

# Nonlinear PDE-Based control of the electron temperature in H-mode tokamak plasmas

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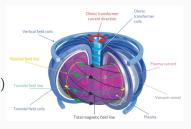
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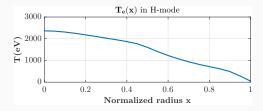
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# INTRODUCTION

### Introduction

- Fixed flux surface geometry, 1D transport.
- Plasma profile control using auxiliary heating sources.
- The High confinement mode (H-mode) is the target (increase  $\tau_E$ ).





#### Introduction

- Data driven model ([Moreau et al., 2008])
- Discretized model for control ([Ou et al., 2008])
- The coupling ([Mavkov et al., 2017])

- First-principle model ([Witrant et al., 2007])
- Distributed control ([Bribiesca Argomedo et al., 2013])
- The nonlinearity of the electron temperature dynamics  $T_e$

Control of the electron temperature profile in TCV<sup>1</sup> H-mode plasmas.

- Infinite-dimensional model: Nonlinear parabolic PDE.
- Lyapunov control function approach.
- Sum of squares technique.

<sup>&</sup>lt;sup>1</sup>Tokamak à Configuration Variable at EPFL in Lausanne, Switzerland.

# SYSTEM DESCRIPTION AND

**CONTROL PROBLEM** 

# **Electron temperature dynamics**

• Electron temperature dynamics  $T_e$ 

$$\frac{3}{2}\frac{\partial(n_eT_e)}{\partial t} = \frac{1}{a^2}\frac{1}{x}\frac{\partial}{\partial x}(xn_e\chi_e\frac{\partial T_e}{\partial x}) + P,\tag{1}$$

with boundary conditions:

$$\frac{\partial T_e}{\partial x}(0,t) = 0, \ T_e(1,t) = T_{e,edge}(t), \forall t \ge t_0$$
 (2)

and the initial condition:

$$T_e(x, t_0) = T_0(x), \forall x \in [0, 1]$$
 (3)

 $\chi_{e}(x,t)$ ; the electron heat diffusivity,

$$P(x,t) = P_{ohm}(x,t) + P_{aux}(x,t) - P_{rad}(x,t)$$

# Electron heat diffusivity model

 Semi-empirical model: the modified bohm/gyrobohm model [Pianroj and Onjun, 2012]:

$$\begin{cases} \chi_{e} = \chi_{e_{c}} \times f_{s}, & \chi_{e_{c}} = (2\chi_{Be} + \chi_{gBe}), \\ \chi_{Be} = 4 \times 10^{-5} R \left| \frac{\nabla (n_{e} T_{e})}{n_{e} B_{\phi_{0}}} \right| q^{2} \left( \frac{T_{e,0.8} - T_{e,1.0}}{T_{e,1.0}} \right), \\ \chi_{gBe} = 5 \times 10^{-6} \sqrt{T_{e}} \left| \frac{\nabla T_{e}}{B_{\phi_{0}}^{2}} \right|, \\ f_{s} = \frac{1}{1 + k \left( \frac{\omega_{E \times B}}{\gamma_{ITG}} \right)^{2}} \times \frac{1}{\max(1, (s - s_{thres})^{2})} \end{cases}$$
(4)

 $f_s$ ; taken as polynomial approximation  $f_s(x)$  of the experimental results obtained in the same paper.

# Control problem

To formulate the control system model, we use the assumptions:

- 1. q(x,t) is constant in time and the density dynamics can be neglected.
- 2.  $\frac{\partial T_e}{\partial x} \leq$  0 always and everywhere and  $T_{e,edge} \ll T_{e,center}$  and to be neglected.

We get:

$$\frac{\partial T_e}{\partial t} = \frac{A}{x} \frac{\partial}{\partial x} \left( x \left( B(x) + C(x) \sqrt{T_e} \right) \left( \frac{\partial T_e}{\partial x} \right)^2 \right) + u, \quad \forall x \in [0, 1]$$
 (5)

with boundary conditions,

$$\frac{\partial T_e}{\partial x}(0,t) = 0, \quad T_e(1,t) = 0, \forall t \ge t_0 \tag{6}$$

where: 
$$A = \frac{2}{3a^2}$$
,  $B(x) = \frac{-8 \times 10^{-5} R L_{T_e}}{B_{\phi_0}} q^2(x) f_s(x)$ ,  $C(x) = \frac{-5 \times 10^{-6}}{B_{\phi_0}^2} f_s(x)$ .

# DISTRIBUTED CONTROL

STABILITY ANALYSIS AND

Nominal stability of  $(\bar{u}(x), \bar{T}_e(x))$  with Lyapunov function candidate:

$$V(T_e) = \frac{1}{2} \int_0^1 x P_{T_e}(x) (T_e - \bar{T}_e)^2 dx, \quad P_{T_e}(x) > 0$$
 (7)

#### **Theorem**

Suppose that for a given  $\alpha_1 > 0$ , there exist a polynomial  $P_{T_e}(x) > 0$  and a  $5 \times 5$  symmetric polynomial matrix H(x) with  $H(0) \ge 0$ ,  $H_{1,1}(1) \le 0$  and:

$$F(x) + \bar{H}(x) \ge 0, \quad \forall x \in [0, 1]$$
(8)

where (8) is a differential matrix inequality.

Then the time derivative  $\dot{V}$  of V along the solutions of (5)-(6) verifies:

$$\dot{V}(T_e) \le -\alpha_1 V(T_e) + \int_0^1 x P_{T_e}(x) (T_e - \bar{T}_e) \tilde{u} \, dx \tag{9}$$

where  $\tilde{u}$  is defined as  $\tilde{u} = u - \bar{u}$ .

proof: along [Valmorbida et al., 2015]

Perform the time derivative of the Lyapunov function, integration by part:

$$\dot{V}(T_e) + \alpha_1 V(T_e) = -\int_0^1 G(x, \sqrt{T_e}, D^1 T_e) dx + \int_0^1 x P_{T_e}(x) (T_e - \bar{T}_e) \tilde{u} dx$$
(10)

Change of variable  $\tau = \sqrt{T_e}$ , we write G in terms of  $\tau$  as:

$$G(x,D^1\tau) = \xi(D^1\tau)^T F(x)\xi(D^1\tau)$$

$$\xi(D^{1}\tau) := \left[1, \tau, \frac{\partial \tau}{\partial x}, \tau^{2}, \tau \frac{\partial \tau}{\partial x}, \left(\frac{\partial \tau}{\partial x}\right)^{2}, \tau^{3}, \tau^{2} \frac{\partial \tau}{\partial x}, \dots\right]$$
$$\tau \left(\frac{\partial \tau}{\partial x}\right)^{2}, \left(\frac{\partial \tau}{\partial x}\right)^{3}, \tau^{4}, \tau^{3} \frac{\partial \tau}{\partial x}, \tau^{2} \left(\frac{\partial \tau}{\partial x}\right)^{2}, \tau \left(\frac{\partial \tau}{\partial x}\right)^{3}\right]^{T}$$

We also have:

$$\int_{0}^{1} \xi(D^{1}\tau)^{T} F(x) \xi(D^{1}\tau) dx = \int_{0}^{1} \left[ \xi(D^{1}\tau)^{T} F(x) \xi(D^{1}\tau) + \frac{d}{dx} \left( \mu(\tau)^{T} H(x) \mu(\tau) \right) \right] dx - \left( \mu(\tau)^{T} H(x) \mu(\tau) \right) \Big|_{0}^{1}$$
(11)

With  $\mu(\tau) := \left[1, \tau, \tau^2, \tau^3, \tau^4\right]^T$ , and finally:

$$\dot{V}(T_e) + \alpha_1 V(T_e) = -\int_0^1 \xi(D^1 \tau)^T \underbrace{(F(x) + \bar{H}(x))}_{\geq 0} \xi(D^1 \tau)$$

$$+ \underbrace{H_{1,1}(1)}_{\leq 0} - \mu(\tau(0))^T \underbrace{H(0)}_{\geq 0} \mu(\tau(0)) + \int_0^1 x P_{T_e}(x) (T_e - \bar{T}_e) \tilde{u} \ dx$$

С

# **Calculation of the weighting Function**

To solve the differential matrix inequality  $F(x) + \bar{H}(x) \ge 0$  over  $\{x | x(1-x) \ge 0\}$  for  $P_{T_e}(x)$ , we cast the problem into a Sum Of Squares problem (SOSP):

Find  $P_{T_e}(x)$ , H(x), N(x) subject to:

- $F(x) + \bar{H}(x) N(x)x(1-x) \in \Sigma^{14\times 14}[x]$ ,
- $N(x) \in \Sigma^{14 \times 14}[x]$ ,
- $H(0) \geq 0, H_{1,1}(1) \leq 0.$

The problem is then formulated as an SDP and solved using Yalmip and the SOS module. The resulting  $P_{T_e}(x)$  is a decreasing polynomial strictly positive on [0,1].

### Distributed control

### Corollary

If the conditions of Theorem 1 are verified, we choose the control input  $u_{ctrl} = \bar{u} + \tilde{u}$ , where  $\tilde{u}$  is calculated to verify the equality:

$$\int_{0}^{1} x P_{T_{e}}(x) (T_{e} - \bar{T}_{e}) \tilde{u} \ dx = -\alpha_{2} V(T_{e}) \tag{12}$$

- Exponential stability of the closed-loop system with convergence rate  $\alpha_1+\alpha_2$
- An explicit control law from (12) is the proportional controller:

$$u_{ctrl} = \bar{u} - \frac{\alpha_2}{2} (T_e - \bar{T}_e) \tag{13}$$

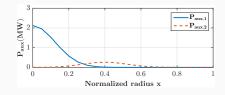
# AND RESULTS

**CONTROL IMPLEMENTATION** 

# **RAPTOR** simulator and Input constraints

- RAPTOR<sup>2</sup> (EPFL) is a lightweight, simplified transport physics code used as a nonlinear plasma simulator and as a real time state observer.
- Radio-Frequency heating: ECRH<sup>3</sup> and ECCD<sup>4</sup> are gaussian distributed.

$$P_{aux,i}(x,t) = \frac{P_i(t)}{\int_0^1 e^z V' dx},$$
  
 $z = \frac{-4(x - x_{dep,i})^2}{w_{dep,i}^2}$ 



<sup>&</sup>lt;sup>2</sup>Rapid Plasma Transport simulatOR.

<sup>&</sup>lt;sup>3</sup>Electron Cyclotron Resonance Heating

<sup>&</sup>lt;sup>4</sup>Electron Cyclotron Current Drive

# Tracking control

• Implementation with the engineering input  $u_{ac} = [P_1(t), P_2(t)]^T$ 

$$u_{ac}^{*} = arg \min_{u_{ac}} \int_{0}^{1} \left[ \underbrace{x P_{T_{e}}(x) \Big( (P_{aux}(u_{ac}) - u_{ref}) + \frac{\alpha_{2}}{2} (T_{e} - T_{e,ref}) \Big)}_{Cf \ Corollary} \right]^{2} dx$$

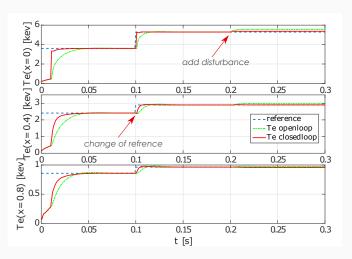
$$(14)$$

subject to:  $0 \le P_1 \le 1 \mathrm{MW}, \ 0 \le P_2 \le 1 \mathrm{MW}.$  where  $(T_{e,ref}, u_{ref})$  are the reference temperature profile and its corresponding input.

• In RAPTOR, we use TCV configuration with H-mode transport model [Kim et al., 2016], with ECRH at  $x_{dep}=0$  and ECCD at  $x_{dep}=0.4$ .

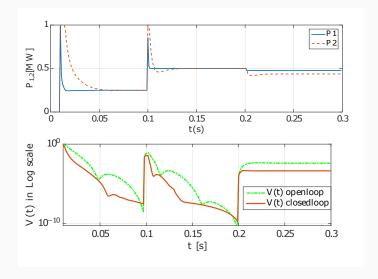
### Simulation results

• Changing the reference profile at  $t=0.1\ s$  and adding a disturbance at  $t=0.2\ s$  (third ECCD source at  $x_{dep}=0.2$  with  $w_{dep}=0.35$  and  $P_3=0.1\ \text{MW}$ ).



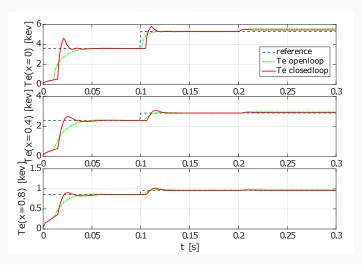
### Simulation results

• Time-evolution of  $P_{1,2}$  and Lyapunov function evaluation:



### Simulation results

 Considering the computation and the transportation of the control signal as a time-delay of 5 ms.



# WORK

**CONCLUSION AND FUTURE** 

### Conclusion

- Stability analysis and control of the electron temperature profile in H-mode tokamak plasmas.
- Using Lyapunov approach and sum of squares techniques.
- RAPTOR plasma simulator was used in H-mode and the control input constraints were taken into account.
- The robustness of the controller was investigated with respect to changing operating point, input disturbances and time-delays.

### **Future work**

- Include the ohmic heating power density  $P_{ohm}(x,t)$  and the radiated power density  $P_{rad}(x,t)$  dynamics to the system.
- Explore central-core-edge diffusivity models that enables internal transport barrier (ITB) [Kim et al., 2016].
- Add the integral action to the analysis to handle the static error.

# THANK YOU!

$F(x) + \tilde{H}(x) =$													
$\lceil f_9 + h_1 \rceil$	h <sub>2</sub>	h <sub>3</sub>	$f_7 + h_4$	h <sub>5</sub>	0	$h_6$	h <sub>7</sub>	0	0	$f_8 + h_8$	hg	$f_5$	0٦
•	$f_7 + h_{10}$	h <sub>11</sub>	h <sub>12</sub>	h <sub>13</sub>	0	$f_8 + h_{14}$	h <sub>15</sub>	$f_5$	0	h <sub>16</sub>	h <sub>17</sub>	f <sub>6</sub>	0
•	•	0	h <sub>18</sub>	0	0	h <sub>19</sub>	f <sub>5</sub>	Õ	0	h <sub>20</sub>	f <sub>6</sub>	ō	0
•	•	•	$f_8 + h_{21}$	h <sub>22</sub>	$f_5$	h <sub>23</sub>	h <sub>24</sub>	$f_6$	0	h <sub>25</sub>	$2h_{27}$	$f_1$	f <sub>3</sub>
•	•	•	•	$f_5$	Õ	h <sub>26</sub>	f <sub>6</sub>	Ō	0	h <sub>27</sub>	$f_1$	f <sub>3</sub>	ő
•	•	•	•	•	0	f <sub>6</sub>	ō	0	0	$f_1$	f <sub>3</sub>	Õ	0
•	•	•	•	•	•	h <sub>28</sub>	h <sub>29</sub>	$f_1$	$f_3$	h <sub>30</sub>	h <sub>31</sub>	$f_2$	f <sub>4</sub>
•	•	•	•	•	•	•	$f_1$	f <sub>3</sub>	Õ	$\frac{3}{4}h_{31}$	f <sub>2</sub>	f <sub>4</sub>	0
	•	•	•	•	•	•	•	0	0	f <sub>2</sub>	f <sub>4</sub>	0	0
•	•	•	•	•	•	•	•	•	0	$f_4$	ō	0	0
	•	•	•	•	•	•	•	•	•	h <sub>32</sub>	h33	0	0
•	•	•	•	•	•	•	•	•	•	•	0	0	0
•	•	•	•	•	•	•	•	•	•	•	•	0	0
	•	•	•	•	•	•	•	•	•	•	•	•	٥٦
(15)													

$$f_{1} = \frac{4}{9}(ABxP'_{T_{e}}), \ f_{2} = \frac{2}{3}(ACxP'_{T_{e}}), \ f_{3} = \frac{4}{5}(ABxP_{T_{e}}), \ f_{4} = (ACxP_{T_{e}}),$$

$$f_{5} = -\frac{4}{9}(ABx(P'_{T_{e}}\bar{T}_{e} + P_{T}\bar{T}_{e}')), \ f_{6} = -\frac{2}{5}(ACx(P'_{T_{e}}\bar{T}_{e} + P_{T}\bar{T}_{e}')),$$

$$f_{7} = \frac{1}{3}(\alpha_{1}xP_{T_{e}}\bar{T}_{e} - xP_{T_{e}}\bar{u}), \ f_{8} = -\frac{1}{10}(\alpha_{1}xP_{T_{e}}),$$

$$f_{9} = (xP_{T_{e}}\bar{T}_{e}\bar{u} - \frac{1}{2}\alpha_{1}xP_{T_{e}}\bar{T}_{e}^{2}). \ h_{1} = \frac{dH_{1,1}}{dx}, \ h_{2} = \frac{dH_{1,2}}{dx}, \ h_{3} = H_{1,2}, \ h_{4} = \frac{2}{3}\frac{dH_{1,3}}{dx},$$

$$h_{5} = H_{1,3}, \ h_{6} = \frac{1}{2}\frac{dH_{1,4}}{dx}, \ h_{7} = H_{1,4}, \ h_{8} = \frac{2}{5}\frac{dH_{1,5}}{dx}, \ h_{9} = H_{1,5}, \ h_{10} = \frac{2}{3}\frac{dH_{1,3}}{dx} + \frac{dH_{2,2}}{dx},$$

$$h_{11} = H_{1,3} + H_{2,2}, \ h_{12} = \frac{1}{2}\frac{dH_{1,4}}{dx} + \frac{dH_{2,3}}{dx}, \ h_{13} = H_{1,4} + 2H_{2,3},$$

$$h_{14} = \frac{2}{5}\frac{dH_{1,5}}{dx} + \frac{dH_{2,4}}{dx}, \ h_{15} = H_{1,5} + 3H_{2,4}, \ h_{16} = \frac{dH_{2,5}}{dx}, \ h_{17} = 4H_{2,5},$$

$$h_{18} = H_{1,4} + H_{2,3}, \ h_{19} = H_{1,5} + H_{2,4}, \ h_{20} = H_{2,5}, \ h_{21} = \frac{2}{5}\frac{dH_{1,5}}{dx} + \frac{dH_{3,3}}{dx},$$

$$h_{22} = H_{1,5} + 2H_{3,3}, \ h_{23} = \frac{dH_{3,4}}{dx}, \ h_{24} = 3H_{3,4}, \ h_{25} = \frac{dH_{3,5}}{dx}, \ h_{26} = 2H_{3,4}, \ h_{27} = 2H_{3,5},$$

$$h_{28} = \frac{dH_{4,4}}{dx}, \ h_{29} = 3H_{4,4}, \ h_{30} = \frac{dH_{4,5}}{dx}, \ h_{31} = 4H_{4,5}, \ h_{32} = \frac{dH_{5,5}}{dx}, \ h_{33} = 4H_{5,5}.$$

$$(16)$$

### References i



Bribiesca Argomedo, F., Witrant, E., Prieur, C., Brémond, S., Nouailletas, R., and Artaud, J.-F. (2013).

Lyapunov-based distributed control of the safety-factor profile in a tokamak plasma.

Nuclear Fusion, 53(3):033005.



Kim, D., Merle, A., Sauter, O., and Goodman, T. (2016). Simple predictive electron transport models applied to sawtoothing plasmas.

Plasma Physics and Controlled Fusion, 58(5):055002.



Mavkov, B., Witrant, E., and Prieur, C. (2017).

Distributed control of coupled inhomogeneous diffusion in tokamak plasmas.

IEEE Transactions on Control Systems Technology, (99):1–8.

### References ii



Moreau, D., Mazon, D., Ariola, M., De Tommasi, G., Laborde, L., Piccolo, F., Sartori, F., Tala, T., Zabeo, L., Boboc, A., et al. (2008).

A two-time-scale dynamic-model approach for magnetic and kinetic profile control in advanced tokamak scenarios on JET. Nuclear Fusion, 48(10):106001.



Ou, Y., Xu, C., Schuster, E., Luce, T., Ferron, J., Walker, M., and Humphreys, D. (2008).

Design and simulation of extremum-seeking open-loop optimal control of current profile in the diii-d tokamak.

Plasma Physics and Controlled Fusion, 50(11):115001.

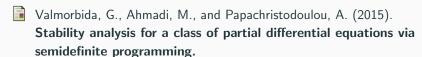


Pianroj, Y. and Onjun, T. (2012).

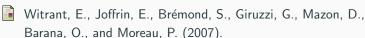
Simulations of H-mode plasmas in tokamak using a complete core-edge modeling in the BALDUR code.

Plasma Science and Technology, 14(9):778.

### References iii



IEEE Transactions on Automatic Control, 61(6):1649–1654.



A control-oriented model of the current profile in tokamak plasma.

Plasma Physics and Controlled Fusion, 49(7):1075.