## Phase Transitions

## Michael Brodskiy

Professor: A. Stepanyants

November 28, 2023

- First-Order Phase Transitions
  - Three Phases: Solid, Liquid, Gas (Vapor)
  - Liquid and Gas/Vapor are known as fluids
  - From liquid to gas, solid to gas, and solid to liquid, the process would be isothermal
  - Phase Changes:
    - \* Solid to Liquid  $\rightarrow$  Melting
    - \* Liquid to Solid  $\rightarrow$  Crystalization
    - \* Solid to Gas  $\rightarrow$  Sublimation
    - \* Gas to Solid  $\rightarrow$  Deposition
    - \* Liquid to Gas  $\rightarrow$  Boiling
    - \* Gas to Liquid  $\rightarrow$  Condensation
- Coexistence of two phases (l and q)

$$G = G_l + G_g = N_l \mu_l(P, \tau) + N_g \mu_g(P, \tau)$$
$$N_g + N_l = N \quad \text{(constant)}$$

- This expression may be rewritten as:

$$G = N_l(\mu_l - \mu_g) + N\mu_g$$

- If l and g coexist, then:
  - \*  $\tau_l = \tau_g$  (thermal equilibrium)
  - \*  $P_l = P_g$  (mechanical equilibrium)
  - \*  $\mu_l = \mu_g$  (chemical equilibrium)
- Thus, for coexistence of l and g, we may write:

$$\mu_l(\tau, P) = \mu_g(\tau, P)$$

– The pressure differential with respect to  $\tau$  may be written as:

$$\frac{dP}{d\tau} = \frac{\left(\frac{\partial \mu_l}{\partial \tau}\right)_P - \left(\frac{\partial \mu_g}{\partial \tau}\right)_P}{\left(\frac{\partial \mu_g}{\partial \tau}\right)_{\tau} - \left(\frac{\partial \mu_l}{\partial \tau}\right)_{\tau}}$$

- Furthermore, implementing pressure and volume per molecules (s and v), we may write:

$$\frac{dP}{dt} = \frac{s_g - s_l}{v_q - v_l}$$

• We can recall that, in a reversible process:

$$\delta Q = \tau \, d\sigma$$
 
$$\Delta Q = \tau \Delta \sigma$$
 
$$\Delta Q = \tau (s_g - s_l)$$

- Thus, we may described the heat gained from moving 1 molecule from g to l
- This  $\Delta Q$  term is referred to as L, or the latent heat of vaporization per molecule
- This leads us to the Clausius-Clapeyron Equation:

$$\frac{dP}{d\tau} = \frac{L}{\tau(v_q - v_l)}$$

- This is used in liquid-gas coexistence
- We assume the gas is ideal
- We also assume  $v_l \ll v_g$
- Also, L is assumed to be constant
- In assuming this, we may write:

$$\frac{dP}{d\tau} = \frac{LP}{\tau^2}$$

- Solving this, we get:

$$P(\tau) = P_0 e^{-\frac{L}{\tau}}$$

• Van der Waals Equation of state

- As a reminder, the free energy for an ideal gas is:

$$F_{ideal} = -N\tau \left( \ln \left( \frac{n_Q}{n} \right) + 1 \right)$$
$$P = -\left( \frac{\partial F}{\partial V} \right)_{\tau N} \to PV = N\tau$$

- In the above, V must be replaced with V bN, where b is the volume of a single molecule, and N is the total number of molecules
- To find the interaction between particles, we use a mean field approximation

$$\Delta U = \frac{1}{2} \sum_{i \neq j=1}^{N} \phi(|\bar{r}_i - \bar{r}_j|) = \frac{1}{2} (N^2 - N) \langle \phi(r) \rangle$$

- Assume that the molecules are uniformly distributed in the accessible volume, V-bn, and average overall the positions of all molecules
- We can approximate it to:

$$\frac{1}{2}(N^2 - N) \int_V \frac{\phi(r)}{V - bN} d^3r$$
$$\Delta U = -\frac{N^2 a}{V}, \quad a > 0$$

- from here, we can write the free energy:

$$F_{VdW} = -N\tau \left( \ln \left( \frac{n_Q(V - bN)}{N} \right) + 1 \right) - \frac{N^2 a}{V}, \quad a, b > 0$$

$$P = \frac{\tau N}{V - bN} - \frac{N^2 a}{V^2}$$

- The equation of state may be simplified as:

$$\left(P + \frac{N^2 a}{V}\right)(V - bN) = N\tau$$

- \* The term dependent on a is the reduction in pressure due to attractive interactions between molecules
- \* The term dependent on b is the excluded volume due to molecule volume
- Critical points may be defined as:

$$\tau_c = \frac{8a}{27b}$$

$$P_c = \frac{a}{27b^2}$$

$$V_c = 3Nb$$

- The Gibbs Free Energy for this Van der Waals gas is:

$$G = F + PV = F + P(V - bN)$$
$$G = -N\tau \ln \left(\frac{n_Q(V - bN)}{N}\right) - \frac{2N^2a}{V}$$

– We find the following relations:

$$\begin{split} \left(\frac{\partial G}{\partial P}\right)_{\tau,N} &= \frac{V}{N} \\ \left(\frac{\partial G}{\partial \tau}\right)_{P,N} &= -\frac{\sigma}{N} = -S \end{split}$$

- Ferromagnetism
  - There are N magnetic moments  $\mu$  in a binary-spin system. From this, we get:

$$U = -$$

- Assume that M is the magnetization of the volume — net magnetic moment per volume

$$B_E = \lambda M$$

- Is the effective magnetic field due to M (mean field approximation)