

# Homework 4

Michael Brodskiy

Professor: G. Fiete

February 22, 2025

1. (a) To normalize the state, we can take the magnitudes of each component to get:

$$\sqrt{A^2 + (-A)^2 + (A)^2} = 1$$

This gives us:

$$\sqrt{3A^2} = 1$$

$$A\sqrt{3} = 1$$

$$A = \frac{1}{\sqrt{3}}$$

Thus, we may write the initial state as:

$$|\psi(t=0)\rangle = \frac{1}{\sqrt{3}}(|\phi_1\rangle - |\phi_2\rangle + i|\phi_3\rangle)$$

- (b) We know that the possible measurements are the eigenvalues of the corresponding eigenstates. Let us refer to them as the respective energies of each state,  $E_n$ . Therefore, we can measure any of the following observables:

$$|\phi_1\rangle \rightarrow E_1, |\phi_2\rangle \rightarrow E_2, |\phi_3\rangle \rightarrow E_3$$

Note that, for an infinite square well, the possible energy potentials for a Spin-1 system are:

$$E_1 = \frac{\pi^2 \hbar^2}{2mL^2}, E_2 = \frac{2\pi^2 \hbar^2}{mL^2}, E_3 = \frac{9\pi^2 \hbar^2}{2mL^2},$$

We can find the probabilities of each as:

$$P_n = |\langle \phi_n | \psi(t) \rangle|^2$$

We may observe that, because the magnitude of each coefficient is same, they all occur with the same probability, or:

$$P_{E_n} = 1/3$$

(c) We can find the average energy by using the probability as weights:

$$\langle E \rangle = P_{E_1} E_1 + P_{E_2} E_2 + P_{E_3} E_3$$

This gives us:

$$\langle E \rangle = \frac{1}{3} [E_1 + E_2 + E_3]$$

Using the energy values from above, we write:

$$\begin{aligned} \langle E \rangle &= \frac{1}{3} \left[ \frac{\pi^2 \hbar^2}{2mL^2} + \frac{2\pi^2 \hbar^2}{mL^2} + \frac{9\pi^2 \hbar^2}{2mL^2} \right] \\ \langle E \rangle &= \frac{1}{3} \left[ \frac{7\pi^2 \hbar^2}{mL^2} \right] \end{aligned}$$

$$\langle E \rangle = \frac{7\pi^2 \hbar^2}{3mL^2}$$

(d) Using our time evolution formula in tandem with our  $\phi_n$  eigenstates and corresponding  $E_n$  eigenvalues, we may write the time evolution as:

$$|\psi(t)\rangle = \frac{1}{\sqrt{3}} \left[ e^{-\frac{iE_1 t}{\hbar}} |\phi_1\rangle - e^{-\frac{iE_2 t}{\hbar}} |\phi_2\rangle + ie^{-\frac{iE_3 t}{\hbar}} |\phi_3\rangle \right]$$

(e) Plugging this in to the above equation, we find:

$$\left| \psi \left( \frac{\hbar}{E_1} \right) \right\rangle = \frac{1}{\sqrt{3}} \left[ e^{-i} |\phi_1\rangle - e^{-\frac{iE_2}{E_1}} |\phi_2\rangle + ie^{-\frac{iE_3}{E_1}} |\phi_3\rangle \right]$$

We may observe, however, that energy states are stationary, and, therefore, the probabilities remain the same.

2. Given that we can infer the well is of length  $L$  based on the formula, we may normalize by writing:

$$\int_0^L |\psi|^2 dx = 1$$

This gives us:

$$\int_0^L (AxL - Ax^2)^2 dx = 1$$

We continue to solve:

$$\begin{aligned}
\int_0^L A^2 x^4 - 2A^2 x^3 L + A^2 x^2 L^2 dx &= 1 \\
\left[ \frac{A^2 x^5}{5} - \frac{A^2 x^4 L}{2} + \frac{A^2 x^3 L^2}{3} \right] \Big|_0^L &= 1 \\
\left[ \frac{A^2 L^5}{5} - \frac{A^2 L^5}{2} + \frac{A^2 L^5}{3} \right] &= 1 \\
\frac{A^2 L^5}{30} &= 1 \\
\boxed{A = \pm \sqrt{\frac{30}{L^5}}}
\end{aligned}$$

Thus, we write the equation of state as:

$$\boxed{|\psi(x, t = 0)\rangle = x \sqrt{\frac{30}{L^5}} (L - x)}$$

We know that, for a particle of mass  $m$  in an infinite square well of length  $L$ , the eigenstates may be written as:

$$\Psi(x, 0) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

And the energies are written as:

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

We may write the time evolution of the system as:

$$\psi(x, t) = \sum_n c_n \Psi(x, 0) e^{-iE_n t/\hbar}$$

We can find  $c_n$  as:

$$c_n = \int_0^L \psi(x, 0) \Psi(x, 0) dx$$

We expand to write:

$$c_n = \frac{2}{L^3} \sqrt{15} \int_0^L x(L - x) \sin\left(\frac{n\pi x}{L}\right) dx$$

Entering this into a solver, we get:

$$c_n = \frac{2}{L^3} \sqrt{15} \left[ \frac{2L^3}{(n\pi)^3} [\cos(n\pi) - 1] \right]$$

This is equivalent to:

$$c_n = \frac{4\sqrt{15}}{n^3\pi^3} (1 - (-1)^n)$$

We may see that, for even  $n$  values,  $c_n = 0$ , and for odd values, we find:

$$c_n = \frac{8\sqrt{15}}{n^3\pi^3}, \quad n = 1, 3, 5, \dots$$

Therefore, we may conclude that the time-evolving state may be written as:

$$\psi(x, t) = \sum_{n=1,3,5,\dots} \left( \frac{8\sqrt{15}}{n^3\pi^3} \right) \left( \sqrt{\frac{2}{L}} \sin \left( \frac{n\pi x}{L} \right) \right) e^{-\frac{in^2\pi^2\hbar}{2mL^2}}$$

Pulling constants out front gives us:

$$\psi(x, t) = \frac{8}{\pi^3} \sqrt{\frac{30}{L}} \sum_{n=1,3,5,\dots} \left( \frac{1}{n^3} \right) \left( \sin \left( \frac{n\pi x}{L} \right) \right) e^{-\frac{in^2\pi^2\hbar}{2mL^2}}$$

Finally, we may calculate the expectation value as:

$$\langle x \rangle = \int_0^L \psi^*(x, t) \cdot x \cdot \psi(x, t) dx$$

We may observe that the factors independent of  $x$  (let us call them  $I_x$ ), in tandem with the cancellation of the exponential, allow this to be simplified to:

$$\langle x \rangle = I_x \int_0^L x \sin^2 \left( \frac{n\pi x}{L} \right) dx$$

Plugging this into a solver, we find:

$$\langle x \rangle = \frac{I_x L^2}{4}$$

We expand our factor to get:

$$\langle x \rangle = \frac{30}{L} \cdot \left( \frac{2}{\pi} \right)^6 \cdot \left( \sum_{n=1,3,5,\dots} n^{-2} \right)^2 \frac{L^2}{4}$$

Finally, this simplifies to:

$$\boxed{\langle x \rangle = \frac{15L}{2\pi^2}}$$

3. We know that the wave function may be written as:

$$\phi(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

Now with  $L \rightarrow 3L$ , we find:

$$\phi(x) = \sqrt{\frac{2}{3L}} \sin\left(\frac{n\pi x}{3L}\right)$$

This means that, for the ground and first excited states, we have:

$$\begin{aligned}\phi_{n=1}(x) &= \sqrt{\frac{2}{3L}} \sin\left(\frac{\pi x}{3L}\right) \\ \phi_{n=2}(x) &= \sqrt{\frac{2}{3L}} \sin\left(\frac{2\pi x}{3L}\right)\end{aligned}$$

From here, we may calculate the probabilities of each state as:

$$P_{n=1} = |\langle \phi_1(x) | \psi(x) \rangle|^2$$

Given that the wave function is real, we know  $\phi_{n=1}^*(x) = \phi_{n=1}(x)$ . This gives us:

$$\begin{aligned}P_{n=1} &= \left( \int_0^L \phi_{n=1}(x) \psi(x) dx + \int_L^{3L} \phi_{n=1}(x) (0) dx \right)^2 \\ P_{n=1} &= \left( \frac{2}{\sqrt{3L}} \int_0^L \sin\left(\frac{\pi x}{3L}\right) \sin\left(\frac{\pi x}{L}\right) dx \right)^2\end{aligned}$$

Entering this into a solver, we get:

$$P_{n=1} = \left( \frac{9}{8\pi} \right)^2$$

$$\boxed{P_{n=1} = \frac{81}{64\pi^2}}$$

Similarly, we may find:

$$P_{n=2} = |\langle \phi_2(x) | \psi(x) \rangle|^2$$

We proceed in a similar manner to the ground state, to write:

$$P_{n=2} = \left( \int_0^L \phi_{n=2}(x) \psi(x) dx \right)^2$$

$$P_{n=2} = \left( \frac{2}{\sqrt{3}L} \int_0^L \sin\left(\frac{2\pi x}{3L}\right) \sin\left(\frac{\pi x}{L}\right) dx \right)^2$$

Once again, we use a solver to get:

$$P_{n=2} = \left( \frac{9}{5\pi} \right)^2$$

$$P_{n=2} = \frac{81}{25\pi^2}$$