# The Hydrogen Atom in Wave Mechanics

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## Contents

1	A One-Dimensional Atom	3
2	Angular Momentum in the Hydrogen Atom	3
3	The Hydrogen Atom Wave Functions	4

#### 1 A One-Dimensional Atom

- Analyzing a Hydrogen Atom in Quantum Mechanics
  - The potential energy, derived from the Coulomb force, is:

$$U(r) = -\frac{e^2}{4\pi\varepsilon_o r}$$

- This may be converted to be in terms of x, and plugged into the Schrödinger equation:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} - \frac{e^2}{4\pi\varepsilon_0 x}\psi(x) = E\psi(x)$$

- To keep this finite, the following two conditions need to be met:

$$\begin{cases} x \to 0, & \psi(x) = 0 \\ x \to \infty, & \psi(x) = 0 \end{cases}$$

– Using  $\psi(x) = Axe^{-bx}$  with the Schrödinger equation, solving would obtain:

$$b = \frac{me^2}{4\pi\varepsilon_o\hbar^2} = \frac{1}{\alpha_o} \Rightarrow \text{Bohr Radius}$$

– Using the normalization condition, and  $\psi(x)=Axe^{-bx}=Axe^{\frac{x}{\alpha_o}},$  we get:

$$A = \frac{2}{\alpha_o^{\frac{3}{2}}}$$

- Conclusions:
  - \* For the ground state, there is an uncertainty in the location of  $e^-$
  - \* The most probable region to find  $e^-$  is near  $x = \alpha_o$  (consistent with Bohr model)
  - \* But, it's possible to find  $e^-$  anywhere (which is very difference from Bohr)

### 2 Angular Momentum in the Hydrogen Atom

- The angular momentum in a planetary system is constant, and vector  $\vec{L}(L_x, L_y, L_z)$  has three components
  - \* l is the angular momentum quantum number; it determines the length of the vector
  - \*  $m_l$  is the magnetic number; it determines one of the components of the vector

\* We have:

$$|\vec{L}| = \sqrt{l(l+1)}\hbar, \qquad l = 0, 1, 2 \cdots$$
  
 $\vec{L}_z = m_l \hbar, \qquad m_l = 0, \pm 1, \pm 2, \cdots \pm l$ 

- \* We know l and  $m_l$  for  $\vec{L}_z$ , but what is the direction?
- \* Applying the uncertainty principle, we obtain:

$$\Delta \vec{L}_z \Delta \phi \ge \hbar$$

### 3 The Hydrogen Atom Wave Functions

• Moving from a unidimensional case to a tridimensional case, we try to apply the Schrödinger equation (in Cartesian coordinates):

$$-\frac{\hbar^2}{2m} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) + U(x, y, z) \psi(x, y, z) = E\psi(x, y, z)$$

• The potential energy is:

$$U(x, y, z) = -\frac{e^2}{4\pi\varepsilon_o} \frac{1}{\sqrt{x^2 + y^2 + z^2}}$$

• The wave function may be defined as:

$$\psi(x, y, z) = X(x)Y(y)Z(z)$$

• Converting to polar coordinates to make calculations simpler, the Schrödinger equation becomes:

$$-\frac{\hbar^2}{2m} \left( \frac{\partial^2 \psi}{\partial r^2} + \frac{2}{r} \frac{\partial \psi}{\partial r} + \frac{1}{r^2 \sin(\theta)} \frac{\partial}{\partial \theta} \left( \sin(\theta) \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2(\theta)} \frac{\partial^2 \psi}{\partial \phi^2} \right) + U(r) \psi(r, \theta, \phi) = E \psi(r, \theta \phi)$$

• The wave function then becomes

$$\psi(r, \theta, \phi) = R(r)\Theta(\theta)\Phi(\phi)$$

- Where R is the radial function,  $\Theta$  is the polar function, and  $\Phi$  is the azimuthal function
- This breaks  $\psi$  into 3 equations, making the wave function easier to solve for
- The defined quantum numbers are as follows:

Number	Name	Values
n	Principle	$1, 2, 3, \cdots$
l	Angular Momentum	$0,1,2,\cdots,n-1$
$m_l$	Magnetic	$0,\pm 1,\pm 2,\cdots,\pm l$