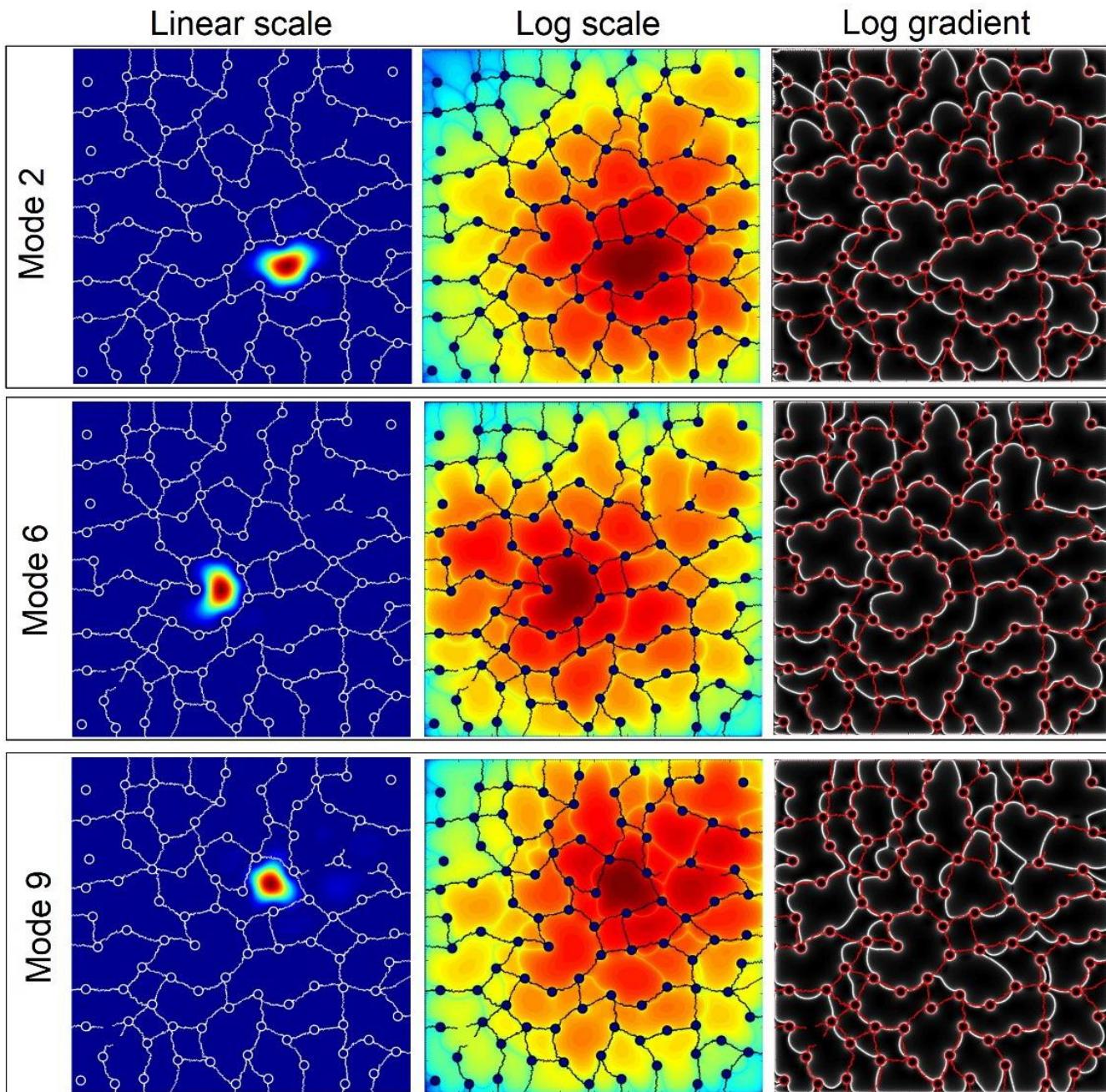
freq=7940 Hz



PREDICTION OF THE LOW FREQUENCY SPECTRUM





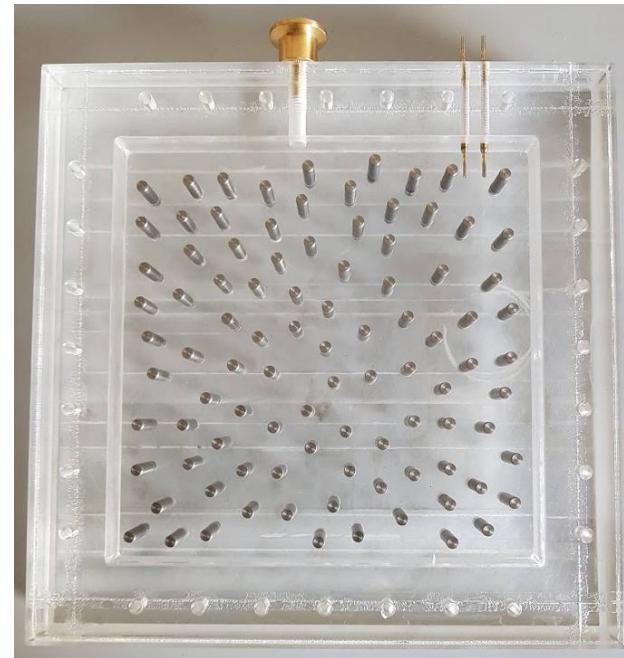
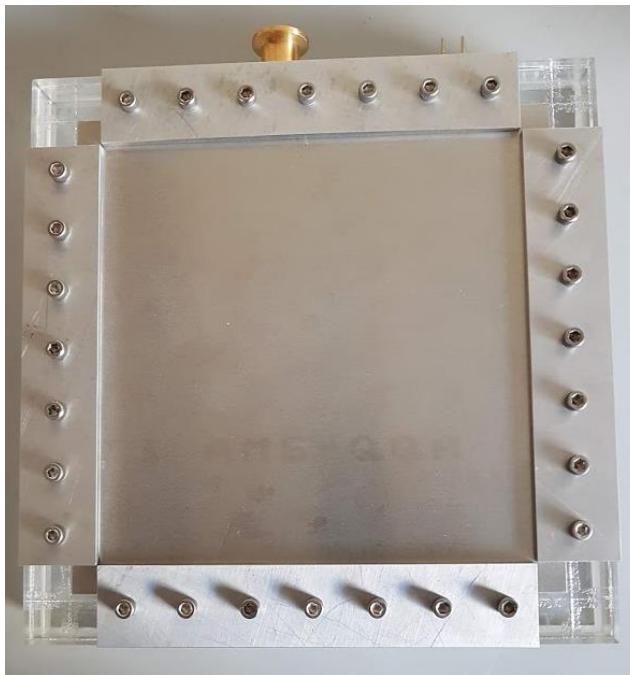


IN SUMMARY

- Disordered-induced localized modes have been measured in a randomly-pinned plate
- The extension of the modes increases with frequency
- The landscape describes the eigenmodes
 Gives their position, described how they decay, gives the low frequency spectrum

COMING NEXT ...

- Measuring the landscape in the pinned plate



II - LOCALIZATION BY RESONANT SCATTERERS

MUTIPLE SCATTERING OF ELASTIC WAVES



Available online at www.sciencedirect.com



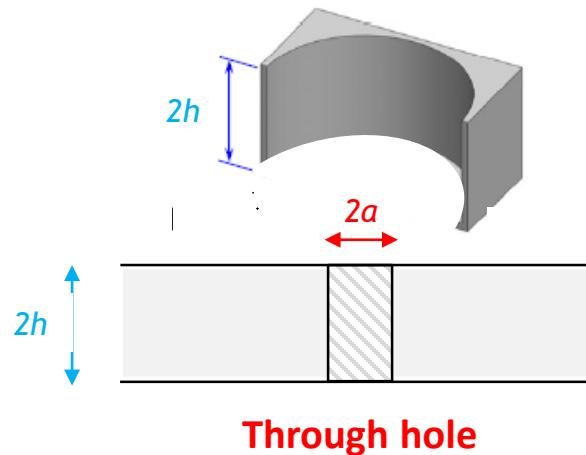
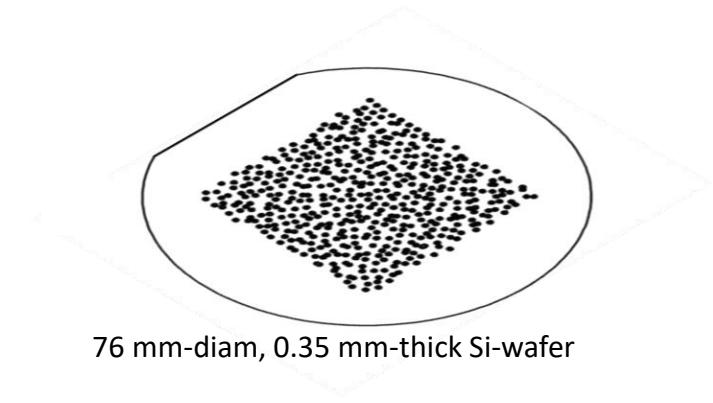
Ultrasonics 42 (2004) 775–779

Ultrasonics

www.elsevier.com/locate/ultras

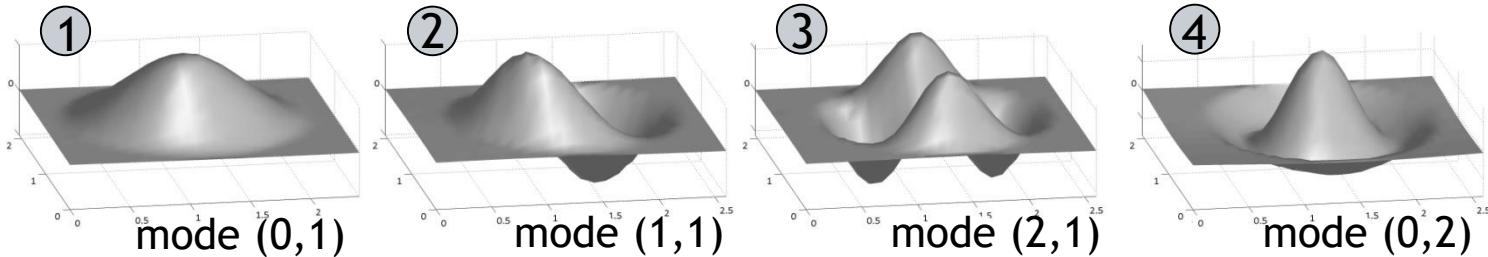
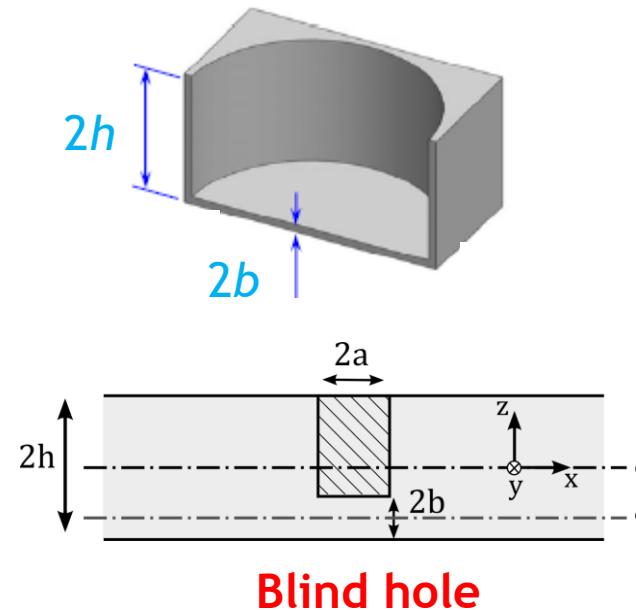
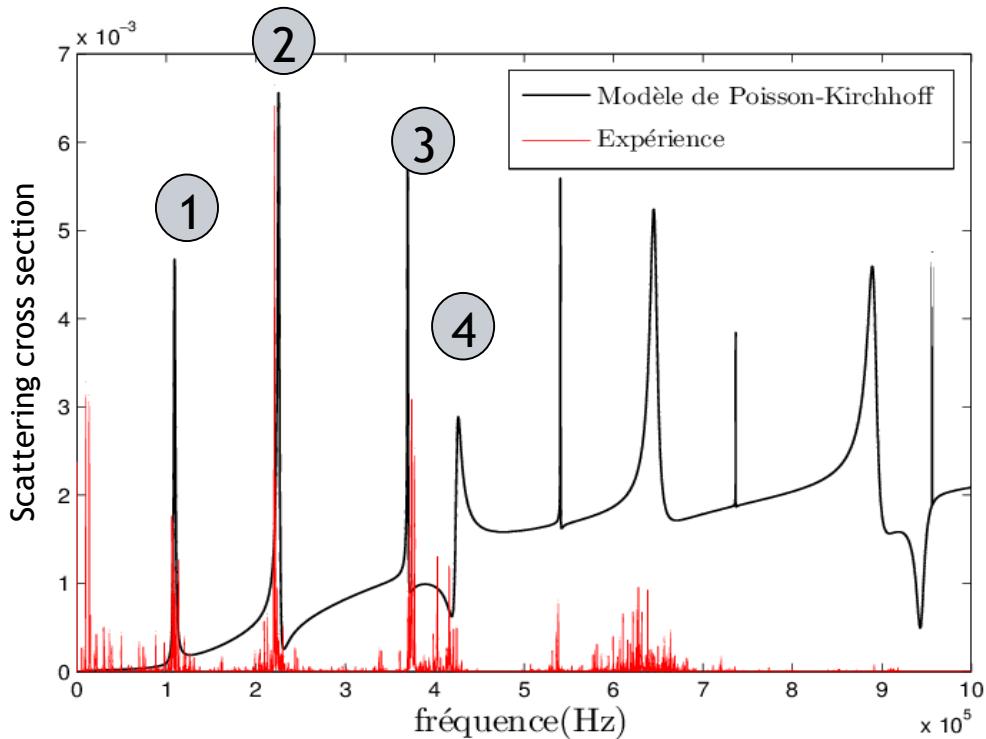
Time reversed wave propagation experiments in chaotic
micro-structured cavities

Rudolf Sprik ^{a,*}, Arnaud Tourin ^b

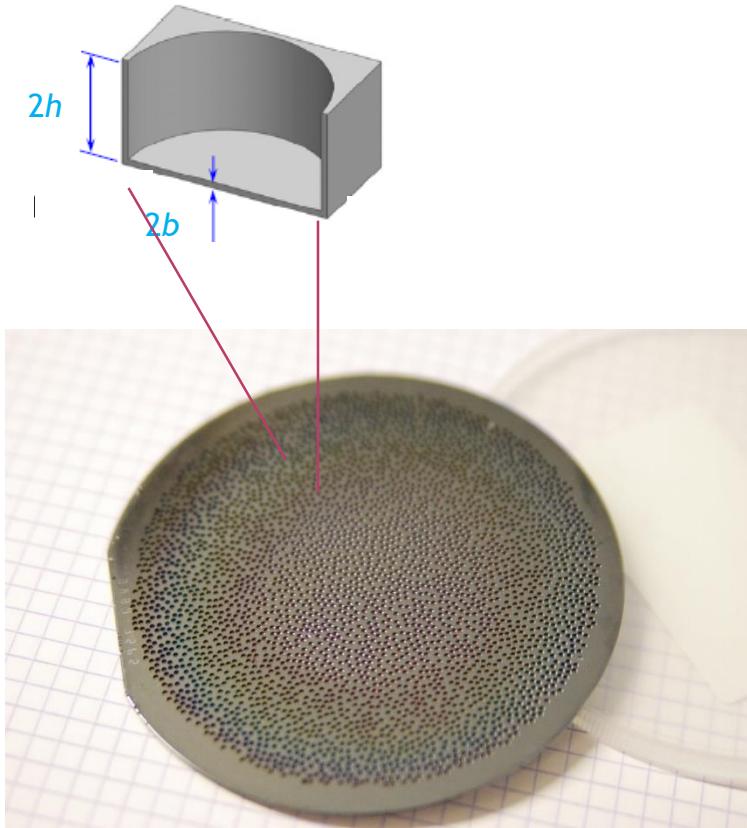


**Disordered-induced localization
was not observed**

RESONANT SCATTERERS



DISORDERED THIN PLATE



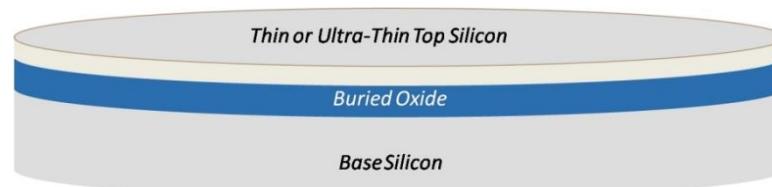
3400 blind holes randomly distributed

SOI wafer : 560 μm -thick

Blind holes : 530 μm -depth, $\emptyset=1 \text{ mm}$

Minimal distance : 100 μm

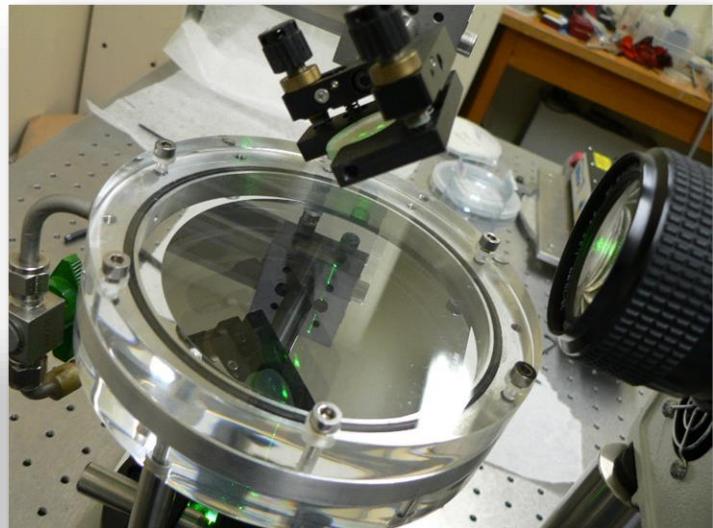
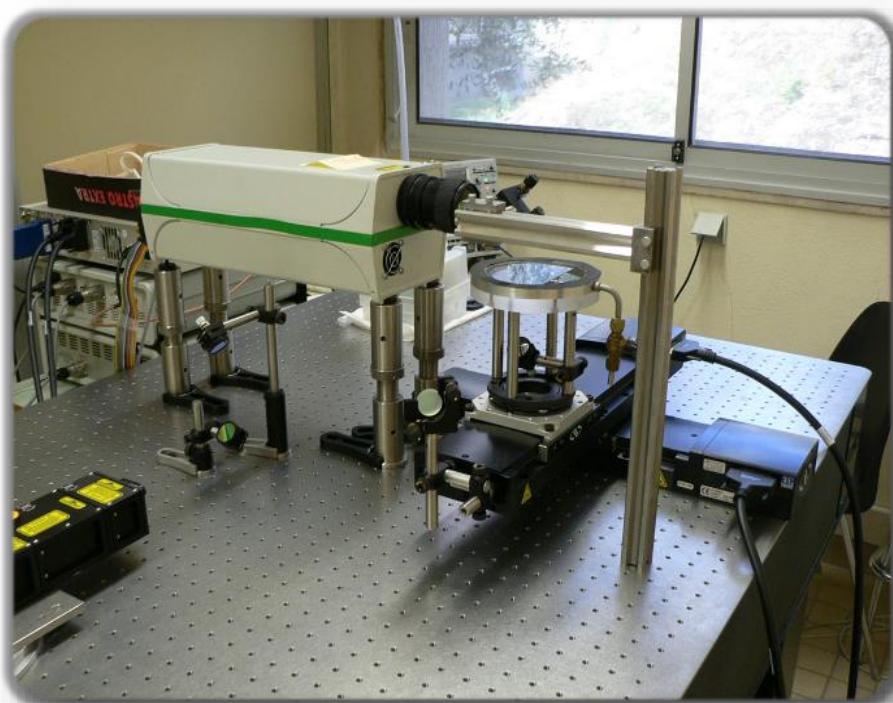
Challenge : reducing blind
hole-depth dispersion



- SOI (Silicon On Insulator)
- Photolithography
- Ion-beam etching

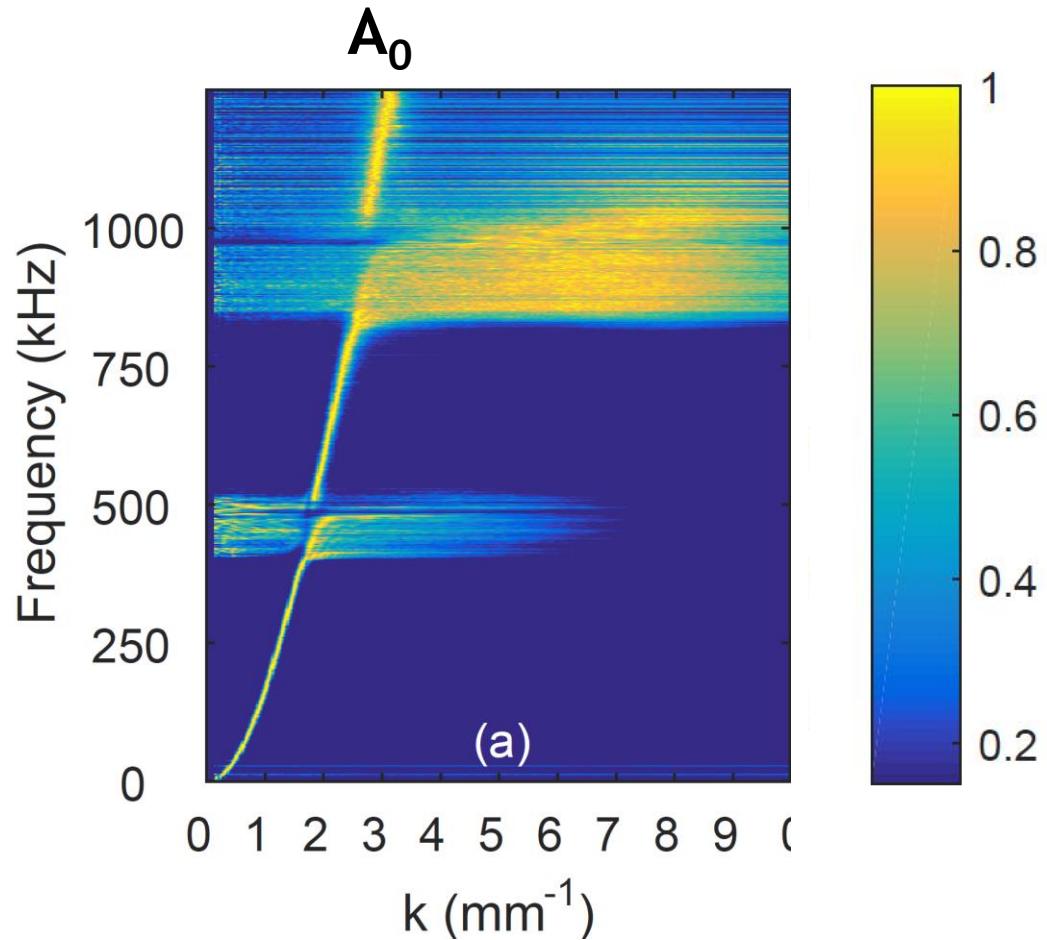
*Fabrication by Etienne Herth,
FemtoST, Besançon*

EXPERIMENTAL SETUP



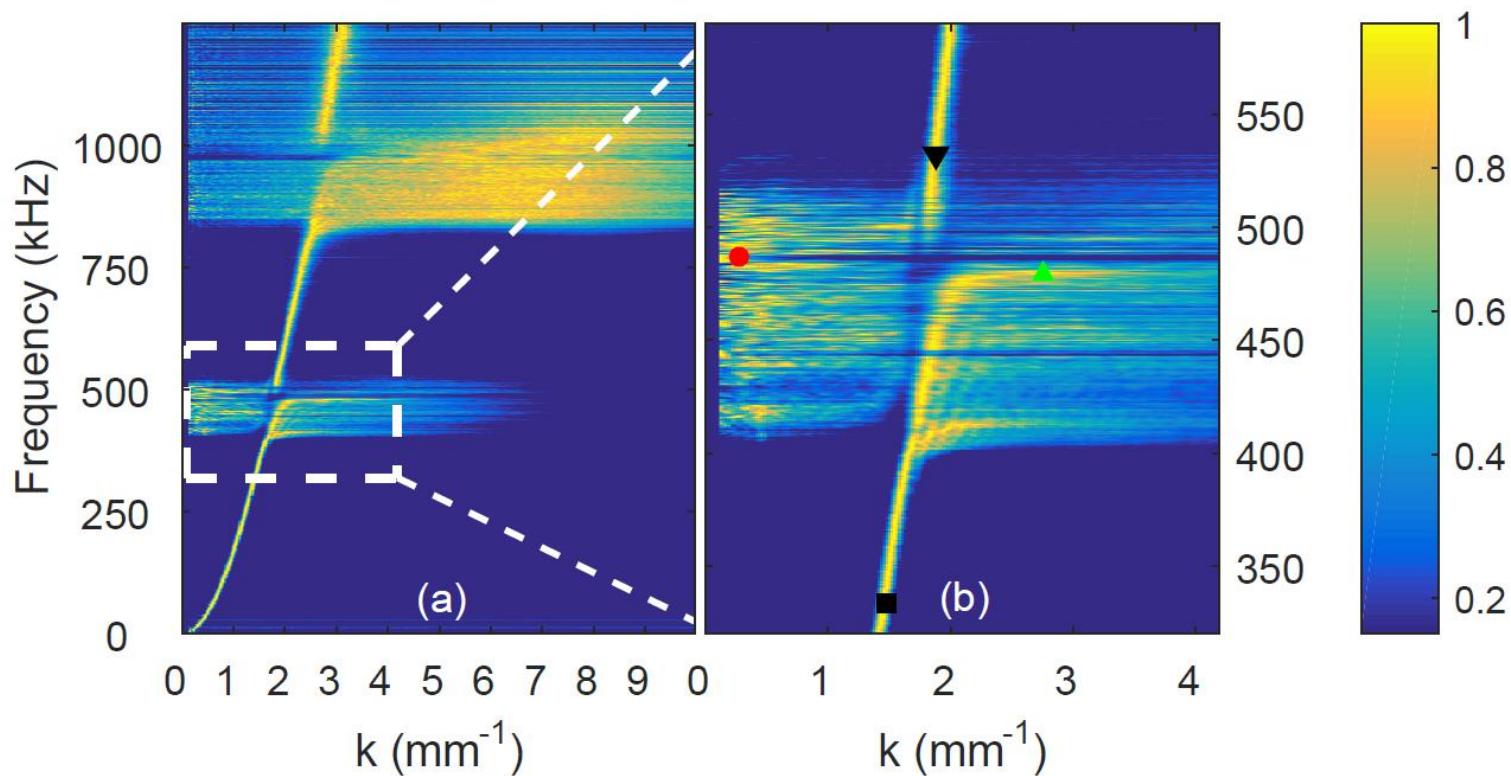
Heterodyne laser interferometer

DISPERSION DIAGRAM

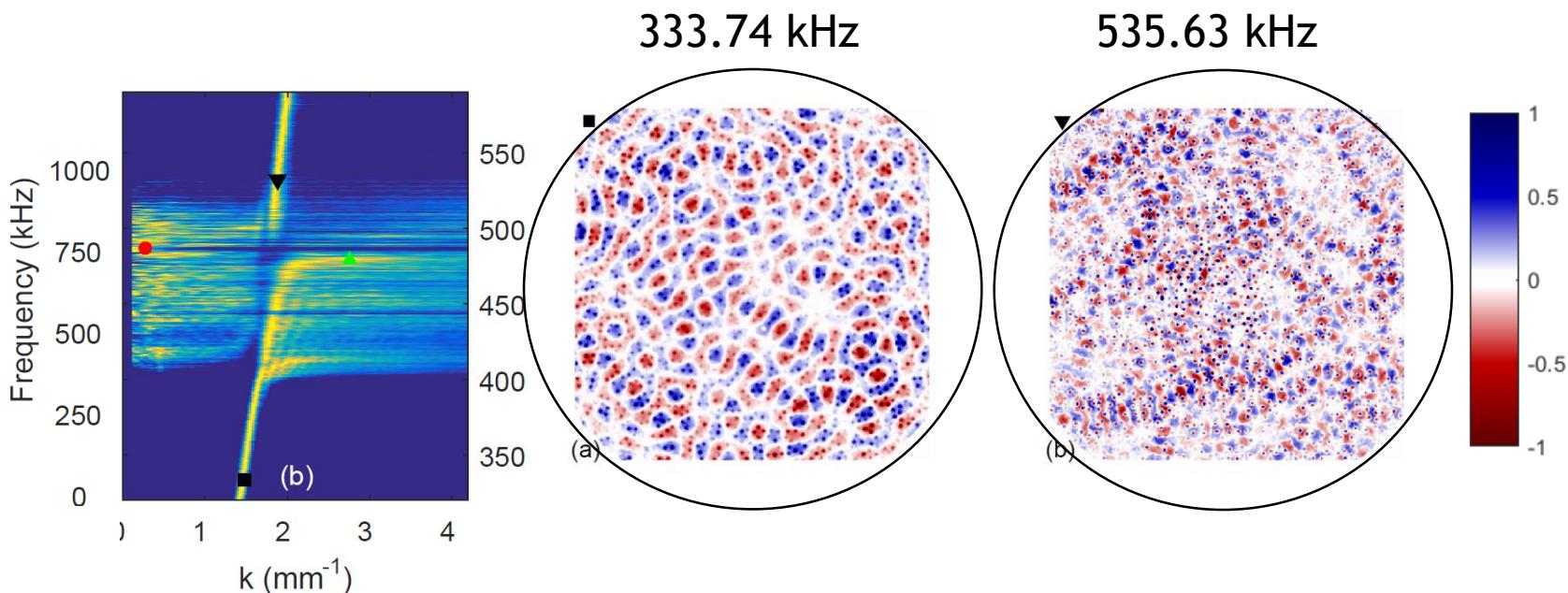


Hybridization gap

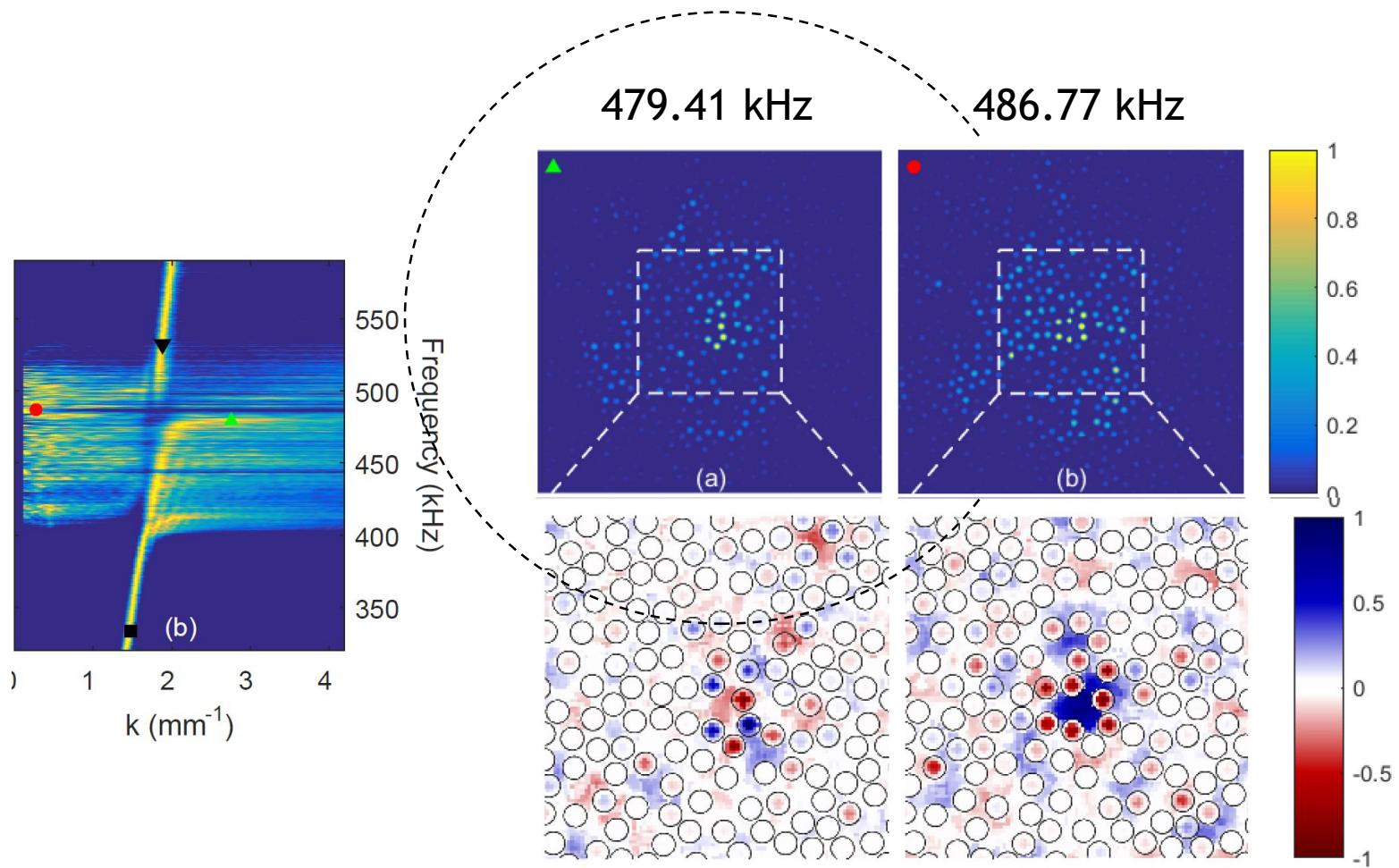
DISPERSION DIAGRAM



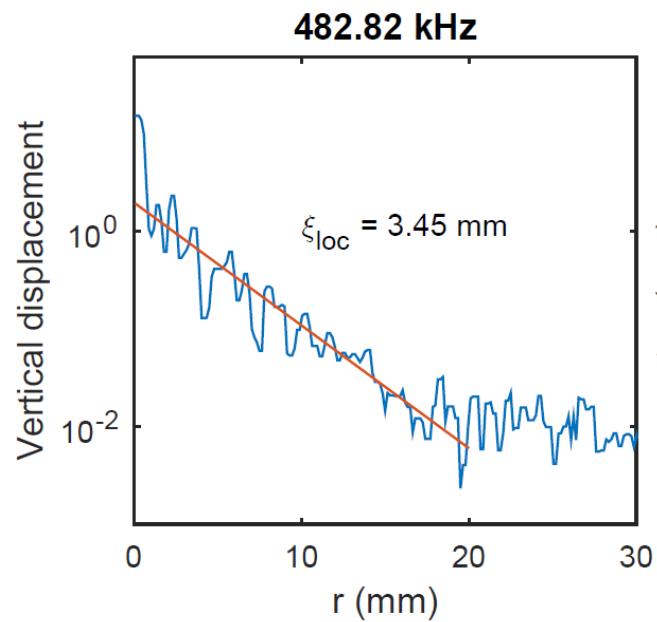
EXTENDED STATED



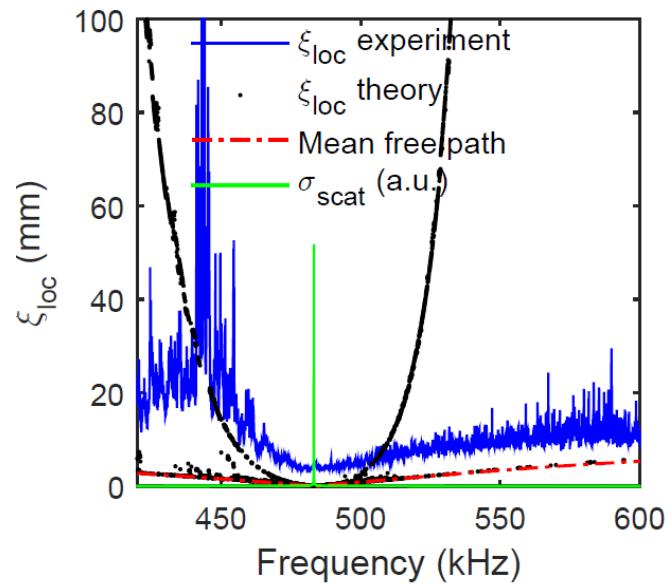
LOCALIZED STATES



LOCALIZATION LENGTH



Mode profile @ 482.82 kHz
(semilog)



Localization length vs.
frequency

SUMMARY

- Resonant scattering is an interesting route toward localization
- It is demonstrated here for flexural waves in a thin plate with blind holes
- The localization length is measured and shows a minimum near the resonance of the blind hole.
- Two types of mode are identified on both sides of the hybridization gap.

CONCLUSION

We have demonstrated modal confinement and localization in 2 particular situations:

- 1 - Localization by randomly pinned plate
- 2 - Localization by resonant scatterers

They correspond to 2 different situations :

- 1 - Many microcavities possessing their own resonant frequency.
At higher frequencies, they coupled to one another.
- 2 - Off-resonance: weak scattering. Near resonance: hybridisation gap and strong confinement of the wavefunction.

The Landscape theory of localization works in the pinned plate.
No theory exists yet in the second case.

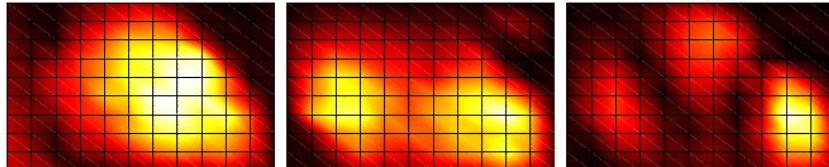
G. Lefebvre *et al.*, One single static measurement predicts wave localization in complex structures, Phys. Rev. Lett. 117, 074301 (2016).

G. Lefebvre *et al.*, Localization induced by resonant scatterers in thin plates, in preparation.



THANK YOU FOR
YOUR ATTENTION

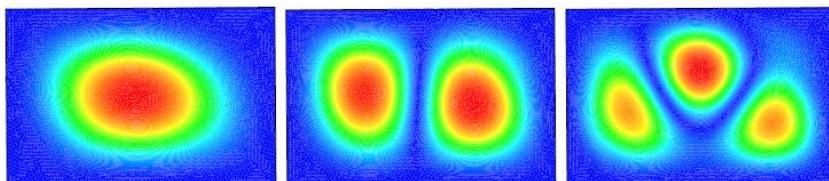
LOW AND HIGH FREQUENCY MODES



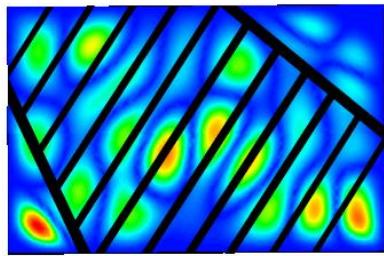
(a) (1,1)-mode

(b) (2,1)-mode

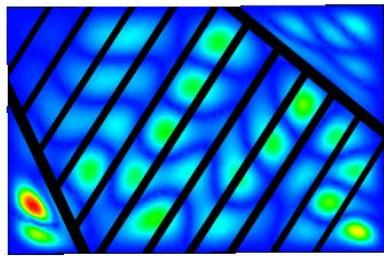
(c) (3,1)-mode



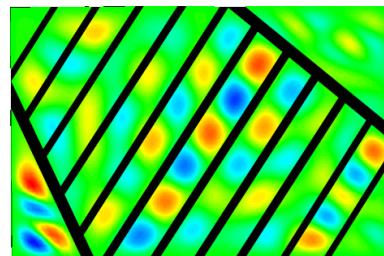
First three modal shapes of the upright piano soundboard



$f_n = 776 \text{ Hz}$



$f_n = 1089 \text{ Hz}$



$f_n = 1593 \text{ Hz}$

Intensity (and field) distribution of higher modes of the upright piano soundboard

OTHER STRING INSTRUMENTS

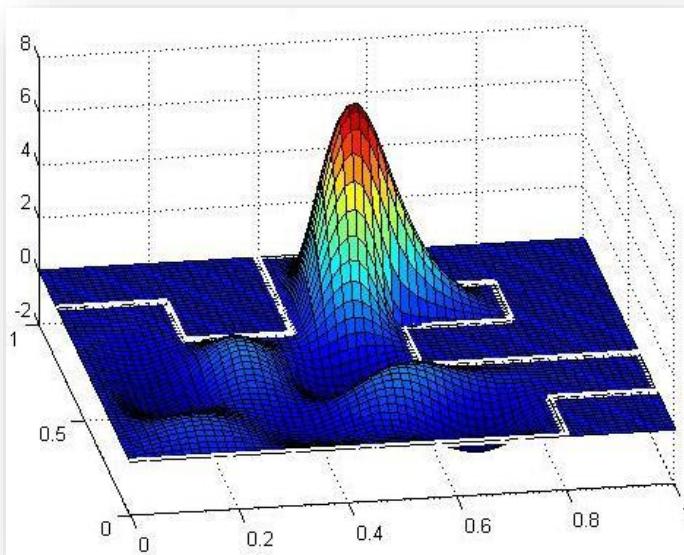


MOTIVATION

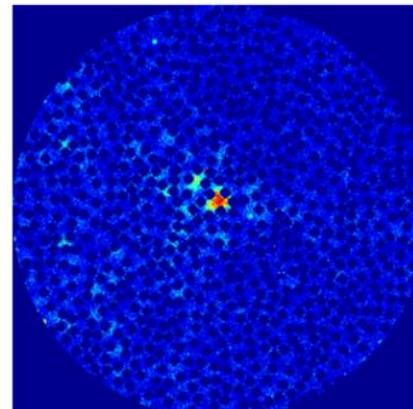
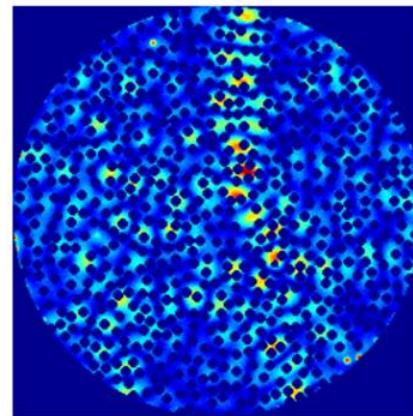
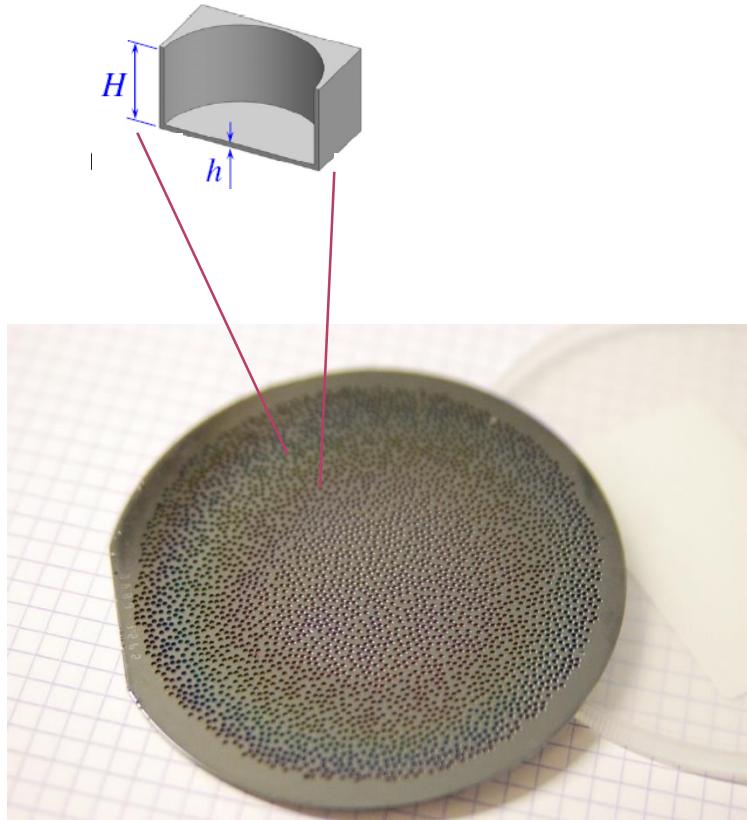
- Eigenmodes and eigenvalues : Key ingredients to understand the physical properties of most vibratory systems
- Yet, **the prediction of the spatial and spectral distributions of the stationary waves** remains an open problem in many cases.
- And **the design of the structure with predefined modal characteristics** remains a challenge

WAVE LOCALIZATION : A BROAD DEFINITION

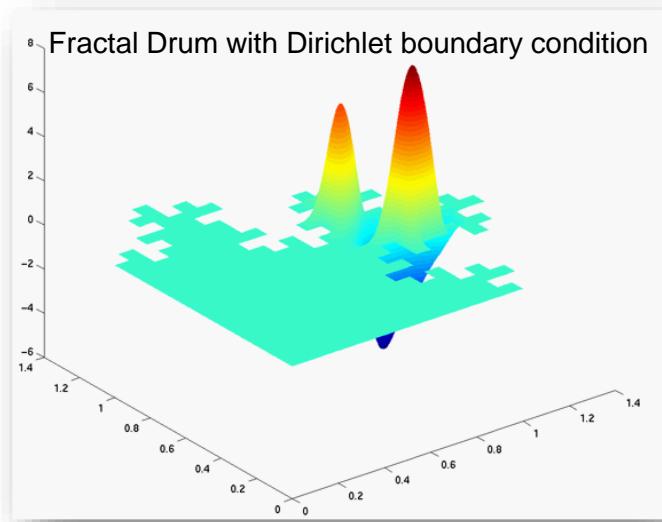
A state is said to be “localized” if the vibration amplitude is large only in a restricted region of the total domain



MODE CONFINEMENT IN A DISORDERED THIN PLATE



WAVE LOCALIZATION IN COMPLEX STRUCTURES



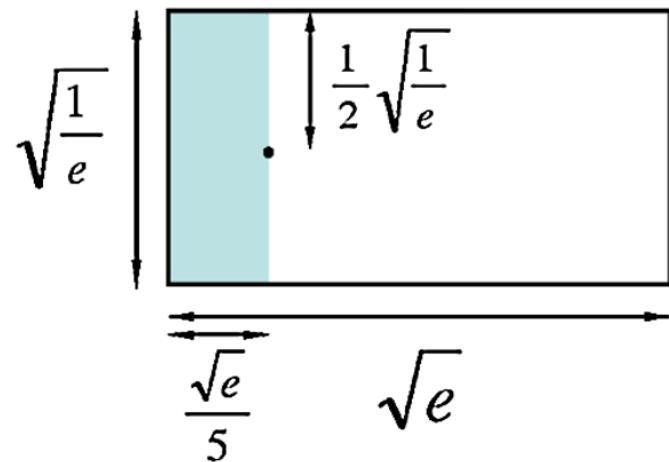
Félix et al., J Sound Vib, 2007



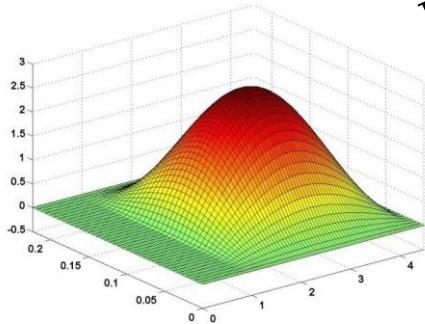
Heilman & Strichartz, Not Am Math Soc, 2010

*"Localized eigenfunctions: here you see them,
there you don't"*

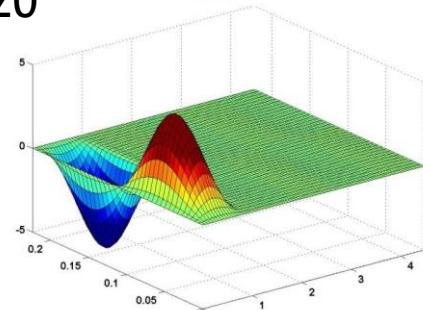
WAVE LOCALIZATION IN A CLAMPED PLATE



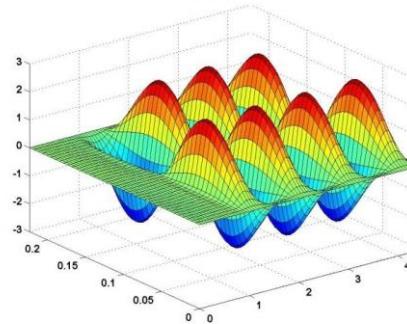
Mode 1



Mode 40

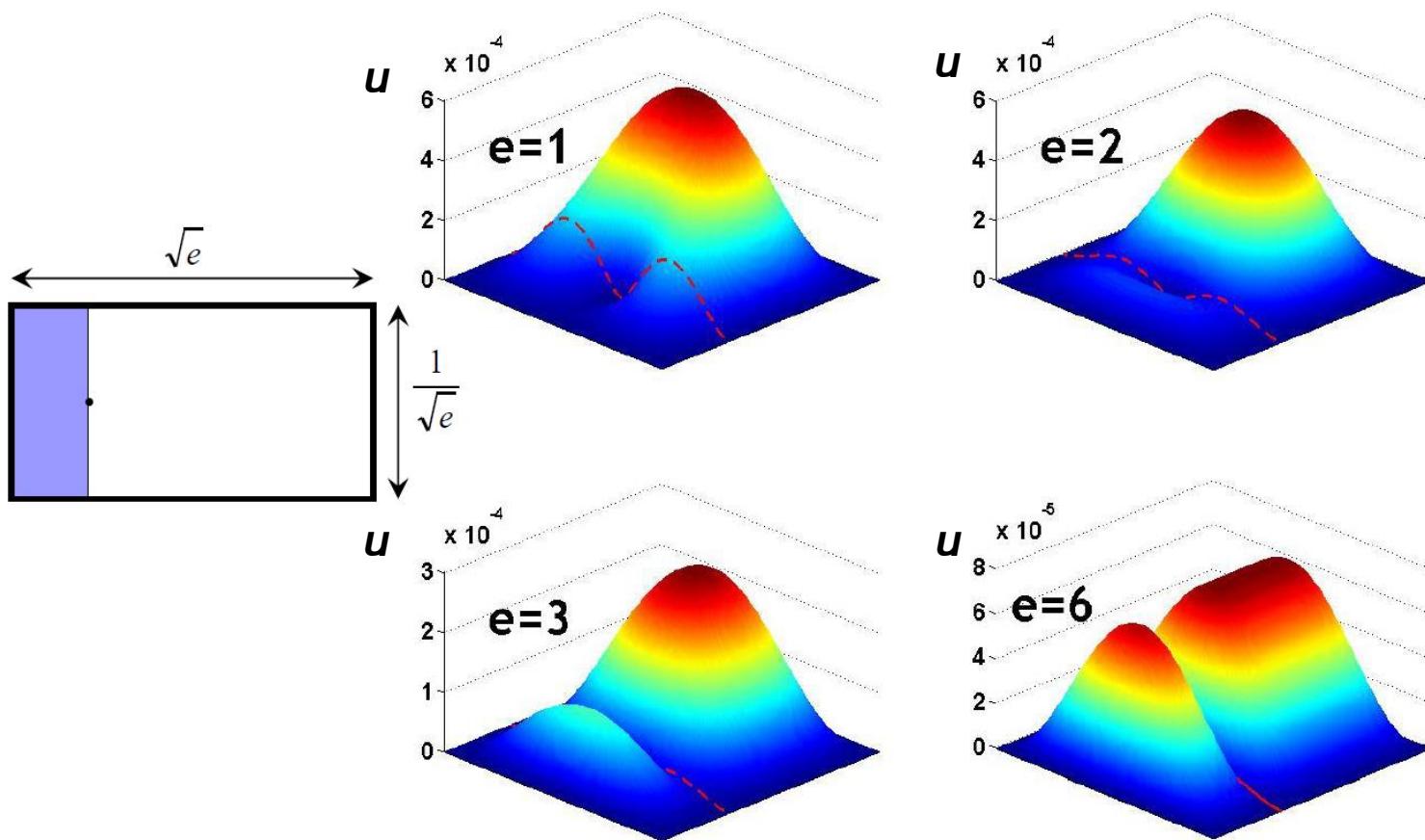


Mode 44

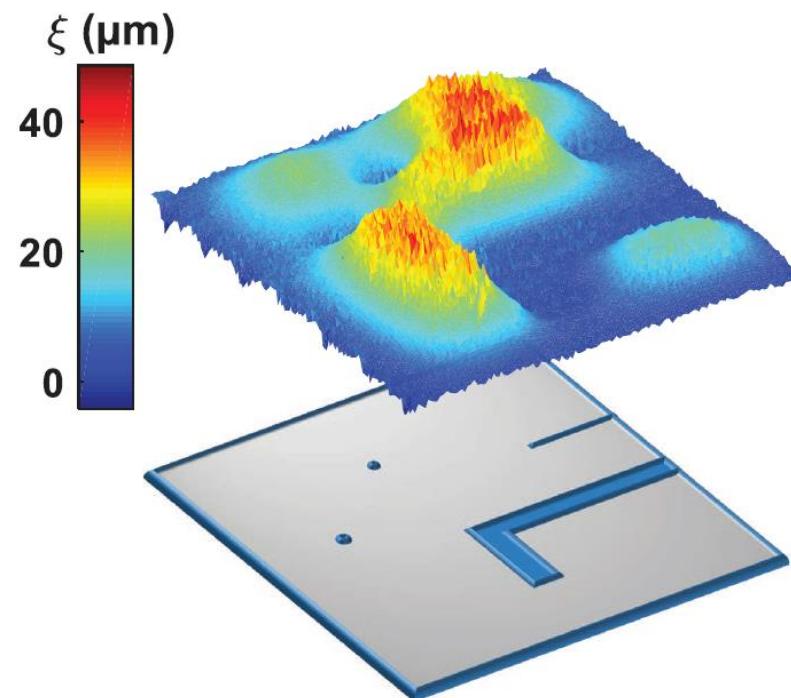
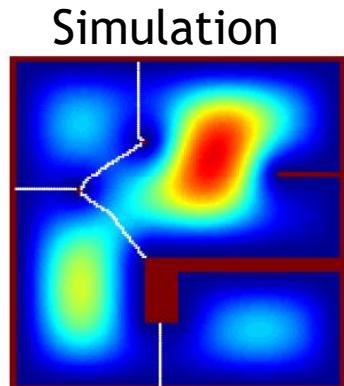
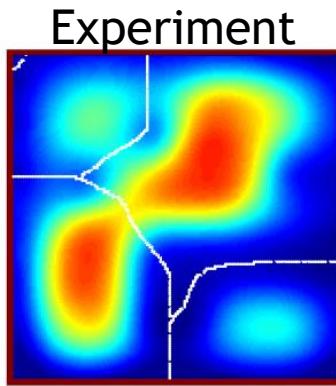


M.Filoche & Mayboroda, *PRL*103, 254301 (2009)

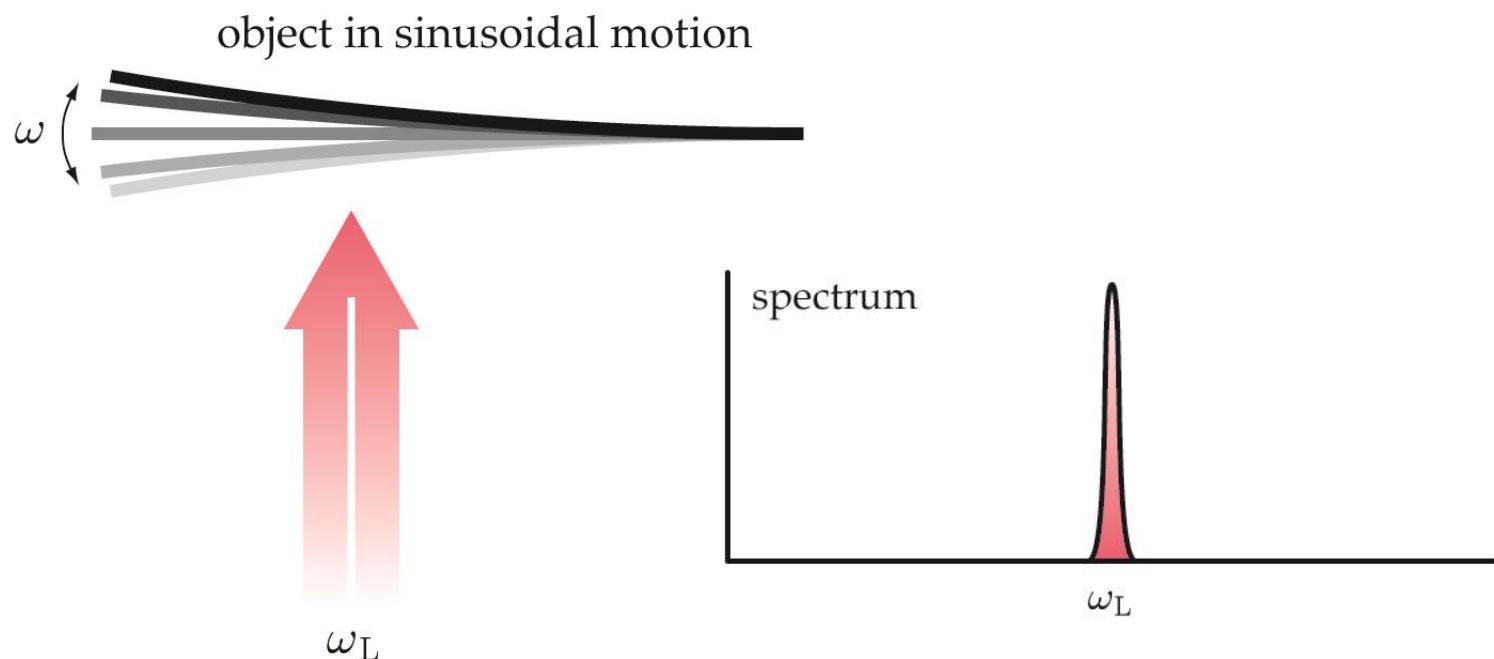
THE LANDSCAPE IN THE CLAMPED PLATE



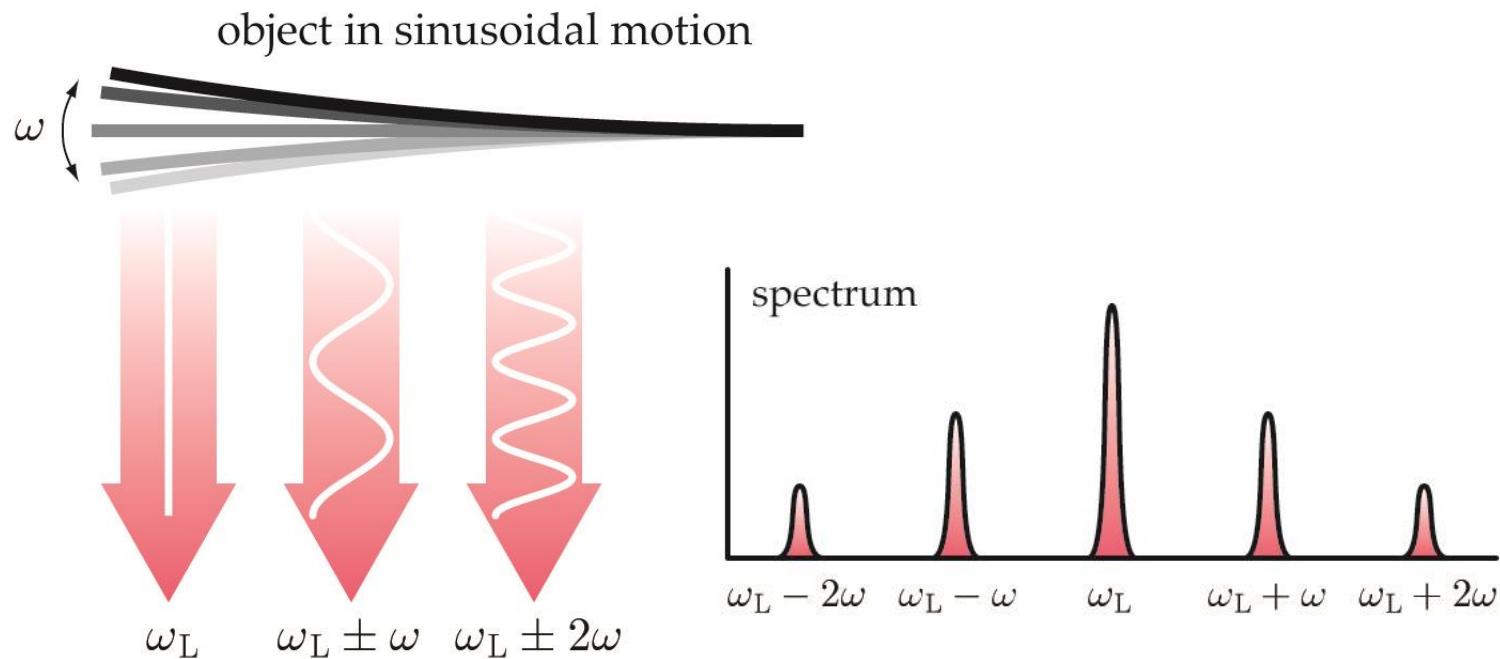
MEASURING THE LANDSCAPE



LIGHT SCATTERED BY AN OBJECT IN SINUSOIDAL MOTION

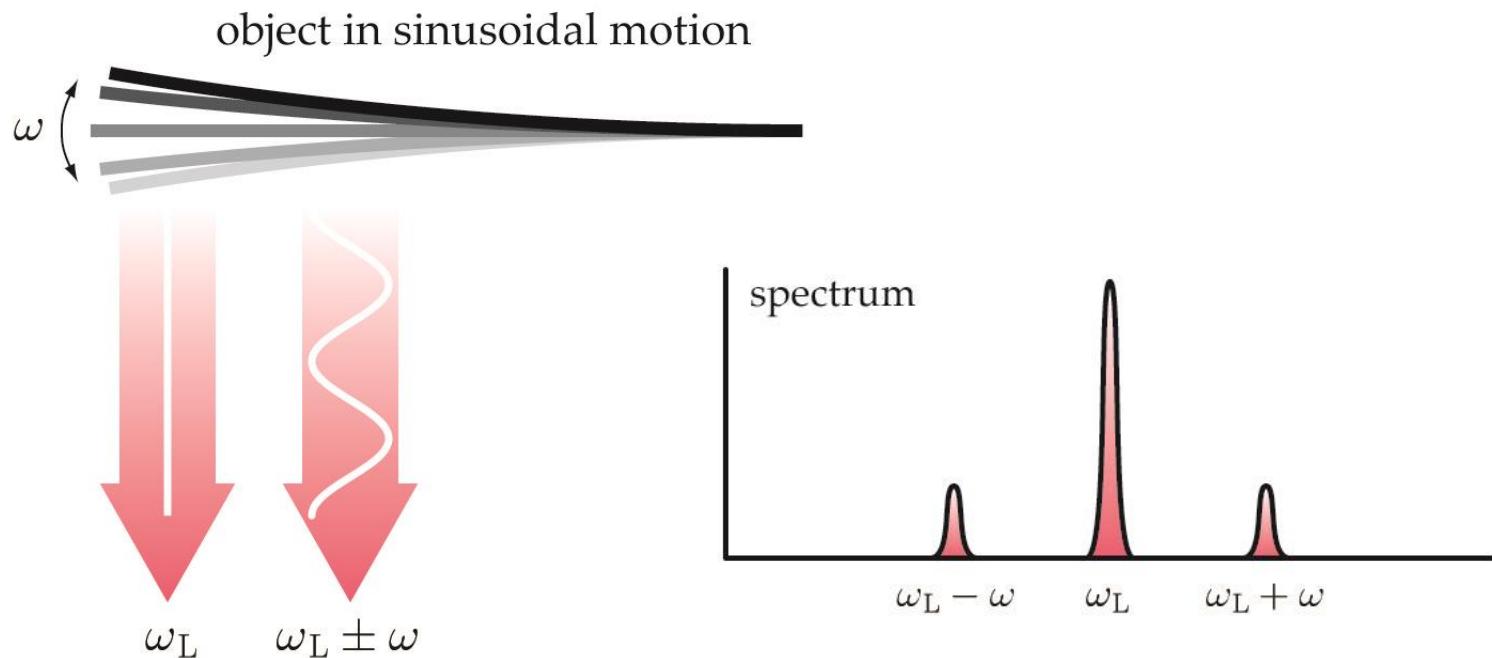


LIGHT SCATTERED BY AN OBJECT IN SINUSOIDAL MOTION



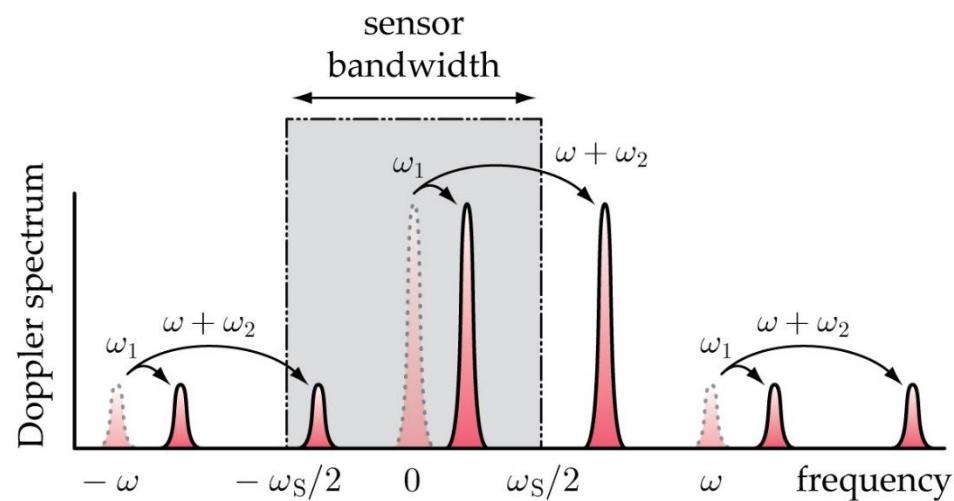
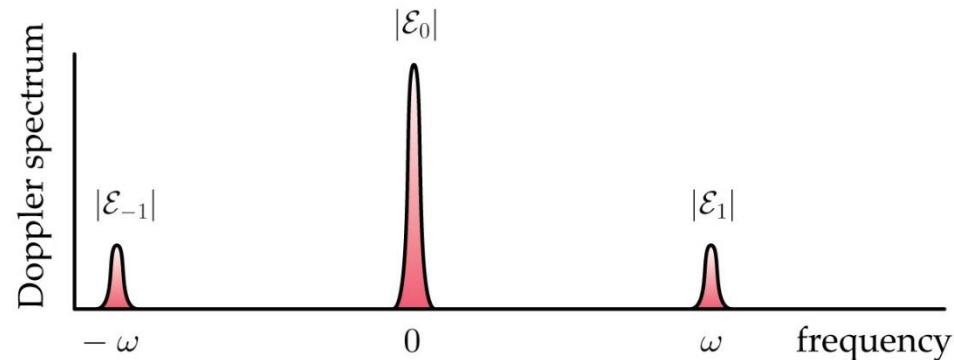
Heterodyne holographic interferometry (M. Atlan *et al.* Opt. Lett. 39, 2014)

SMALL VIBRATIONS



Heterodyne holographic interferometry (M. Atlan *et al.* Opt. Lett. 39, 2014)

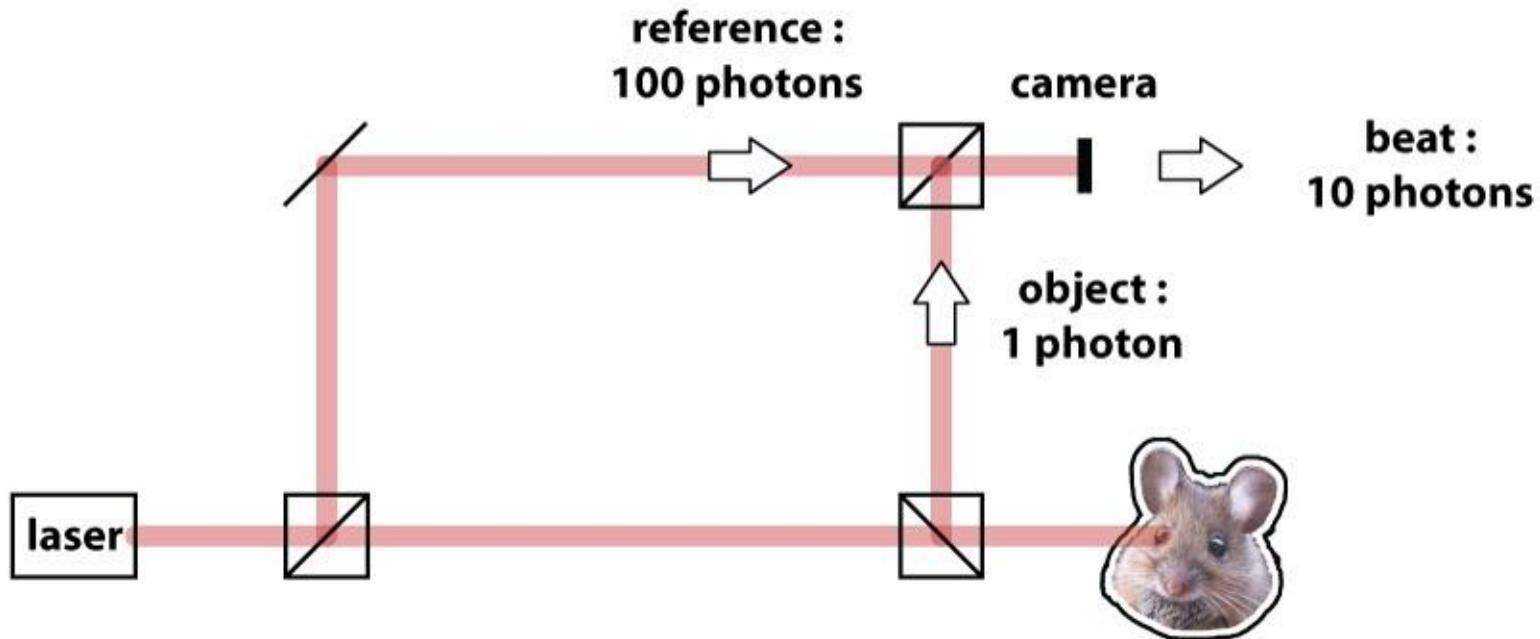
DOUBLE SIDEBAND MODULATION



PRE-DETECTION OPTICAL SIGNAL ENHANCEMENT

HIGH HETERODYNE GAIN REGIME

The weak object wave beats against the reference wave.
Signal modulation depth is enhanced



EXPERIMENTAL VALIDATION OF THE LANDSCAPE THEORY

- Show that the landscape function controls the amplitude of the localized waves = *effective potential*
- Show that the landscape function predicts the eigenvalues of the vibratory modes
- Show that the landscape function gives the evolution of the mode extension vs. frequency = *transition from localized to extended*

MAX_I(U)

$$u = \sum_i \alpha_i \varphi_i \quad \text{with} \quad \alpha_i = (u, \varphi_i) = \int u \phi_i \, dx$$

$$(Lu, \varphi_i) = (1, \varphi_i) = \int \varphi_i \, dx$$

$$(Lu, \varphi_i) = (u, L\varphi_i) = \lambda_i (u, \varphi_i) = \lambda_i \alpha_i$$

$$\alpha_i = \frac{1}{\lambda_i} \int \varphi_i \, dx .$$

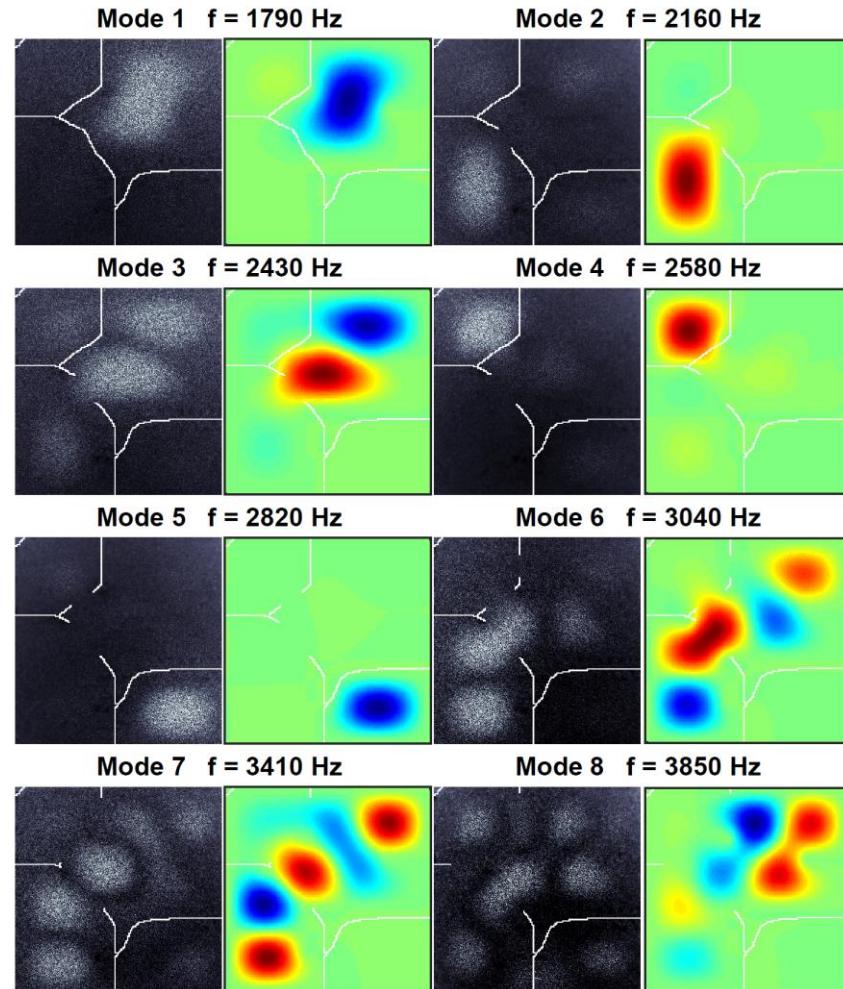
For the fundamental mode of a localization subregion

$$\lambda_0 = \frac{\int \varphi_0 \, dx}{\alpha_0} = \frac{\int \varphi_0 \, dx}{(u, \varphi_0)}$$

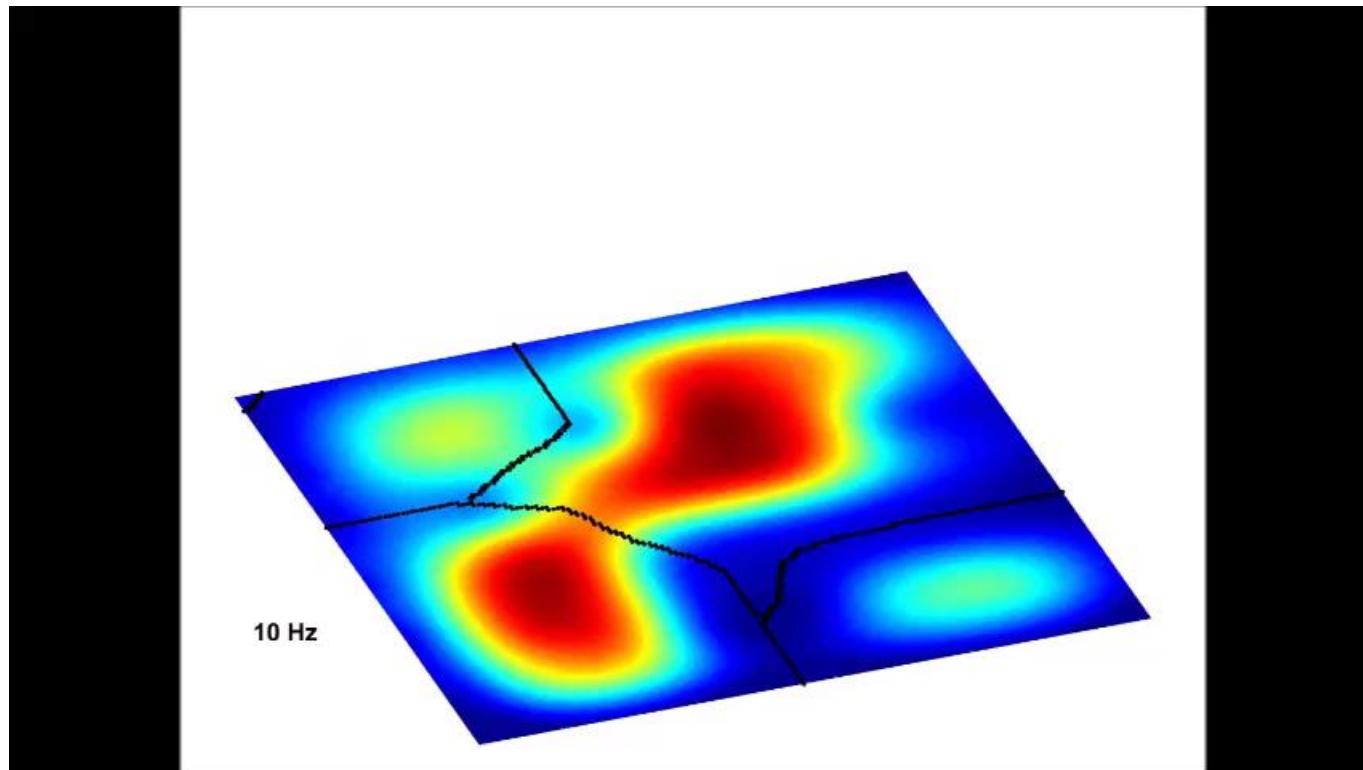
u is essentially the fundamental mode

$$\lambda_0 \approx \frac{\int \frac{u}{\|u\|} \, dx}{\left(u, \frac{u}{\|u\|}\right)} = \frac{\int u \, dx}{\int u^2 \, dx}$$

MEASUREMENTS VS. SIMULATIONS



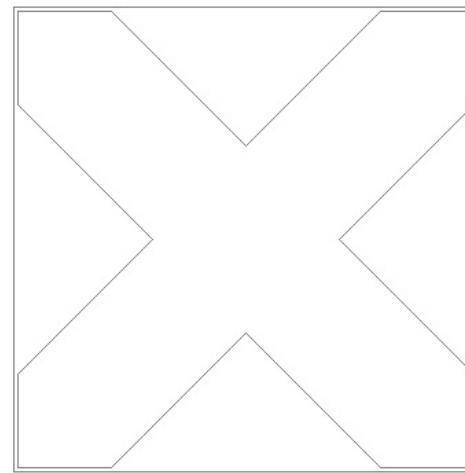
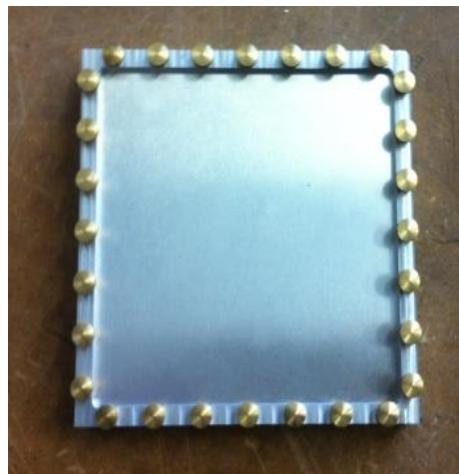
4 - EFFECTIVE VALLEY NETWORK



THE LOCALIZATION LANDSCAPE

- Describes the sub-regions of localization
- Gives, in a single measurement
 - the eigenfrequencies
 - The eigenmode profile
 - The frequency dependence of the mode spatial extension
- This is a static measurement, which does not require a precise knowledge of the microscopic geometry
- It can be used to address the inverse problem and design the system geometry corresponding that a desired spatial distribution of the modes

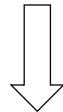
HOW TO MAKE AN X VIBRATE ?



PROBLEM FORMULATION

Kirchhoff-Love equation
of flexural waves

$$\frac{\partial^2 w}{\partial t^2} + \frac{Eh^2}{12\rho(1-\nu^2)} \nabla^4 w = 0$$



Steady-state
equation of the eigenmodes

$$\begin{cases} LW_m = \omega_m^2 W_m & \text{on } \Omega \\ W_m|_{\partial\Omega} = 0 \\ \mathbf{n} \cdot \nabla W_m|_{\partial\Omega} = 0 \end{cases}$$

$$L = \frac{1}{\alpha^2} \nabla^4.$$



Minimization of the distance

$$\|u - u^*\|$$



Derivation of the
experimental landscape with
N clamped points

$$\begin{cases} Lu = 1 & \text{on } \Omega \\ u|_{\partial\Omega \cup A_1 \cup \dots \cup A_N} = 0 \\ \mathbf{n} \cdot \nabla u|_{\partial\Omega \cup A_1 \cup \dots \cup A_N} = 0 \end{cases}$$



Derivation of the
Ideal landscape

$$\begin{cases} Lu^* = 1 & \text{on } \Omega \\ u^*|_{\partial\Omega \cup \partial\Omega_1} = 0 \\ \mathbf{n} \cdot \nabla u^*|_{\partial\Omega \cup \partial\Omega_1} = 0 \end{cases}$$

FREQUENCY RESPONSE OF A PLATE CLAMPED WITH 16 PINS

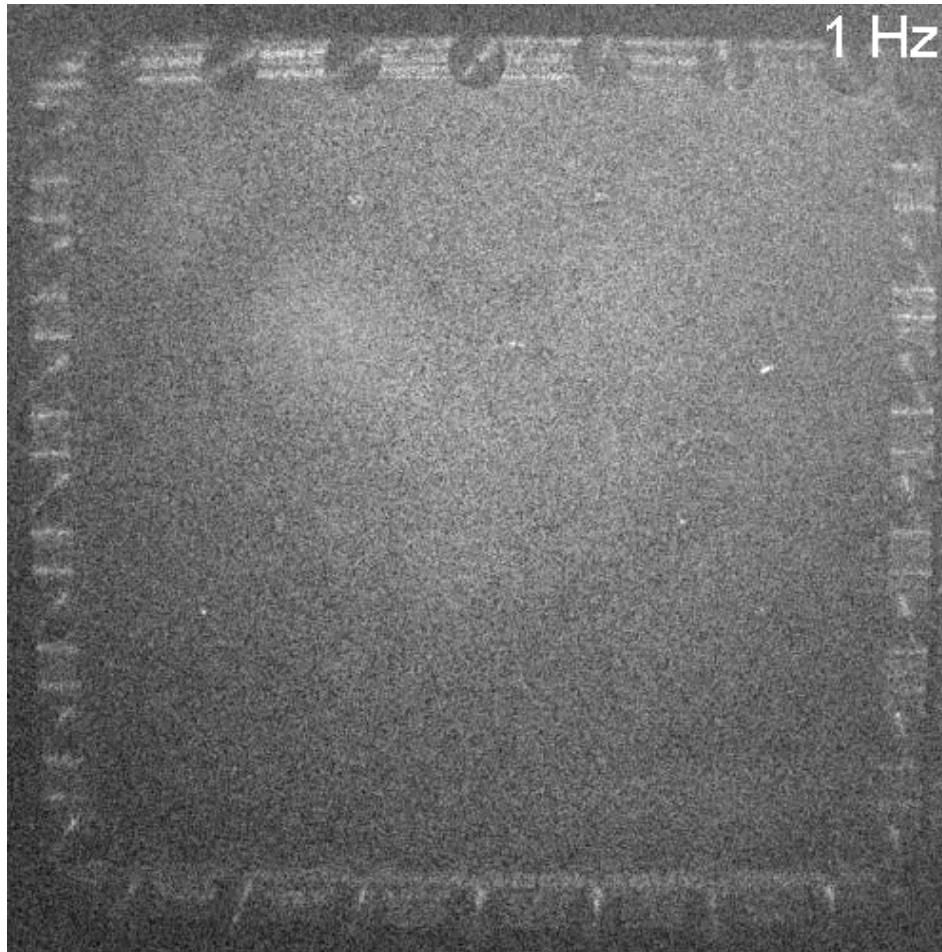
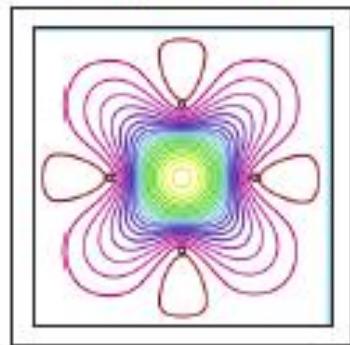


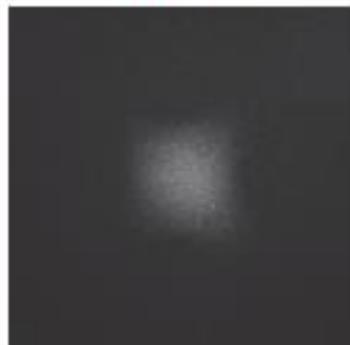
PLATE WITH 4 PINS

$m = 0, 1474$ Hz



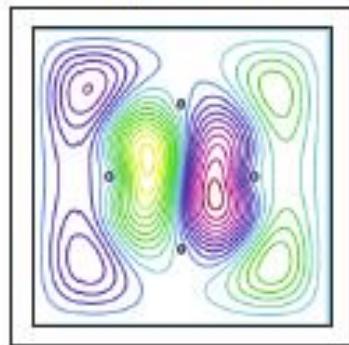
(a)

1100 Hz



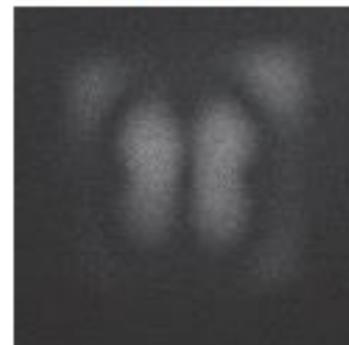
(b)

$m = 6, 2754$ Hz



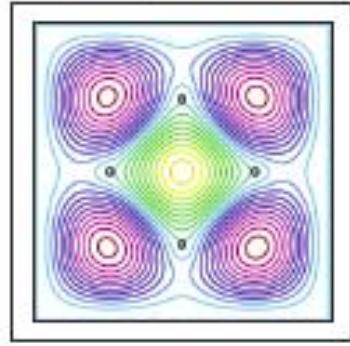
(e)

2630 Hz



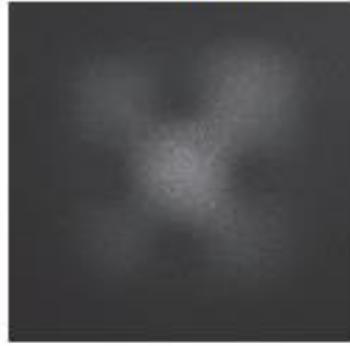
(f)

$m = 4, 1782$ Hz



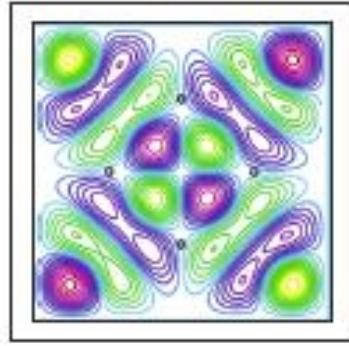
(c)

1450 Hz



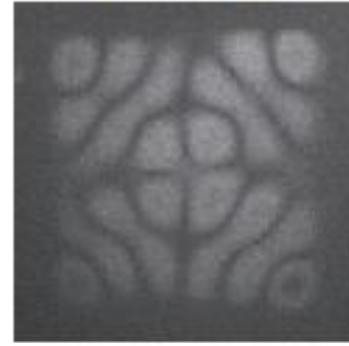
(d)

$m = 33, 7290$ Hz



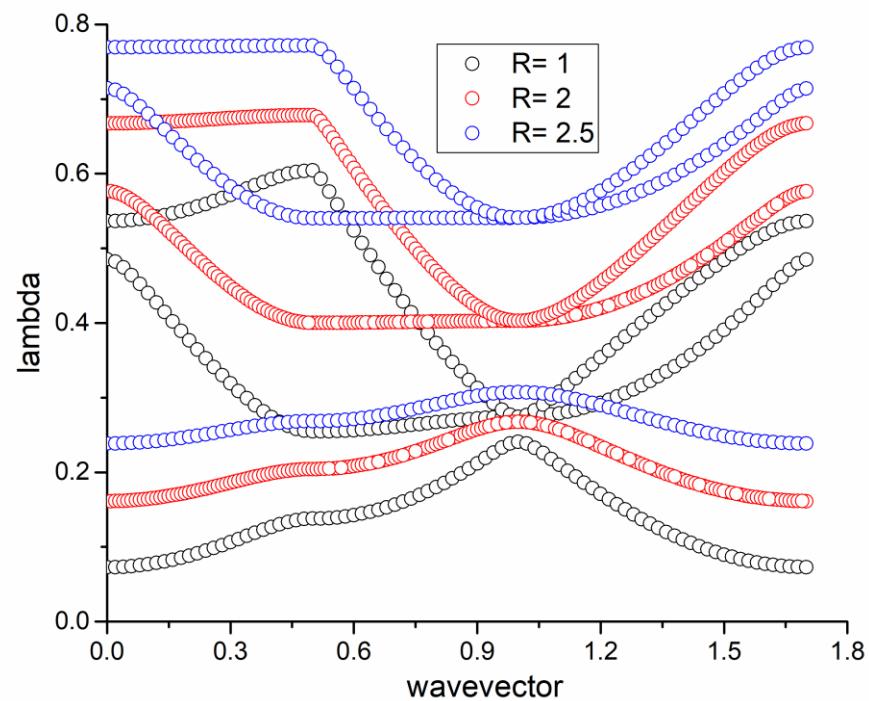
(g)

6850 Hz



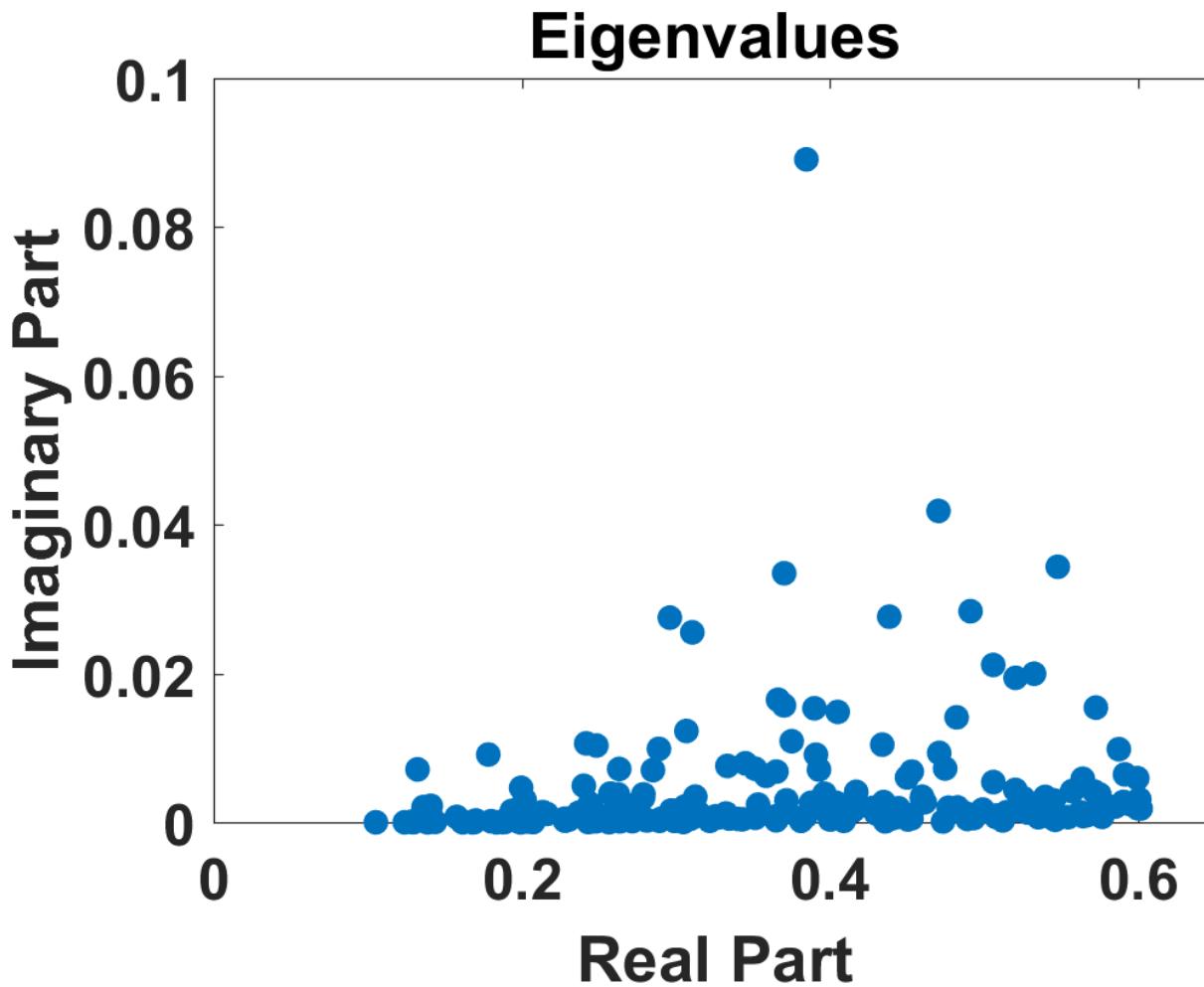
(h)

BAND DIAGRAM

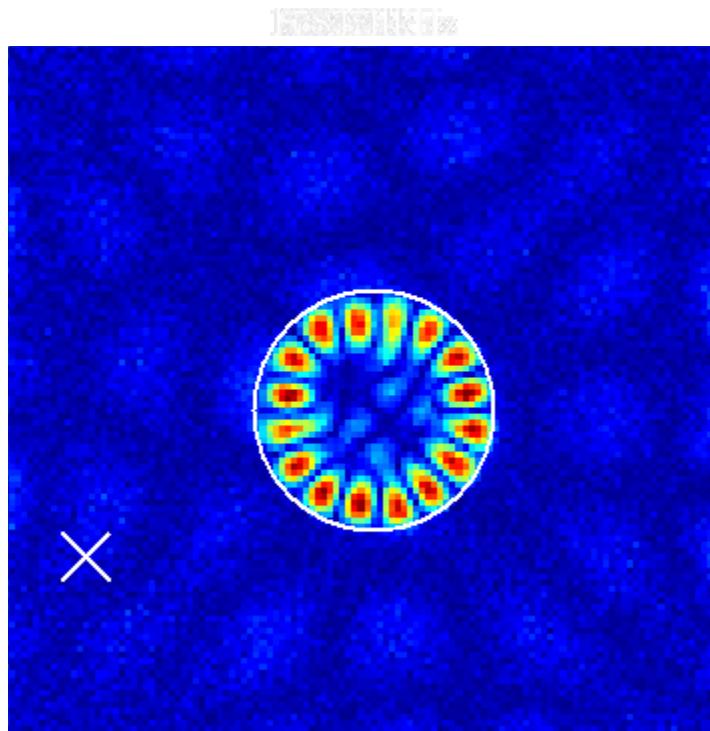


Low frequency gap

DISTRIBUTION OF RESONANCES



BLIND HOLE RESONANCES



COLLABORATORS

- ⦿ **PhD students :**

- ⦿
 - Marc DUBOIS (Institut Langevin, Paris, France)
 - Olivier XERIDAT (LPMC, Nice, France)

- ⦿ **Post Doc :**

- ⦿
 - Gautier LEFEBVRE (Institut Langevin, Paris, France)
 - Kun TANG (Bar Ilan University, Israel)

- ⦿ **Fabrication :**

- ⦿
 - Etienne HERTH (FemtoST, Besançon, France)

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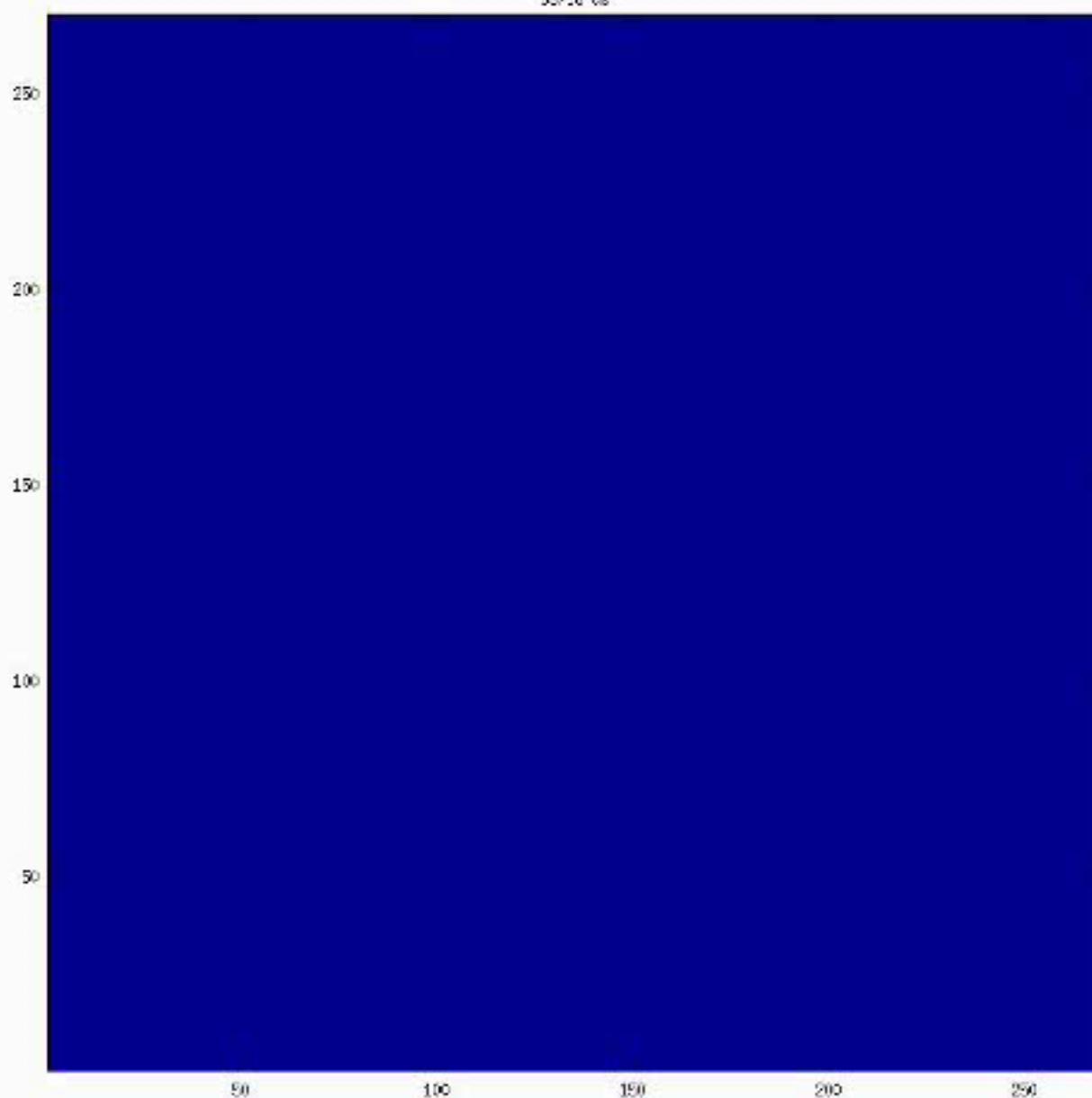
United States-Israel Binational Science Foundation (NSF/BSF 2015694).

Papers :

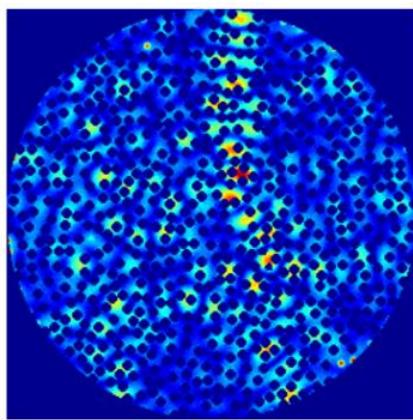
G. Lefebvre *et al.*, One single static measurement predicts wave localization in complex structures, Phys. Rev. Lett. 117, 074301 (2016).

G. Lefebvre *et al.*, Localization induced by resonant scatterers in thin plates, in preparation.

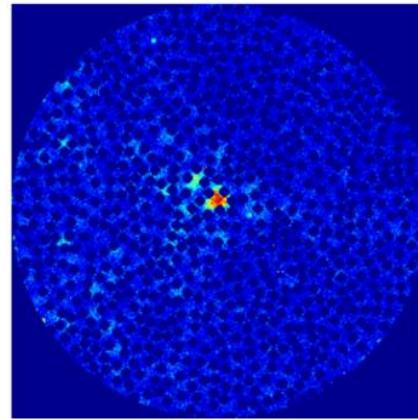
997.6 vs



LOCALIZATION OF FLEXURAL WAVES



Extended mode



Localized mode

PERSPECTIVES

- Can be generalized to any complex system, where the structural or microscopic information is not accessible
- Towards the prediction of the full dynamics from a static measurement
- Extension to other systems, other wave equations
- Inverse problem: towards engineering of modes