Last week we generalised the explicit finite difference scheme. Additionally, the stability of these schemes were analysed using a *von Neumann* analysis. From here boundary conditions were studied, in particular, how to build these into the general matrix scheme that was derived.

# 1 Backwards difference scheme for parabolic PDEs

So far we have looked at using a *forwards in time, centred in space* approach sometimes referred to as FTCS scheme. However, as this leads to an **explicit** scheme, it has stability criterion - i.e. the CFL number in which the  $\sigma$  coefficient (following lecture notation) must be less than or equal to 1/2.

If we consider the diffusion equation,

$$\partial_t u = D \partial_{xx} u$$

and discretise this using a backwards difference we get,

$$\frac{u_j^m - u_j^{m-1}}{\Delta t} = \frac{D(u_{j+1}^m - 2u_j^m + u_{j-1}^m)}{(\Delta x)^2}$$
$$u_j^{m-1} = u_j^m - \frac{D\Delta t}{(\Delta x)^2} \left( u_{j+1}^m - 2u_j^m + u_{j-1}^m \right)$$

Now, taking the normal approach we set  $\sigma = \frac{D\Delta t}{(\Delta x)^2}$  and write,

$$u_j^{m-1} = (1+2\sigma)u_j^m - \sigma u_{j+1}^m - \sigma u_{j-1}^m$$

In order to solve this, we need to write our system of equations in matrix form so that we have,

$$A\mathbf{u}^{m} = \mathbf{u}^{m-1}$$

$$\implies \mathbf{u}^{m} = A^{-1}\mathbf{u}^{m-1}$$

Therefore, we are able to find our timestep m from the previous, m-1, and update our system. In doing this, be careful about the top and bottom row of your matrix A, as this must incorporate boundary conditions!

### 1.1 Boundary conditions

For Neumann (or flux) boundary conditions, we can treat these as first or second order,

$$\frac{\partial u}{\partial x} \approxeq \begin{cases} \frac{u_{j+1}^{m+1} - u_j^{m+1}}{\Delta x} = c \\ \frac{-3u_j^{m+1} + 4u_{j+1}^{m+1} - u_{j+2}^{m+1}}{2\Delta x} = c \end{cases} \implies u_{j+1}^{m+1} - u_j^{m+1} = c\Delta x \text{ , first order }$$

#### 1.2 Characteristics of implicit scheme

- The backward difference scheme is unconditionally stable permits any value of  $\sigma$ ;
- The local truncation error is  $O(\Delta x^2) + O(\Delta t)$ , the same as for the forward difference scheme
- Unconditional stability permits much larger time steps, but they attract larger truncation error
- If the time step is chosen as roughly equivalent to the grid spacing, then  $O(\Delta t) \ll O(\Delta x^2)$ , and so the error of the backward difference scheme becomes  $\sim O(\Delta t)$
- It is possible to use a mixed approach, which maintains unconditional stability, and reduces the leading error term of the local truncations

## 2 Crank-Nicolson (CN) scheme for diffusion

Okay, so we achieved unconditional stability - but what if we want to reduce error? This is where the CN scheme comes in. The idea here is that we use a mixed central difference, namely we take the average central difference between the current and previous timestep,

$$\partial_{xx} u \approx \frac{1}{2} \left( \frac{u_{j+1}^m - 2u_j^m + u_{j-1}^m}{\Delta x^2} + \frac{u_{j+1}^{m-1} - 2u_j^{m-1} + u_{j-1}^{m-1}}{\Delta x^2} \right)$$

In addition to this, we use our backwards difference in time for the temporal derivative. Substituting these into the diffusion equation and arranging terms of matching timestep on LHS and RHS (*try this yourself!*) we get,

$$-\frac{\sigma}{2}u_{j-1}^{m}+(1+\sigma)u_{j}^{m}-\frac{\sigma}{2}u_{j+1}^{m}=\frac{\sigma}{2}u_{j-1}^{m-1}+(1-\sigma)u_{j}^{m-1}+\frac{\sigma}{2}u_{j+1}^{m-1}$$

This, we can then write into matrix form and look to solve. In the end, this gives us a scheme that has error  $O(\Delta x^2) + O(\Delta t^2)$ .

## 2.1 Boundary conditions

These can be applied in the same way as our previous schemes, namely, discretise the if a Neumann boundary condition and form this equation in the top/bottom rows of the matrices. For the CN scheme, the matrix associated with the previous timestep may have all zeros for the top and bottom row. This is a result of the boundary conditions not being dependent (necessarily) on the previous time.