ISM 26-09-2019 Know difference between statistical and thermodynamic equilibrium = detailed balance : each process is reversed by He inverse process SE: sum of all processes into one level is equal to the sum of all processes at · of that level $\frac{n_0}{n_0} = \frac{g_0}{g_0} \exp\left(-\frac{hv_{01}}{hT_{ex}}\right)$ We often simplify by considering systems as two-level systems. Given Einstein A coefficient we can calculate all other coefficients. [AU] = 5 (average # lifetime) $\left(\frac{\partial n}{\partial t}\right) = \left(\frac{\partial n}{\partial t}\right) = B_{iv} n_i \int_{\Gamma} v \varphi_i dv$ density of photons

(as by Draine)

Lightman / Rybich uses Iv

Bri C Bornere ce will look at the profile Jordv=1 + Ine is somewhere Cross - section ow (1)

R = (ku) Ne No ; @ collision [cm 35-1] 6 dow(r) cdt dry) = ne for (v) c hiver dv a no four wood because of narrow = Ben ne U podr spectral line Bu & Bu BUL S AUL Jozu (V) dv = ge on vuz Aux variable Toonstant Ow (V) = ge dry 2 Aur Pr OU (V) = gu O20 W Harmonic oscillator line profile Pr = 48012 16n2 (v-vo)2 + xu2. DVANHO = 801 -> [5] SEST +7 t/2

 $y_{ij} = \sum_{j \in \mathcal{J}} A_{ij} + \sum_{j \in \mathcal{L}} A_{ij}$ Some numeric values -> Du ~ 0.1 kms in applical (3) All ISM lines are much vider than this so in most cases we can ignore the natural broadening. In X-ray, au no kms-, comparable to ISM linewidth. In X-ray, due to T the thermal linewidth is still dominating the broadening 0 0 $P_{V} = \sqrt{2\pi} \sigma_{V} \exp\left(-\frac{(V-V_{0})^{2}}{2\sigma_{V}^{2}}\right)$ 0 e.g. the fraction of particles with velocity between v and of v+dv. Combining natural and gaussian line profile from the turbulent as broadening (convolution) is the Voigt profile -> Voigt = Natural & Gaussian (Draine 6.34) 0 0 Profile has vings more prominent than Gaussian. At some point the natural broadening vings will thus start to dominate. Cross-over point 0 happens at V-Vo 1, 4.5 ov Normally we do not measure signal here Cobe to low SNR). What happens with an incoming beam of radiation Iv+dIv

S S+dS

remissivity tem dIv = - Iv kvds + Jvds [eg stan3tht fractional absorption & per unit length [cm-1] colone = hvan A.u. pars/volume/ss = 40 No AUL pover/volume/sr/ 4 j= hvu no Au Pu emissivity Kr = no ou w) - no ou w) pure absorption shim emission netto absorption $= \int_{\mathcal{U}} \mathcal{O}_{\mathcal{U}} \left(V \right) \left(1 - \frac{\int_{\mathcal{U}} \int_{\mathcal{U}}}{g_{0}/q_{0}} \right)$ = nov(v) 1- exp(-hvu/ktex) -hulletex) of > On low = ne go onvoi Auc (1-e Optical depth: du = krds

We can now write equation of radiative transfer in terms of optical depth. 4 dIv=-Iv dov + (by) dov -> de = (Sv-Iv) dov (1) etr (dIv + to da) = etr Sudtr d(e Iv) = et Sudon 1 N II multiply by et to to State of the State of t Iv (Tv) = Iv (o) e + f e (Tv - Tv') Sudti intensity emission in slab, but corrected for Sv = Vc = c2 & e hylkiex -1 = Br (Tex) Ly Kirchhoff's law.

Iv = Iv (0) e + Br (Tex) (1-e-tr) & TAN Z Infe Limiting cases: Absorption only I) o $T \rightarrow cn$ $T_{v(e)e^{-tor}} \Rightarrow 0$ $P_{1-e^{-tr}} \rightarrow P_{anck}$ $T_{v} = B_{v}(T_{ex})$ Planck 2) = T > 0 for with source Function abside medium. I-e-to 20 } See right through,
Iv = Io(0) 3) Abs only: Iv= Iv(o)e-Tr 3) Emission only Iv

Total

Tv = Br(Tex)(1-e-v) Iv = By [tv] = TuBV(Tex) emission proportional to optical Str= ne ge znvor Al (1-exp(hvor/texk) c/ods = M₂ $\frac{g_{\nu}}{g_{1}} \frac{e^{2}}{2nv_{\nu}r^{2}} A_{\nu} \left(1-e^{-hv_{\nu}/kTex}\right) \phi_{\nu}$ coldens

Tu = N_L $\frac{c^2}{2\pi v_{al}^2}$ Aul (e hu/kitex -1) ϕ_V To solve, we need to know excitation temperature -additional measurements - use guesses Tex can become negative if ne) ge, as T is just an expression for head values, this is no problem. 1 know something of Masers -Tex 40 - Rotation - etc

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