## Transmission Lines

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- Transmission lines connect inputs to loads
  - -l is the length of the transmission lines
  - If l is not much smaller than  $\lambda$ , we need detailed analysis
  - If l is comparable to  $\lambda$ , then we can not use the lumped parameter model
  - We can, however, partition the transmission lines into segments where  $l \ll \lambda$ , then we can apply Kirchoff's circuit laws to each subdivided segment
  - For an imperfect dielectric, there is some loss
  - Some per unit-length properties:
    - 1. Resistance per unit length: R' (ohm per meter)

$$R = R' \Delta z$$

2. Inductance per unit length: L' (Henry per meter)

$$L = L'\Delta z$$

3. Capacitance per unit length: C' (Farad per meter)

$$C = C'\Delta z$$

4. Conductance per unit length: G' (Siemens per meter)

$$G = G'\Delta z$$

• Using this, we obtain the Helmholtz Equation:

$$\frac{d^2\tilde{V}(z)}{dz^2} - \gamma^2\tilde{V}(z) = 0$$

– Where 
$$\gamma = \sqrt{(j\omega C' + G')(j\omega L' + R')}$$

- For a Unidimensional Wave Equation:
  - The characteristic equation becomes:  $m^2 \gamma^2 = 0 \rightarrow m = \pm \gamma$
  - The solutions to the above differential equation are superpositions of  $\{e^{\gamma z}, e^{-\gamma z}\}$
  - $\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')}$ , where:
    - \* R' is the resistance per unit length
    - \* L' is the impedance per unit length
    - \*  $i = \sqrt{-1}$
    - $* \gamma = R_e \gamma + I_m \gamma$
    - \*  $R_e \gamma = \alpha$ , the attenuation constant
    - \*  $I_m \gamma = \beta$ , the phase constant (sometimes, notation used is  $k = \beta$ , like the wave number)

$$\widetilde{V}(z) = V_o^+ e^{-\gamma z} + V_o^- e^{\gamma z}$$

- We know from the definitions above that  $\gamma = \alpha + j\beta$
- For  $\alpha > 0$ , we are dealing with a (passive) lossy material ( $\alpha = 0$ ) is a loss less material
- For  $\alpha < 0$ , we are dealing with a gainy material
- The wave coming from the source is known as the "incident" wave, and the wave combing from the load is known as "reflected"

$$-\frac{dV}{dz} = (R' + j\omega L')\widetilde{I}(z)$$

– Solving this by incorporation the equation for  $\widetilde{V}$  above, we obtain:

$$\widetilde{I} = \left(V_o^+ e^{-\gamma z} - V_o^- e^{\gamma z}\right) \left(\frac{\gamma}{R' + j\omega L'}\right)$$

– We then define  $z_o = \frac{R' + j\omega L}{\gamma}$  as our characteristic impedance:

$$\widetilde{I} = \frac{1}{z_o} \left( V_o^+ e^{-\gamma z} - V_o^- e^{\gamma z} \right)$$

- There are thus two unknowns:  $V_o^+$ , which depends on source, and  $V_o^-$ , which is the reflected wave amplitude (depends on load)
  - \* The incident wave (in phase domain):  $\widetilde{V} = V_o^+ e^{-\alpha z} e^{-j\beta z}$
  - \* In time domain, this becomes:  $v(z,t) = |V_o^+|e^{-\alpha z}\cos(\omega t \beta z + \phi_+)^1$

 $<sup>^{1}\</sup>beta$  represents the wave number

- For a lossless line (R' = 0, G' = 0, and a pure inductor and capacitor)

$$\gamma = \sqrt{(jC'\omega)(jL'\omega)} \to \beta = \omega\sqrt{L'C'} = \frac{\omega}{U_{ph}}$$

$$U_{ph} = \frac{1}{\sqrt{L'C'}} = \frac{1}{\sqrt{\mu\varepsilon}} = c$$
 in dielectric

– The reflection coefficient  $(\Gamma)$  is given by:

$$\Gamma = \frac{V_o^-}{V_o^+}$$
 
$$V_{load} = V_o^+(1+\Gamma)$$
 
$$I_{load} = \frac{V_o^+}{z_o}(1-\Gamma)$$

– The normalized load impedance:  $\hat{z_L} = \frac{z_L}{z_o}$ 

$$\Gamma = \frac{z_l - z_o}{z_l + z_o} \to \Gamma = \frac{\hat{z_L} - 1}{\hat{z_L} + 1}$$

- Special Cases:
  - \* Short Circuit:

$$\Gamma_{sc} = -1$$

\* Open Circuit:

$$\Gamma_{oc} = 1$$

- · Reactive load, no real absorption
- A Phase-Shifted  $\Gamma$  would look as follows:

$$\Gamma_l = \left(\frac{z_L - z_o}{z_L + z_o}\right) e^{-j(2\beta l)}$$

- Standing Waves
  - $\tilde{V}(d) = V_o^+ e^{-j\beta d} + V_o^+ \Gamma e^{j\beta d}$

$$|\tilde{V}(d)| = |V_o^+| [1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta d - \theta_r)]^{\frac{1}{2}}$$

- \* This fluctuates between a minimum  $(V_{min})$  and maximum  $(V_{max})$  value
- \* Maximum/minimum possible value is determined by the cos term and is  $\pm 2|\Gamma$
- From the formula for  $\Gamma$ , we see that  $\Gamma=0$  when the load impedance matches the internal impedance
  - \* This is known as load matching

- Thus, when  $\Gamma=0,$  we know  $|\tilde{V}(d)|=|V_0^+|$
- The standing wave ratio is defined as follows:

$$SWR = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

- This measures the matching level of load to the line; ideally,  $\Gamma=0$ , and SWR=1; worst case,  $|\Gamma|=1$ , and  $SWR=\infty$
- In reference to the maximums and minimums, these occur when  $2\beta d \theta_r = 2n\pi$  and  $2\beta d \theta_r = (2n+1)\pi$ , respectively. This results in:

$$\begin{cases} d_{max} = \frac{2n\pi + \theta_r}{2\beta} \\ d_{min} = \frac{(2n+1)\pi + \theta_r}{\beta} \end{cases}$$