

1 Introduction

Let A denote the one-dimensional Schrödinger operator with a periodic delta potential at $(x_k)_{k \geq 1}$ and let Ω be a fundamental domain of periodicity associated with the periodic delta potential. Now assuming w.l.o.g. that $x_0 = 0 \in \Omega$ and $|\Omega| = 1$ yields

$$A_0 u := A|_{\Omega} u = -\Delta u + \delta(x_0)u, \quad \text{for } u \in \text{dom} A_0$$

For fixed $k \in \mathbb{R}$, we are interested in the spectral problem of the weak formulation of our differential equation such that the main focus point is going to be the following equation

$$\int_{\Omega} u' \varphi' - \mu \int_{\Omega} u \varphi + cu(x_0) \varphi(x_0) = \int_{\Omega} f \varphi \quad \forall \varphi \in H_k^1$$

for a $u \in H_k^1 := \{u \in H^1(\Omega) : u(\frac{1}{2}) = e^{ik}u(-\frac{1}{2})\}$ and for all $\varphi \in H_k^1$

$$\begin{aligned} \text{dom} A_0 = \left\{ u \in H^2(-\frac{1}{2}, 0) \cap H^2(0, \frac{1}{2}), u(-0) = u(+0), \right. \\ \left. u'(-0) - u'(+0) + cu(0) = 0, u(\frac{1}{2}) = e^{ik}u(-\frac{1}{2}), u'(\frac{1}{2}) = e^{ik}u'(-\frac{1}{2}) \right\} \end{aligned}$$

The boundary conditions on $\partial\Omega$, so-called semi-periodic boundary conditions, are stated here to later extended eigenfunctions on the whole of \mathbb{R} by those conditions.

Subsequently it is shown that A_k is well defined on its domain, that given the differential equation with a bounded $f \in L^2$ there exists a unique solution u in the domain for all $f \in L^\infty$ and that there is an orthonormal basis of eigenfunction to the corresponding eigenvalue problem.

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Lemma 2.1. *The operator A is well-defined on its domain.*

Proof. *To show that A_0 is well-defined on its domain consider $u \in \text{dom}A$ and an arbitrary $v \in H_k^1$*

$$\int_{\Omega} (A_0 u - \mu u) \bar{v} = \int_{\Omega} u' \bar{v}' - \mu \int_{\Omega} u \bar{v} + cu(x_0) \overline{v(x_0)} \quad (*)$$

Now choosing a specific v such that $\text{supp } v = (-\frac{1}{2}, x_0)$ alters $()$ to*

$$\int_{-\frac{1}{2}}^{x_0} u' v' dx = \int_{-\frac{1}{2}}^{x_0} Au \bar{v} \iff \int_{-\frac{1}{2}}^{x_0} u \bar{v}'' dx = \int_{-\frac{1}{2}}^{x_0} -Au \bar{v} dx$$

$\Rightarrow u'' = -Au \in L^2$ on $(-\frac{1}{2}, x_0)$ and an analogously on $(x_0, \frac{1}{2})$. One can therefore fix

$$\text{dom}A \subset \{u \in H^2(-\frac{1}{2}, x_0), u \in H^2(x_0, \frac{1}{2})\}$$

Next, integrating $()$ on both sides of x_0 by parts yields*

$$\begin{aligned} - \left(\int_{-\frac{1}{2}}^{x_0} + \int_{x_0}^{\frac{1}{2}} \right) u'' \bar{v} + (u'(x_0-0)v(x_0-0) - u'(x_0+0)v(x_0+0)) + cu(x_0)\bar{v}(x_0) \\ = - \int_{-\frac{1}{2}}^{x_0} u'' v - \int_{x_0}^{\frac{1}{2}} u'' v \end{aligned}$$

But as $v \in H_k^1 \subseteq C$ this is equivalent

$$u'(x_0-0) - u'(x_0+0) + cu(x_0) = 0$$

Such that

$$\text{dom}A \subset \{u \in H^2(-\infty, x_0) \cap H^2(x_0, \infty), u'(x_0-0) - u'(x_0+0) + cu(x_0) = 0\} =: B$$

The opposite inclusion one shows by taking an arbitrary $u \in B$ and proving that u is in the domain of A_0 . However¹, as $\text{range}R_{\mu,0} = \text{dom}A_0$ one can also show that $u = \mathcal{R}(R_{\mu,0})$. As $\text{dom}R_{\mu,0} = L^2$ define

$$f := A_0 u = \begin{cases} -u'', & (-\frac{1}{2}, x_0) \\ -u'', & (x_0, \frac{1}{2}) \end{cases}$$

¹notice, $R_{\mu,0} := (A_0 - \mu I)^{-1}$ denotes the resolvent

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and show that $u = R_\mu(f - \mu u)$:

$$\begin{aligned} \int_{\Omega} (f - \mu u)v &= \int_{\Omega} u'v' + cu(x_0)v(x_0) - \mu \int_{\Omega} uv \\ \iff - \int_{-\frac{1}{2}}^{x_0} u''v - \int_{x_0}^{\frac{1}{2}} u''v &= \int_{\Omega} u'v' + cu(x_0)v(x_0) \end{aligned}$$

Partial integration with $\text{supp } v = (-\frac{1}{2}, \frac{1}{2})$ yields

$$\begin{aligned} \int_{-\frac{1}{2}}^{x_0} u'v' + \int_{x_0}^{\frac{1}{2}} u'v' - u'(x_0-0)v(x_0) + u'(x_0+0)v(x_0) &= \int_{\Omega} u'v' + cu(x_0)v(x_0) \\ \iff u(x_0+0)v(x_0) - u(x_0-0)v(x_0) - cu(x_0)v(x_0) &= 0 \end{aligned}$$

Such that $u \in B$.

2.1 A_k is a self-adjoint operator

Last but not least, to show that A_k is self-adjoint, we focus first on $R_{\mu,k}^{-1}$ which is given by

$$R_{\mu,k}(A)^{-1} = (A - \lambda I)$$

First one has to notice that $R_{\mu,k}^{-1}$ is symmetric, as $\forall v \in H_k^1$:

$$\begin{aligned} \langle R_{\mu,k}^{-1}u, v \rangle &= \langle (A - \lambda I)u, v \rangle \\ &= \int (A - \lambda I)(u)v dx \\ &= \int u'v' dx - \int \lambda uv dx + cu(x_0)v(x_0) \\ &= \int u(A - \lambda I)(v) dx \\ &= \langle u, (A - \lambda I)v \rangle = \langle u, R_{\mu,k}^{-1}v \rangle \end{aligned}$$

Now as $\text{dom} R_{\mu,k} = L^2(\mathbb{R})$ and $\text{range} R_{\mu,k} = \text{dom} R_{\mu,k}^{-1}$, we want to show that for each $f, g \in L^2$

$$\langle R_{\mu,k}f, g \rangle - \langle f, R_{\mu,k}g \rangle = \gamma$$

$\gamma = 0$. Now there are $u, v \in \text{dom} A_k$ with $Rf = u, Rg = v$ applying to A_k to u, v one gets for all $\varphi, \psi \in H_k^1$

$$\begin{aligned} \int u'\varphi' + cu(0)\varphi(0) - \mu \int u\varphi &= \int f\varphi \\ \int v'\psi' + cv(0)\psi(0) - \mu \int v\psi &= \int g\psi \end{aligned}$$

As it has to hold for all $\varphi, \psi \in H_k^1$ the special choice of $\varphi = v$ and $\psi = u$ yields $\gamma = 0$ and $R_{\mu,k}$ is therefore symmetric.

All in all we can use this to show that $\mathbb{R}_{\mu,k}$ is self-adjoint, as we get for an arbitrary $v^* \in \text{domain} R_{\mu,k}^{-1}$ there exists a $v \in \text{dom} R_{\mu,k}$:

$$\begin{aligned} \langle u, v^* \rangle &= \langle R_{\mu,k}^{-1}R_{\mu,k}u, v^* \rangle \\ &= \langle R_{\mu,k}u, (R_{\mu,k}^{-1})v^* \rangle \\ &= \langle R_{\mu,k}u, v \rangle = \langle u, R_{\mu,k}v \rangle \end{aligned}$$

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So $v^* \in \text{range} R_{\mu,k} = \text{dom} R_{\mu,k}^{-1}$ with that also $R_{\mu,k}^{-1}$ is self-adjoint and as A_k is simply $R_{\mu,k}^{-1}$ shifted by the real constant μ , $R_{\mu,k}^{-1}$ is self-adjoint as well.

□

2.2 A_k being compact

Let $B_{H_k^1} = \{f \in H_k^1(\Omega) : \|f\| \leq 1\}$. We want to show that $\forall \epsilon > 0 \exists g_1, \dots, g_{n_\epsilon}$:

$$\forall f \in B \exists g \in \{g_1, \dots, g_{n_\epsilon}\} : \|f - g\| \leq \epsilon$$

Together with the closure of H_k^1 this yields the compact embedding. Now, as $H^1(\Omega) \subset C(\Omega)$:

$$|f(x) - f(y)| \leq c|x - y|^{\frac{1}{2}} \text{ for some } c > 0 \quad (*)$$

Now, for a $f \in B_{H^1}$ follows from $(*)$ that

$$|f(x)|^2 \leq 2\|f\|_{L^2}^2 + 2 \leq 4 \quad \forall x \in \Omega$$

And with that we can approximate a $f \in B$ by simple functions through partitioning Ω into n_ϵ equidistant intervals. As our simple function is constant on each subinterval, we chose this constant c_k such that

$$|f(\frac{k}{n}) - c_{k+1}| < \frac{1}{n}$$

such that

$$\begin{aligned} \|f - g\|_{L^2}^2 &= \sum_{k=0}^{n-1} \int_{\frac{k}{n}}^{\frac{k+1}{n}} |f - c_{k+1}|^2 dx \\ &= 2 \sum_{k=0}^{n-1} \int_{\frac{k}{n}}^{\frac{k+1}{n}} |f - f(\frac{k}{n})|^2 dx + 2 \sum_{k=0}^{n-1} \int_{\frac{k}{n}}^{\frac{k+1}{n}} |f(\frac{k}{n}) - c_{k+1}|^2 dx \\ &\leq 2 \sum_{n=0}^{n-1} \frac{1}{n^2} + 2 \sum_{n=0}^{n-1} \frac{1}{n^3} = \frac{2}{n} + \frac{2}{n^2} < \epsilon^2 \text{ for } n \text{ small enough.} \end{aligned}$$

2.3 Existence of a unique solution

Next, one has to show that $R_{\mu,k}$ is well-defined, which means that for all $u \in \text{dom}(A_k)$ there exists a unique solution. Lets assume that f , as the righthand-side of the given differential equation, is a bounded linear functional. Lax-Milgram's theorem² would indeed guarantee the existence and uniqueness to prove the existence of a unique solution and whereby $R_{\mu,k}$ being well-defined, but one has to show that H_k^1 is a Hilbert space, for

$$\begin{aligned} B(u, \varphi) &:= \langle \nabla u, \nabla \varphi \rangle + cu(x_0)\varphi(x_0) - \mu \langle u, \varphi \rangle \\ l(\varphi) &:= \langle f, \varphi \rangle \end{aligned}$$

B is bounded and $B[u, u]$ is coercive.

2.3.1 H_k^1 being a Hilbert space

H_k^1 is evidently a subspace of the Hilbert space $H^1(\Omega)$, but additionally H_k^1 is also closed, as for an arbitrary sequence $(\psi_j)_{j \geq 1} \in H_{1,k}$ the value on the boundary coincides. Define $f := \psi_j - \lim \psi_j$ and then

$$\begin{aligned} |f(-\frac{1}{2})|^2 &= 2|f(x)|^2 + 2 \left(\int_{-\frac{1}{2}}^x f'(\tau) d\tau \right)^2 \\ &\leq 2|f(x)|^2 + 2 \int_{-\frac{1}{2}}^{\frac{1}{2}} |f'|^2 d\tau \\ &\leq 2\|f\|_{H^1(-\frac{1}{2}, \frac{1}{2})}^2 \end{aligned}$$

With that $\psi \in H_k^1$ as

$$\psi(-\frac{1}{2}) = \lim_{j \rightarrow \infty} \psi_j(-\frac{1}{2}) = \lim_{j \rightarrow \infty} e^{ik} \psi_j(\frac{1}{2}) = e^{ik} \psi(\frac{1}{2})$$

2.3.2 The bilinear form B is bounded

$$\begin{aligned} |B(u, \varphi)| &:= |\langle \nabla u, \nabla \varphi \rangle + cu(x_0)\varphi(x_0) - \mu \langle u, \varphi \rangle| \\ &\stackrel{\text{Schwarz's}}{\leq} \|\nabla u\| \cdot \|\nabla \varphi\| + cu(x_0)\varphi(x_0) - \mu \|u\| \cdot \|\varphi\| \\ &\stackrel{\text{Inequality}}{\leq} \end{aligned}$$

²formulation and prove in appendix A

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Again we require $H^1(\mathbb{R}) \Subset C(\mathbb{R})$, we can estimate the modulus of $v(x_0) \in \{u(x_0), \varphi(x_0)\}$ over the periodicity cell I_k :

$$\begin{aligned} |v(x_0)|^2 &= \left| v(x) + \int_x^{x_0} \nabla v(\tau) d\tau \right|^2 \quad \text{for an arbitrary } x \in (\inf I, x_0) \\ &\stackrel{\text{convexity}}{\leq} 2|v(x)|^2 + 2 \left| \int_x^{x_0} \nabla v(\tau) d\tau \right|^2 \\ &\stackrel{\text{trace}}{\leq} 2|v(x)|^2 + 2 \int_I |\nabla v(\tau)|^2 d\tau \cdot (x_0 - x) \end{aligned}$$

Integrating both sides over the interval I yields:

$$\begin{aligned} |v(x_0)|^2 \cdot |I| &= 2 \int_I |v(x)|^2 dx + 2 \int_I |\nabla v(\tau)|^2 d\tau \cdot |I| \cdot (x_0 - x) \\ \Rightarrow |v(x_0)|^2 &= \frac{2}{|I|} \int_I |v(x)|^2 dx + 2 \int_I |\nabla v(\tau)|^2 d\tau \cdot \underbrace{(x_0 - x)}_{\leq |I|} \end{aligned}$$

and results in the following

$$\begin{aligned} |B(u, \varphi)| &\leq \| \nabla u \| \cdot \| \nabla \varphi \| + c \cdot u(x_0) \varphi(x_0) - \mu \| u \| \cdot \| \varphi \| \\ &\leq \| \nabla u \| \cdot \| \nabla \varphi \| + c \left(u(x_0)^2 \varphi(x_0)^2 \right)^{\frac{1}{2}} - \mu \| u \| \cdot \| \varphi \| \\ &= \| \nabla u \| \cdot \| \nabla \varphi \| + 2c \left(\frac{1}{|I|} \| u \|^2 + \| \nabla u \|^2 \cdot |I| \right)^{\frac{1}{2}} \\ &\quad \cdot \left(\frac{1}{|I|} \| \varphi \|^2 + \| \nabla \varphi \|^2 \cdot |I| \right)^{\frac{1}{2}} - \mu \| u \| \cdot \| \varphi \| \\ &= \left| (1 + 2c \cdot |I|) \cdot \| \nabla u \| \cdot \| \nabla \varphi \| + \left(\frac{2c}{|I|} - \mu \right) \cdot \| u \| \cdot \| \varphi \| \right. \\ &\quad \left. + 2c (\| u \| \cdot \| \nabla \varphi \| + \| \nabla u \| \cdot \| \varphi \|) \right| \\ &\leq \alpha \cdot \| u \|_{H^1} \cdot \| \varphi \|_{H^1} \end{aligned} \quad \square$$

2.3.3 $B[u, u]$ is coercive

Next, the coercivity for $c \geq 0$ and as assumed at the start μ is small enough, here $\mu < -1$

$$\begin{aligned} B(u, u) &= \langle \nabla u, \nabla u \rangle + cu(x_0)^2 - \mu \langle u, u \rangle \\ &\geq \langle \nabla u, \nabla u \rangle - \mu \langle u, u \rangle \\ &\geq \langle \nabla u, \nabla u \rangle + \langle u, u \rangle \\ &= \|u\|_{H^1}^2 \end{aligned}$$

and for $c < 0$

$$\begin{aligned} B(u, u) &= \langle \nabla u, \nabla u \rangle + c|u(x_0)|^2 - \mu \langle u, u \rangle \\ &= \langle \nabla u, \nabla u \rangle + c \left(\frac{2}{I} \int_I |u(x)|^2 dx + 2I \int_I |\nabla u(\tau)|^2 d\tau \right) - \mu \langle u, u \rangle \\ &= (1 + 2cI) \|\nabla u\|^2 + (-1 + c\frac{2}{I}) \|u\|^2 \\ &\geq \beta \|u\|_{H^1}^2 \end{aligned} \quad \square$$

All in all, Lax-Milgram's theorem now guarantees a unique element $u \in H$ such that

$$B(u, v) = l(\varphi)$$

for all $\varphi \in H_k^1$

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3.1 The inverse of a self-adjoint operator

If $T \in B(X, Y)$ is invertible, where X, Y are Hilbert spaces, then T^* has an inverse and $(T^*)^{-1} = (T^{-1})^*$

Proof. Let $T \in B(X, Y)$ be invertible, notice that $\langle Tv, u \rangle = \langle v, T^*u \rangle$ for all $v \in X, u \in Y$. Then $\langle T^*(T^{-1})^*v, u \rangle = \langle (T^{-1})^*v, Tu \rangle = \langle v, T^{-1}Tu \rangle = \langle u, v \rangle$.

Therefore $T^*(T^{-1})^* = I$, hence $(T^{-1})^* = (T^*)^{-1}$

3.2 Lax-Milgram

Let H be a real Hilbert space, with norm $\|\cdot\|$ and inner product $\langle \cdot, \cdot \rangle$ as well as the pairing of H with its dual space. Assume that

$$B: H \times H \rightarrow \mathbb{R}$$

is a bilinear mapping, for which there exist constant $\alpha, \beta > 0$ such that

$$|B[u, v]| \leq \alpha \|u\| \|v\| \quad (u, v \in H)$$

and

$$\beta \|u\|^2 \leq B[u, u] \quad (u \in H)$$

Finally, let $f: H \rightarrow \mathbb{R}$ be a bounded linear functional on H .

Then there exists a unique element $u \in H$ such that

$$B[u, v] = \langle f, v \rangle$$

for all $v \in H$.

Proof. For each fixed element $u \in H$, the mapping $v \mapsto B[u, v]$ is a bounded linear functional on H ; whence the Riesz Representation Theorem asserts the existence of a unique element $w \in H$ satisfying

$$B[u, v] = \langle w, v \rangle \quad (*)$$

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Let us write $Au = w$ whenever $(*)$ holds; so that

$$B[u, v] = \langle Au, v \rangle \quad (u, v \in H)$$

We first claim $A: H \rightarrow H$ is a bounded linear operator. Indeed if $\lambda_1, \lambda_2 \in \mathbb{R}$ and $u_1, u_2 \in H$, we see for each $v \in H$ that

$$\begin{aligned} \langle A(\lambda_1 u_1 + \lambda_2 u_2), v \rangle &= B[\lambda_1 u_1 + \lambda_2 u_2, v] \quad (\text{by } (*)) \\ &= \lambda_1 B[u_1, v] + \lambda_2 B[u_2, v] \\ &= \lambda_1 \langle Au_1, v \rangle + \lambda_2 \langle Au_2, v \rangle \quad (\text{by } (*) \text{ again}) \\ &= \langle \lambda_1 Au_1 + \lambda_2 Au_2, v \rangle. \end{aligned}$$

This equality obtains for each $v \in H$, and so A is linear. Furthermore

$$\|Au\|^2 = \langle Au, Au \rangle = B[u, Au] \leq \alpha \|u\| \|Au\|.$$

Consequently $\|Au\| \leq \alpha \|u\|$ for all $u \in H$, and so A is bounded.

Next we assert

$$\begin{cases} A \text{ is one-to-one, and} \\ R(A), \text{ the range of } A, \text{ is close in } H. \end{cases} \quad (+)$$

To prove this, let us compute

$$\beta \|u\|^2 \leq B[u, u] = \langle Au, u \rangle \leq \|Au\| \|u\|$$

Hence $\beta \|u\| \leq \|Au\|$. This inequality easily implies $(+)$.

We demonstrate now

$$R(A) = H \quad (-)$$

For if not, then, since $R(A)$ is closed, there would exist a nonzero element $w \in H$ with $w \in R(A)^\perp$. But this fact in turn implies the contradiction $\beta \|w\|^2 \leq B[w, w] = \langle Aw, w \rangle = 0$.

Next, we observe once more from the Riesz' Representation Theorem that

$$\langle f, v \rangle = \langle w, v \rangle \text{ for all } v \in H$$

for some element $w \in H$. We then utilise $(+)$ and $(-)$ to find $u \in H$ satisfying $Au = w$. Then

$$B[u, v] = \langle Au, v \rangle = \langle w, v \rangle = \langle f, v \rangle (v \in H)$$

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and this is the claim.

Finally, we show there is at most one element $u \in H$ verifying the claim. For if both $B[u, v] = \langle f, v \rangle$ and $B[\tilde{u}, v] = \langle f, v \rangle$, then $B[u - \tilde{u}, v] = 0$ ($v \in H$). We set $v = u - \tilde{u}$ to find $\beta \|u - \tilde{u}\|^2 \leq B[u - \tilde{u}, u - \tilde{u}] = 0$.

3.3 Sobolev Embedding

For $s > d/2$, the follow holds

$$H^s \subset C_b(\mathbb{R}^d)$$

Proof. For $u \in S(\mathbb{R}^d)$.

$$\begin{aligned} u(x) &= \int_{\mathbb{R}^d} e^{2i\pi x\xi} (1 + |\xi|^2)^{-s/2} \hat{u} (1 + |\xi|^2)^{s/2} d\xi \\ &\leq \left(\int (1 + |\xi|^2)^{-s} d\xi \right)^{1/2} \left(\int |\hat{u}|^2 (1 + |\xi|^2)^s d\xi \right)^{1/2} \end{aligned}$$

with $(\int (1 + |\xi|^2)^{-s} d\xi)^{1/2} \leq \infty$ und $(\int |\hat{u}|^2 (1 + |\xi|^2)^s d\xi)^{1/2} = \|u\|_{H^s}$.

3.4 $\text{dom } \mathbf{A}_k = \text{range } \mathbf{R}_{\mu,k}$

Proof. As we introduced the variational problem

$$\forall v \in H^1(\mathbb{R}) : \quad \int \nabla u \overline{\nabla v} dx - \mu \int u \bar{v} dx + \alpha u(x_0) v(x_0) = \int f \bar{v} dx \quad (1)$$

$\exists_1 u \in H^1(\mathbb{R})$ satisfying (1)

$$L^2(\mathbb{R}) \ni f \mapsto u =: R_\mu f$$

For $f_1 \neq f_2 \Rightarrow u_1 \neq u_2$, since:

$$\text{Suppose } u_1 = u_2 \Rightarrow \int (f_1 - f_2) \bar{v} = 0 \quad \forall v \in H^1(\mathbb{R}) \Rightarrow f_1 = f_2$$

$\underbrace{\hspace{10em}}_{\text{and therefore } \forall v \in L^2(\mathbb{R})}$

Since H^1 is dense in $L^2 \Rightarrow f_1 = f_2$

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$$\Rightarrow \left. \begin{array}{l} f = R_{\mu}^{-1}u \\ Au - \mu u \end{array} \right\} Au = R_{\mu}^{-1}u + \mu u$$

$$\Rightarrow \operatorname{dom} A = \operatorname{range} R_{\mu}$$