## Magnetostatics

## Michael Brodskiy

Professor: D. Wood

October 19, 2023

## • Magnetostatics:

$$\vec{F}_{lorentz} = q\vec{v} \times \vec{B}$$

- This is "static" in the sense of steady flow of magnetic field
- The units of  $\vec{B}$  are Teslas [T]
- In the simple case where  $\vec{B} = B_o \hat{\mathbf{z}}$  and  $v_z = 0$ , the force could be described as:

$$|\vec{F}| = qvB$$

- With inward direction. We can then write:

$$qvB = \frac{mv^2}{R}$$
$$qB = \frac{mv}{R}$$
$$qB = \frac{p}{R}$$
$$R = \frac{p}{qB}$$

- The frequency of rotation may be written:

$$\omega = \frac{qB}{m}$$

- Total force may be written as:

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B} = q(\vec{E} + \vec{v} \times \vec{B})$$

\* Special Case:  $\vec{F} = 0$ 

$$\vec{E} + \vec{v} \times \vec{B} = 0$$

\* Given  $\vec{E} = E_o \hat{\mathbf{z}}$  and  $\vec{B} = B_o \hat{\mathbf{x}}$ , we would get:

$$E_o\mathbf{\hat{z}} + v_z B_o\mathbf{\hat{y}} - v_y B_o\mathbf{\hat{z}} = 0$$

- \* From this, we get  $v_z=0$  and  $v_y=\frac{E_o}{B_o}$
- \* This can be used to construct a velocity selector
- \* In the event that  $\vec{E} \parallel \vec{B}$ , an expand or contracting helix about the fields as a pole would be constructed
- We now go back to a general case:  $\vec{E} \perp \vec{B}$ ,  $\vec{F}_{net} \neq 0$ ,  $\vec{E} = E\hat{\mathbf{z}}$ , and  $\vec{B} = B\hat{\mathbf{x}}$

$$\frac{1}{m}\frac{d\vec{v}}{dt} = q(\vec{E} + \vec{v} \times \vec{B})$$

\* The solutions for this would look as follows:

$$y(t) = c_1 \cos(\omega t) + c_2 \sin(\omega t) + \frac{E}{B}t + c_3$$
$$z(t) = c_2 \cos(\omega t) + c_1 \sin(\omega t) + c_4$$
$$\omega = \frac{qB}{m}$$

- \* This shape is known as a cycloid
- \* Special case: starting from rest at (0,0); this would give us:

$$c_1 = 0 c_2 = -\frac{E}{\omega B} c_3 = -c_1 c_4 = -c_2$$

$$y(t) = -\frac{E}{\omega B} \sin(\omega t) + \frac{E}{B} t = R(\omega t - \sin(\omega t))$$

$$z(t) = R(1 - \cos(\omega t))$$

$$\vec{v}_{cent} = \left(\frac{\vec{E} \times \vec{B}}{B^2}\right)$$

\* Work:

$$dW = \vec{F} \cdot d\vec{l} = (a\vec{v} \times \vec{B}) \cdot \vec{v} dt = a dt (\vec{v} \times \vec{B}) \cdot \vec{v} = 0$$

- \* Critical note: magnetic fields do no work
- \* Magnetic fields may induce electric fields to do work, but do not do work themselves
- Continuous Systems
  - Given a wire in space, with a bit of charge, dq, moving with velocity  $\vec{v}$  shaped in rectangle, and placed in a magnetic field  $\vec{B} = \frac{A}{z} \hat{\mathbf{x}}$ :

$$dq \, \vec{v} = I \, d\vec{l}$$
 
$$d\vec{F} = I \, d\vec{l} \times \vec{B}$$
 
$$\vec{F}_{mag} = \int I \, d\vec{l} \times \vec{B}$$

- Solving would give us:

$$F_{tot} = -IL\left(\frac{A}{a}\right)\hat{\mathbf{z}} + \hat{\mathbf{y}} \int_{a}^{b} \frac{A}{z} dz + \hat{\mathbf{z}}IL\left(\frac{A}{b}\right) - \hat{\mathbf{y}}I \int_{a}^{b} \frac{A}{z} dz$$
$$= -\hat{\mathbf{z}}ILA\left(\frac{1}{a} - \frac{1}{b}\right)$$

- There are several types of densities:

Linear:  $I d\vec{l}$ Surface:  $\vec{K} da$ Bulk:  $\vec{J} d\tau$ 

• The following is an important continuity equation:

$$\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$$

- This is another form of charge conservation
- The Biot-Savart law is akin to Coulomb's law, but for magnetism:

$$d\vec{B} = \frac{\mu_o I d\vec{l} \times \hat{\mathbf{R}}}{4\pi R^2}$$

- Not the expression can be modified:

$$I d\vec{l} \iff \vec{K} da \iff \vec{J} d\tau$$

- The magnetic permeability of free space is:

$$\mu_o = 4\pi \cdot 10^{-7} \left[ \frac{\mathrm{N}}{\mathrm{A}^2} \right]$$

- \* This term "defines the amp", and the current then "defines the coulomb"
- \* Also note: the magnetic field is defined in Newtons per amp-meter
- Ampère's Law

$$\vec{\nabla} \times \vec{B} = \mu_o J \Longleftrightarrow \vec{B} = \frac{\mu_o}{4\pi} \int \frac{\vec{J} \times \hat{\mathbf{R}} \, d\tau}{R^2}$$

- Notice, we can then use Stokes' Theorem to write:

$$\oint \vec{B} \, dl = \int_{V} (\vec{\nabla} \times \vec{B}) \, d\vec{a} = \int \mu_{o} J^{2} \cdot \, d\vec{a} = \mu_{o} I_{enc}$$