1 Unitary Operations and RBMs

The state of an N-spin system is described by a wavefunction $\Psi(s)$ with $s \in \{0,1\}^N$. $\Psi(s)$ can be encoded in a so-called Boltzmann Machine, a simple two-layer ANN. It consists of N_v visible and N_h hidden nodes and is specified by parameters $a \in \mathbb{R}^{N_v}$, $b \in \mathbb{R}^{N_h}$ and $w \in \mathbb{R}^{N_h \times N_v}$. The energy function

$$E(v,h) = v^T a + h^T b + h^T w v \tag{1}$$

allows one to express probability distributions defined on \mathbb{R}^{N_v} according to

$$P(v) = \frac{\sum_{h} e^{-E(v,h)}}{\sum_{v} \sum_{h} e^{-E(v,h)}}.$$
 (2)

Similarly, if we allow complex $a \in \mathbb{C}^{N_v}$, $b \in \mathbb{C}^{N_h}$ and $w \in \mathbb{C}^{N_h \times N_v}$ we can map each spin configuration $s \in \{0,1\}^{N_S}$ of system with N_S spins to a complex amplitude. This is then the wavefunction

$$\Psi(s) = \frac{\sum_{h} e^{-E(s,h)}}{\sum_{s} \sum_{h} e^{-E(s,h)}}.$$
 (3)

Consider now a 1-body unitary operator $O \in U(2)$. It is completely specified by its four matrix elements $O_{ss'}$ where $s, s' \in \{0, 1\}$. Equivalently, we can express this Operator in terms of an exponential function:

$$O(s, s') = A \exp(\alpha s + \beta s' + \omega s s'). \tag{4}$$

To do this we simply associate

$$\begin{split} A &= O_{00} \\ \alpha &= \ln \left(\frac{O_{10}}{A} \right) \\ \beta &= \ln \left(\frac{O_{01}}{A} \right) \\ \omega &= \ln \left(\frac{O_{11}}{A} \right) - (\alpha + \beta). \end{split}$$

Clearly the function O(s, s') defined this way behaves like the operator. The wavefunction $\Psi'(s)$ after O has acted upon the initial state Psi(s) is given by

$$\Psi'(s) = \sum_{s'} O_{ss'} \Psi(s_1, \dots, s', \dots, s_N)$$

$$= \sum_{s'} \exp(\alpha s + \beta s' + \omega ss') \Psi(s_1, \dots, s', \dots, s_N)$$
(5)

Expressing the RHS wave-function now in terms of our RBM (cf. 3), the sum over spin s' can be associated with another hidden node.

2 Unitary Operations and POVM Distributions

Suppose now the RBM encodes some probability distribution $P(\vec{a})$ of an informationally complete POVM measurement. This distribution then uniquely specifies some quantum state. Let $O_{\vec{a},\vec{b}}$ now denote a single-body unitary operator acting on POVM distribution vectors. Is it still possible to rewrite this operator as an exponential function and describe the effect of the unitary operation as the insertion of an additional node into the RBM?