

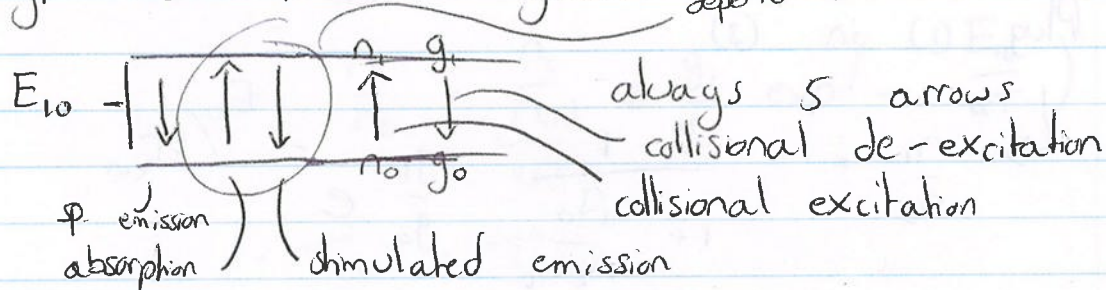
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## Collisional excitation and nebular diagnostics

- Why is eg O III as b almost as bright as H $\alpha$  in many H II regions, whereas the abundance is much lower?

Setting up excitation problem. Book goes into more detail.

Hypothetical two level system depend on radiation field



In some cases, collisions where not important. Generally need 5 however. Excitations and de-excitations are caused by collisions with other particles

$n_c$  = density of collision partner. What that is depends on physics:  $H_2$ ,  $HI$ ,  $e^-$

Upwards collisional excitation  $k_{01} n_c$   
Down " " " " de- " "  $k_{10} n_c$

Excitation costs energy which comes from kinetic energy of gas.

QM : cross section for process.

For most things, cross sections and collision coefficients are well known.

$k_{10}$  is going to depend on temperature, naturally  
 since low  $T$  has low  $E_{kin}$ .

$$(1) \quad k_{01} = k_{10} e^{-E_{10}/k_B T} \frac{g_1}{g_0}$$

$$(2) \quad \frac{dn_1}{dt} = -n_1 A_{10} + n_0 n_c k_{01} - n_1 n_c k_{10}$$

$$(3) \rightarrow \frac{n_1}{n_0} = \frac{n_c k_{01}}{n_c k_{10} + A_{10}}$$

Plug (1) in (3)

$$\rightarrow \frac{n_1}{n_0} = \frac{1}{1 + \frac{A_{10}}{n_c k_{10}}} \frac{g_1}{g_0} e^{-E_{10}/k_B T_{kin}}$$

If no spontaneous decays, this is exactly the Boltzmann distribution.

$$n_{crit} = \frac{A_{10}}{k_{10}} \quad [s^{-1}]$$

$\downarrow [cm^3 s^{-1}]$

$$\frac{n_1}{n_0} = \frac{1}{1 + \frac{n_{crit}}{n_c}} e^{-E_{10}/k_B T_{kin}}$$

Property of molecule but depends on collision partner  
~~Dreue~~ Refers to transition and partner and Temp,  
 but not on density.



Limiting cases

$$n_c \gg n_{\text{crit}} : \frac{n_1}{n_0} = \frac{g_1}{g_0} \exp\left(-\frac{E_{10}}{kT}\right)$$

- $kT_{\text{ex}} = T_{\text{kin}}$
- levels are thermalized
- independent of  $n_c$

~~$n_c \gg n_{\text{crit}}$~~

$$n_c \ll n_{\text{crit}}$$

$$\frac{n_1}{n_0} = \underbrace{\frac{n_c}{n_{\text{crit}}}}_{\ll 1} \frac{g_1}{g_0} \exp\left(-\frac{E_{10}}{kT}\right)$$

- $T_{\text{ex}} < T_{\text{kin}}$
- subthermal excitation
- depends on density of collision partner

If we know  $n_c$  and  $n$ , we know in what regimes we are.

$n_{\text{crit}}$  ranges from small (e.g. C II  $2 \times 10^4$ ) to large (He V  $1.8 \cdot 10^7$ ). These densities are however found in the ISM.

Now we add radiation to this, since there are situations where this is essential.

$$n_\gamma = \frac{c^2}{2h\nu^3} I_\nu$$

direction - averaged,

$$\bar{n}_\gamma = \frac{c^2}{8\pi h\nu^3} u_\nu$$

Black body : Planck function

photon occupation number

$$n_\gamma = \frac{1}{\exp(h\nu/kT_r) - 1}$$

Define  $\rightarrow$  this temperature as our  $T_r$ .

$$\frac{dn_i}{dt} = \underbrace{\left[ n_0 \bar{n}_\gamma \frac{g_i}{g_0} \right]}_{\text{absorption}} - \underbrace{\left[ n_i (1 + \bar{n}_\gamma) \right]}_{\text{sp. emission}} A_0$$

$$+ \underbrace{n_0 n_c k_{0i}}_{\text{col ex}} - \underbrace{n_i n_c k_{i0}}_{\text{col de-ex.}} \quad \text{stimulated emission}$$

Need this result in mol. clouds (and) when lines are optically thick (because of absorption).

$$= A_0 n_0 \left( n_c k_{0i} + \bar{n}_\gamma \frac{g_i}{g_0} A_{0i0} \right) - n_i \left( n_c k_{i0} + (1 + \bar{n}_\gamma) A_{0i0} \right)$$

Steady state thus above is zero :

$$\rightarrow \frac{n_i}{n_0} = \frac{n_c k_{0i} + \bar{n}_\gamma \frac{g_i}{g_0} A_{0i0}}{n_c k_{i0} + (1 + \bar{n}_\gamma) A_{0i0}}$$



Now the critical density in a more general way

$$n_{\text{crit}} = \frac{(1 + \bar{n}_\gamma) A_{10}}{k_{10}}$$

$$\frac{n_1}{n_0} = \frac{1}{1 + \frac{n_{\text{crit}}}{n_c}} \frac{g_1}{g_0} \exp\left(-\frac{E_{10}}{kT_{\text{kin}}}\right) + \frac{1}{1 + \left(\frac{n_{\text{crit}}}{n_c}\right)^{-1}} \frac{g_1}{g_0} \exp\left(-\frac{E_{10}}{kT_r}\right)$$

Again limiting cases

$\bar{n}_\gamma \gg 1$  (strong radiation field)

↳ if  $n_c \ll n_{\text{crit}}$  :  $T_{\text{ex}} = T_r$

HI (21 cm) line :  $\frac{E_{10}}{k} = 0.0682 \text{ K}$

$$A_{10} = 2.88 \cdot 10^{-15} \text{ s}^{-1}$$

$$k_{10} = 1.2 \cdot 10^{-10} \text{ cm}^3 \text{ s}^{-1} @ 100 \text{ K}$$

$$T_r = 2.73 \text{ K} + \underbrace{1 \text{ K}}_{\substack{\text{milky way / galaxy} \\ \text{background rad field}}} = 3.73 \text{ K}$$

$$\bar{n}_\gamma = 55$$

$$\rightarrow n_{\text{crit}} = 1.7 \cdot 10^{-3} \text{ cm}^{-3}$$

$$T_s = T_{\text{kin}}$$

## Applications of this

### Nebular Diagnostics

- Upper level temperatures (wrt expected temperatures)
- Critical densities (with respect to expected densities)

Using line with  $n_c < \text{expected density}$ , we are in the thermal regime. Then lines are not thermalized so line ratio only depends on temperature.

Temperature of ionizing star differs a lot from H II region to H II region, however the temperature of H II region is almost always near  $T \sim 10^4 \text{ K}$ .

If  $n \gg n_{crit}$ , population of levels are independent of density. Hence, line ratio of thermalized lines is good temperature probe. Opposite,  $n \ll n_{crit}$ ,  $\frac{n_1}{n_0} = \frac{n_c}{n_{crit}}$ . This again is independent of  $T$ . Ratio of two sub-thermal lines again produce good  $T$  estimator.

SII / OII diagrams for density estimation because one transition sub thermal and one thermalized.

Far infrared lines, very easily excited as  $T_{ex}$  is low. So we probe density probe for OII 4IR lines. If critical densities are different, OII IR lines are good density probes.