

# Lecture 5

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- Associated Legendre Functions

- For  $l = 0, 1, 2, 3 \dots$

$$\left[ (1 - z^2) \frac{d^2}{dz^2} - 2z \frac{d}{dz} + l(l + 1) - \frac{m^2}{1 - z^2} \right] P_l(z) = 0$$

- With:

$$P_l^m(z) = P_l^{-m}(z) = (1 - z^2)^{m/2} \frac{d^m}{dz^m} P_l(z)$$

- Since  $P_l(z)$  is an  $l$ -th order polynomial, then  $P_l^m(z)$  vanishes if  $m > l$

- Properties of  $P_l^m(z)$ :

- \*  $P_l^m(z) = 0$  if  $|m| > l$
    - \*  $P_l^m(z) = P_l^{-m}(z)$
    - \*  $P_l^m(\pm 1) = 0$  for  $m \neq 0$
    - \*  $P_l^m(-z) = (-1)^{l-m} P_l^m(z)$
    - \*  $\int_{-1}^1 P_l^m(z) P_q^m(z) dz = \frac{2}{2l + 2} \cdot \frac{(l + m)!}{(l - m)!} \delta_{lq}$

- We can obtain our  $\Theta(\theta)$  function as:

$$\Theta_l^m(\theta) = (-1)^m \frac{(2l + 1)}{2} \frac{(l - m)!}{(l + m)!} P_l^k(\cos(\theta)), \quad m \geq 0$$

And:

$$\Theta_l^{-m}(\theta) = (-1)^m \Theta_l^m(\theta)$$

– Here, we arrive at a point where the associated Legendre Functions are needed:

$$\begin{aligned}
P_0^o &= 1 \\
P_1^o &= \cos(\theta) \\
P_1^1 &= \sin(\theta) \\
P_2^o &= \frac{1}{2}(3\cos^3(\theta) - 1) \\
P_3^o &= \frac{1}{2}(5\cos^5(\theta) - 3\cos(\theta))
\end{aligned}$$

– Spherical Harmonics

$$\begin{aligned}
Y_l^m(\theta, \phi) &= (-1)^{(m+|m|)/2} \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-|m|)!}{(l+|m|)!}} P_l^m(\cos(\theta)) e^{im\phi} \\
Y_l^m(\theta, \phi) &= (-1)(Y_l^m(\theta, \phi))^*
\end{aligned}$$

Our first few values may be written as:

$$\begin{aligned}
Y_0^o(\theta, \phi) &= \frac{1}{\sqrt{4\pi}} \\
Y_1^o(\theta, \phi) &= \sqrt{\frac{3}{4\pi}} \cos(\theta) \\
Y_1^{\pm 1}(\theta, \phi) &= \mp \sqrt{\frac{3}{8\pi}} \sin(\theta) e^{\pm i\phi}
\end{aligned}$$

– Important Properties:

\* Orthonormality:

$$\langle l_1 m_1 | l_2 m_2 \rangle = \delta_{l_1 l_2} \delta_{m_1 m_2}$$

\* Completeness:

$$\begin{aligned}
\psi(\theta, \phi) &= \sum_{l=0}^{\infty} \sum_{m=-l}^l C_{lm} Y_l^m(\theta, \phi) \\
C_{lm} = \langle lm | \psi \rangle &= \int_0^{2\pi} \int_0^{\pi} (Y_l^m(\theta, \phi))^* \psi(\theta, \phi) d\Omega
\end{aligned}$$

\* Parity:

$$Y_l^m(\pi - \theta, \phi + \pi) = (-1)^l Y_l^m(\theta, \phi)$$

• The Radial Component

- We may consider a general case:

$$R_n(\rho) = \rho^l H(\rho) e^{-\gamma \rho}$$

- \* Where  $\rho = r/a$  with  $r$  as the radial distance and  $a$  as some scaling factor
- \* Furthermore, we have:

$$\gamma^2 = \frac{E}{\left(\frac{\hbar^2}{2\mu a^2}\right)} \Rightarrow E_n = -\frac{1}{2n^2} \left(\frac{Ze^2}{4\pi\epsilon_o}\right)^2 \frac{\mu}{\hbar^2}$$

- \* Our constraints then become:

$$\begin{aligned} n &= 1, 2, \dots, \infty \\ l &= 0, 1, \dots, n-1 \\ m &= -l, -l+1, \dots, l-1, l \end{aligned}$$

## • Hydrogen Energies and Spectrum

- The principal quantum number,  $n$ , is also called the shell number
- We may observe that, as  $n \rightarrow \infty$  we find the ionization limit
  - \* Note,  $E$  does not depend on  $m$
- For an electron (with a rest energy of 511[keV]), we may find:

$$E_n = -\frac{13.6}{n^2}$$

- The Bohr Radius becomes:

$$a_o = \frac{4\pi\epsilon_o\hbar^2}{m_e e^2}$$

- We may write the energy scale as:

$$E_n = -\frac{1}{2n^2} \left(\frac{e^2}{4\pi\epsilon_o a_o}\right)$$

- Noteworthy Features:

- \* Infinite number of bound states since the Coulomb potential falls off slowly for  $r \rightarrow \infty$  (finite square well in 3D has only a finite number of bound states)
- \* We compute the degeneracy of  $E_n$  by counting up all possible values:

$$\begin{aligned} \sum_{l=0}^{n-1} (2l+1) &= 2 \sum_{l=0}^{n-1} l + \sum_{l=0}^{n-1} 1 \\ \sum_{l=0}^{n-1} (2l+1) &= n(n-1) + n \end{aligned}$$

$$\sum_{l=0}^{n-1} (2l+1) = n^2$$

- \* If one includes the spin of the  $e^-$  atom, the total degeneracy is  $2n^2$
- \*  $m$  degeneracy is a result of spherical symmetry, and is removed if an electric field or magnetic field is applied
- \*  $l$  degeneracy is a result of the  $1/r$  potential and is removed if this changes
- \* The energies of emitted or absorbed light can be obtained as:

$$E_{\text{photon}} = \Delta E_{fi} = |E_f - E_i|$$

$$|E_f - E_i| = \frac{1}{2}(m_e c^2) \left( \frac{e^2}{4\pi\epsilon_o \hbar c} \right)^2 \left| \frac{1}{n_i^2} - \frac{1}{n_f^2} \right|$$

- \* Furthermore, we know:

$$E_{\text{photon}} = \hbar\omega = hv = \frac{hc}{\lambda}$$

- \* We may thus conclude:

$$\frac{1}{\lambda} = \frac{m_e}{4\pi\hbar^3 c} \left( \frac{e^2}{4\pi\epsilon_o} \right)^2 \left| \frac{1}{n_i^2} - \frac{1}{n_f^2} \right|$$

- \* Not all transitions are allowed in the hydrogen atom; transitions require a non-zero value of  $\langle n_f l_f m_f | V_{int} | n_i l_i m_i \rangle$
- \* For electromagnetic interactions, the selection rules, which follow from conservation of angular momentum, are:

$$\Delta l = l_f - l_i = \pm 1$$

$$\Delta m = m_f - m_i = 0, \pm 1$$

- The Radial Wave Functions

– Using:

$$a = \frac{4\pi\epsilon_o \hbar^2}{m_e Z e^2} = \frac{a_o}{Z}$$

$$\gamma = \frac{1}{n}$$

$$\rho = \frac{r}{a} = \frac{Zr}{a_o} \Rightarrow R_{nl}(r) = \left( \frac{Zr}{a_o} \right)^l e^{-Zr/na_o} H \left( \frac{Zr}{a_o} \right)$$

- The Overall Hydrogen Wave Function

$$|nlm\rangle = \psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_l^m(\theta, \phi)$$

– We then establish the following relationships:

$$H |nlm\rangle = -\frac{13.6}{n^2} |nlm\rangle$$

$$\vec{L}^2 |nlm\rangle = l(l+1)\hbar^2 |nlm\rangle$$

$$L_z |nlm\rangle = m\hbar |nlm\rangle$$