



1 Dual-Spin Spacecraft

Spin stabilization is useful, but problematic for two reasons: 1) major axis rule dictates the body must be oblate; 2) the entire spacecraft must spin. To circumvent these issues we add another spinning body.

Consider a spacecraft composed of a platform, denoted \mathcal{P} , and a wheel, denote \mathcal{W} , shown in Fig. 1. The wheel is spinning relative to the platform. Let \mathcal{B} denote the $\mathcal{P} + \mathcal{W}$ system.

Let \mathcal{F}_a be an inertial frame (which is not shown in Fig. 1).

Let \mathcal{F}_p be the platform frame that is fixed in the platform.

Let \mathcal{F}_w be the wheel frame that is fixed in the wheel. Let the origin of \mathcal{F}_w coincide with the center of mass of the wheel.

Let O^w denote the origin of \mathcal{F}_w .

Let c denote the center of mass of the $\mathcal{P} + \mathcal{W}$ system.

Let the inertia of the platform relative to c resolved in \mathcal{F}_p be $\mathbf{J}_p^{\mathcal{P}c}$.

Let the inertia of the wheel relative to w resolved in \mathcal{F}_w be $\mathbf{I}_w^{\mathcal{W}w}$.

Let the inertia of the wheel relative to c resolved in \mathcal{F}_w be $\mathbf{J}_w^{\mathcal{W}c}$.

Let \underline{a} be the unit-length spin axis of the wheel.

The approach to deriving the equations of motion of a dual-spinning spacecraft are as follows. First, we'll find the angular momentum of \mathcal{P} relative to c w.r.t. \mathcal{F}_a expressed in \mathcal{F}_p . Next, we'll find the angular momentum of \mathcal{W} relative to w w.r.t. \mathcal{F}_a expressed in \mathcal{F}_w . Along the way we'll discuss second order tensors. The angular momentum of \mathcal{W} relative to c w.r.t. \mathcal{F}_a expressed in \mathcal{F}_w will then be computed. We will combine the expressions for angular momentum to get the angular momentum of \mathcal{B} relative to c w.r.t. \mathcal{F}_a expressed in \mathcal{F}_p ; this is the total angular momentum of the system. To find the equations of motion the time-rate-of-change of the total angular momentum w.r.t. \mathcal{F}_a will be found.

2 Angular Momentum of the Platform

The angular momentum of \mathcal{P} relative to c w.r.t. \mathcal{F}_a is

$$\underline{h}^{\mathcal{P}c/a} = \underline{\mathcal{F}}_p^T \mathbf{h}_p^{\mathcal{P}c} = \underline{\mathcal{F}}_p^T \mathbf{J}_p^{\mathcal{P}c} \omega_p^{pa}.$$

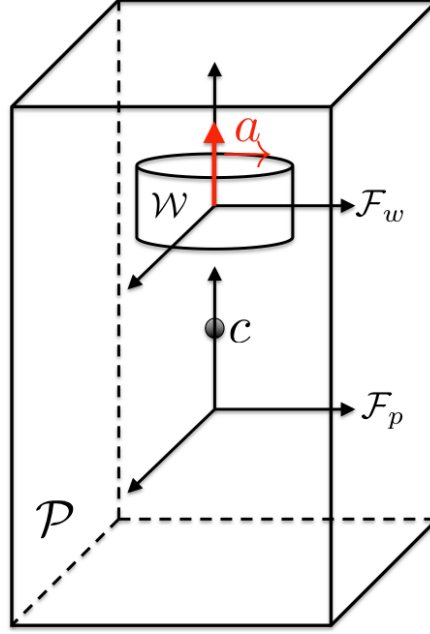


Figure 1: Dual spin spacecraft.

3 Angular Momentum of the Wheel

3.1 Angular Momentum of the Wheel Relative to O^w w.r.t. \mathcal{F}_a

The angular momentum of \mathcal{W} relative to O^w w.r.t. \mathcal{F}_a is

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{W}w/a} &= \int_{\mathcal{V}} \underline{\rho}_{\rightarrow}^{dm w} \times \underline{\rho}_{\rightarrow}^{dm w \bullet} dm \\
 &= \int_{\mathcal{V}} -\underline{\rho}_{\rightarrow}^{dm w} \times (\underline{\rho}_{\rightarrow}^{dm w} \times \underline{\omega}_{\rightarrow}^{wa}) dm \\
 &= \underline{\mathcal{F}}_{\rightarrow w}^T \underline{\mathbf{I}}_w^{\mathcal{W}w} \underline{\omega}_w^{wa} \\
 &= \underline{\mathcal{F}}_{\rightarrow w}^T \underline{\mathbf{h}}_w^{\mathcal{W}w/a},
 \end{aligned}$$

where $\underline{\rho}_{\rightarrow}^{dm w \bullet} = \underline{\omega}_{\rightarrow}^{wa} \times \underline{\rho}_{\rightarrow}^{dm w} = -\underline{\rho}_{\rightarrow}^{dm w} \times \underline{\omega}_{\rightarrow}^{wa}$. Note that

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{W}w/a} &= \underline{\mathcal{F}}_{\rightarrow w}^T \underline{\mathbf{I}}_w^{\mathcal{W}w} \underbrace{\underline{\mathcal{F}}_{\rightarrow w} \cdot \underline{\mathcal{F}}_{\rightarrow w}^T}_{\mathbf{1}} \underline{\omega}_w^{wa} \\
 &= \underbrace{\underline{\mathcal{F}}_{\rightarrow w}^T \underline{\mathbf{I}}_w^{\mathcal{W}w} \underline{\mathcal{F}}_{\rightarrow w}}_{\underline{I}_{\rightarrow}^{\mathcal{W}w}} \cdot \underbrace{\underline{\mathcal{F}}_{\rightarrow w}^T \underline{\omega}_w^{wa}}_{\underline{\omega}_{\rightarrow}^{wa}} \\
 &= \underline{I}_{\rightarrow}^{\mathcal{W}w} \cdot \underline{\omega}_{\rightarrow}^{wa}
 \end{aligned}$$

where $\underline{I}_{\rightarrow}^{\mathcal{W}w}$ is the inertia tensor of \mathcal{W} relative to w .

3.2 Second Order Tensors

An inertia tensor is a second order tensor. A second order tensor, also called a dyadic, maps a first order tensor (i.e., a physical vector) to another first order tensor. Consider $\underline{h}^{\mathcal{W}w} = \underline{I}^{\mathcal{W}w} \cdot \underline{\omega}^{wa}$; $\underline{I}^{\mathcal{W}w}$ maps physical vectors to physical vectors. Just like a physical vector, a second order tensor can be resolved in any frame:

$$\begin{aligned}\underline{I}^{\mathcal{W}w} &= \underline{\mathcal{F}}_{\rightarrow w}^T \mathbf{I}_w^{\mathcal{W}w} \underline{\mathcal{F}}_{\rightarrow w} \\ &= \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{C}_{pw} \mathbf{I}_w^{\mathcal{W}w} \mathbf{C}_{wp} \underline{\mathcal{F}}_{\rightarrow p} \\ &= \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{I}_p^{\mathcal{W}w} \underline{\mathcal{F}}_{\rightarrow p},\end{aligned}$$

where $\mathbf{I}_p^{\mathcal{W}w} = \mathbf{C}_{pw} \mathbf{I}_w^{\mathcal{W}w} \mathbf{C}_{wp}$ and $\underline{\mathcal{F}}_{\rightarrow w}^T = \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{C}_{pw}$. Note that $\mathbf{I}_w^{\mathcal{W}w}$ is the inertia matrix of \mathcal{W} relative to O^w resolved in \mathcal{F}_w , while $\mathbf{I}_p^{\mathcal{W}w}$ is the inertia matrix of \mathcal{W} relative to O^w resolved in \mathcal{F}_p .

3.3 Angular Momentum of the Wheel Relative to the Center of Mass w.r.t. \mathcal{F}_a

The angular momentum of \mathcal{W} relative to c w.r.t. \mathcal{F}_a is

$$\begin{aligned}\underline{h}^{\mathcal{W}c/a} &= \int_{\mathcal{V}} (\underline{r}^{wc} + \underline{\rho}^{dmw}) \times (\underline{r}^{wc\bullet} + \underline{\rho}^{dmw\bullet}) dm \\ &= \int_{\mathcal{V}} \left(\underline{r}^{wc} \times \underline{r}^{wc\bullet} + \underline{r}^{wc} \times \underline{\rho}^{dmw\bullet} + \underline{\rho}^{dmw} \times \underline{r}^{wc\bullet} + \underline{\rho}^{dmw} \times \underline{\rho}^{dmw\bullet} \right) dm \\ &= m \underline{r}^{wc} \times \underline{r}^{wc\bullet} + \int_{\mathcal{V}} \underline{r}^{wc} \times \underline{\rho}^{dmw\bullet} dm + \int_{\mathcal{V}} \underline{\rho}^{dmw} \times \underline{r}^{wc\bullet} dm + \int_{\mathcal{V}} \underline{\rho}^{dmw} \times \underline{\rho}^{dmw\bullet} dm \\ &= m \underline{r}^{wc} \times \underline{r}^{wc\bullet} + \int_{\mathcal{V}} -\underline{r}^{wc} \times (\underline{\rho}^{dmw} \times \underline{\omega}^{wa}) dm \\ &\quad + \int_{\mathcal{V}} \underline{\rho}^{dmw} \times \underline{r}^{wc\bullet} dm + \int_{\mathcal{V}} -\underline{\rho}^{dmw} \times (\underline{\rho}^{dmw} \times \underline{\omega}^{wa}) dm \\ &= m \underline{r}^{wc} \times \underline{r}^{wc\bullet} + \underline{h}^{\mathcal{W}w/a}.\end{aligned}$$

3.4 Expressing the Angular Momentum of the Wheel in the Platform Frame

Consider the angular velocity of \mathcal{F}_w relative to \mathcal{F}_a :

$$\begin{aligned}\underline{\omega}^{wa} &= \underline{\omega}^{wp} + \underline{\omega}^{pa} \\ &= \underline{\mathcal{F}}_{\rightarrow w}^T \mathbf{a}_w \omega_{w,s}^{wp} + \underline{\mathcal{F}}_{\rightarrow p}^T \omega_p^{pa} \\ &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{C}_{pw} \mathbf{a}_w \omega_{w,s}^{wp} + \omega_p^{pa}) \\ &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{a}_p \omega_{p,s}^{wp} + \omega_p^{pa}) \\ &= \underline{\mathcal{F}}_{\rightarrow p}^T \omega_p^{wa}\end{aligned}$$

where $\underline{a}_w = \underline{\mathcal{F}}_{\rightarrow w}^T \mathbf{a}_w = \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{a}_p$ is the unit-length (i.e., $\underline{a}_w \cdot \underline{a}_w = 1$) spin axis of the wheel, $\mathbf{a}_p = \mathbf{C}_{pw} \mathbf{a}_w$, and $\omega_{w,s}^{wp} = \omega_{p,s}^{wp}$ is the wheel speed about the spin axis relative to \mathcal{F}_p expressed in \mathcal{F}_w . Then, the angular

momentum of \mathcal{W} relative to O^w w.r.t. \mathcal{F}_a is

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{W}w/a} &= \underline{\mathcal{F}}_w^T \mathbf{I}_w^{\mathcal{W}w} \omega_w^{wa} \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{C}_{pw} \mathbf{I}_w^{\mathcal{W}w} \mathbf{C}_{wp} \mathbf{C}_{pw} \omega_w^{wa} \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{I}_p^{\mathcal{W}w} \omega_p^{wa} \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{I}_p^{\mathcal{W}w} \mathbf{a}_p \omega_{w,s}^{wp} + \mathbf{I}_p^{\mathcal{W}w} \omega_p^{pa}) \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{I}_p^{\mathcal{W}w} \omega_p^{pa} + \mathbf{a}_p h_{p,s})
 \end{aligned}$$

where $\mathbf{I}_p^{\mathcal{W}w} = \mathbf{C}_{pw} \mathbf{I}_w^{\mathcal{W}w} \mathbf{C}_{wp}$, $\mathbf{C}_{wp} \mathbf{C}_{pw} = \mathbf{1}$, and $\omega_p^{wa} = \mathbf{a}_p \omega_{p,s}^{wp} + \omega_p^{pa}$. Additionally, $h_{p,s} = I_{p,s}^{\mathcal{W}w} \omega_{p,s}^{wp}$ where and $I_{p,s}^{\mathcal{W}w}$ is the inertia of the wheel relative to O^w resolved in \mathcal{F}_p along the spin axis defined by \underline{a}_p . It follows that

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{W}c/a} &= m \underline{r}_{\rightarrow}^{wc} \times \underline{r}_{\rightarrow}^{wc\bullet} + \underline{h}_{\rightarrow}^{\mathcal{W}w/a} \\
 &= \underline{\mathcal{F}}_p^T (-m \mathbf{r}_p^{wc \times} \mathbf{r}_p^{wc \times} \omega_p^{pa}) + \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{I}_p^{\mathcal{W}w} \omega_p^{pa} + \mathbf{a}_p h_{p,s}) \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T ((\mathbf{I}_p^{\mathcal{W}w} - m \mathbf{r}_p^{wc \times} \mathbf{r}_p^{wc \times}) \omega_p^{pa} + \mathbf{a}_p h_{p,s}) \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{J}_p^{\mathcal{W}c} \omega_p^{pa} + \mathbf{a}_p h_{p,s})
 \end{aligned}$$

where $\underline{r}_{\rightarrow}^{wc\bullet} = -\underline{r}_{\rightarrow}^{wc} \times \underline{\omega}_{\rightarrow}^{pa}$ and $\mathbf{J}_p^{\mathcal{W}c} = \mathbf{I}_p^{\mathcal{W}w} - m \mathbf{r}_p^{wc \times} \mathbf{r}_p^{wc \times}$.

4 The Total Angular Momentum of the Body

Consider the angular momentum of \mathcal{B} relative to c w.r.t. \mathcal{F}_a :

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{B}c/a} &= \underline{h}_{\rightarrow}^{\mathcal{P}c} + \underline{h}_{\rightarrow}^{\mathcal{W}c} \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T \mathbf{J}_p^{\mathcal{P}c} \omega_p^{pa} + \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{J}_p^{\mathcal{W}c} \omega_p^{pa} + \mathbf{a}_p h_{p,s}) \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T ((\mathbf{J}_p^{\mathcal{P}c} + \mathbf{J}_p^{\mathcal{W}c}) \omega_p^{pa} + \mathbf{a}_p h_{p,s}) \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{I}_p^{\mathcal{B}c} \omega_p^{pa} + \mathbf{a}_p h_{p,s})
 \end{aligned}$$

where $\mathbf{I}_p^{\mathcal{B}c} = \mathbf{J}_p^{\mathcal{P}c} + \mathbf{J}_p^{\mathcal{W}c}$ is the moment of inertia of \mathcal{B} relative to c resolved in \mathcal{F}_p .

5 The Equations of Motion

The time-rate-of-change of $\underline{h}_{\rightarrow}^{\mathcal{B}c}$ w.r.t. \mathcal{F}_a is

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{B}c/a\bullet} &= \underline{h}_{\rightarrow}^{\mathcal{B}c/a\circ} + \underline{\omega}_{\rightarrow}^{pa} \times \underline{h}_{\rightarrow}^{\mathcal{B}c/a} \\
 &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{I}_p^{\mathcal{B}c} \dot{\omega}_p^{pa} + \mathbf{a}_p \dot{h}_{p,s} + \omega_p^{pa \times} (\mathbf{I}_p^{\mathcal{B}c} \omega_p^{pa} + \mathbf{a}_p h_{p,s})) \\
 &= \underline{\tau}_{\rightarrow}^{\mathcal{B}c},
 \end{aligned}$$

where $(\)^\circ$ denotes time differentiation w.r.t. \mathcal{F}_p . For a constant wheel speed $\dot{h}_{p,s} = 0$, thus

$$\begin{aligned}
 \underline{h}_{\rightarrow}^{\mathcal{B}c/a\bullet} &= \underline{\mathcal{F}}_{\rightarrow p}^T (\mathbf{I}_p^{\mathcal{B}c} \dot{\omega}_p^{pa} + \omega_p^{pa \times} (\mathbf{I}_p^{\mathcal{B}c} \omega_p^{pa} + \mathbf{a}_p h_{p,s})) = \underline{\mathcal{F}}_{\rightarrow p}^T \underline{\tau}_p^{\mathcal{B}c}, \\
 \mathbf{I}_p^{\mathcal{B}c} \dot{\omega}_p^{pa} + \omega_p^{pa \times} (\mathbf{I}_p^{\mathcal{B}c} \omega_p^{pa} + \mathbf{a}_p h_{p,s}) &= \underline{\tau}_p^{\mathcal{B}c}
 \end{aligned}$$

Consider the case where

$$\begin{aligned}
 \boldsymbol{\omega}_p^{pa} &= \text{diag} \{ \omega_1, \omega_2, \omega_3 \}, \\
 \mathbf{I}_p^{\mathcal{B}c} &= \text{diag} \{ I_{p,1}^{\mathcal{B}c}, I_{p,2}^{\mathcal{B}c}, I_{p,3}^{\mathcal{B}c} \} = \text{diag} \{ I_1, I_2, I_3 \}, \\
 \mathbf{I}_p^{\mathcal{W}w} &= \text{diag} \{ I_{p,1}^{\mathcal{W}w}, I_{p,2}^{\mathcal{W}w}, I_{p,3}^{\mathcal{W}w} \} = \text{diag} \{ I_{p,t}^{\mathcal{W}w}, I_{p,t}^{\mathcal{W}w}, I_{p,s}^{\mathcal{W}w} \}, \\
 \mathbf{a}_p &= [0 \ 0 \ 1]^\top, \\
 \boldsymbol{\tau}_p^{\mathcal{B}c} &= \mathbf{0},
 \end{aligned}$$

so that

$$\mathbf{I}_p^{\mathcal{W}w} \mathbf{a}_p = [0 \ 0 \ I_{p,3}^{\mathcal{W}w}]^\top$$

and $I_{p,3}^{\mathcal{W}w} = I_{p,s}^{\mathcal{W}w}$. Then, the dual-spin spacecraft equations of motion become

$$\begin{aligned}
 I_1 \dot{\omega}_1 + (I_3 - I_2) \omega_3 \omega_2 + h_{p,s} \omega_2 &= 0, \\
 I_2 \dot{\omega}_2 + (I_1 - I_3) \omega_1 \omega_3 - h_{p,s} \omega_1 &= 0, \\
 I_3 \dot{\omega}_3 + (I_2 - I_1) \omega_1 \omega_2 &= 0.
 \end{aligned}$$