



GIT Electronics Engineering

ELEC 232 and ELEC 331

Some Derivations for the Analysis of Simple Electronic Circuits

Updated February 16, 2016 - 22:54

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Abstract

In English: The contents of this document are derivations that are frequently used and mostly common to the topics of the courses ELEC 232 and ELEC 331 (Electronic Circuits 1 and 2, respectively).

In Turkish: Bu dokümanın içerisinde ELEC 232 ve ELEC 331 (Elektronik Devreler 1 ve 2) derslerinde sıkça kullanılan ve her iki dersin konularında ortak olan çıkarımlar bulunmaktadır.

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1. Introduction and Essentials

1.1 Resistances

Significance of Equivalent Parallel Resistances

Worked on this on =
January 15, 2016 / Fri

Parallel resistance $R_p = R_1 // R_2$

$$\frac{1}{R_1} + \frac{1}{R_2} = \frac{1}{R_p}$$

$$\frac{R_2}{R_1 R_2} + \frac{R_1}{R_1 R_2} = \frac{1}{R_p}$$

$$\frac{R_1 + R_2}{R_1 R_2} = \frac{1}{R_p}$$

then
$$R_p = \frac{R_1 R_2}{R_1 + R_2}$$

Note the identity

$$\frac{1}{R_1} + \frac{1}{R_2} = \frac{1}{R_1 // R_2}$$

where $R_1 // R_2 = R_p$

Also note the identity

$$R_p = \frac{1}{\frac{1}{R_1} + \frac{1}{R_2}} \Rightarrow \frac{1}{\frac{1}{R_1 // R_2}} = R_1 // R_2 = R_p$$

Note

$$R_p = \frac{R_1 R_2}{R_1 + R_2} = R_1 \cdot \frac{R_2}{R_1 + R_2} < R_1$$

Since $\frac{R_2}{R_1 + R_2} < 1$

Assuming
 $R_1, R_2 > 0$

Also

$$R_p = R_2 \cdot \frac{R_1}{R_1 + R_2} < R_2$$

Since $\frac{R_1}{R_1 + R_2} < 1$

Then $\overbrace{R_p = R_1 // R_2} < \max(R_1, R_2)$

and $R_p = R_1 // R_2 < \min(R_1, R_2)$

We have $R_p < \min(R_1, R_2) \leq \max(R_1, R_2)$

↓
since it may be
that
 $R_1 = R_2$

Now assume that for some $\alpha > 0$

$$R_2 = \alpha R_1$$

$$R_p = R_1 // R_2 = \frac{R_1 R_2}{R_1 + R_2} = \frac{\alpha R_1^2}{R_1 + \alpha R_1} = \frac{\alpha R_1^2}{(1+\alpha) R_1}$$
$$= \frac{\alpha}{1+\alpha} R_1$$

For an α that is very big ($\alpha > 10$)

$$R_p = \frac{\alpha}{1+\alpha} R_1 \approx R_1 \text{ since } \lim_{\alpha \rightarrow +\infty} \frac{\alpha}{1+\alpha} = 1$$

Note that for $\alpha > 10$

$$R_p < \min(R_1, R_2) = \min(R_1, \alpha R_1) = \cancel{R_1}$$

since
 $\alpha R_1 > R_1$

Then

$$R_p < R_1$$

But $R_p \approx R_1$

For two resistances that are in parallel, if there is a big difference between the two, the equivalent parallel resistance will almost be equal to the minimum of the two.

Now assume

$$R_2 = \alpha R_1 \quad \text{but } \alpha \ll 1$$

$$R_p = R_1 // R_2 = \frac{\alpha R_1^2}{(1+\alpha)R_1} = \frac{\alpha}{1+\alpha} R_1 \approx \alpha R_1$$

$$R_p < \min(R_1, R_2) = \min(R_1, \alpha R_1) = \underbrace{\alpha R_1}_{\text{since } \alpha R_1 < R_1}$$

And

$$R_p < \alpha R_1$$

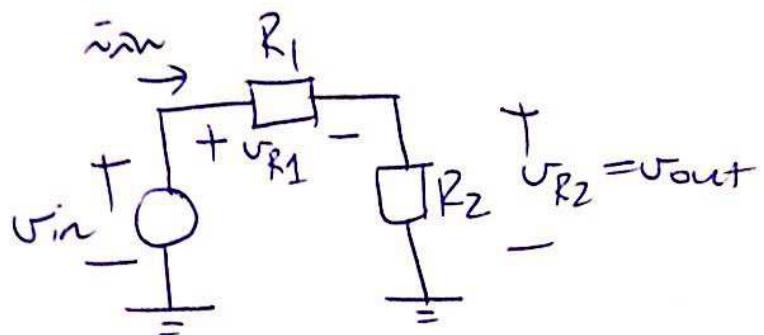
But

$$R_p \approx \alpha R_1$$

Read the note at the end of page 3.

1.2 Dividers

Voltage divider



Compute V_{out} in terms of V_{in} .

KVL

$$+V_{in} - V_{R1} - V_{R2} = 0$$

Ohm's law

$$V_{R1} = R_1 \bar{i}_{in}$$

$$V_{R2} = R_2 \bar{i}_{in}$$

Then

$$+V_{in} - R_1 \bar{i}_{in} - R_2 \bar{i}_{in} = 0$$

$$V_{in} = (R_1 + R_2) \bar{i}_{in}$$

$$\frac{V_{in}}{\bar{i}_{in}} = R_1 + R_2$$

And

$$V_{out} = V_{R2} = R_2 \bar{i}_{in}$$

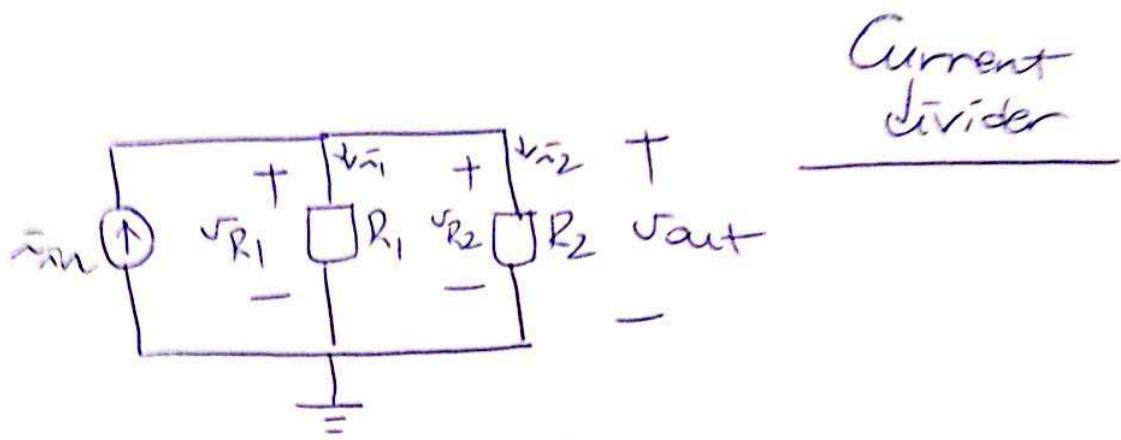
$$\frac{V_{out}}{\bar{i}_{in}} = R_2$$

Note that

$$\frac{V_{out}}{V_{in}} = \frac{R_2}{R_1 + R_2}$$

$$\frac{R_2}{R_1 + R_2} = \frac{V_{out}}{V_{in}}$$

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Note that

$$v_{R_1} = v_{R_2} = v_{out}$$

$$v_{R_1} = \bar{\pi}_1 R_1$$

$$v_{R_2} = \bar{\pi}_2 R_2$$

$$\bar{\pi}_{in} = \bar{\pi}_1 + \bar{\pi}_2 \text{ by KCL}$$

Then

$$\bar{\pi}_1 R_1 = \bar{\pi}_2 R_2 \Rightarrow \bar{\pi}_2 = \bar{\pi}_1 \frac{R_1}{R_2}$$

$$\bar{\pi}_1 + \bar{\pi}_1 \frac{R_1}{R_2} = \bar{\pi}_{in}$$

$$\bar{\pi}_1 \left(1 + \frac{R_1}{R_2}\right) = \bar{\pi}_{in}$$

$$\frac{\bar{\pi}_1}{\bar{\pi}_{in}} = \frac{1}{1 + \frac{R_1}{R_2}} = \frac{R_2}{R_1 + R_2} \quad \left. \begin{array}{l} \text{current} \\ \text{division} \\ \text{factor} \end{array} \right\}$$

Note also that

$$\frac{\bar{\pi}_1}{\bar{\pi}_{in}} + \frac{\bar{\pi}_2}{\bar{\pi}_{in}} = \frac{\bar{\pi}_1 + \bar{\pi}_2}{\bar{\pi}_{in}} = \frac{\bar{\pi}_{in}}{\bar{\pi}_{in}} \Rightarrow \frac{\bar{\pi}_2}{\bar{\pi}_{in}} = \frac{R_1}{R_1 + R_2} \quad \left. \begin{array}{l} \text{another} \\ \text{current} \\ \text{division} \\ \text{factor} \end{array} \right\}$$

(1)

Current Divider

$$\frac{U_{\text{out}}}{i_{\text{in}}} = ?$$

$$U_{\text{out}} = U_{R_1} = R_1 \cdot \frac{R_2}{R_1 + R_2} i_{\text{in}}$$

$$\frac{U_{\text{out}}}{i_{\text{in}}} = R_1 \cdot \frac{R_2}{R_1 + R_2} = \frac{R_1 R_2}{R_1 + R_2} = R_1 // R_2$$

Could also compute this as:

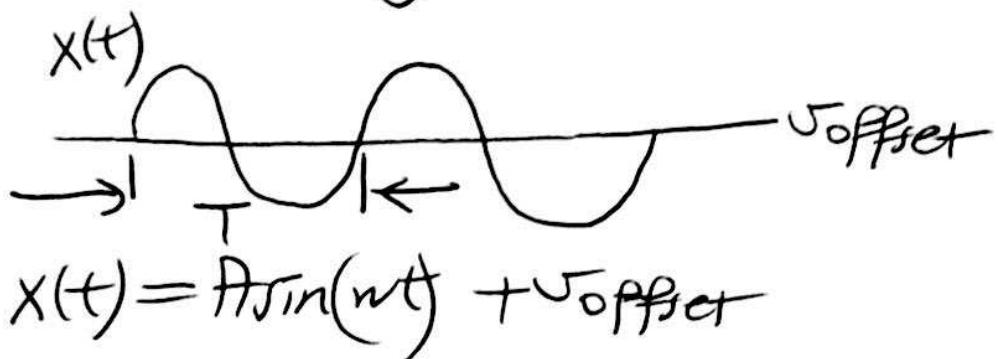
$$U_{\text{out}} = U_{R_2} = R_2 \cdot \frac{R_1}{R_1 + R_2} i_{\text{in}}$$

$$\frac{U_{\text{out}}}{i_{\text{in}}} = \frac{R_1 R_2}{R_1 + R_2} = R_1 // R_2$$

1.3 AC and DC Coupling

AC and DC coupling

For a sinusoidal signal with offset



$$x(t) = A \sin(\omega t) + V_{\text{offset}}$$

On the oscilloscope; the DC coupled waveform is $x(t)$ itself.

The DC component of $x(t)$ is computed as in

$$x_{\text{DC}} = \frac{1}{T} \int_0^T x(t) dt$$

where T is the period of the signal.

The DC component is invariant under time-shifts in the signal, i.e.

$$x_{\text{DC}} = \frac{1}{T} \int_{t_0}^{t_0+T} x(t) dt \quad \text{for } t_0 \in \mathbb{R}$$

For this signal $x(t)$ note that

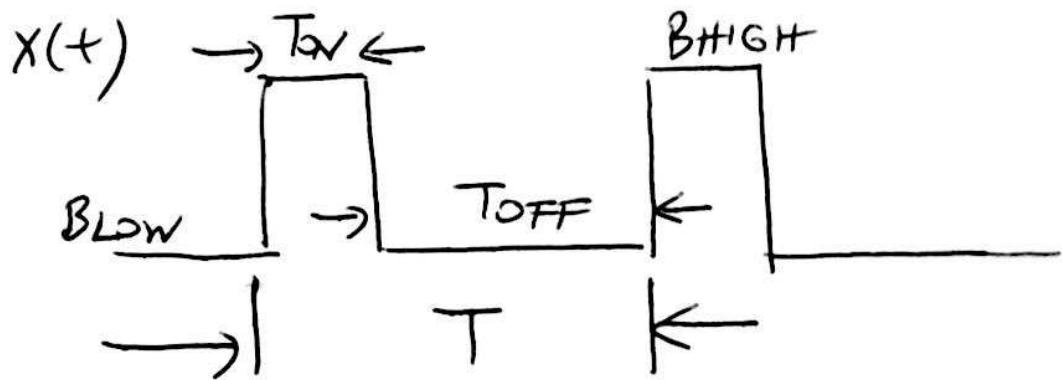
$$T = \left(\frac{\omega}{2\pi}\right)^{-1} = \left(\frac{2\pi f}{2\pi}\right)^{-1} = \frac{1}{f}$$

and

$$\begin{aligned}x_{DC} &= \frac{1}{T} \int_{t_0}^{t_0+T} [A \sin(\omega t) + v_{offset}] dt \\&= \frac{A}{T} \int_{t_0}^{t_0+T} \sin(\omega t) dt + \frac{1}{T} \int_{t_0}^{t_0+T} v_{offset} dt \\&= 0 \quad \frac{1}{T} \pi A \cos(\omega t) \Big|_{t_0}^{t_0+T} \\&\quad + \frac{1}{T} T v_{offset}\end{aligned}$$

$$x_{DC} = v_{offset}$$

AC and DC coupled waveforms
for a periodic pulse signal



Note that $\underbrace{T}_{\text{period}} = T_{on} + T_{off}$

DC coupled waveform $\bar{x}(t)$ itself.

DC component of $x(t)$ \bar{x}

$$\begin{aligned}x_{DC} &= \frac{1}{T} \int_0^T x(t) dt \\&= \frac{1}{T} [T_{on} B_{high} + T_{off} B_{low}]\end{aligned}$$

Note now that

$$x_{DC, \text{coupled}}(t) = x(t)$$

x_{DC} as computed on pg 1

$$x_{AC, \text{coupled}}(t) = x(t) - x_{DC}$$

Q: When do we have

$$x_{DC, \text{coupled}}(t) = x_{AC, \text{coupled}}(t) ?$$

A: When $x_{DC} = 0$

$$\frac{1}{T} [T_{ON} B_{HIGH} + T_{OFF} B_{LOW}] = 0$$

$$\frac{T_{ON}}{T} B_{HIGH} + \frac{T - T_{ON}}{T} B_{LOW} = 0$$

$$\frac{T_{ON}}{T} [B_{HIGH} - B_{LOW}] = -B_{LOW}$$

$$\frac{T_{ON}}{T} = -\frac{B_{LOW}}{B_{HIGH} - B_{LOW}}$$

(2)

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The duty cycle of a periodic pulse waveform is

$$\text{duty cycle} = \frac{T_{ON}}{T} \times 100\%$$

Note that on pg 1 we have computed

$$\begin{aligned} X_{DC} &= \frac{T_{ON}}{T} B_{HIGH} + \frac{T-T_{ON}}{T} B_{LOW} \\ &= \frac{T_{ON}}{T} [B_{HIGH} - B_{LOW}] + B_{LOW} \end{aligned}$$

then

$$\text{duty cycle} = \frac{T_{ON}}{T} \times 100\% = \frac{X_{DC} - B_{LOW}}{B_{HIGH} - B_{LOW}} \times 100\%$$

Also note that

X_{DC} changes if B_{HIGH} and B_{LOW} remain constant but the duty cycle changes.

This means since X_{DC} changes with the duty cycle, $X_{AC, \text{coupled}}(t)$ will change with the duty cycle as well.

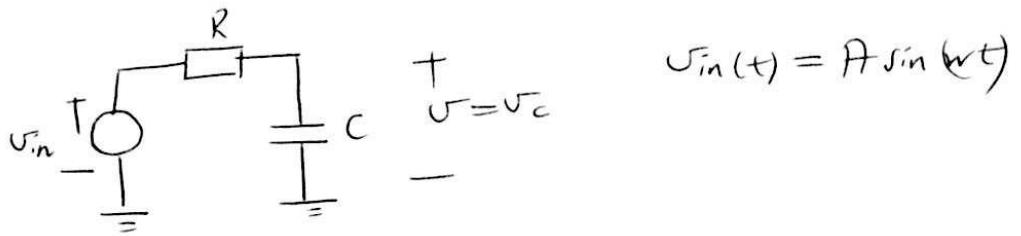
Observe this with a signal generator and oscilloscope.

(4)

1.4 RC Low-Pass Filter

Response of an RC
Low-Pass Filter to
a sinusoidal input

Worked on this on: Dec. 24, 2015/Th



$$\frac{v - v_{in}}{R} + C \frac{dv}{dt} = 0$$

$$\frac{v - v_{in}}{RC} + \frac{dv}{dt} = 0$$

$$\frac{dv}{dt} + \frac{1}{RC} v = \frac{1}{RC} v_{in}(t)$$

$$\begin{aligned}
 sV(s) - v(0^-) + \frac{1}{RC} V(s) &= \frac{1}{RC} V_{in}(s) \\
 &= \frac{A}{RC} \frac{\omega}{s^2 + \omega^2}
 \end{aligned}$$

$$\left(s + \frac{1}{RC}\right)V(s) = v(0^-) + \frac{A}{RC} \frac{\omega}{s^2 + \omega^2}$$

$$V(s) = \frac{v(0^-)}{s + \frac{1}{RC}} + \frac{A}{RC} \frac{\omega}{(s^2 + \omega^2)(s + \frac{1}{RC})}$$

Should now compute the periodic part
of the soln $v(t)$ in the time domain.

Take care of

$$\frac{1}{(s^2 + w^2)(s + \frac{1}{RC})} = \frac{Fs + G}{s^2 + w^2} + \frac{H}{s + \frac{1}{RC}}$$

$$(Fs + G)\left(s + \frac{1}{RC}\right) + H(s^2 + w^2) = 1$$

$$Fs^2 + \frac{G}{RC} + Gs + \frac{F}{RC}s + \underline{Hs^2 + Hw^2} = 1$$

$$F + H = 0 \implies F = -H$$

$$G + \frac{F}{RC} = 0 \implies G - \frac{H}{RC} = 0 \Rightarrow G = \frac{H}{RC}$$

$$\underbrace{\frac{G}{RC} + Hw^2}_m = 1$$

$$\frac{H}{(RC)^2} + Hw^2 = 1$$

$$H \left[1 + (RCw)^2 \right] = (RC)^2$$

$$H = \frac{(RC)^2}{1 + (RCw)^2}$$

$$G = \frac{H}{RC} = \frac{RC}{1 + (RCw)^2}$$

$$F = -\frac{(RC)^2}{1 + (RCw)^2}$$

Note that the periodic part has the Laplace Transform

$$V_{\text{per}}(s) = \frac{Fw}{RC} \left[\frac{Fs + G}{s^2 + w^2} \right]$$
$$= \frac{AFw}{RC} \frac{s}{s^2 + w^2} + \frac{AG}{RC} \frac{w}{s^2 + w^2}$$

$$v_{\text{per}}(t) = \frac{AFw}{RC} \cos(wt) + \frac{AG}{RC} \sin(wt)$$

$$K \cos \phi = \frac{AG}{RC}$$

$$K \sin \phi = \frac{AFw}{RC}$$

Then

$$\tan \phi = \frac{\cancel{AFw}}{\cancel{RC}} \frac{\cancel{RC}}{\cancel{AG}} = \frac{Fw}{G}$$

$$= - \frac{(RC)^2}{1 + (RCw)^2} w \frac{1 + (RCw)^2}{RC}$$

$$\tan \phi = -RCw$$

and

$$K^2 = \left(\frac{AFw}{RC} \right)^2 + \left(\frac{AG}{RC} \right)^2$$
$$= \left(\frac{A}{RC} \right)^2 \left[(Fw)^2 + G^2 \right]$$
$$= \left(\frac{A}{RC} \right)^2 \left[\frac{(RC)^4 w^2}{[1 + (RCw)^2]^2} + \frac{(RC)^2}{[1 + (RCw)^2]^2} \right]$$

$$K^2 = \left(\frac{A}{RC}\right)^2 \left[(RC)^2\right] \left[\frac{1 + (RCw)^2}{[1 + (RCw)^2]^2} \right]$$

$$= \frac{A^2}{1 + (RCw)^2}$$

$$K = \frac{A}{\sqrt{1 + (RCw)^2}}$$

—————

We have

$$\varphi_{per}(t) = K \sin(wt + \phi) = \frac{A}{\sqrt{1 + (RCw)^2}} \sin(wt + \arctan(-RCw))$$

$$\text{at } w = w_0 = \frac{1}{RC}$$

$$K = \frac{A}{\sqrt{1 + (RC \frac{1}{RC})^2}} = \frac{A}{\sqrt{2}}$$

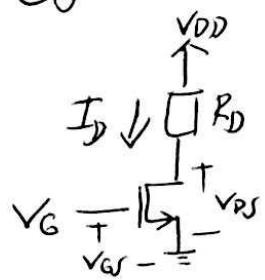
$$\begin{aligned} \phi &= \arctan\left(-RC \frac{1}{RC}\right) = \arctan(-1) \\ &= -\frac{\pi}{4} \end{aligned}$$

—————

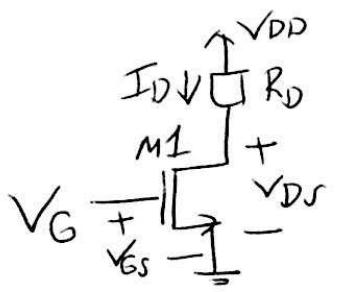
2. Device Models

2.1 Transitions between Regimes

Solution for the Boundary of
Triode and Saturation Regimes
in a ~~CS~~ type stage



Worked on this on: Dec 23, 2015 / Wed



What is the value of $V_G = V_{GS}$
for which

$\Rightarrow M1$ runs ^{just} into triode?
 $\Rightarrow M1$ resides just at the boundary of triode and saturation?

} equivalent questions

$$\begin{aligned}V_G &= V_G - 0V \\&= V_G - V_S \\&= V_{GS}\end{aligned}$$

$$\begin{aligned}V_D &= V_D - 0V \\&= V_D - V_S \\&= V_{DS} \\&= V_{DD} - R_D I_D\end{aligned}$$

The boundary between triode and saturation

determined by

$$\begin{aligned}V_{DS} &= V_{GS} - V_{th} \\&= V_{ov} \\&= V_{overdrive}\end{aligned}$$

$$V_{DS} = V_{DD} - R_D I_D \\ = V_{GS} - V_{th} \quad (1)$$

At The boundary between the triode and saturation

$\Rightarrow I_D$ is still equal to

$$I_D = \frac{1}{2} \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_{th})^2$$

Then we have to solve from (1)

$$V_{DD} - R_D \frac{1}{2} \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_{th})^2 = V_{GS} - V_{th} \quad (2)$$

Plug in
 $K_n = \frac{1}{2} \mu_n C_{ox} \frac{W}{L}$

and $V_{ov} = V_{GS} - V_{th}$

Then (2) becomes

$$V_{DD} - R_D K_n V_{ov}^2 = V_{ov} \\ \Rightarrow R_D K_n V_{ov}^2 + V_{ov} - V_{DD} = 0 \quad (3)$$

In (3), check the units

$$V_{ov} : \text{Volts} \quad R_n : \Omega \\ V_{DD} : \text{Volts} \quad K_n = \frac{\text{A}}{\text{V}^2}$$

$$\text{Then } R_D K_n V_{DD}^2 \Rightarrow \text{unit: } \Omega \frac{A}{V^2} \cdot V^2 = V$$

units check!

Solve (3).

$$V_{OR} = \frac{-1 \pm \sqrt{1 + 4R_D K_n V_{DD}}}{2R_D K_n}$$

Check the units:

$$R_D K_n V_{DD} \Rightarrow \text{unit: } \Omega \frac{A}{V^2} V \Rightarrow \text{units}$$

$$R_D K_n \Rightarrow \text{unit: } \Omega \frac{A}{V^2} \Rightarrow \frac{1}{V} = V^{-1}$$

Units again check!

Note that

$$\sqrt{1 + 4R_D K_n V_{DD}} > 1$$

Then definitely

$$-1 - \sqrt{1 + 4R_D K_n V_{DD}} < 0$$

$$-1 + \sqrt{1 + 4R_D K_n V_{DD}} > 0$$

$$V_{OR} = \frac{-1 + \sqrt{1 + 4R_D K_n V_{DD}}}{2R_D K_n} \quad \text{in units of Volts}$$

is the physically valid soln.

$$V_{GS} = V_{OR} + V_{TH} = V_{TH} + \frac{-1 + \sqrt{1 + 4R_D K_n V_{DD}}}{2R_D K_n}$$

2.2 Small-Signal Equivalent Circuits

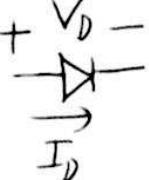
Small-Signal Model
Derivation for the Diode
around a D.C. Operating
Point

$$v_D + \frac{1}{r_D} \downarrow i_D \xrightarrow{s.s.} v_D + \frac{1}{r_D} \downarrow i_D$$

Worked on this on: Oct 19, 2015/Mon

Small Signal Model Derivation

For a Diode



$$\text{Diode model}$$

$$I_D = I_s \left(\exp\left(\frac{V_D}{nV_T}\right) - 1 \right) \quad (*)$$

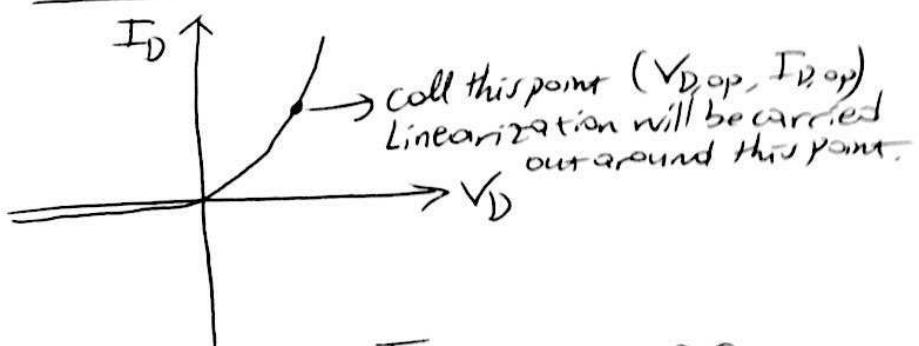
I_s : saturation current

n : emission coef.

Thermal voltage $\leftarrow V_T = \frac{kT}{q} \rightarrow$ Temperature in Kelvin
 Boltzmann constant

charge of a single carrier

I_D vs V_D plot (DC transfer characteristic)



Around $(V_{D,op}, I_{D,op})$ [called the DC operating point, Q-point, or simply DC O.P.]

(*) can be approximated as

$$I_D \approx I_s \exp\left(\frac{V_D}{nV_T}\right)$$

Some Definitions

V_D : DC voltage across the diode

\bar{V}_D : small perturbation on the DC voltage

$$\bar{V}_D = V_D + \bar{v}_D \Rightarrow \text{sum of the two}$$

I_D : DC current through the diode

\bar{I}_D : small perturbation on the DC current

$$\bar{I}_D = I_D + \bar{i}_D$$

—————

Note that

$$I_D = I_S \exp\left(\frac{V_D}{nV_T}\right) \text{ implies } \bar{I}_D = I_S \exp\left(\frac{\bar{V}_D}{nV_T}\right)$$

(**) 

Now Taylor-expand both sides
of the equation in (**).

From (**):

$$(**) I_{D,op} + \bar{i}_d = \sum_{k=0}^{+\infty} \frac{1}{k!} \left. \frac{\partial^k f(x)}{\partial x^k} \right|_{x=V_{D,op}} \cdot (V_D - V_{D,op})^k$$

The expansion is around $V_{D,op}$

The whole expansion above is around the DC operating point $(V_{D,op}, I_{D,op})$.

First note in (**) that

$$V_D - V_{D,op} = \underbrace{v_d}_{\text{small signal perturbation}}$$

Second note that

$$f(x) = I_s \exp\left(\frac{x}{nV_T}\right)$$

Third express the sum in the RHS of (**)

as two terms as in:

$$(**) I_{D,op} + \bar{i}_d = \sum_{k=0}^1 \frac{1}{k!} \left. \frac{\partial^k f(x)}{\partial x^k} \right|_{x=V_{D,op}} v_d^k + \sum_{k=2}^{+\infty} \frac{1}{k!} \left. \frac{\partial^k f(x)}{\partial x^k} \right|_{x=V_{D,op}} \bar{i}_d^k$$

In (***) the terms with $\sum_{k=2}^{+\infty}$

correspond to higher order terms
and can be omitted if \bar{v}_d is small, which it is.

Then from (****) we have:

$$(\text{****}) \quad I_{D,op} + \bar{v}_d \approx \frac{1}{0!} \left. \frac{\partial^0 f(x)}{\partial x^0} \right|_{x=v_{D,op}} \sqrt{d}^0$$

$$+ \frac{1}{1!} \left. \frac{\partial^1 f(x)}{\partial x^1} \right|_{x=v_{D,op}} \sqrt{d}^1$$

Above $\sqrt{d}^0 = 1 \quad \sqrt{d}^1 = \sqrt{d}$
 $0! = 1 \quad 1! = 1$
 $\frac{\partial^0 f(x)}{\partial x^0} = f(x) \quad \frac{\partial f(x)}{\partial x} = \frac{I_s}{nV_T} \exp\left(\frac{x}{nV_T}\right)$
 $= I_s \exp\left(\frac{x}{nV_T}\right)$

Then

$$(\text{****}) \quad I_{D,op} + \bar{v}_d = I_s \exp\left(\frac{x}{nV_T}\right) \Big|_{x=v_{D,op}} \rightarrow I_s \exp\left(\frac{v_{D,op}}{nV_T}\right) = I_{D,op}$$

$$+ \frac{I_s}{nV_T} \exp\left(\frac{x}{nV_T}\right) \Big|_{x=v_{D,op}} \cdot \sqrt{d}$$

$$\downarrow$$

$$I_{D,op}$$

We have

~~$I_{D,op} + \bar{v}_d = I_{D,op} + \frac{I_{D,op}}{nV_T} \cdot \sqrt{d}$~~

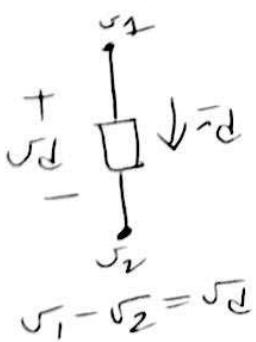
Therefore the fundamental relation for a diode in small-signal around an operating point $(V_{D,op}, I_{D,op})$ is:

$$\bar{v}_d = \frac{I_{D,op}}{nV_T} \cdot v_d$$

manipulate this to obtain

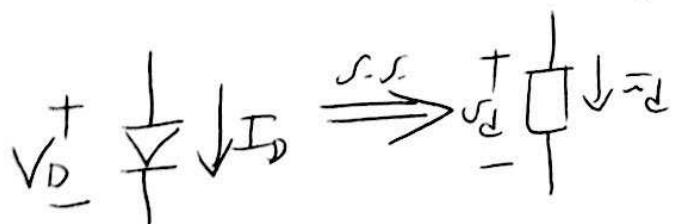
$$(\star\star\star\star) \quad \frac{\bar{v}_d}{\bar{i}_d} = \frac{nV_T}{I_{D,op}}$$

This equation defines the small-signal model for a diode:

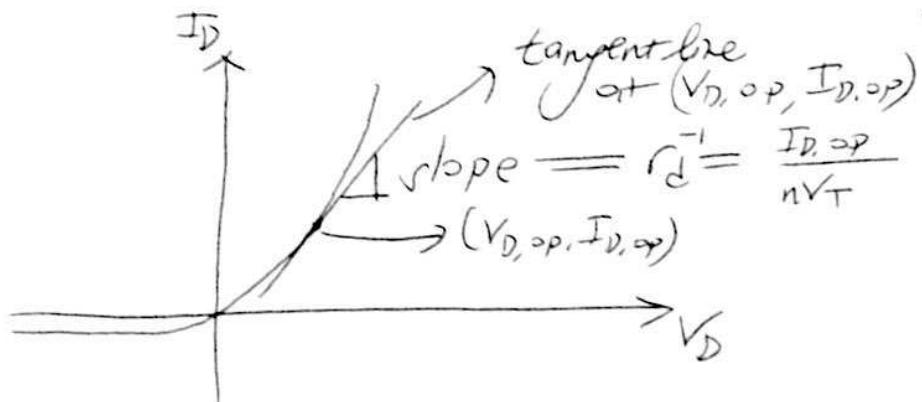


The current for this passive element flows in between the nodes ① and ② whose difference defines v_d . i_s and \bar{i}_d are related by a constant factor (depending on $I_{D,op}$). See $(\star\star\star\star\star\star)$.

Then the small-signal model is indeed a resistance with value $r_d = \frac{\bar{v}_d}{\bar{i}_d} = \frac{nV_T}{I_{D,op}}$.



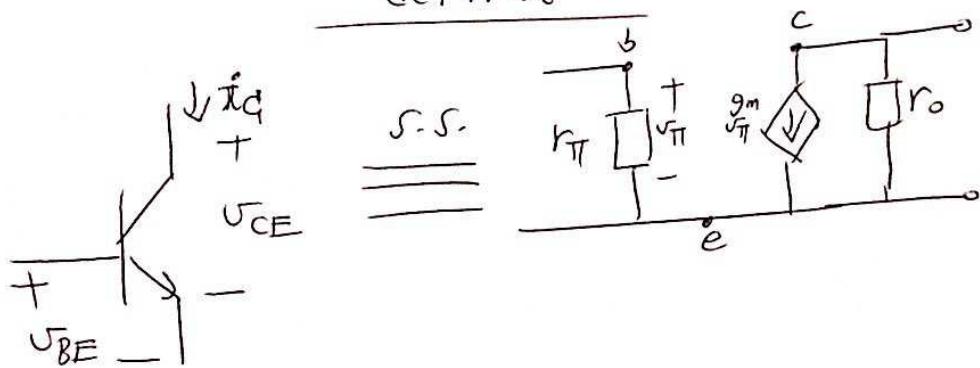
How does this resistance relate to the I_D vs V_D plot (the DC transfer characteristic)?



For higher values of $I_{D,op}$, I_D vs V_D plot is steeper, r_d^{-1} is higher, r_d is lower.
For lower values of $I_{D,op}$, vice versa.

Most of the applications with diodes utilize the large signal model of diodes, these are nonlinear circuits. For Nonlinear applications of diodes, the above derivation is not of the use, since DC op changes with the applied inputs (with a large enough range).

BJT small-signal
equivalent model
derivation



Worked on this on: Nov 5, 2014 / Wed

BJT small-signal model derivation

Fundamental equations:

$$I_C = I_S \exp\left(\frac{V_{BE}}{V_T}\right) \quad (1)$$

Diagram illustrating the components of the fundamental equation:

- Collector current (I_C)
- Saturation current (I_S)
- Thermal voltage (V_T)
- Base-emitter voltage (V_{BE})

$$I_C = \beta I_B \quad (2)$$

Diagram illustrating the components of the fundamental equation:

- Collector current (I_C)
- Base current (I_B)
- Current gain (β)

Defns:

$$v_{BE} = V_{BE} + v_{be} \quad (3)$$

Diagram illustrating the components of the definition equation:

- DC value (V_{BE})
- Small signal value (v_{be})

$$\bar{v}_q = \bar{I}_q + \bar{v}_c \quad (4)$$

Diagram illustrating the components of the definition equation:

- DC value (\bar{I}_q)
- Small signal value (\bar{v}_c)

$$\bar{v}_B = \bar{I}_B + \bar{v}_b \quad (5)$$

Diagram illustrating the components of the definition equation:

- DC value (\bar{I}_B)
- Small signal value (\bar{v}_b)

Apply (3) and (4) on (1) \Rightarrow perturb the variables of (1)
 and obtain a relation between the small signal parameters

$$\bar{i}_c = I_s \exp\left(\frac{v_{BE}}{V_T}\right)$$

$$I_d + \bar{i}_c = I_s \exp\left(\frac{v_{BE} + v_{be}}{V_T}\right)$$

apply Taylor expansions

$$\left[I_s \exp\left(\frac{v}{V_T}\right) \right] \Big|_{v=V_{BE}}$$

$$+ \frac{d}{dv} \left[I_s \exp\left(\frac{v}{V_T}\right) \right] \Big|_{v=V_{BE}} \cdot (v_{BE} - V_{BE})$$

+ Higher Order Terms (H.O.T.)

Note that $\overbrace{I_d = I_s \exp\left(\frac{v_{BE}}{V_T}\right)}$

Now we have (canceling equal terms
and H.O.T.)

$$\bar{z}_c \approx \frac{I_s}{V_T} \exp\left(\frac{V_{BE}}{V_T}\right) \cdot V_{be}$$

$\frac{I_s}{V_T}$ \Rightarrow DC value

call
the
transconductance

$$g_m \triangleq \frac{\bar{z}_c}{V_{be}} = \frac{I_d}{V_T} \quad (6)$$

Apply (4) and (5) on (2).

$$\bar{z}_d = \beta \bar{z}_b$$

$$I_d + \bar{z}_c = \beta (\bar{z}_b + \bar{z}_b) \quad \begin{array}{l} \text{(Taylor expansion is not} \\ \text{necessary, this is a} \\ \text{linear relation)} \end{array}$$

Cancel
equal
terms

$$\bar{z}_c = \beta \bar{z}_b$$

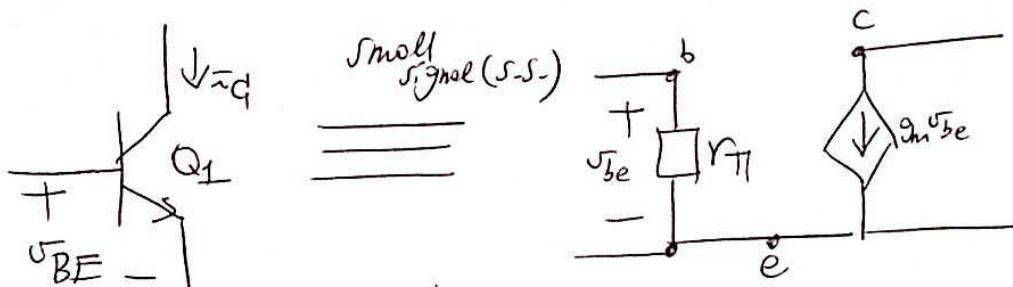
(7)

call the
input
impedance

$$R_{II} = \frac{V_{be}}{\bar{z}_b} = \frac{\frac{\bar{z}_c}{g_m}}{\frac{\bar{z}_c}{\beta}} \xrightarrow{\text{From (6)}} \text{From (7)}$$

$$= \frac{\beta}{g_m} \quad (8)$$

The small-signal equivalent model thus far is



Transistor Q_1 is DC-biased in forward active mode.

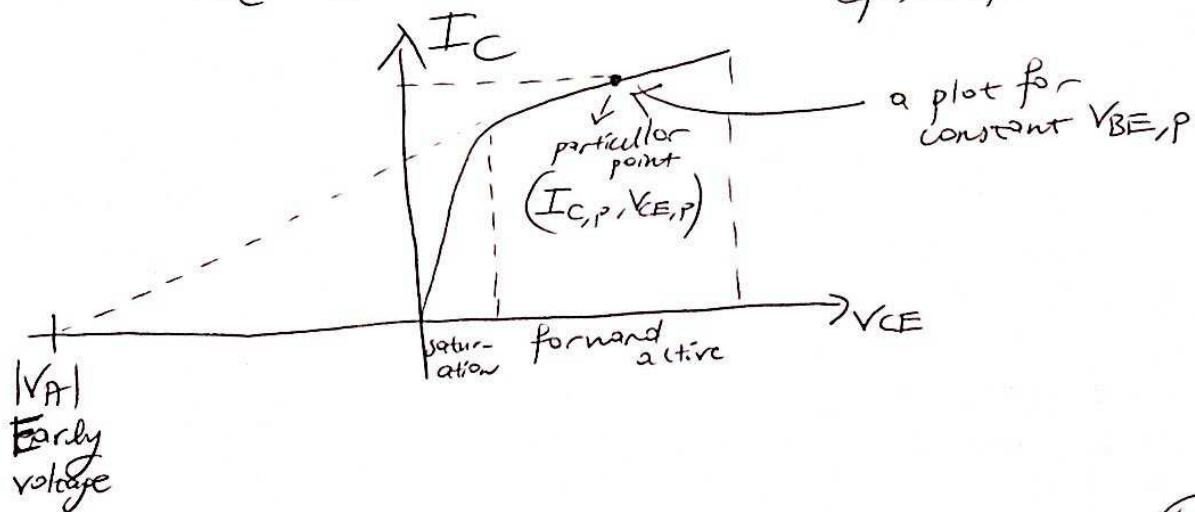
We will next account for the output impedance.

r_o appears due to base-width modulation (otherwise called the Early effect).

Due to this effect, the collector current I_C continues to increase with increasing V_{CE} .

r_o is the parameter that expresses this effect in the small signal model.

$I_C - V_{CE}$ (DC transfer characteristics) plot of the npn BJT



The slope in the $I_C - V_{CE}$ (forward active mode) plot

is given by

$$\text{slope} = \frac{I_{C,P}}{|V_A| + V_{CE,P}}$$

The slope is almost constant in this regime (forward-active).

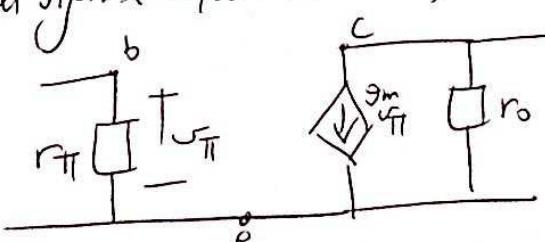
Typically $|V_A| \gg V_{CE,P}$.

Therefore

$$\text{slope} \approx \frac{I_{C,P}}{|V_A|} \quad \text{at a particular DC operating point}$$

$$\text{The s.s. output impedance } r_o = \frac{1}{\text{slope}} = \frac{|V_A|}{I_{C,P}} \quad \text{at the DC o.p. } (V_{BE,P}, V_{CE,P}, I_{C,P})$$

The small signal equivalent model of BJT is then

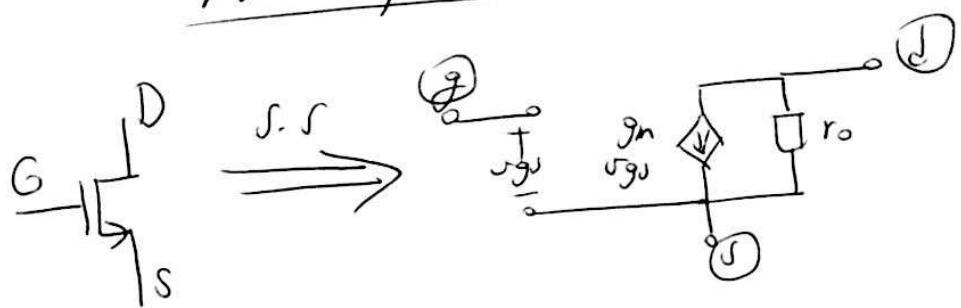


$$\text{where } U_{\pi\pi} = U_{be}$$

$$r_{\pi\pi} = \frac{\beta}{g_m} \quad r_o = \frac{|V_A|}{I_Q}$$

$$g_m = \frac{I_d}{V_T}$$

Derivation of the small signal model for MOS



Worked on this on:

April 9, 2015 / Thursday

Small-signal model derivation for NMOS

The large signal model can be approximated as in

$$I_D = \frac{1}{2} \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_t)^2$$

Some definitions:

I_D : DC value of the drain current

\bar{i}_d : small perturbation on the drain current

$$\bar{i}_D = I_D + \bar{i}_d$$

V_{GS} : DC value of the gate-source voltage

v_{gs} : small perturbation on the gate-source voltage

$$v_{GS} = V_{GS} + v_{gs}$$

Define function f :

$$I_D = f(V_{GS})$$

see here

It is also true that

$$\bar{i}_D = f(v_{GS})$$

Then we have

$$I_D + \bar{I}_D = f(V_{GS} + v_{GS})$$

Expand the RHS around V_{GS} (the DC value)

$$I_D + \bar{I}_D = \sum_{k=0}^{+\infty} \frac{1}{k!} \left. \frac{\partial^k f(x)}{\partial x^k} \right|_{x=V_{GS}} \cdot (v_{GS} - V_{GS})^k$$

We will keep the first two terms on the RHS

and omit the rest (the higher order terms or H.O.T.)

$$I_D + \bar{I}_D \approx \sum_{k=0}^1 \frac{1}{k!} \left. \frac{\partial^k f(x)}{\partial x^k} \right|_{x=V_{GS}} \cdot (v_{GS} - V_{GS})^k$$

$$= \frac{1}{0!} f(x) \Big|_{x=V_{GS}} \cdot (v_{GS} - V_{GS})^0$$

$$+ \frac{1}{1!} \left. \frac{\partial f(x)}{\partial x} \right|_{x=V_{GS}} \underbrace{(v_{GS} - V_{GS})^1}_{\text{see the defns on pg 1}}$$

$$= f(V_{GS})$$

$$+ \left. \frac{\partial f(x)}{\partial x} \right|_{x=V_{GS}} \cdot v_{GS}$$

In all we have

$$\cancel{I_D + \bar{I}_D} = f(V_{GS}) + \left. \frac{\partial f(x)}{\partial x} \right|_{x=V_{GS}} \cdot v_{GS}$$

\rightarrow see pg 1 how these cancel

We now have a relation between \bar{i}_d and v_{gs}
 (the small-signal variables)

$$\bar{i}_d = \left. \frac{\partial f(x)}{\partial x} \right|_{x=V_{GS}} \cdot v_{gs}$$

In this equation

$$\begin{aligned} \left. \frac{\partial f(x)}{\partial x} \right|_{x=V_{GS}} &= \left. \frac{\partial}{\partial x} \left[\frac{1}{2} \mu_n C_{ox} \frac{W}{L} (x - V_t)^2 \right] \right|_{x=V_{GS}} \\ &= \left. \left[\mu_n C_{ox} \frac{W}{L} (x - V_t) \right] \right|_{x=V_{GS}} \end{aligned}$$

the 1st
formula
for the
transconductance
of MOS

$$g_m = \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_t)$$

$$\bar{i}_d = g_m v_{gs}$$

this is the
proportionality
constant between
 \bar{i}_d and v_{gs}

Note that

$$\begin{aligned} \left[2 \mu_n C_{ox} \frac{W}{L} \right] \left[\frac{1}{2} \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_t)^2 \right] &= \\ \left[\mu_n C_{ox} \frac{W}{L} (V_{GS} - V_t) \right]^2 &= g_m^2 \end{aligned}$$

Therefore the second formula for the transconductance of MOS is

$$g_m = \sqrt{2\mu_n C_{ox} \frac{W}{L} I_D}$$

Also note that

$$\frac{2 \left[\frac{1}{2} \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_t)^2 \right]}{(V_{GS} - V_t)} = \mu_n C_{ox} \frac{W}{L} (V_{GS} - V_t)$$

therefore the third formula for the transconductance of MOS is

$$g_m = \frac{2 I_D}{V_{GS} - V_t}$$

Now compare the g_m 's of BJT and MOS (using the 3rd formula for the g_m of MOS)

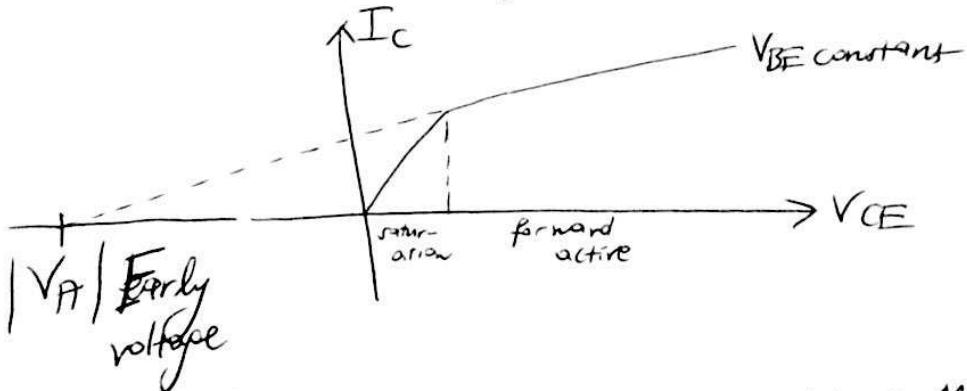
	BJT	MOS
g_m	$\frac{I_C}{V_T}$	$\frac{I_D}{(V_{GS} - V_t)/2}$

V_T is 25-30 mV.

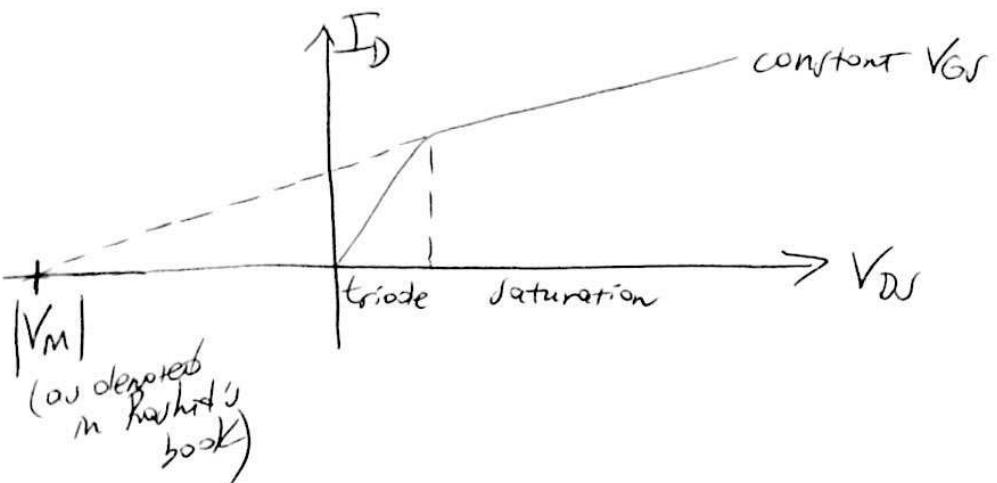
$\frac{V_{GS} - V_t}{2}$ is on the order of 100s of millivolts.

Therefore the transconductance/current ratio of BJTs is higher than that of MOS.
BJT is said to be stronger than MOS for this reason.

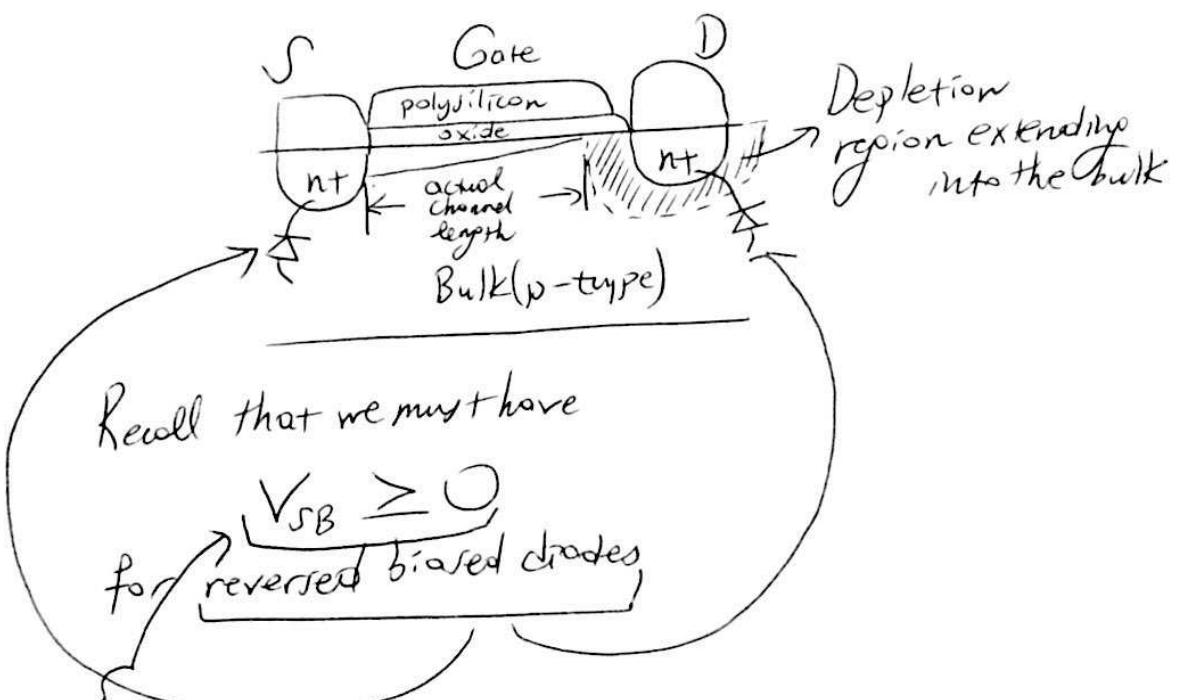
Recall that base-width modulation is responsible in BJTs for the slight increase of I_C with increasing V_{CE} in the forward active regime. Another name for this phenomenon is ^{the} Early effect.



Channel length modulation is responsible in MOS for the slight increase of I_D with increasing V_{DS} in the saturation regime.



The structure of the NMOS,



Recall that we must have

$$V_{SB} \geq 0$$

for reversed biased drains

$V_D > V_S$ makes the junction we call the drain.

This condition ensures

$$V_{DB} \geq 0$$

since we have this.

The B-D pn-junction is more reverse-biased than the B-S pn-junction. The depletion region into the bulk of B-D junction widens into the channel, shortening the effective channel length.

The equation for I_D is

$$I_D = \frac{1}{2} \mu n C_{ox} \frac{W}{L} (V_{GS} - V_t)^2$$

In this eqn, shorter L means larger current.

Therefore, by this effect (channel length modulation) we have the I_D current slightly increasing with V_{DS} .

The I_D equation can be remodeled as

$$I_D = \frac{1}{2} \mu n C_{ox} \frac{W}{L} (V_{GS} - V_t)^2 [1 + \lambda V_{DS}]$$

where λ ($\%/\text{V}$) models the rise in I_D with V_{DS} .

Now I_D is a function of both V_{GS} and V_{DS}

$$I_D = f(V_{GS}, V_{DS})$$

It is also true that

$$I_D + \bar{I}_D = f(V_{GS}, V_{DS})$$

$\downarrow \quad \downarrow \quad \downarrow$

$$V_{GS} + v_{GS} \quad V_{DS} + v_{DS}$$

The first few terms of the Taylor expansion
for the RHS reads:

$$I_D + \bar{I}_D \approx f(V_{GS}, V_{DS})$$
$$+ \frac{\partial f(x,y)}{\partial x} \Big|_{\substack{x=V_{GS} \\ y=V_{DS}}} \cdot \underbrace{(V_{GS} - V_{GS})}_{v_{GS}}$$
$$+ \frac{\partial f(x,y)}{\partial y} \Big|_{\substack{x=V_{GS} \\ y=V_{DS}}} \cdot \underbrace{(V_{DS} - V_{DS})}_{v_{DS}}$$

$$\frac{\partial f(x,y)}{\partial x} \Big|_{\substack{x=V_{GS} \\ y=V_{DS}}} \approx g_m = \text{the three formulas as found before on pg 3-4}$$

$$\left. \frac{\partial f(x,y)}{\partial y} \right|_{\substack{x=V_{GS} \\ y=V_{DS}}} = \frac{\partial}{\partial y} \left[\frac{1}{2} \mu_n C_o x \frac{W}{L} (x - V_t)^2 [1 + \lambda y] \right]$$

$$= \lambda \left[\frac{1}{2} \mu_n C_o x \frac{W}{L} (V_{GS} - V_t)^2 \right] \approx I_D$$

$$\cong \lambda I_D$$

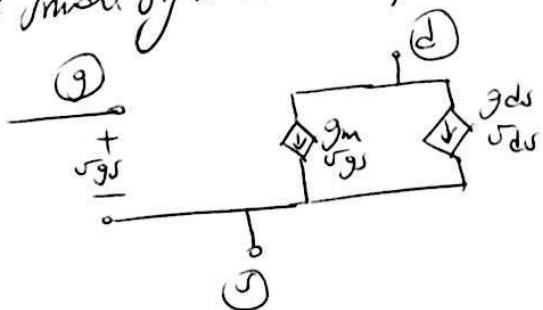
In all

$$I_D + i_d = f(V_{GS}, V_{DS}) + g_m v_{GS} + \lambda I_D v_{DS}$$

~~$I_D + i_d = f(V_{GS}, V_{DS})$~~

these two cancel since they are the same

In the small signal model of MOS



The element

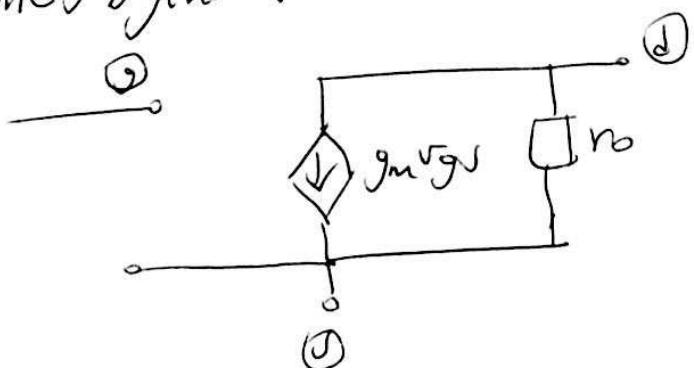
controlled by the voltages of the nodes it is connected to. This can be modelled by a linear resistor since it is linear.

The resistance value r_o is

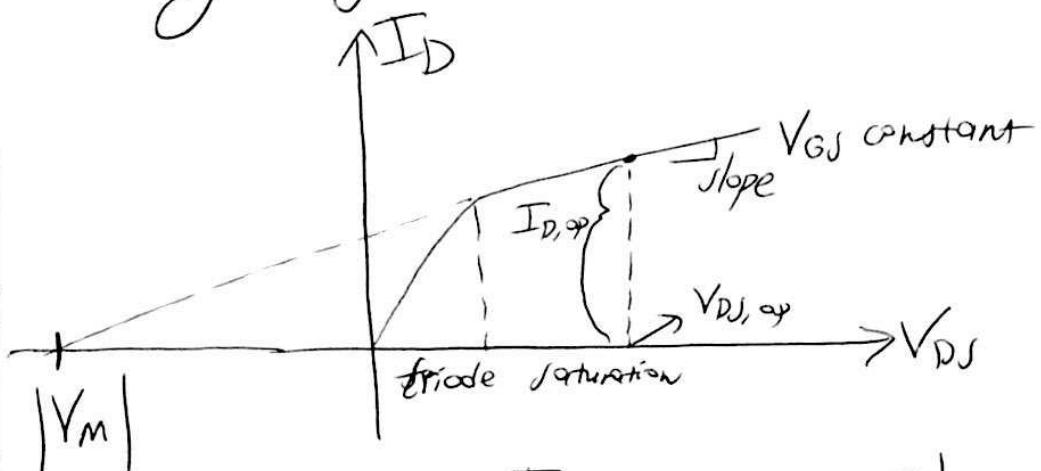
$$r_o = \frac{1}{g_{ds}} = \frac{1}{\lambda I_D} \Rightarrow \left(\frac{V}{A} \right) \text{ in units}$$

\downarrow
 $(\frac{1}{V})$ (A)

Therefore the simplified small signal model of MOS is given as:



In some books λ is not given, but the channel length modulation is modelled with a voltage that corresponds to the Early voltage in BJTs.



$$\text{slope} = \frac{I_{D,op}}{|V_m| + V_{DS,op}} = (r_o)^{-1}$$

$$r_o = \frac{|V_m| + V_{DS,op}}{I_{D,op}}$$

Since usually $|V_m| \gg V_{DS,op}$

$$r_o \approx \frac{|V_m|}{I_{D,op}}$$

3. Diode Applications

3.1 Clamping Circuits

Diode Applications

3.52 Figure P3.52 shows a clamping circuit and its input voltage waveform.

- Carefully sketch the output waveform if the diode is ideal.
- Sketch the output waveform if the diode has an offset of 0.6 V.
- Repeat part (a) after reversing the diode orientation.
- Repeat part (b) after reversing the diode orientation.

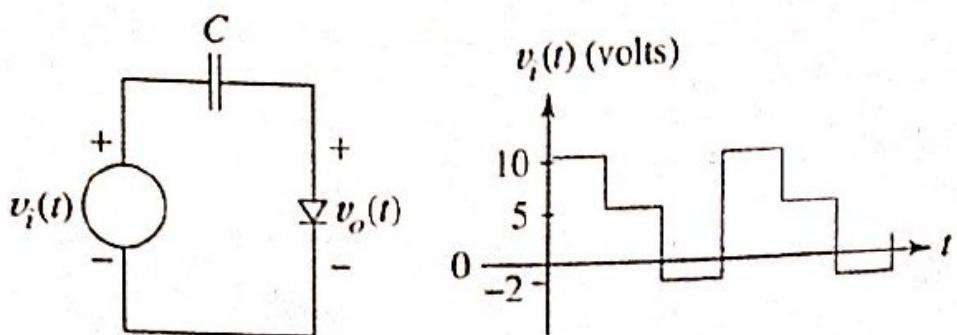
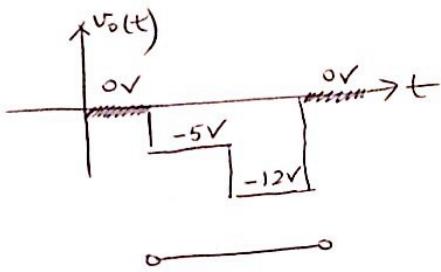


Figure P3.52

capacitor voltage charged to
 $v_{\text{z,peak}}$ (positive peak)

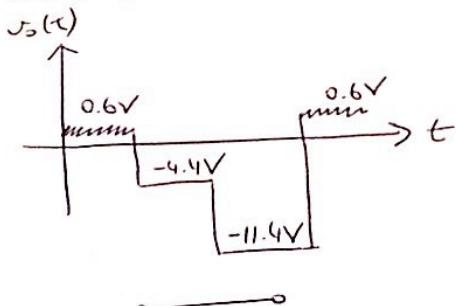
$$= 10 \text{ V}$$

After reaching the first positive peak in the periodic waveform $\Rightarrow v_o(t) = v_z(t) - v_{\text{z,peak}}$
 $= v_z(t) - 10 \text{ V}$



The cap charges to $v_{\text{z,peak}} - v_{\text{o,ON}}$
 $v_{\text{z,peak}} - 0.6 \text{ V} = 9.4 \text{ V}$

After the init. transient $\Rightarrow v_o(t) = v_z(t) - 9.4 \text{ V}$



In the initial case \Rightarrow cap has no charge on it
 \rightarrow the diode turns on when

$$v_z(t) = 0$$

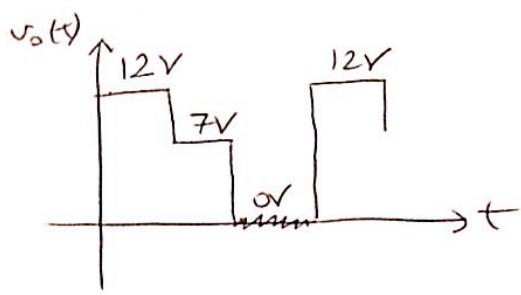
\rightarrow then the cap voltage is charged

$$\rightarrow -2 \text{ V}$$

\rightarrow the diode turns off and never turns on again

\Rightarrow after the init transient $v_o(t) = v_z(t) + 2 \text{ V}$

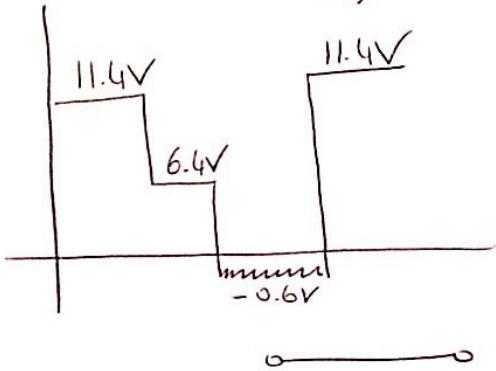
c)



Mohit 3.52
cont.

c)
cont-

Similar to (c) but the diode turns ON
when $v_o(t) = -0.6V$
the cap charges to $-1.4V$
after the init. transient
 $v_o(t) = v_{in}(t) + 1.4V$



4. Amplifiers

4.1 Types

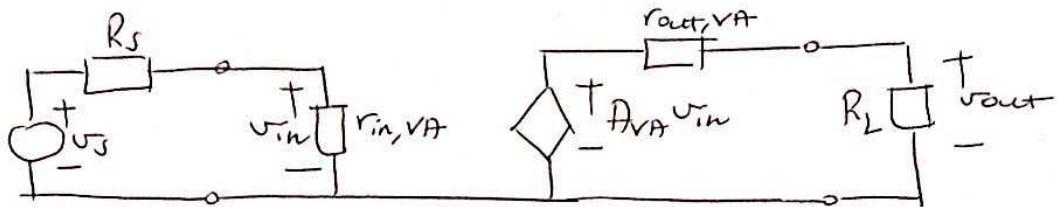
Properties of Voltage Amplifiers (VA)

⇒ Input/output impedances (for an ideal VA)

⇒ Load/source impedances (for an ideal VA)

Worked on this on : Oct 30, 2014 / Th

Voltage amplifiers - properties



Realized voltage gain

$$A_{VA,realized} = \frac{r_{in,VA}}{r_{in,VA} + R_s} A_{VA} \frac{R_L}{R_L + r_{out,VA}}$$

if $\begin{cases} r_{out,VA} \rightarrow 0 \\ r_{in,VA} \rightarrow \infty \end{cases} \Rightarrow A_{VA,realized} \rightarrow A_{VA}$

These are the properties of an ideal voltage amplifier with finite gain (A_{VA}).

if $\begin{cases} R_L \rightarrow \infty \\ R_s \rightarrow 0 \end{cases} \Rightarrow A_{VA,realized} \rightarrow A_{VA}$

These are the ideal load

and ideal source for a voltage ampli.

In the case that the V_{in} is loaded and sourced by these quantities, $r_{out,VA}$ and $r_{in,VA}$ will have no effect. The V_{in} will act as if it is ideal with finite gain (A_{VA}).

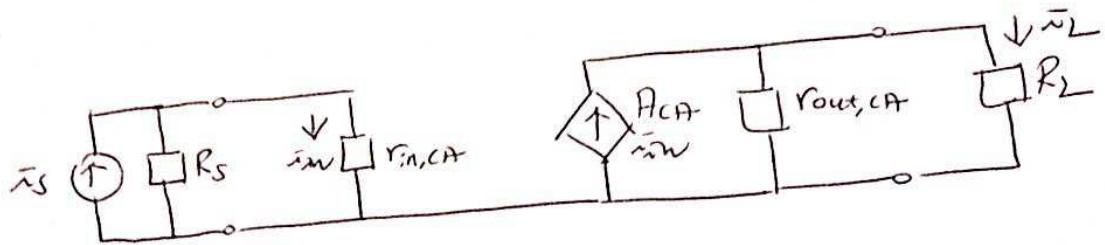
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Properties of Current Amplifiers

- ⇒ Input/Output impedances for an ideal CA
- ⇒ Ideal source and load impedances for any CA

Worked on this on: Oct 30, 2014 / Th

Current amplifiers - properties



Realized current gain

$$A_{CA, \text{realized}} = \frac{R_s}{R_s + r_{in, CA}} A_{CA} \frac{r_{out, CA}}{r_{out, CA} + R_L}$$

if $r_{in, CA} \rightarrow 0$ } $\Rightarrow A_{CA, \text{realized}} = A_{CA}$
 $r_{out, CA} \rightarrow \infty$ } for any source
 and load

↓
 properties of
 an ideal current
 amplifier with
 finite gain (A_{CA}).
 ↓

if $R_s \rightarrow \infty$ } $\Rightarrow A_{CA, \text{realized}} = A_{CA}$
 $R_L \rightarrow 0$ }

↓
 ideal source and
 load for current
 amplifiers, respectively.

If a CA is sourced and loaded by these impedances, the CA will act as an ideal CA with finite gain (A_{CA}).

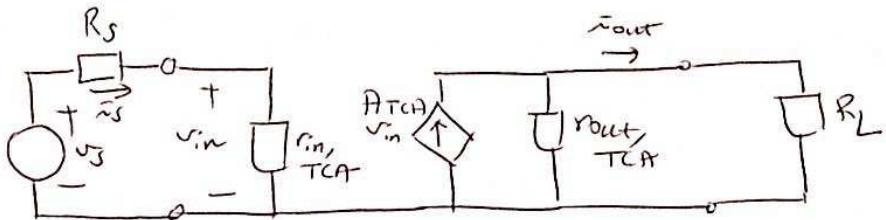
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Properties of the Transconductance Amplifier

- ⇒ properties of an ideal amplifier
- ⇒ ideal input/output impedances
- ⇒ ideal source and load
- ⇒ how to compute input/output impedances and gain

Worked on this on: Oct 19, 2014 Sun

Properties of the transconductance amplifier



$$\frac{v_{out}}{v_s} = \frac{r_{in,TCA}}{r_{in,TCA} + R_s} A_{TCA} \frac{r_{out,TCA}}{r_{out,TCA} + R_L} = A_{TCA, \text{realized}}$$

If $r_{in,TCA} \rightarrow +\infty$ and $r_{out,TCA} \rightarrow +\infty$ then the transconductance amplifier is ideal with finite gain, and $A_{TCA, \text{realized}} = A_{TCA}$

If $R_s = 0$ and $R_L = 0$

Computing A_{TCA}

⇒ Connect v_s without R_s so that $v_s = v_{in}$

⇒ set R_L to short circuit

⇒ Compute $\frac{v_{out}}{v_s} = A_{TCA}$

Computing $r_{in,TCA}$

⇒ Connect v_s without R_s so that $v_s = v_{in}$

⇒ set R_L to short circuit

⇒ Compute $\frac{v_s}{v_{in}} = r_{in,TCA}$

Computing $r_{out, TCA}$

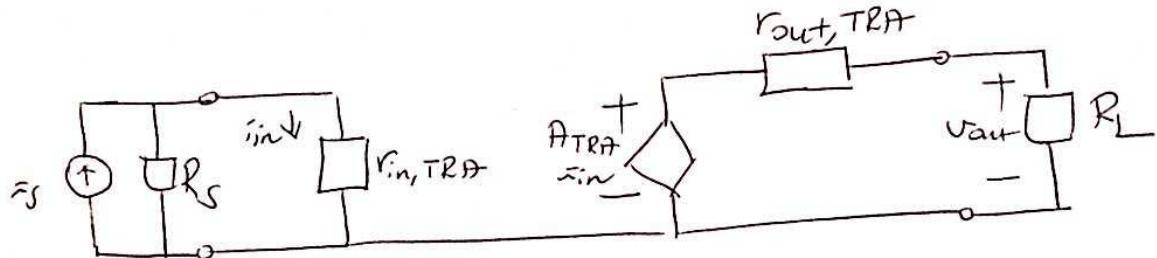
- ⇒ set $v_s = 0$ and connect no R_s
- ⇒ connect a v_{test} instead of R_L
(i_{test} flows out of the positive terminal of v_{test})
- ⇒ compute $\frac{v_{test}}{i_{test}} = r_{out, TCA}$

Properties of Transresistance Amplifiers

- ⇒ Input/output impedances
for the ideal TRA
- ⇒ Ideal source/load impedances
for any TRA

Worked on this on: Oct 30, 2014 / Th

Transresistance amplifier - properties



$A_{TRA, \text{realized}} \Rightarrow \text{realized gain}$

$$A_{TRA, \text{realized}} = \frac{R_s}{R_s + r_{in, TRA}} A_{TRA} \frac{R_L}{R_L + r_{out, TRA}}$$

if $r_{in, TRA} \rightarrow 0$ } $\Rightarrow A_{TRA, \text{realized}} = A_{TRA}$
 $r_{out, TRA} \rightarrow 0$ }

ideal input/output impedances for the TRA with finite gain (A_{TRA})

if $R_L \rightarrow +\infty$ } $\Rightarrow A_{TRA, \text{realized}} = A_{TRA}$
 $R_s \rightarrow +\infty$ }

ideal load/source impedances
(respectively) for the TRA

if the TRA is loaded/sourced
by these impedances, the TRA
will act as an ideal TRA with finite gain (A_{TRA}).

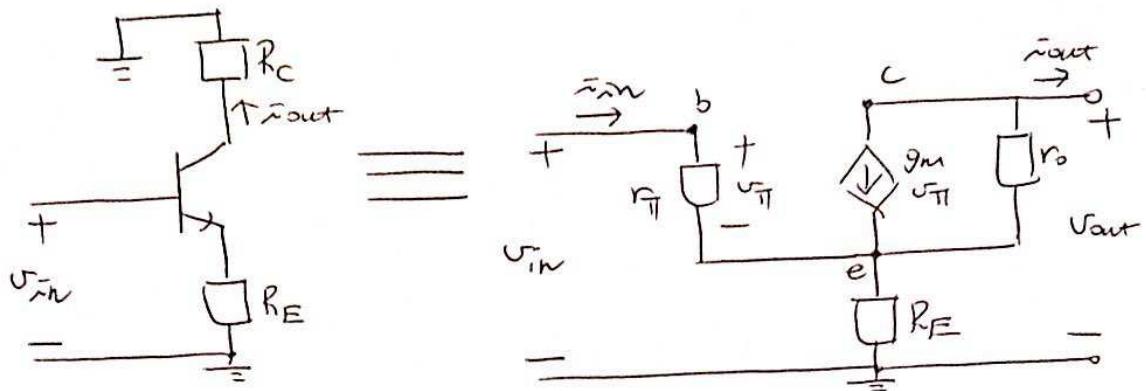
4.2 Modeling Examples

Modeling the Common Emitter
Amplifier as a
Transconductance Amplifier

⇒ A_{TCA} , $r_{in,TCA}$, $r_{out,TCA}$
exact and approximate results

Worked on this on: Oct 31, 2014/F

Analysis of the common emitter amplifier
as a transconductance amplifier



⇒ Computation of
 $A_{TCA} = \frac{\bar{v}_{out}}{\bar{v}_{in}} \Big|_{v_{out}=0}$

$$v_e = v_{in} - v_\pi$$

KCL at (e)

$$\frac{v_{in} - v_\pi - v_e}{r_\pi} + \frac{v_{in} - v_\pi}{R_E} - g_m v_\pi + \frac{v_{in} - v_\pi}{r_o} = 0 \quad (1)$$

Ignore this term
assuming $r_o \rightarrow \infty$

$$-g_m v_\pi = \bar{v}_{out} \quad (\text{assuming } r_o \rightarrow \infty) \quad (2)$$

From (1) we have

$$\frac{v_{in}}{R_E} = \frac{v_\pi}{r_\pi \parallel R_E \parallel \frac{1}{g_m}} \Downarrow \frac{-\bar{v}_{out}}{g_m} = \frac{v_\pi}{r_\pi \parallel R_E \parallel \frac{1}{g_m}}$$

From (2)

(3)

From (3), we have

$$\frac{v_{out}}{v_{in}} = - \frac{g_m}{R_E} \left[r_\pi \parallel R_E \parallel \frac{1}{g_m} \right]$$

$$= - \frac{g_m}{R_E} \left[\frac{\frac{r_\pi}{R_E/g_m}}{r_\pi + \frac{R_E/g_m}{R_E + 1/g_m}} \right]$$

$$= - \left[\frac{r_\pi \frac{g_m}{1+g_m R_E}}{r_\pi + \frac{R_E}{1+g_m R_E}} \right]$$

$$A_{TCA} = - \left[\frac{r_\pi g_m}{R_E + (1+g_m R_E) r_\pi} \right] \quad (4)$$

$$A_{TCA} = - \left[\frac{\beta}{r_\pi + (1+\beta) R_E} \right] \quad (5)$$

From (4)
we have $A_{TCA} \approx - \frac{g_m}{1+g_m R_E}$

(6) approximate answer

$$\Rightarrow \text{computation of } \left. \frac{v_{in}}{\tilde{v}_{in}} \right|_{v_{out}=0} = r_{in, TCA}$$

Note that

$$\tilde{v}_{in} = \frac{v_{\pi}}{r_{\pi}} \quad \frac{v_{in}}{\tilde{v}_{in}} = \frac{v_{in}}{v_{\pi}} r_{\pi} = - \frac{v_{in}}{v_{out}} g_m r_{\pi} \quad (7)$$

$$\text{From (2)} \Rightarrow v_{\pi} = - \frac{v_{out}}{g_m}$$



From (7) and (5)

$$r_{in, TCA} = \frac{v_{in}}{\tilde{v}_{in}} = - \left[- \frac{r_{\pi} + (1+\beta) R_E}{\beta} \right] \cancel{\beta} = r_{\pi} + (1+\beta) R_E \quad (8)$$

exact answer
(for $r_o \rightarrow \infty$)

From (7) and (6)

$$r_{in, TCA} = \frac{v_{in}}{\tilde{v}_{in}} \approx - \left[- \frac{1 + g_m R_E}{g_m} \right] \cancel{g_m} r_{\pi} \\ = r_{\pi} [1 + g_m R_E] \quad (9)$$

approximate answer (for $r_o \rightarrow \infty$)



\Rightarrow Computation of

$$r_{out, TCA} = - \frac{v_{out}}{i_{out}} \Big|_{v_{in}=0}$$

$$v_e = -\sqrt{\pi}$$

$$(KVL) \Rightarrow v_{out} = (-i_{out} - g_m v_{\pi}) r_o - \sqrt{\pi} \\ = -i_{out} r_o - \sqrt{\pi} (1 + g_m r_o) \quad (10)$$

and we have

$$-i_{out} \frac{R_E}{R_E + r_{\pi}} \cdot r_{\pi} = -\sqrt{\pi} \\ i_{out} (r_{\pi} // R_E) = \sqrt{\pi} \quad (11)$$

From (10) and (11)

$$v_{out} = -i_{out} \left[r_o + (r_{\pi} // R_E) (1 + g_m r_o) \right]$$

$$r_{out, TCA} = - \frac{v_{out}}{i_{out}} \Big|_{v_{in}=0} = r_o + (r_{\pi} // R_E) (1 + g_m r_o) \quad (12)$$

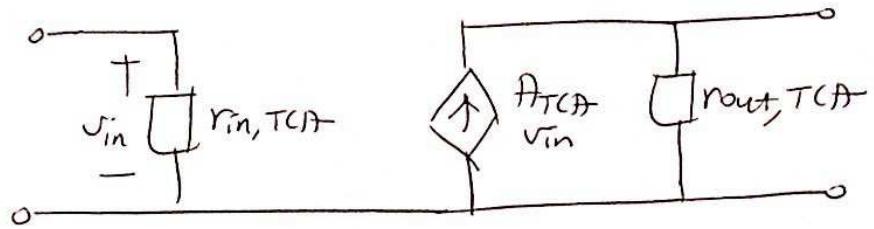
if $r_{\pi} \gg R_E$

$$r_{out, TCA} \approx r_o + R_E (1 + g_m r_o) \\ = R_E + r_o (1 + g_m R_E) \approx r_o (1 + g_m R_E) \quad (13)$$

approx answer

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Transconductance amplifier model



$A_{TCA} \Rightarrow$ see (4), (5) (6)

$r_{in, TCA} \Rightarrow$ see (8), (9)

$r_{out, TCA} \Rightarrow$ see (12), (13)

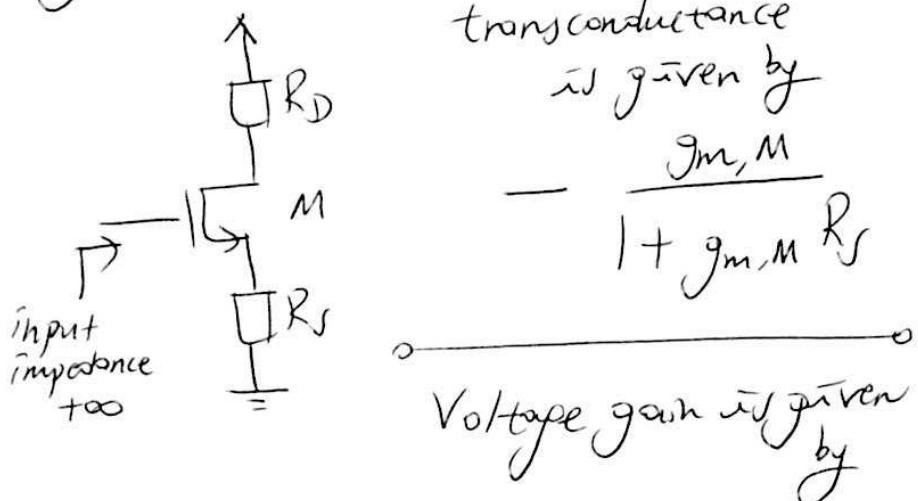
4.3 Design Considerations

How to construct an
amplifier with large input
impedance and big voltage gain?

C_S-CE cascaded stages

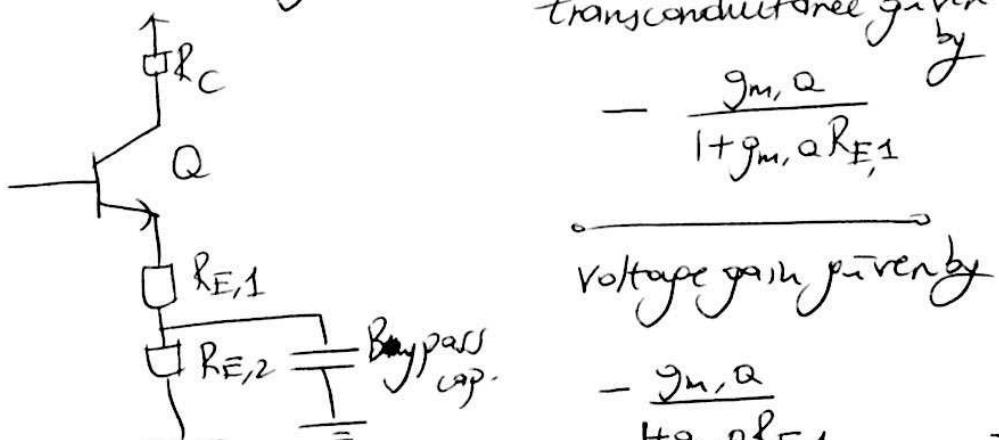
Worked on this on-
April 9, 2015 / Thursday

Consider using a CS (common source) stage at the input

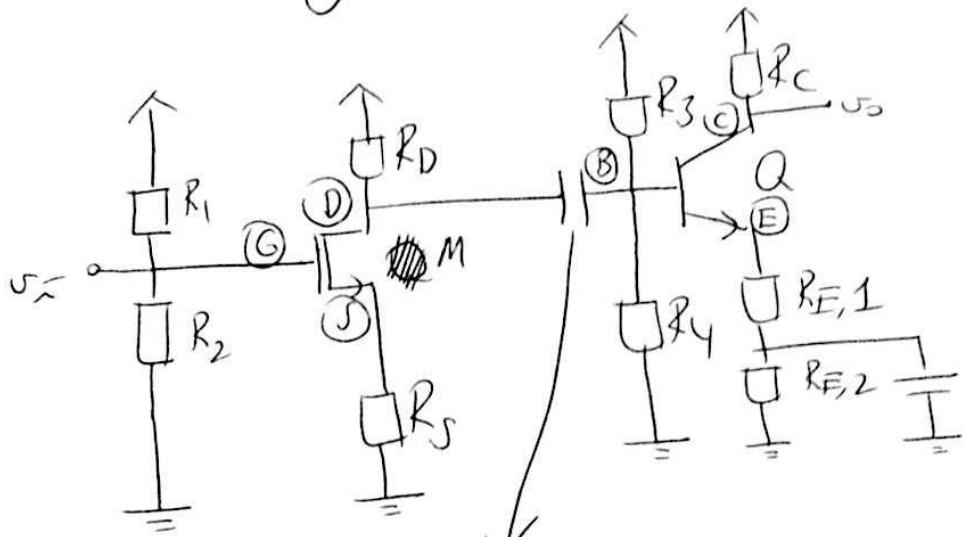


$$- \frac{g_{m,M}}{1 + g_{m,M} R_S} \left[R_D \parallel [r_{o,M}(1 + g_{m,M} R_S)] \right]$$

Cascade this stage with a CE (common emitter)



Two two stages cascaded with complete bias detail:



The DC bias voltage of nodes D and B may be different, then we need this ~~coupling cap.~~ coupling

We desire $R_{E,1} + R_{E,2}$ to establish bias point stability. $R_{E,1}$ is better designed as small since we do not want to lose much of the gain of the second stage, though feedback is retained in AC as well. But then $R_{E,1}$ cannot be too small since the small signal resistance

$$R_{in, base, Q} \approx r_{\pi, Q} (1 + g_{m, Q} R_{E,1})$$

comes in parallel with $R_3 // R_4 // R_D // [r_{o, M} (1 + g_{m, M} R_S)]$

Having $R_{in,base,Q}$ too small due to
 small $R_{E,1}$ would kill the voltage gain of
 the first stage.

The overall expression for the voltage gain is

$$\frac{V_o}{V_i} = \left[-\frac{g_{m,M}}{1 + g_{m,M} R_S} \right] \left[\begin{array}{l} [r_{\pi,Q}(1 + g_{m,Q} R_{E,1})] \\ // R_3 // R_4 // R_D \\ // [r_{o,M}(1 + g_{m,M} R_S)] \end{array} \right]$$

$$\left[-\frac{g_{m,Q}}{1 + g_{m,Q} R_{E,1}} \right] R_C // [r_{o,Q}(1 + g_{m,Q} R_{E,1})]$$


5. Small-Signal Impedances

5.1 Impedances into the Three Nodes of a BJT

Computation of Impedance

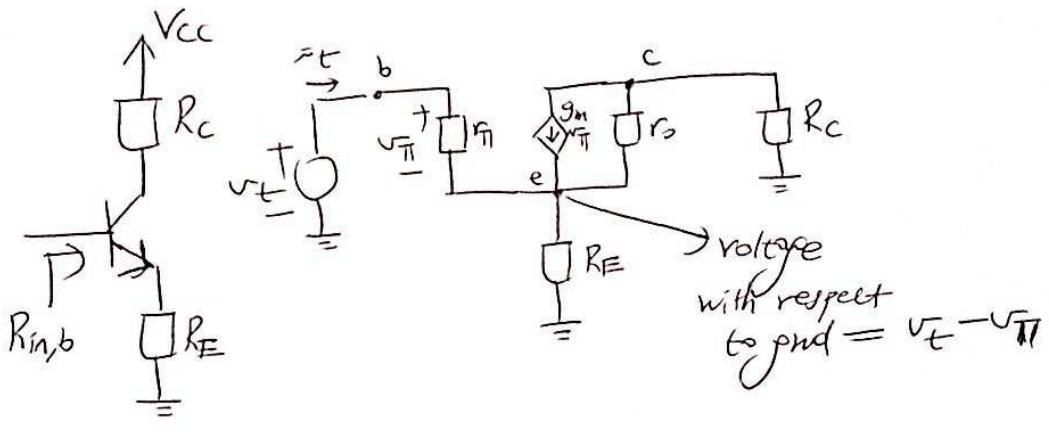
$R_{in,b}$

Looking into the base of
an NPN BJT

worked on this on:-

Oct 19, 2014 / Sun

Computation of R looking into the base



KCL at (e)

$$(1) \quad \frac{v_e - v_\pi}{R_E} + \frac{v_e - v_\pi - v_t}{r_\pi} - g_m v_\pi + \frac{v_t - v_\pi - v_c}{r_o} = 0$$

KCL at (c)

$$(2) \quad \frac{v_c}{R_C} + \frac{v_c - v_t - v_\pi}{r_o} + g_m v_\pi = 0$$

and we have

$$v_t r_\pi = v_\pi$$

From (1) and (2)

$$(3) \quad \frac{v_t}{R_E // r_o} = \frac{v_\pi}{R_E // r_\pi // \frac{1}{g_m} // r_o} + \frac{v_c}{r_o}$$

$$(4) \quad \frac{v_c}{R_C // r_o} + \frac{v_\pi}{r_o // \frac{1}{g_m}} = \frac{v_t}{r_o}$$

\Rightarrow in (3) and (4), solve for v_C :

$$(5) \left[\frac{v_t}{R_E // r_o} - \frac{v_{\pi}}{R_E // r_{\pi} // \frac{r_o}{1+g_m r_o}} \right] r_o$$

$$\left[\frac{v_t}{r_o} - \frac{v_{\pi}}{\frac{r_o}{1+g_m r_o}} \right] (R_C // r_o) = \left[\frac{v_t}{r_o} - \frac{(1+g_m r_o)v_{\pi}}{R_C + r_o} \right] \frac{R_C r_o}{R_C + r_o}$$

$$v_t \frac{R_E + r_o}{R_E} \cdot \cancel{r_o} - v_t \frac{R_C}{R_C + r_o}$$

$$= v_t \left[1 + \frac{r_o}{R_E} - \frac{R_C}{R_C + r_o} \right] = v_t \left[\frac{r_o}{R_E} + \frac{r_o}{R_C + r_o} \right]$$

$$= v_t \left[\frac{r_o}{R_E // (R_C + r_o)} \right]$$

Note that

$$r_o \frac{1}{(R_E // r_{\pi}) // \left(\frac{r_o}{1+g_m r_o} \right)} = \frac{(1+g_m r_o)[R_E // r_{\pi}] + r_o}{(R_E // r_{\pi}) \cancel{r_o}}$$

$$= (1+g_m r_o) + \frac{r_o}{R_E // r_{\pi}}$$

$$v_{\pi} \left[(1+g_m r_o) + \frac{r_o}{(R_E // r_{\pi})} - (1+g_m r_o) \frac{R_C}{R_C + r_o} \right] r_o$$

$$= v_{\pi} \left[(1+g_m r_o) \frac{r_o}{R_C + r_o} + \frac{r_o}{(R_E // r_{\pi})} \right] = v_{\pi} \frac{\frac{r_o}{R_C + r_o} // (R_E // r_{\pi})}{\cancel{(R_E // r_{\pi}) // (R_E // r_{\pi})}}$$

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Therefore from (5), we have

$$\begin{aligned}
 \frac{v_t}{(R_c + r_o) // R_E} &= \frac{\pi_t r_{\pi}}{\left(\frac{(R_c + r_o)}{1 + g_m r_o}\right) // (R_E // r_{\pi})} \\
 R_{in,b} &= \frac{v_t}{\pi_t} = \frac{\frac{(R_c + r_o) R_E}{R_c + r_o + R_E} \cancel{\times}}{\frac{(R_c + r_o)}{1 + g_m r_o} \cdot \frac{R_E \cancel{r_{\pi}}}{R_E + r_{\pi}}} \\
 &\quad \frac{\frac{R_c + r_o}{1 + g_m r_o} + \frac{R_E r_{\pi}}{R_E + r_{\pi}}}{R_c + r_o + R_E} \\
 &= \frac{(R_c + r_o)(R_E + r_{\pi}) + R_E r_{\pi}(1 + g_m r_o)}{R_c + r_o + R_E} \\
 &= \frac{r_{\pi}(R_c + r_o + R_E) + R_E(R_c + r_o) + g_m r_o R_E r_{\pi}}{R_c + r_o + R_E} \\
 &= r_{\pi} + \frac{\tilde{R}}{R_c + r_o + R_E}
 \end{aligned}$$

where $\tilde{R} = \frac{R_E R_c + r_o R_E (1 + g_m r_{\pi})}{R_c + r_o + R_E}$

if we let $r_o \rightarrow \infty$

$$\begin{aligned}
 (6) \quad R_{in,b} &\approx r_{\pi} + R_E (1 + \tilde{R})^{\beta} \rightarrow \approx R_E (1 + \beta) \\
 &= R_E + (1 + g_m R_E) r_{\pi} \rightarrow \approx r_{\pi} (1 + g_m R_E)
 \end{aligned}$$

Note that we could have gotten
the same approx result if we had let
 $r_0 \rightarrow \infty$ in (3) and (4).
then (3) and (4) become uncoupled eqns.

From the approximate form of (3), we have

$$\frac{v_t}{R_E} \underset{\approx}{=} \frac{v_{\pi}}{R_E \parallel r_{\pi} \parallel \frac{1}{g_m}} = \bar{v}_t r_{\pi}$$

$$\frac{v_t}{\bar{v}_t} = \frac{R_E r_{\pi}}{r_{\pi} \parallel \frac{R_E}{1 + g_m R_E}} = \frac{\frac{R_E r_{\pi}}{R_E r_{\pi}}}{\frac{R_E r_{\pi}}{R_E r_{\pi}}} [R_E + (1 + g_m R_E) r_{\pi}]$$

From which $\xrightarrow{\text{follow}}$ the results of (6).

$R_{in,e}$ derivation

Impedance (small signal) looking
into the emitter of an NPN BJT

$\Rightarrow (9)$ is the exact formula

$$\Rightarrow (10) \quad R_{in,e} \approx \frac{R_B + r_\pi}{1 + \beta} \quad (\text{the usual approx formula for BJTs})$$

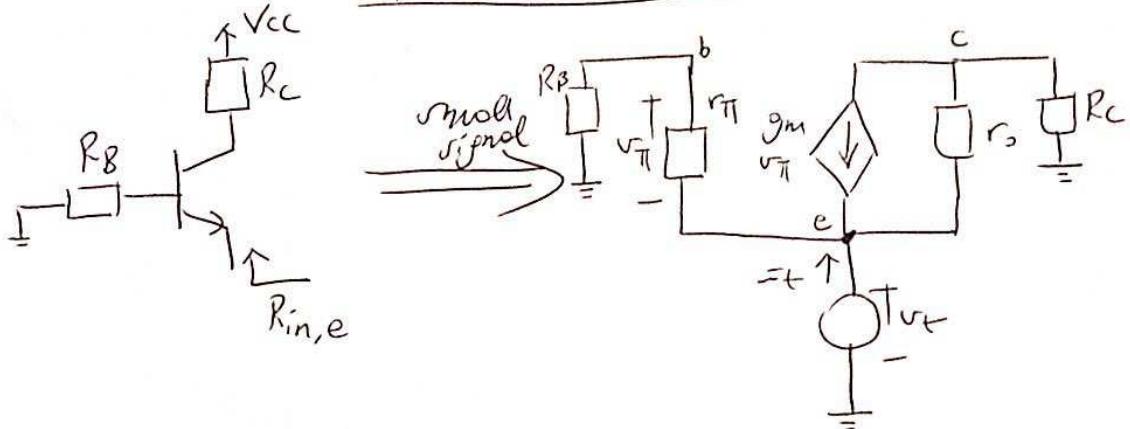
$$\Rightarrow (11) \quad R_{in,e} \approx \frac{1}{g_m} \quad (\text{a more approximate formula in the case } r_\pi \gg R_B \text{ or } R_B = 0)$$

$$\Rightarrow (13) \quad R_{in,e} \approx \frac{R_C + r_o}{1 + g_m r_o} \quad (\text{approx formula for MOS})$$

$$\Rightarrow (14) \quad R_{in,e} \approx \frac{1}{g_m} \quad (\text{a more approx formula for MOS in the case } r_o \gg R_C \text{ or } R_C = 0)$$

Worked on this on: Oct 19, 2014/Sun

$R_{in,e}$: impedance looking into the emitter
of an NPN BJT



KCL at (e)

$$\frac{v_t}{R_B + r_\pi} - \bar{i}_t - g_m v_\pi + \frac{v_t - v_c}{r_0} = 0 \quad (1)$$

$$\frac{KCL \text{ at } (c)}{g_m v_\pi + \frac{v_c - v_t}{r_0} + \frac{v_c}{R_C} = 0} \quad (2)$$

and we have

$$r_\pi = v_t \left[\frac{-r_\pi}{r_\pi + R_B} \right] \quad (3)$$

From (1) and (3)

$$\frac{v_t}{R_B + r_\pi} + \frac{g_m r_\pi v_t}{R_B + r_\pi} - \bar{i}_t + \frac{v_t}{r_0} = \frac{v_c}{r_0} \quad (4)$$

From (2) and (3)

$$\frac{v_c}{R_C // r_0} = \frac{v_t}{r_0} + \frac{g_m r_\pi v_t}{r_\pi + R_B} \quad (5)$$

From (4) and (5)

$$\begin{aligned}
 & \cancel{K_o} \left[\frac{\frac{v_t}{R_B + r_{\pi}} // r_o}{\beta} + \frac{\frac{v_t}{R_B + r_{\pi}} - \bar{v}_t}{\beta} \right] \\
 &= \left[\frac{\frac{v_t}{R_B + r_{\pi}} // r_o}{\beta} \right] \frac{\frac{R_c r_o}{R_c + r_o}}{\beta} \quad (6)
 \end{aligned}$$

From (6)

$$\frac{\frac{v_t}{R_B + r_{\pi}} // r_o}{\beta} \frac{r_o}{R_c + r_o} + \frac{\frac{v_t}{R_B + r_{\pi}} - \bar{v}_t}{\beta} = \bar{v}_t \quad (7)$$

From (7)

$$\frac{\frac{\beta v_t}{R_B + r_{\pi}} // r_o}{\beta} \frac{r_o}{R_c + r_o} + \frac{\frac{v_t}{R_B + r_{\pi}} - \bar{v}_t}{\beta} = \bar{v}_t \quad (8)$$

From (8)

$$\begin{aligned}
 R_{in,e} &= \frac{\frac{v_t}{R_B + r_{\pi}} // r_o}{\bar{v}_t} = \frac{(R_B + r_{\pi})(R_c + r_o)}{\beta r_o + R_c + r_o + R_B + r_{\pi}} \\
 &= \frac{(R_B + r_{\pi})R_c + (R_B + r_{\pi})r_o}{R_c + R_B + r_{\pi} + r_o(1+\beta)} \quad (9)
 \end{aligned}$$

From (9), with $r_o \rightarrow +\infty$

$$R_{in,e} \approx \frac{v_t}{i_t} = \frac{R_B + r_\pi}{1 + \beta} \quad \begin{array}{l} \text{(the usual} \\ \text{approximate} \\ \text{formula for} \\ \text{BJTs,} \\ i_t \text{ is justified} \\ \text{since } r_o \text{ is} \\ \text{in reality big)} \end{array} \quad (10)$$

From (10), if further we have

$$R_B = 0 \text{ or } r_\pi \gg R_B \quad \begin{array}{l} \text{(and} \\ \text{with } \beta \gg 1 \text{)} \end{array}$$

$$R_{in,e} \approx \frac{r_\pi}{1 + \beta} \approx \frac{r_\pi}{\beta} = \frac{r_\pi}{g_m r_\pi} = \frac{1}{g_m} \quad (11)$$

Reorganize (9)

$$R_{in,e} = \frac{r_\pi (R_c + r_o) + R_B (R_c + r_o)}{r_\pi (1 + g_m r_o) + r_o + R_c + R_B} \quad (12)$$

Let $r_\pi \rightarrow +\infty$ (as for MOS)

$$R_{in,e} \approx \frac{R_c + r_o}{1 + g_m r_o} \quad \begin{array}{l} \text{(the usual} \\ \text{formula for MOS)} \end{array} \quad (13)$$

Note that R_c and r_o may, in cases, be comparable, especially with an active load instead of R_c .

If $R_c = 0$ or $r_o \gg R_c$ (and with
 $g_m r_o \gg 1$,
as the intrinsic
 g_m is much
bigger than 1)

From (B)

$$R_{in,e} \underset{r_o}{\approx} \frac{r_o}{1 + g_m r_o} \underset{g_m r_o}{\approx} \frac{r_o}{g_m r_o} = \frac{1}{g_m} \quad (14)$$

→

(4)

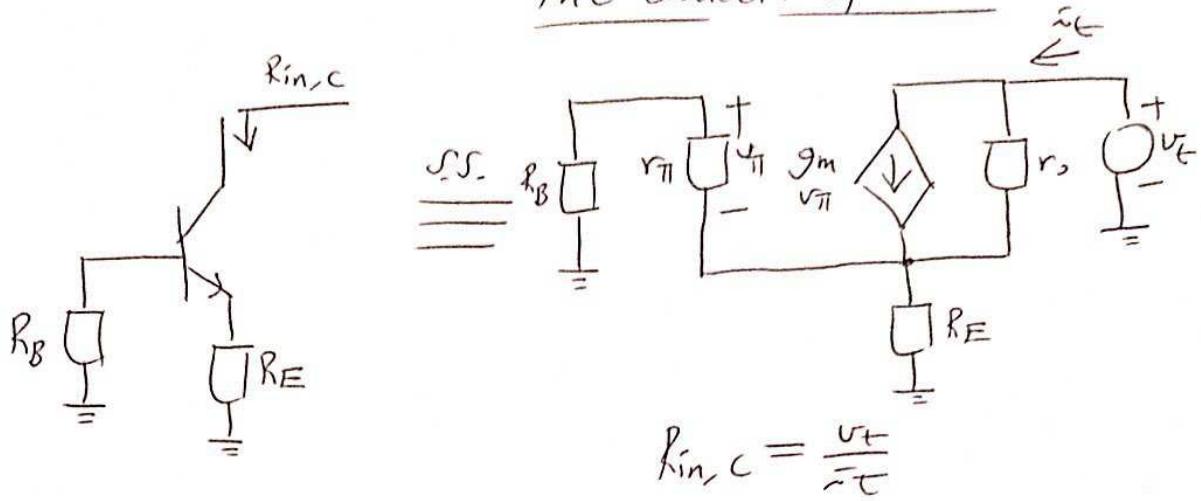
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$R_{in,C}$ calculation:
impedance into the collector
of a BJT

- $\Rightarrow (1)$ exact formula
- $\Rightarrow (2)$ when $R_B = 0$ or $R_B \ll r_\pi$
- $\Rightarrow (3)$ when $r_\pi \rightarrow +\infty$ (the case for MOS)

Worked on this on: Oct 22, 2014/Wed

Calculation of $R_{in,C}$: Impedance looking into the collector of a BJT



$$(KVL) \quad v_t = (\bar{i}_t - g_m v_{\pi}) r_o + \bar{i}_t [R_E // (r_{\pi} + R_B)]$$

$$v_{\pi} = -\bar{i}_t \frac{R_E}{R_E + r_{\pi} + R_B} \cdot r_{\pi}$$

$$(1) \quad \frac{v_t}{\bar{i}_t} = r_o \left[1 + \frac{g_m R_E r_{\pi}}{R_E + r_{\pi} + R_B} \right] + R_E // (r_{\pi} + R_B)$$

Also
can be
written as

$$= r_o + (1 + g_m r_o) \frac{R_E r_{\pi}}{R_E + r_{\pi} + R_B} + \frac{R_E R_B}{R_E + r_{\pi} + R_B}$$

$$(2) \quad = (1 + g_m r_o) \frac{R_E r_{\pi}}{R_E + r_{\pi} + R_B} + r_o + \frac{R_E R_B}{R_E + r_{\pi} + R_B}$$

From (1), if $R_B = 0$:

$$\begin{aligned}
 \frac{v_t}{i_t} &= r_o \left[1 + \frac{g_m R_E r_{\pi}}{R_E + r_{\pi}} \right] + (R_E // r_{\pi}) \\
 &= r_o \left[1 + g_m (R_E // r_{\pi}) \right] + (R_E // r_{\pi}) \\
 (3) \quad &= (R_E // r_{\pi}) \left[1 + g_m r_o \right] + r_o
 \end{aligned}$$

Note that (3) is also approximately correct

if $R_B \ll r_{\pi}$.

If $r_{\pi} \rightarrow \infty$, there will be no current on R_B ,

then $v_b \rightarrow 0$ and $v_{\pi} \rightarrow -v_e$.

In this case, we will have: (from (1) and (3))

$$\begin{aligned}
 \frac{v_t}{i_t} &\approx r_o \left[1 + g_m R_E \right] + R_E \\
 (4) \quad &= R_E \left[1 + g_m r_o \right] + r_o
 \end{aligned}$$

(4) is $R_{in,c}$ when we have a MOS transistor instead of a BJT.

6. Power Amplifiers

6.1 Miscellaneous Analyses

Analysis of the Implications of the Load Resistor Value on the performance of a Class-AB Power amplifier

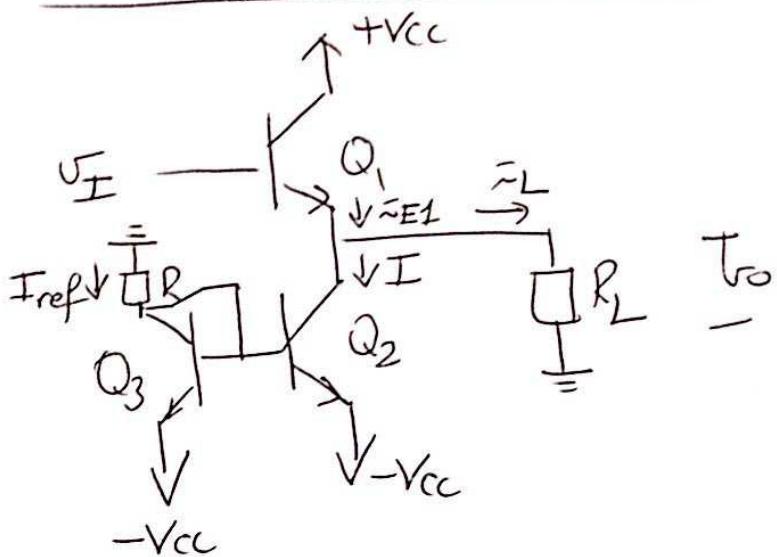
η_{\max} : max efficiency

$R_{L,\eta_{\max}}$: the load resistor
value which yields
the max efficiency

R_L : actual load
resistor value

Worked on this on: May 14, 2015 / Th

Class - A Power Amplifier



Calculations show that:
(if Q_2 and Q_3 are matched and $\beta \rightarrow +\infty$)

$$I = I_{ref} = \frac{0 - (-V_{cc} + V_{BE})}{R}$$

$$I = \frac{V_{cc} - V_{BE}}{R}$$

Also

$R_{L,\eta_{max}}$: the value for R_L which yields the maximum value for the efficiency η ($\eta_{max} = 25\%$)

$$R_{L,\eta_{max}} = \frac{V_{cc} - V_{CE,max}}{I}$$

Note that

$$R = \frac{V_{CC} - V_{BE}}{I}$$

so that

$$\frac{R_{L,\eta_{max}}}{R_{\eta}} = \frac{V_{CC} - V_{CE,sat}}{V_{CC} - V_{BE}}$$

Consider again the two possibilities

\rightarrow
 Q_1 cuts off
before Q_2 runs into
saturation

$$\bar{v}_{E1} = 0 = I + \bar{i}_L$$

$$\bar{i}_L = -I$$

$$v_{out} = -IR_L$$

(*) For this assumption to be true we must have $V_{CE,2} > V_{CE,sat}$ at the moment Q_1 cuts off

$$\begin{aligned} Q_2 \text{ runs into saturation before } Q_1 \text{ cuts off} \\ v_{out} &= -V_{CC} + V_{CE,sat} \\ \bar{i}_L &= \frac{v_{out}}{R_L} \\ &= -\frac{V_{CC} - V_{CE,sat}}{R_L} \end{aligned}$$

(**) For this assumption to be true, we must have $\bar{v}_{E1} = I + \bar{i}_L > 0$ at the moment Q_2 runs into saturation.

Analyze (*) on pg 2

$$V_{CE,2} = V_{out} - (-V_{CC})$$

$$= -IR_L - (-V_{CC})$$

$$= -R_L \frac{V_{CC} - V_{BE}}{R} + V_{CC} > V_{CE,sat}$$

$$\Rightarrow V_{CC} - V_{CE,sat} > \frac{R_L}{R} (V_{CC} - V_{BE})$$

$$\Rightarrow \frac{R_L}{R} < \frac{V_{CC} - V_{CE,sat}}{V_{CC} - V_{BE}} = \frac{R_{L,\eta_{max}}}{R}$$

$$\Rightarrow R_L < R_{L,\eta_{max}}$$

Analyze (***) on pg 2:

$$I + i_L > 0$$

$$\frac{V_{CC} - V_{BE}}{R} - \frac{V_{CC} - V_{CE,sat}}{R_L} > 0$$

$$\Rightarrow \frac{R_L}{R} > \frac{V_{CC} - V_{CE,sat}}{V_{CC} - V_{BE}} = \frac{R_{L,\eta_{max}}}{R}$$

$$\Rightarrow R_L > R_{L,\eta_{max}}$$

Therefore

$$R_L < R_{L,\eta_{\max}} \Rightarrow Q_1 \text{ cuts off before } Q_2 \text{ runs into saturation and } Q_2 \text{ consequently never runs into saturation}$$

$$R_L > R_{L,\eta_{\max}} \Rightarrow Q_2 \text{ runs into saturation before } Q_1 \text{ cuts off}$$

Then

$$R_L = R_{L,\eta_{\max}} \Rightarrow Q_1 \text{ cuts off and } Q_2 \text{ runs into saturation at the same moment}$$

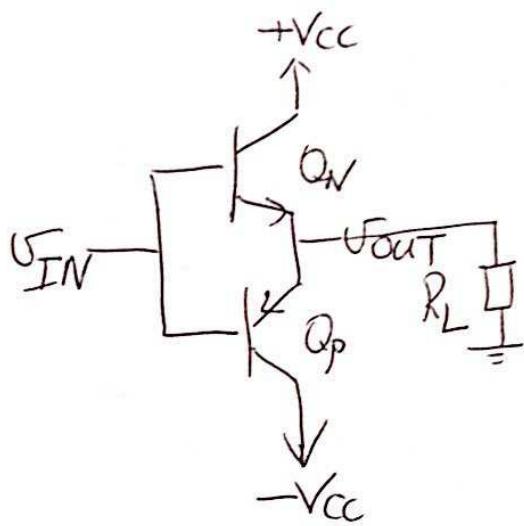
Moreover, $\eta = \eta_{\max} = 25\%$
for the class A power amplifier.

4

Derivation of the
Power Efficiency
for Class B
Power Amplifier

Worked on this on = May 18, 2015/M

Power efficiency of a Class B Amplifier



$$U_{\text{OUT},\text{max}} = +V_{\text{CC}} - V_{\text{CE},\text{sat}}$$

$$U_{\text{out},\text{min}} = -V_{\text{CC}} + V_{\text{CE},\text{sat}}$$

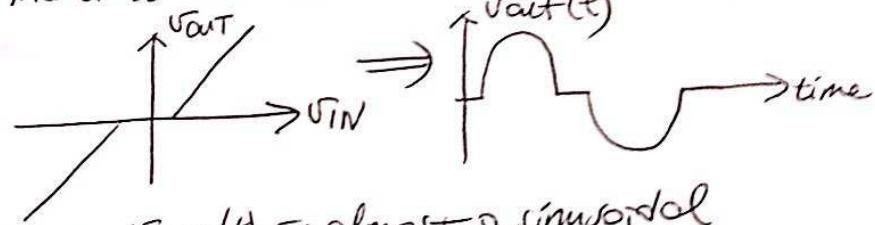
assuming $|V_{\text{CC}}| \gg |V_{\text{CE},\text{sat}}|$

$$\begin{aligned} \text{and } V_{\text{CE},\text{sat}} &= V_{\text{CE},N,\text{sat}} \\ &= V_{\text{EC},P,\text{sat}} \end{aligned}$$

max
peak-to-
peak
value
for
 U_{OUT}

$$U_{\text{out},\text{pp}} = U_{\text{out},\text{max}} - U_{\text{out},\text{min}}$$

Neglecting the crossover distortion effect



We may assume $U_{\text{out}(t)}$ is almost a sinusoidal

given by $U_{\text{out}(t)} = \hat{V}_o \sin(\omega t)$ w: angular freq.
 \hat{V}_o : magnitude

Note that $\hat{V}_o = \frac{U_{\text{out},\text{pp}}}{2} \approx V_{\text{CC}}$

power efficiency $\eta = \frac{\text{average power delivered to the load}}{\text{average power supplied by the DC voltage sources}}$

average power delivered to the load

$$P_{L,av} = \left(\frac{\hat{V}_o}{\sqrt{2}} \right)^2 \frac{1}{R_L}$$

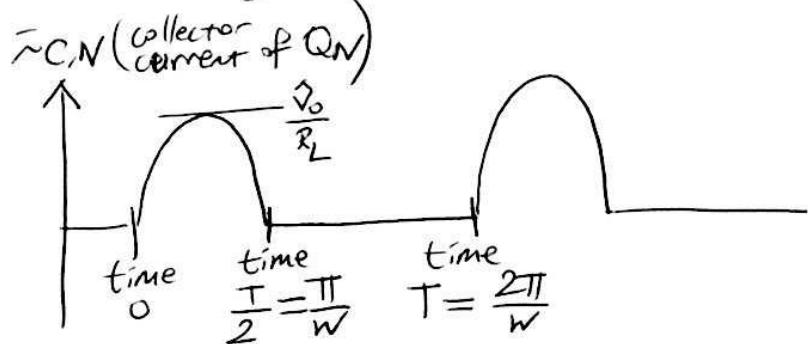
$$= \left(\frac{\text{rms value of the output voltage}}{\sqrt{2}} \right)^2 \frac{1}{R_L}$$

Average power supplied by the voltage sources

$$P_S = P_{S+} + P_{S-}$$

by the (+V_{cc}) source by the (-V_{cc}) source

Plot of the current that comes out of the (+V_{cc}) source
[neglecting the crossover distortion]



$$P_{S+} = (+V_{CC}) \frac{1}{T} \int_0^{T/2} \frac{\hat{V}_o}{R_L} \sin(\omega t) dt$$

DC voltage value instantaneous current
 average current

Note that Q_N is OFF during

$$\frac{T}{2} = \frac{\pi}{\omega} \leq t < \frac{2\pi}{\omega} = T$$

Compute

$$\begin{aligned}
 \int_0^{\frac{\pi}{\omega}} \frac{\hat{V}_o}{R_L} \sin(\omega t) dt &= \frac{\hat{V}_o}{R_L} \frac{1}{\omega} \left[-\cos(\omega t) \right] \Big|_0^{\frac{\pi}{\omega}} \\
 &= \frac{\hat{V}_o}{R_L} \frac{\pi}{2\omega} \left[-\cos(\pi) + \cos(0) \right] \\
 &= \frac{\hat{V}_o}{R_L} \frac{\pi}{2}
 \end{aligned}$$

$$P_{S+} = (+V_{CC}) \frac{1}{T} \frac{\hat{V}_o}{R_L} \frac{\pi}{2} = \frac{V_{CC} \hat{V}_o}{\pi R_L}$$

it happens that $P_{S+} = P_{S-}$

The total average supplied power is

$$P_s = P_{s+} + P_{s-} = \frac{2V_{cc}\hat{V}_o}{\pi R_L}$$

Therefore the efficiency

$$\eta = \frac{P_{L,av}}{P_s} = \frac{\hat{V}_o^2 \frac{1}{2R_L}}{\frac{2V_{cc}\hat{V}_o}{\pi R_L}} = \frac{\hat{V}_o}{V_{cc}} \frac{\pi}{4}$$

See the bottom of fig ①

max value for $\hat{V}_o \approx V_{cc}$

Then $\eta_{max} = \frac{\pi}{4} \approx 78.5\%$

7. Differential Amplifiers

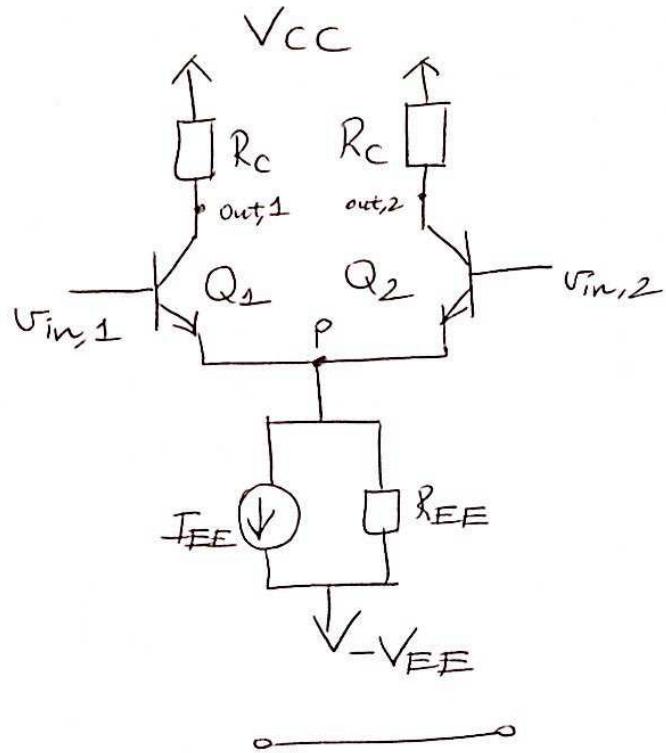
7.1 Fundamentals

Simple Differential Amplifier Analysis

- ⇒ Differential mode half circuit
(virtual ground concept in symmetrical diff. amp.)
- ⇒ Common mode half circuit
- ⇒ Gain, I/O impedances
- ⇒ CMRR

Worked on this on: NOV 5, 2014 / Wed

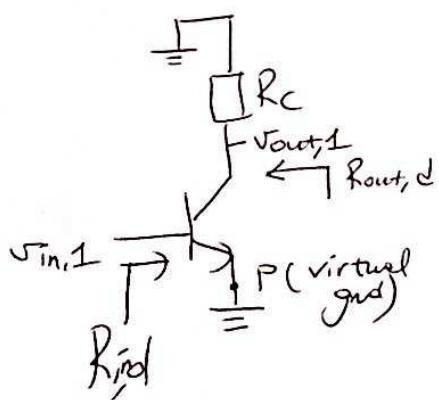
Simple Differential Amplifier Analysis



Q_1 and Q_2
are identical
(the small signal parameters are the same since DC o.p. is the same)

Differential mode half circuit

node P can be proved to be virtual ground (in small signal)



$$\frac{v_{out,1}}{v_{in,1}} = -g_{m1}(R_C // r_{o1})$$

similarly we have

$$\frac{v_{out,2}}{v_{in,2}} = -g_{m2}(R_C // r_{o2})$$

Note that $g_{m1} = g_{m2} \triangleq g_m$

$$r_{o1} = r_{o2} \triangleq r_o$$

$$r_{\pi1} = r_{\pi2} \triangleq r_\pi$$

①

Observe also that

$$\underset{\text{diff mode gain}}{\text{diff mode gain}} A_D \triangleq -g_m (R_C // r_o)$$

$$v_{out,1} = A_D v_{in,1}$$

$$v_{out,2} = A_D v_{in,2}$$

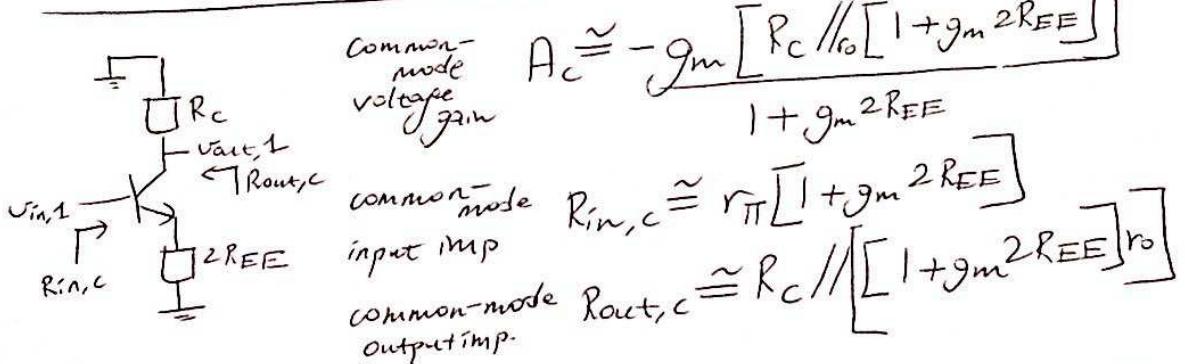
$$v_{out,1} - v_{out,2} = A_D (v_{in,1} - v_{in,2})$$

Then $\frac{v_{out,1} - v_{out,2}}{v_{in,1} - v_{in,2}} = A_D$

Diff mode input imp $\Rightarrow R_{in,d} = r_\pi$

Diff mode output imp $\Rightarrow R_{out,d} = R_C // r_o$

Common mode half circuit



Common-mode Rejection Ratio (CMRR)

$$\frac{|\mathcal{A}_d|}{|\mathcal{A}_c|} = CMRR = \left| \frac{1 - g_m (R_c // r_o)}{1 - \frac{g_m}{1 + g_m^2 R_{EE}} [R_c // [r_o [1 + g_m^2 R_{EE}]]]} \right|$$
$$= (1 + g_m^2 R_{EE}) \frac{R_c // r_o}{R_c // [r_o [1 + g_m^2 R_{EE}]]}$$

if $r_o \gg R_c \approx (1 + g_m^2 R_{EE})$

$\xrightarrow{\hspace{1cm}}$
CMRR is expected / designed to be
much greater than 1.
 $\xrightarrow{\hspace{1cm}}$

8. Frequency Response

8.1 Fundamentals

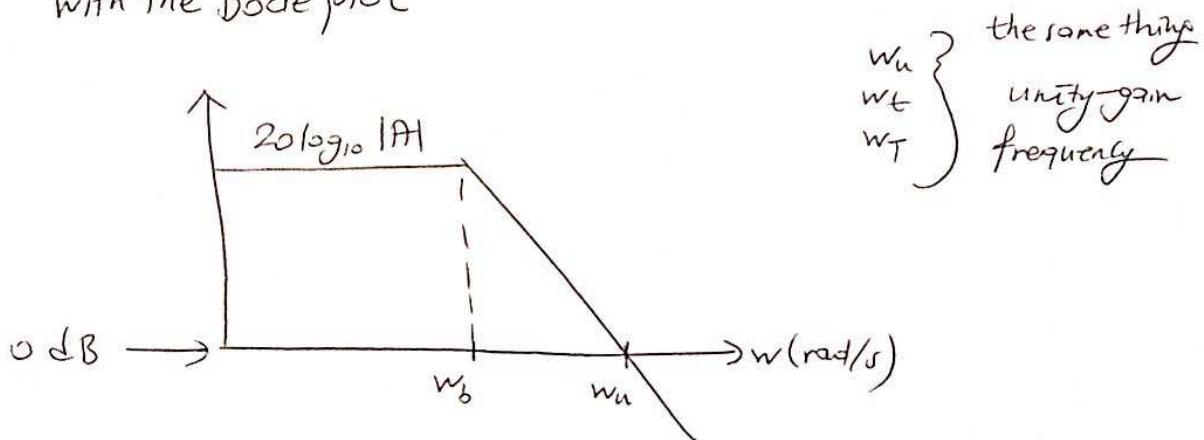
A Note on the
Gain-Bandwidth Product
of $A(jw) = \frac{A}{1 + \frac{jw}{w_b}}$

Worked on this on: Oct 23, 2014/Th

Gain-Bandwidth product (see Sedra & Smith pp 90-93)

Consider $A(j\omega) = \frac{A}{1 + \frac{j\omega}{w_b}}$

with the Bode plot



Note that $w \gg w_b$ for

$$A(j\omega) \approx \frac{A}{\frac{j\omega}{w_b}}$$

then $20 \log_{10} |A(j\omega)| = ① \text{dB}$

$$|A(j\omega)| = 1 = \left| \frac{A}{\frac{j\omega}{w_b}} \right| \quad \text{since } w_a \gg w_b$$

$$\Rightarrow A w_b = w_a \Rightarrow \begin{array}{l} \text{the unity gain freq is} \\ \text{equal to the gain-bandwidth} \\ \text{product.} \end{array}$$

↓ gain
 ↓ bandwidth
 (midband)

OCTC - Open Circuit Time Constants Method

$$\frac{1}{\left(1 + \frac{jw}{w_1}\right)\left(1 + \frac{jw}{w_2}\right)} \xrightarrow{\text{OCTC}} \frac{1}{1 + \frac{jw}{w_0}}$$

Worked on this on: Oct 27, 2015 / Tues

Open Circuit Time Constant Method (OCTC)

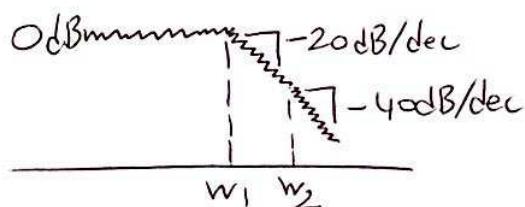
for computing the high frequency cut-off
of an amplifier circuit

OCTC is an approximation. In some cases
it is applied after Miller's approximation.

An example of a high-frequency response:

$$A(j\omega) = \frac{1}{(1 + \frac{j\omega}{\omega_1})(1 + \frac{j\omega}{\omega_2})}$$

The corresponding magnitude Bode Plot is



How do we approximate this frequency response
so that there is only one pole ω_0 to account for?

Consider $A(j\omega) = \frac{1}{1 + \frac{j\omega}{\omega_1} + \frac{j\omega}{\omega_2} + \left(\frac{j\omega}{\omega_1}\right)\left(\frac{j\omega}{\omega_2}\right)}$

Now omit this term

Then approximately

$$H(j\omega) \approx \frac{1}{1 + \frac{j\omega}{\omega_0}}$$

Note that $\frac{1}{\omega_0} = \frac{1}{\omega_1} + \frac{1}{\omega_2}$

ω_1 and ω_2 correspond to time constants τ_1 and τ_2 , respectively.

$$\omega_1 = \frac{1}{\tau_1} \quad \omega_2 = \frac{1}{\tau_2}$$

Then $\frac{1}{\omega_0} = \tau_1 + \tau_2$

$$\omega_0 = \frac{1}{\tau_1 + \tau_2}$$

Note that the time constants may be given as

$$\tau_1 = R_1 C_1 \quad \tau_2 = R_2 C_2$$

In order to make $\tau_2 = 0$ when computing τ_1 ,

we let $C_2 = 0$. This corresponds to $\frac{1}{j\omega C_2} \rightarrow +\infty$.

Effectively an open circuit.

The procedure goes like this:

→ set $C_2 = 0$, make C_2 an open circuit

compute the resistance R_1 that C_1 "sees"

compute $\tau_1 = R_1 C_1$

→ set $C_1 = 0$, make C_1 an open circuit

compute the resistance R_2 that C_2 "sees"

compute $\tau_2 = R_2 C_2$

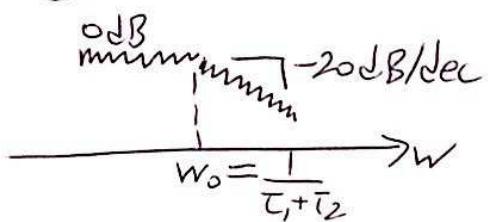
Then compute

$$\omega_0 = \frac{1}{\tau_1 + \tau_2}$$

$H(j\omega)$ can be modelled (approximated) as

$$H(j\omega) \approx \frac{1}{1 + j\frac{\omega}{\omega_0}}$$

The corresponding Bode plot is:



8.2 Intrinsic Responses of Transistors

BJT Frequency Response

$$\Rightarrow \beta(j\omega) = \frac{\bar{r}_c(j\omega)}{\bar{r}_b(j\omega)} \text{ modeling}$$

\Rightarrow Bode plot for $\beta(j\omega)$

$\Rightarrow w_B = -3dB$ freq for $\beta(j\omega)$

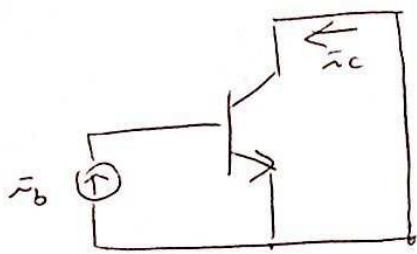
$\Rightarrow w_T$ or $f_T = \frac{w_T}{2\pi}$ is called the transition freq

$$w_T = \frac{g_m}{C_{\mu} + C_{\pi}}$$

Worked on this on: Oct 23, 2014 / Th

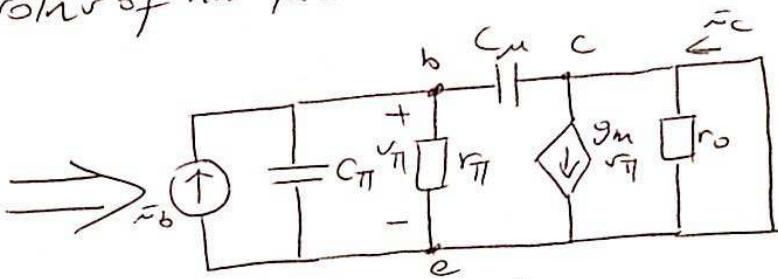
BJT Frequency Response

See Rashed pg 504-506,
excepting some useful facts to be
utilized in the soln of hw questions.



BJT s.s. model

(\bar{i}_b and \bar{i}_c are
small signal
currents,
there are no DC
sources in the
schem.)



BJTss model
replaced by the high
freq model

(Note that there is
no current through r_o ,
and C_{π} and C_{μ} come in
parallel)

$$v_{\pi} = \bar{i}_b \left[r_{\pi} \parallel \frac{1}{sC_{\pi}} \parallel \frac{1}{sC_{\mu}} \right]$$

$\underbrace{\frac{1}{s(C_{\pi} + C_{\mu})}}$ (capacitors in parallel added)

KCL at node b \Rightarrow yields the above formula

$$\underline{\text{KCL at node c}} \Rightarrow \frac{v_{\pi}}{1/sC_{\mu}} + \bar{i}_c = g_m v_{\pi}$$

$$\Rightarrow sC_{\mu} v_{\pi} + \bar{i}_c = g_m v_{\pi}$$

$$\Rightarrow \bar{i}_c = v_{\pi} (g_m - sC_{\mu})$$

calculate

$$\begin{aligned}\left(\frac{\bar{z}_b(s)}{\bar{z}_c(s)}\right)^{-1} &= \frac{\bar{z}_c(s)}{\bar{z}_b(s)} = \frac{\frac{s}{r_{\pi}}(g_m - sC_m)}{\frac{1}{r_{\pi} + \frac{1}{s(C_m + C_{\pi})}}} \\ &= \frac{(g_m - sC_m) \frac{r_{\pi}}{s(C_m + C_{\pi})}}{r_{\pi} + \frac{1}{s(C_m + C_{\pi})}} \\ &= \frac{(g_m - sC_m) r_{\pi}}{1 + s(C_m + C_{\pi}) r_{\pi}}\end{aligned}$$

Now we are given $g_m \gg sC_m$ at the frequencies of interest and we know $\beta = g_m r_{\pi}$.

Then

$$\beta(jw) = \frac{\bar{z}_c(jw)}{\bar{z}_b(jw)} \approx \frac{g_m r_{\pi}}{1 + jw(C_m + C_{\pi}) \frac{\beta}{g_m}} = \frac{\beta}{1 + \frac{jw}{\frac{g_m}{\beta(C_m + C_{\pi})}}}$$

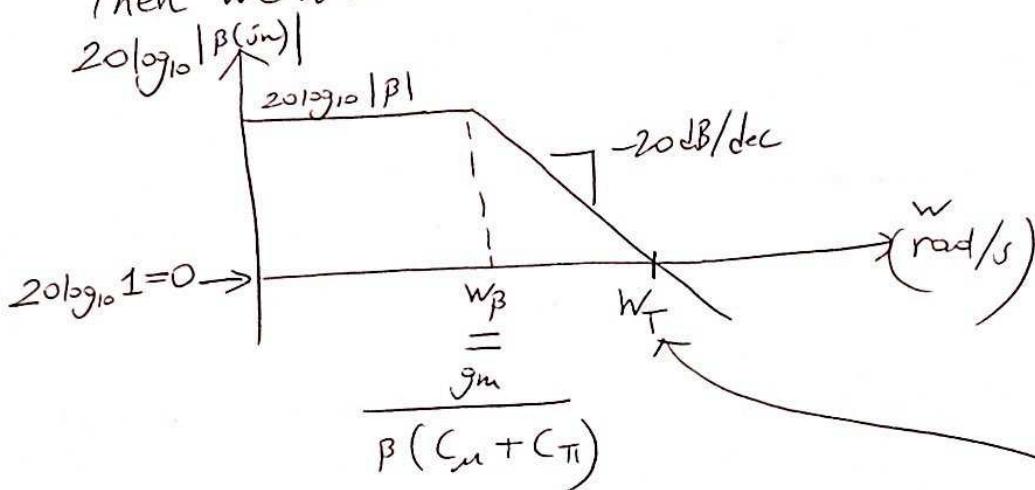
This is the frequency model for the current gain β of the BJT.

Bode plot for $\beta(j\omega)$

Note that we can write

$$\beta(j\omega) = \frac{\beta}{1 + \frac{j\omega}{w_p}} \quad \text{where } w_p \text{ is the } -3\text{dB frequency for } \beta(j\omega)$$

Then we have



Note that for $\omega \gg w_p \Rightarrow \beta(j\omega) \approx \frac{\beta}{j\omega}$

(this is an asymptotic approximation)

Now $20\log_{10} |\beta(j\omega)| = 0$
then
when $\omega = w_T$

then we need to have

$$\left| \frac{\beta}{jw_T - w_B} \right| = 1$$

$$\Rightarrow \frac{\beta w_B}{w_T} = 1 \Rightarrow w_T = \beta w_B$$

$$w_T = \beta w_B = \beta \frac{g_m}{\beta(C_m + C_{Tl})} = \frac{g_m}{C_m + C_{Tl}}$$

↓
called
transition
frequency

for $w > w_T$, $|\beta(jw)| < 1$
then the BJT is not amplifying
it anymore, it is attenuating it.

Note that you will be given either

$$w_T \text{ or } f_T = \frac{w_T}{2\pi} \text{ in the questions.}$$



9. Feedback

9.1 Benefits

Effect of feedback on low frequency response

$\Rightarrow \omega_L$ lowered \longrightarrow By the same factor
 $\Rightarrow A_m$ (midband gain) lowered \longrightarrow By the same factor

Worked on this on: Oct 27, 2014/Mon

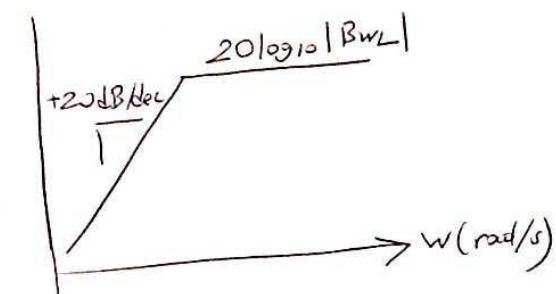
Effect of feedback on low-freq response

$$\text{T.F. model} \Rightarrow A(j\omega) = \beta \frac{j\omega}{1 + \frac{j\omega}{w_L}}$$

$$\text{for } \omega \gg w_L \Rightarrow A(j\omega) \approx \beta \frac{j\omega}{\frac{j\omega}{w_L}}$$

$$A_m = |A(j\omega)| = |\beta w_L| \quad (\text{midband gain})$$

Magnitude
Bode plot \Rightarrow



T.F. model
with feedback
(where β is the
feedback factor)

$$A_f(j\omega) = \frac{A(j\omega)}{1 + \beta A(j\omega)}$$

$$= \frac{\beta \frac{j\omega}{1 + \frac{j\omega}{w_L}}}{1 + \beta \frac{j\omega}{1 + \frac{j\omega}{w_L}}}$$

$$= \frac{\beta j\omega}{1 + \frac{j\omega}{w_L} + \frac{j\omega}{\frac{1}{\beta}}} = \frac{\beta j\omega}{1 + \frac{j\omega}{w_L} + \frac{j\omega}{\frac{1}{\beta \beta}}}$$

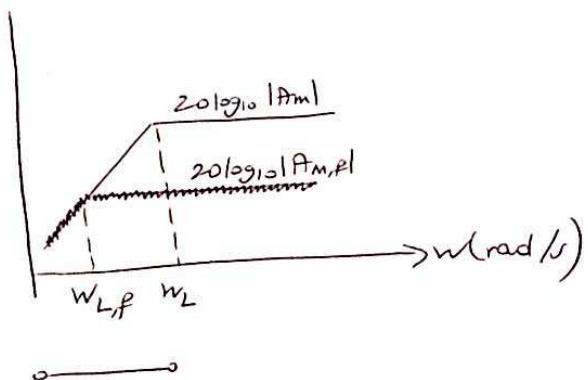
$$= \beta \frac{j\omega}{1 + \frac{j\omega}{w_L} + \frac{j\omega}{1 + \beta w_L}}$$

then $w_{L,f} = \frac{w_L}{1 + \beta B w_L}$ lowered by this factor

w_L low freq cutoff with feedback

new midband gain $\Rightarrow \frac{B w_L}{1 + \beta B w_L} = A_{m,f}$

Bode plots



Note that $w_L > 0$.
 Then the signs of B and β must be the same
 for negative feedback.

(2)

Effects of feedback on High-Frequency Cut-off

$\Rightarrow w_T \left(\cancel{G \text{BW}} \right)$ remains the same
gain-bandwidth product

\Rightarrow bandwidth increased
gain reduced

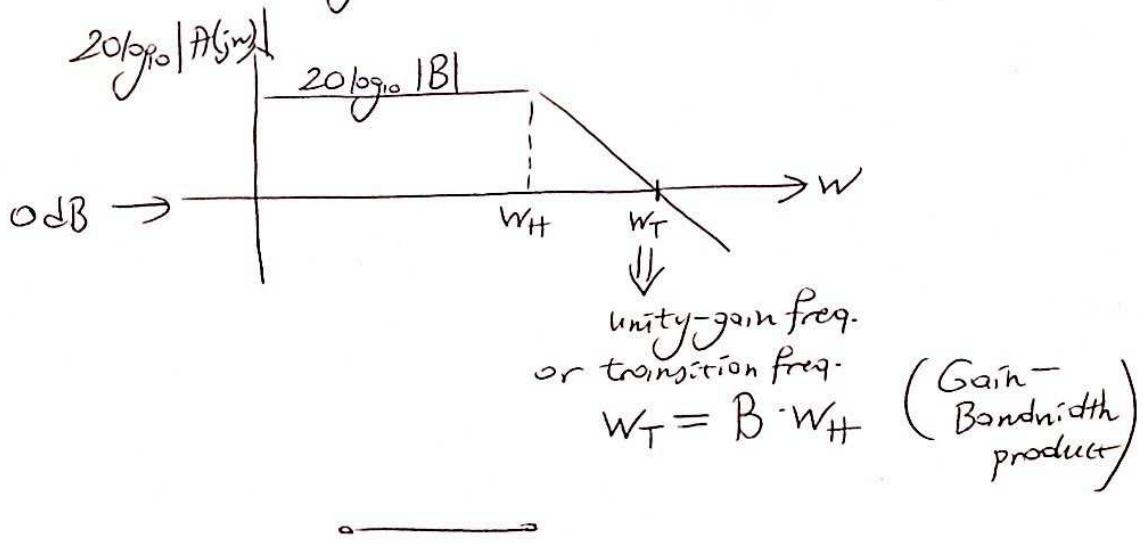
Worked on this on: Nov 5, 2014 / Wed

Effects of feedback on high-frequency cut-off

Our model is $A(j\omega) = \frac{B}{1 + \frac{j\omega}{w_H}}$

Feedback factor is β

Bode Plot (magnitude) of $A(j\omega)$ in asymptotic form



System transfer function with feedback

$$A_f(j\omega) = \frac{A(j\omega)}{1 + \beta A(j\omega)}$$

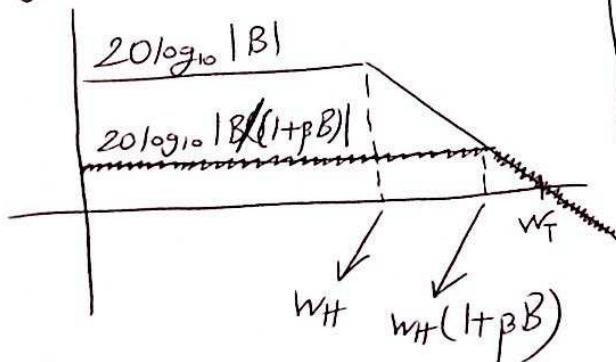
$$A_p(jw) = \frac{\frac{B}{1 + \frac{jw}{w_H}}}{1 + \beta \frac{\frac{B}{1 + \frac{jw}{w_H}}}{1 + \frac{jw}{w_H}}}$$

$$= \frac{B}{1 + \frac{jw}{w_H} + \beta B}$$

$$= \frac{\frac{B}{1 + \beta B}}{1 + \frac{jw}{w_H(1 + \beta B)}}$$

Bode plots of $A(jw)$ and $A_p(jw)$

(Magnitude-asymptotic)



$A(jw) \rightarrow A_p(jw)$
Bandwidth increased but gain reduced

w_T (GBW) for both plots is the same

$$w_T = B \cdot w_H = \frac{B}{1 + \beta B} \cdot w_H(1 + \beta B)$$

2

Scanned by CamScanner

Feedback decreases the
sensitivity with respect to
a parameter

Worked on this on: Nov 6, 2014/Thu

Given an amplifier with gain A and feedback factor β , we examine sensitivities with respect to a parameter P (β is assumed to be independent of P).

$$S_p^A = \frac{\partial A}{\partial P} \frac{P}{A}$$

$$A_f = \frac{A}{1 + \beta A}$$

How is $S_p^{A_f}$ related to S_p^A ?

$$\begin{aligned} S_p^{A_f} &= \frac{\partial A_f}{\partial P} \frac{P}{A_f} = \frac{P}{A_f} \frac{\partial}{\partial P} \left[\frac{A}{1 + \beta A} \right] \\ &= \frac{P}{A_f} \frac{\frac{\partial A}{\partial P} (1 + \beta A) - A \left[\beta \frac{\partial A}{\partial P} \right]}{(1 + \beta A)^2} \\ &= \frac{P}{A_f} \frac{\frac{\partial A}{\partial P}}{(1 + \beta A)^2} \\ &= \frac{P}{A} \frac{1 + \beta A}{(1 + \beta A)^2} \frac{\partial A}{\partial P} \\ &= \frac{S_p^A}{1 + \beta A} \end{aligned}$$

Therefore negative feedback decreased the sensitivity with respect to a parameter by the improvement factor $(1 + \beta A)$.

Effects of Feedback on Distortion

D8.20 A particular amplifier has a nonlinear transfer characteristic that can be approximated as follows:

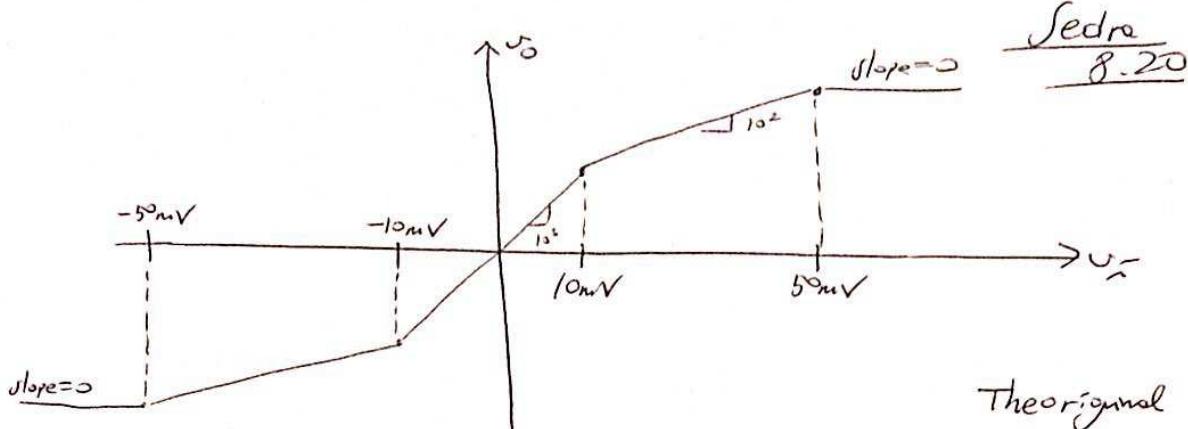
- (a) For small input signals, $|v_I| \leq 10 \text{ mV}$, $v_O/v_I = 10^3$
- (b) For intermediate input signals, $10 \text{ mV} \leq |v_I| \leq 50 \text{ mV}$, $v_O/v_I = 10^2$
- (c) For large input signals, $|v_I| \geq 50 \text{ mV}$, the output saturates

If the amplifier is connected in a negative-feedback loop, find the feedback factor β that reduces the factor-of-10 change in gain (occurring at $|v_I| = 10 \text{ mV}$) to only a 10% change. What is the transfer characteristic of the amplifier with feedback?

Notes: None.

Additional Tasks: None.

Necessary Knowledge and Skills: Gain calculations for a feedback amplifier.



The original
VTC curve
is highly
nonlinear

$$\text{Let } A_1 = 10^3$$

$$A_2 = 10^2$$

with feedback

$$\left. \begin{aligned} A_{1,f} &= \frac{A_1}{1+A_1\beta} \\ A_{2,f} &= \frac{A_2}{1+A_2\beta} \end{aligned} \right\} \text{the same feedback factor } \beta \text{ needs to be used.}$$

With feedback, we will still have

$$A_{1,f} > A_{2,f}$$

It is required that we have

$$A_{1,f} = A_{2,f} (1 + 10\%)$$

$$\frac{A_1}{1+A_1\beta} = \frac{A_2}{1+A_2\beta} (1.1) \Rightarrow \text{solve for } \beta$$

$$10(1+100\beta) = (1+1000\beta)1.1$$

Sedra
8-20
contin-

$$10 + 1000\beta = 1.1 + 1100\beta$$

$$8.9 = 100\beta$$

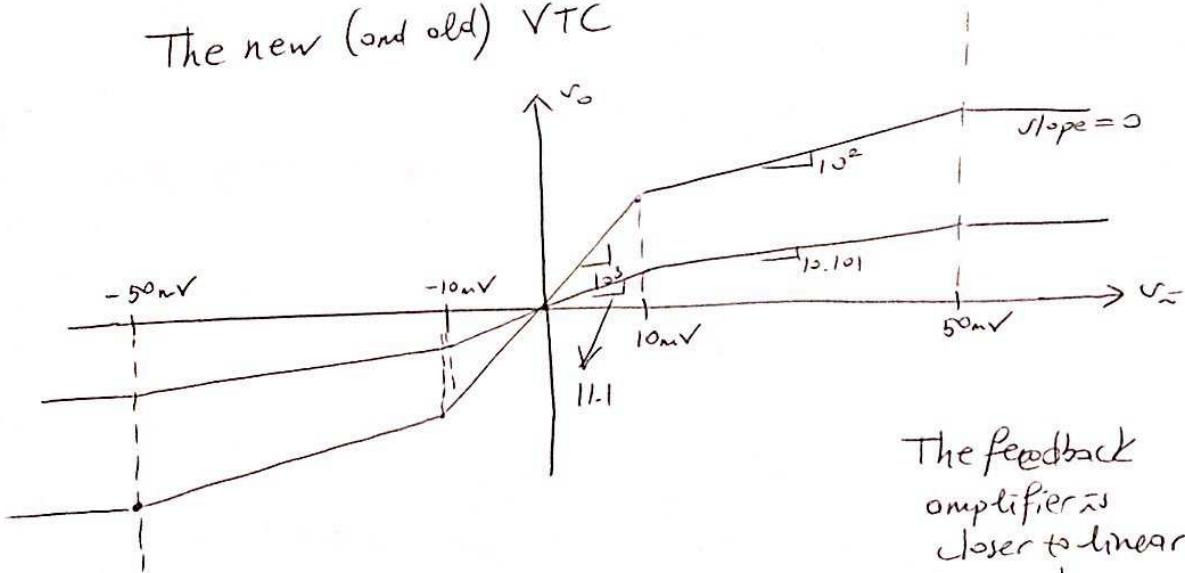
$$\beta = \frac{8.9}{100}$$

then $A_{1,f} = \frac{10^3}{1 + 10^3 \frac{8.9}{10^2}} = \frac{1000}{90} = 11.1$

10% difference

$$A_{2,f} = \frac{10^2}{1 + 10^2 \frac{8.9}{10^2}} = \frac{100}{9.9} \approx 10.101$$

The new (and old) VTC



The feedback amplifier is closer to linear (almost has a uniform gain at all ranges of input.)

Note: Such VTC graphs as above need to be interpreted as in:

→ The component of the input $\leq 10mV$ gets $A_1 (A_{1,f})$ gain

→ The component of the input in excess of $10mV$ but $\leq 50mV$ obtains $A_2 (A_{2,f})$ gain

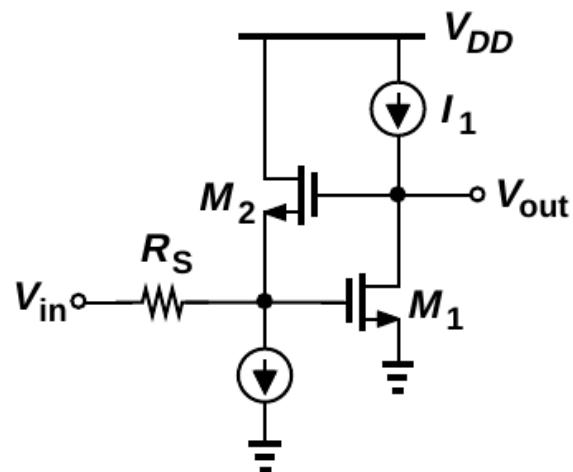
→ The component of the input in excess of $50mV$ obtains no gain.

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9.2 Polarity Determination

Polarity of Feedback

22. Determine the polarity of feedback in each of the stages illustrated in Fig. 12.87.



(a)

Notes: None.

Additional Tasks: None.

Necessary Knowledge and Skills: KCL, KVL, amplifier types, feedback polarity determination.

Razaw
12.22
a)

Assume

$$(*) \Rightarrow v_{g1} \text{ inc.}$$
$$\Rightarrow v_{out} \text{ dec.}$$

$$(**) \Rightarrow v_{s2} = v_{g1} \text{ dec. (source follower)}$$

(negative feedback)

Alternatively.

$$\Rightarrow v_{g1} \text{ inc.}$$
$$\Rightarrow \frac{v_{in} - v_{g1}}{R_s} \text{ dec.}$$
$$(*) \Rightarrow \bar{v}_{d2} \cancel{\text{inc.}} \text{ inc. (due to the}\}$$

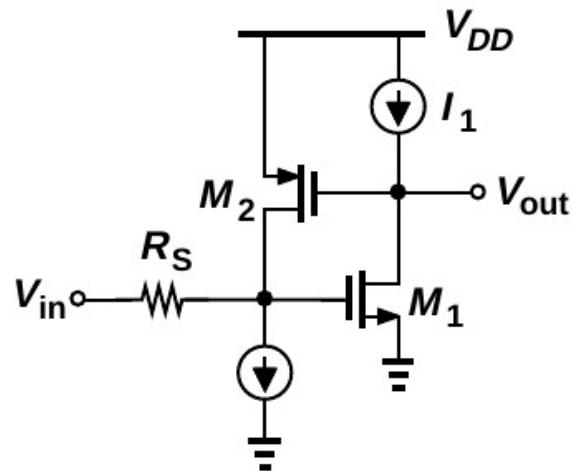
*const. cur. source
at the bottom)*

$$\Rightarrow v_{out} \text{ dec.}$$
$$\Rightarrow v_{gs2} \text{ dec.}$$
$$(**) \Rightarrow \bar{v}_{d2} \text{ dec.}$$

(again negative feedback)

Polarity of Feedback

22. Determine the polarity of feedback in each of the stages illustrated in Fig. 12.87.



(b)

Notes: None.

Additional Tasks: None.

Necessary Knowledge and Skills: KCL, KVL, amplifier types, feedback polarity determination.

Razani
12-22
b)

Assume

$$\Rightarrow v_{g1} \text{ inc}$$

$$\Rightarrow \frac{v_{in} - v_{g1}}{R_s} \text{ dec.}$$

(*) $\Rightarrow i_{d2}$ ~~inc.~~ mcr. (due to the
constant cur-
source at the bottom)

$$\Rightarrow v_{out} \text{ dec.}$$

$$\Rightarrow v_{g2} \text{ inc}$$

(**) $\Rightarrow i_{d2}$ mcr.

(positive feedback)

Positive feedback example

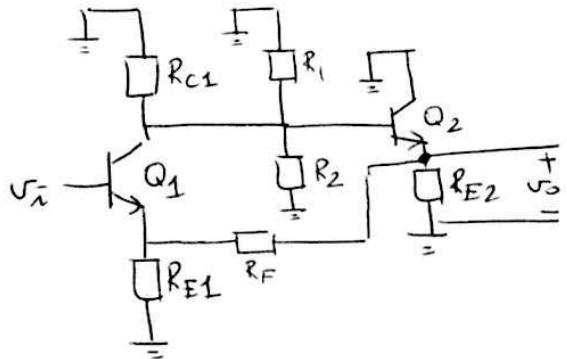
⇒ Voltage amplifier with
cascaded
CE - CC stages

⇒ Compute A_V , f_l and β

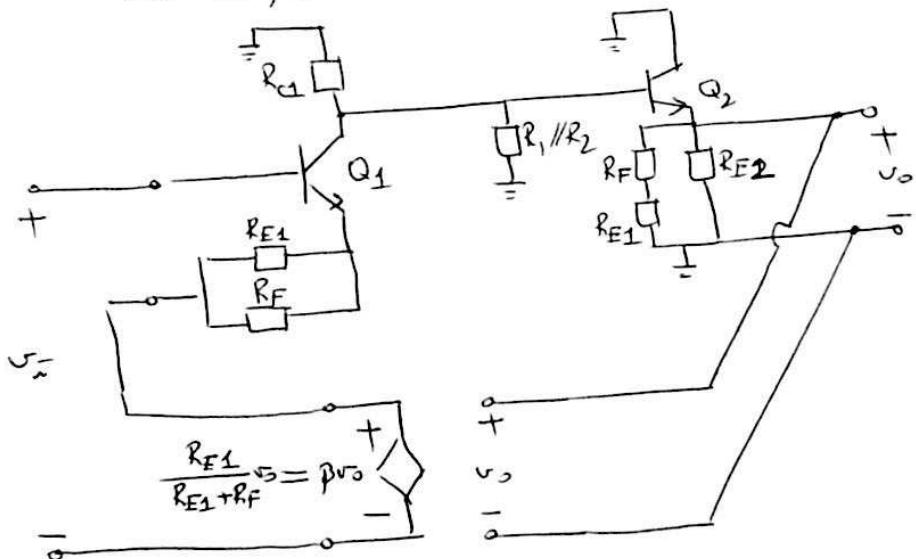
Worked on this on: Nov 18, 2014 / Tues

A positive feedback example

Consider the following CE-CC ampli:



Above is a voltage amplifier with voltage series feedback.
Analyze the V_{TH}, f_L (feedback-loaded voltage amplifier)
and the feedback network:



Kill β ($\beta=0$) above and compute the gain of the V_{TH}, f_L:

1

$$A_{VA, PE} \approx \left[\frac{-g_m 1}{1 + g_m 1 (R_{E1} // R_F)} \right] \cdot \left[R_{C1} // R_1 // R_2 // \left[R_{T2} (1 + g_m 2 (R_{E2} // (R_{E1} + R_F))) \right] \right]$$

$\cdot \frac{R_{E2} // (R_F + R_{E1})}{\frac{1}{g_m 2} + R_{E2} // (R_F + R_{E1})}$

Note that $A_{VA, PE} < 0$

but $\beta = \frac{R_{E1}}{R_{E1} + R_F} > 0$

Therefore $A_{VA, PE} \beta < 0$

so there is positive feedback in this circuit.

Instead consider constructing a
CE-CE-CC stage of VA.



9.3 Types

Voltage Series Feedback

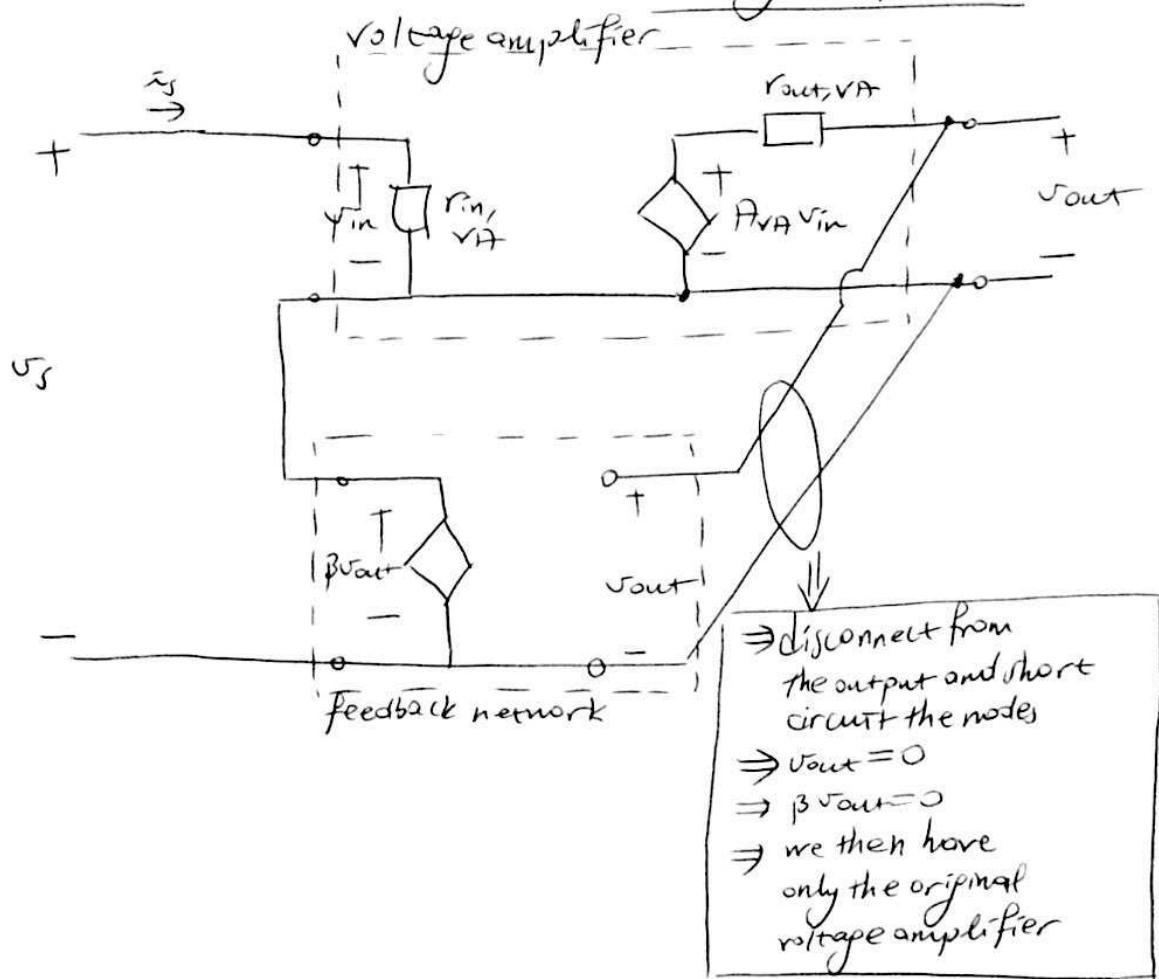
Voltage Amplifiers
on

- ⇒ Gain
- ⇒ Input impedance
- ⇒ Output impedance

Calculations. Comments on the effect
of feedback.

Worked on this on: Oct 20, 2014/Mon

Voltage series feedback on voltage amplifiers



Compute the gain A_f of the feedback amplifier

$$\left. \begin{array}{l} v_{in} = v_s - \beta v_{out} \\ A_{VA} v_{in} = v_{out} \end{array} \right\} \quad \begin{aligned} v_{in} + A_{VA} \beta v_{in} &= v_s \\ \frac{v_{in}}{v_s} &= \frac{1}{1 + A_{VA} \beta} \end{aligned}$$

$$\Rightarrow \frac{v_{out}}{v_s} = \frac{A_{VA}}{1 + A_{VA} \beta}$$

Compute the input impedance of the feedback amplifier

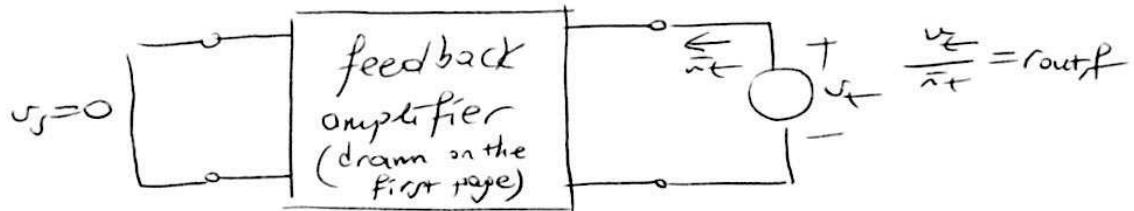
$$r_{in,f} = \frac{v_s}{\bar{z}_s} = ? \quad \frac{v_s - \beta v_{out}}{r_{in,VA}} = \bar{z}_s$$

(output is open
circuit)

$$\frac{v_s - \beta \frac{A_{VA}}{1 + A_{VA}\beta} v_s}{r_{in,VA}} = \bar{z}_s$$

$$\frac{v_s}{\bar{z}_s} = r_{in,VA} (1 + A_{VA}\beta)$$

Compute the output impedance of the feedback amplifier



$$\frac{v_t - A_{VA} v_{in}}{r_{out,VA}} = \bar{z}_t$$

$$v_{out} = v_t$$

$$-v_{in} = \beta v_{out} = \beta v_t$$

$$v_{in} = -\beta v_t$$

$$\frac{v_t - A_{VA}(-\beta v_t)}{r_{out,VA}} = \bar{z}_t \Rightarrow r_{out,f} = \frac{v_t}{\bar{z}_t} = \frac{r_{out,VA}}{(1 + A_{VA}\beta)}$$

	<u>without feedback</u>	<u>with feedback</u>	<u>change</u>
<u>Gain</u>	A_{VTF}	$\frac{A_{VTF}}{1 + A_{VTF}\beta}$	decreased
<u>Input Impedance</u>	$r_{in,VTF}$	$r_{in,VTF}(1 + A_{VTF}\beta)$	increased
<u>Output Impedance</u>	$r_{out,VTF}$	$\frac{r_{out,VTF}}{1 + A_{VTF}\beta}$	decreased

All by the
 same factor.
 $(1 + A_{VTF}\beta)$

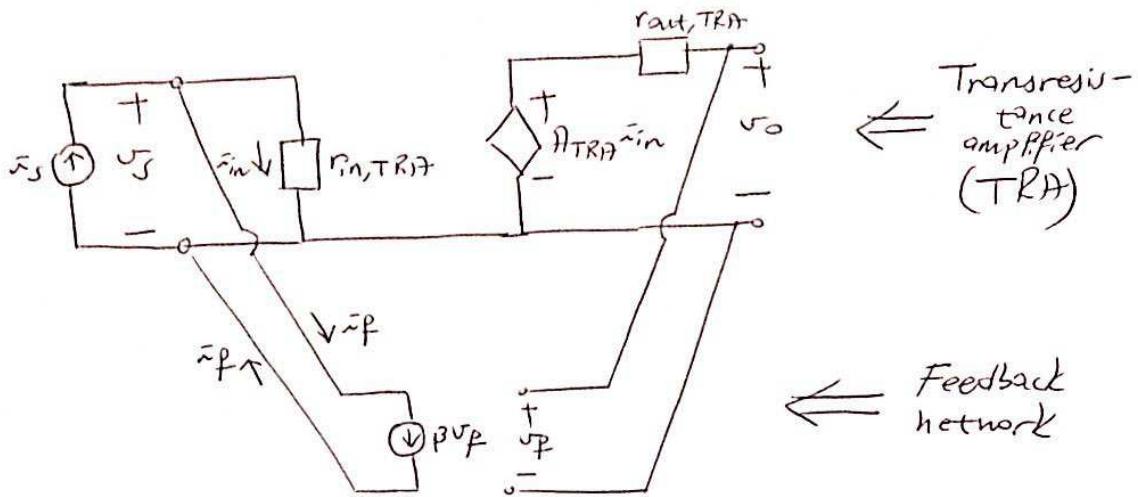
⇒ The voltage amplifier with feedback has properties that are closer to those of an ideal voltage ampl. (high input impedance, low output impedance)

Voltage-Shunt
Feedback
on a
Trans-Resistance amplifier

⇒ Gain, input/output impedance calculations

Worked on this on: Oct 27, 2014 / Mon

Voltage shunt feedback on a transresistance amplifier



equations

$$\tilde{v}_s = \tilde{v}_f + \tilde{v}_{in}$$

$$v_f = v_o$$

$$A_{TRA} \tilde{v}_{in} = v_o$$

$$\tilde{v}_f = \beta v_o$$

Gain of the feedback amplifier (A_f)

$$A_f = ? \quad A_f = \frac{v_o}{\tilde{v}_s}$$

$$\tilde{v}_s = \beta v_o + \tilde{v}_{in}$$

$$A_{TRA} \tilde{v}_{in}$$

$$\tilde{v}_s = \frac{\tilde{v}_{in}}{1 + \beta A_{TRA}}$$

$$A_f = \frac{v_o}{\tilde{v}_s} = \frac{A_{TRA}}{1 + \beta A_{TRA}}$$

input resistance of the feedback amplifier ($r_{in,f}$)

$$r_{in,f} = \frac{v_s}{\tilde{v}_s}$$

$$v_s = \tilde{v}_{in} r_{in,TRA}$$

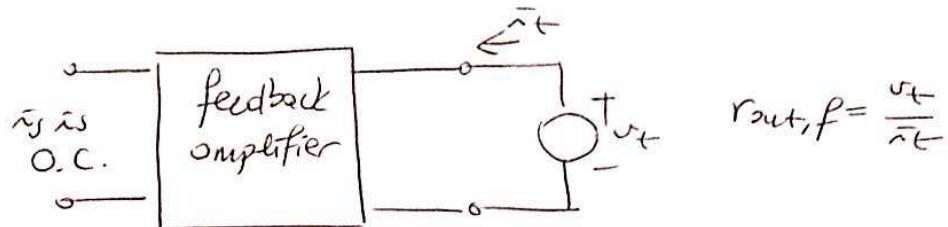
$$\tilde{v}_{in} = \frac{\tilde{v}_s}{1 + \beta A_{TRA}}$$

$$r_{in,f} = \frac{v_s}{\tilde{v}_s} = \frac{r_{in,TRA}}{1 + \beta A_{TRA}}$$

1

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$r_{out,f}$ (Output resistance of the feedback amplifier)



$$\bar{v}_f + \bar{v}_m = 0$$

$$\bar{v}_f = \beta v_o \Rightarrow \bar{v}_m = -\beta v_o$$

$$\frac{v_o - A_{TRA} \bar{v}_m}{r_{out,TRA}} = \bar{v}_o \Rightarrow r_{out,f} = \frac{v_o}{\bar{v}_o} = \frac{r_{out,TRA}}{1 + \beta A_{TRA}}$$

	TRA	Feedback amplifier	
Gain	A_{TRA}	$\frac{A_{TRA}}{1 + \beta A_{TRA}}$	decreased
r_{in}	$r_{in,TRA}$	$\frac{r_{in,TRA}}{1 + \beta A_{TRA}}$	decreased
r_{out}	$r_{out,TRA}$	$\frac{r_{out,TRA}}{1 + \beta A_{TRA}}$	decreased

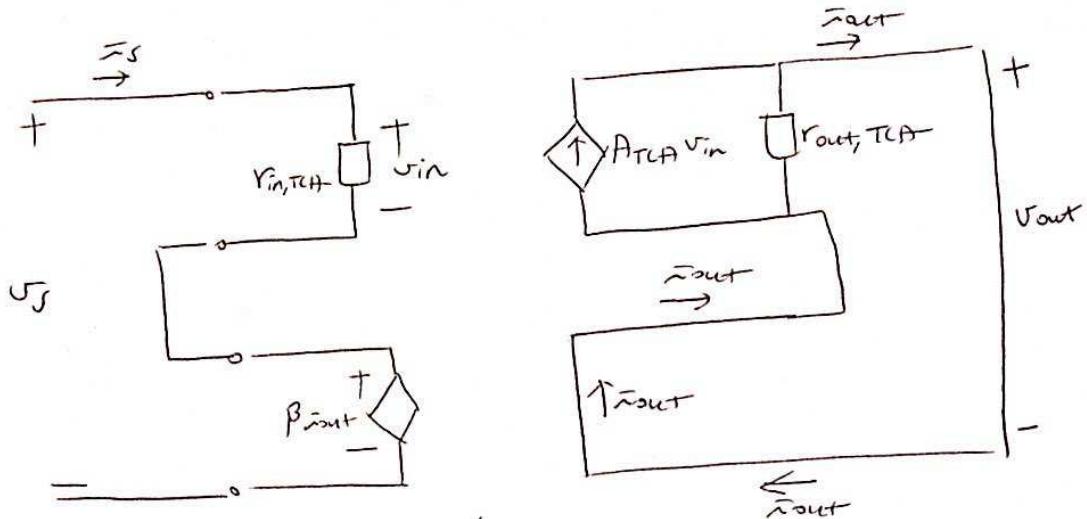
A transresistance amplifier should have low input/output impedances. Feedback decreases both input and output impedances.

Current-Series Feedback
on a Trans-Conductance
Amplifier

⇒ Gain, I/O impedances
computation

Worked on this on: Nov 5, 2014/Wed

Current-Series Feedback on TCA



Note: β in ohms

A_{TCA} is Siemens

Equations

$$\begin{aligned} v_s &= v_{in} + \beta \bar{I}_{out} \\ A_{TCA} v_{in} &= \bar{I}_{out} \end{aligned} \quad \left. \begin{aligned} v_s &= v_{in} (1 + \beta \frac{\bar{I}_{out}}{A_{TCA}}) \\ v_{in} &= \frac{\bar{I}_{out}}{A_{TCA}} \end{aligned} \right.$$



$$v_s = \frac{\bar{I}_{out}}{A_{TCA}} (1 + \beta \frac{\bar{I}_{out}}{A_{TCA}})$$

$$\Rightarrow A_f = \frac{\bar{I}_{out}}{v_s} = \frac{A_{TCA}}{1 + \beta A_{TCA}}$$

(gain of
the
feedback
ampli)

$$r_{in,f} = \frac{v_s}{\bar{I}_{out}} = ?$$

$$v_{in} = \bar{I}_{out} \cdot r_{in,TCA} \Rightarrow v_s = \bar{I}_{out} \cdot r_{in,TCA} (1 + \beta \frac{\bar{I}_{out}}{A_{TCA}})$$

$$r_{in,f} = \frac{v_s}{\bar{I}_{out}} = r_{in,TCA} (1 + \beta \frac{\bar{I}_{out}}{A_{TCA}})$$

$r_{out,f}$ Computation

$r_{out,f} = -\frac{v_{out}}{i_{out}}$ when $v_f = 0$
and v_{act} is connected at the output port.

$$\left(v_{out} - A_{TCA} v_{in} \right) r_{out,TCA} = -v_{out} \quad \left. \begin{array}{l} \\ \end{array} \right\} i_{out} [1 + \beta f] r_{out,TCA} = -v_{out}$$

$$v_{in} = -\beta v_{out}$$

$$r_{out,f} = -\frac{v_{out}}{i_{out}} = \frac{r_{out,TCA}}{(1 + \beta f)}$$

	Without feedback	With feedback
Gain	A_{TCA}	$\frac{A_{TCA}}{1 + \beta A_{TCA}}$
Input impedance	$r_{in,TCA}$	$r_{in,TCA}(1 + \beta A_{TCA})$
Output impedance	$r_{out,TCA}$	$r_{out,TCA}(1 + \beta A_{TCA})$

Gain decreased by the improvement factor but
 impedances multiplied by the same factor.
 We have obtained a better GCA.

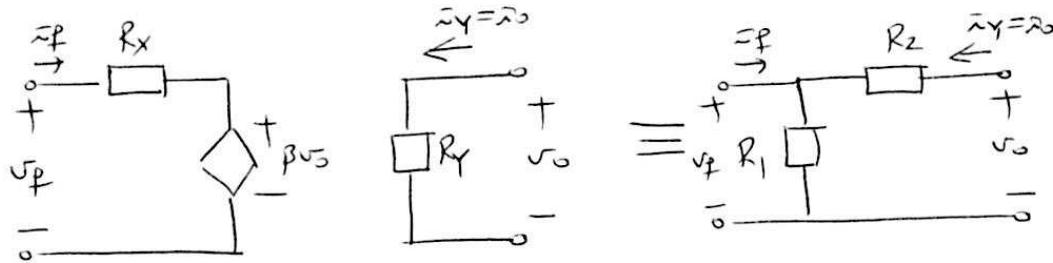
9.4 Modeling of Feedback Networks

Modelling a voltage series
feedback network

used in voltage amplifiers

Worked on this on: Nov 19, 2014/Wed

Modeling a voltage series feedback network



Model

$$\beta = \frac{v_o}{v_f} \Big|_{\bar{v}_f=0} = \frac{R_y}{R_x + R_y}$$

$$R_y = \frac{v_o}{\bar{v}_f} \Big|_{\bar{v}_f=0} = R_1 + R_2$$

$$R_x = \frac{v_f}{v_o} \Big|_{v_o=0} = R_1 // R_2$$

Example

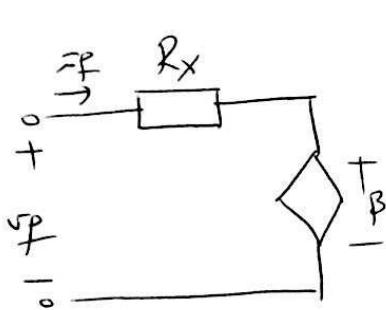
$$\frac{R_1}{R_1 + R_2}$$



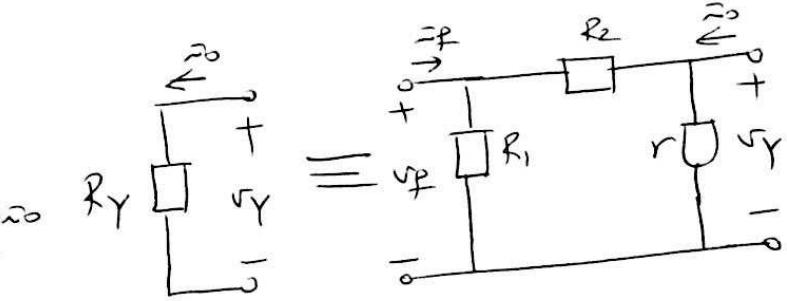
Modelling a current-series
feedback network
(used in transconductance
(TCA) amplifiers)

Worked on this on: Nov 19, 2014/Wed

Modeling a current series feedback network



Model



Example

$$\beta = \frac{v_f}{v_o} \Big|_{i_f=0}$$

$$= \frac{r}{r+R_1+R_2} \cdot R_1$$

$$R_X = \frac{v_f}{i_f} \Big|_{v_o=0}$$

$$= R_1 // (R_2 + r)$$

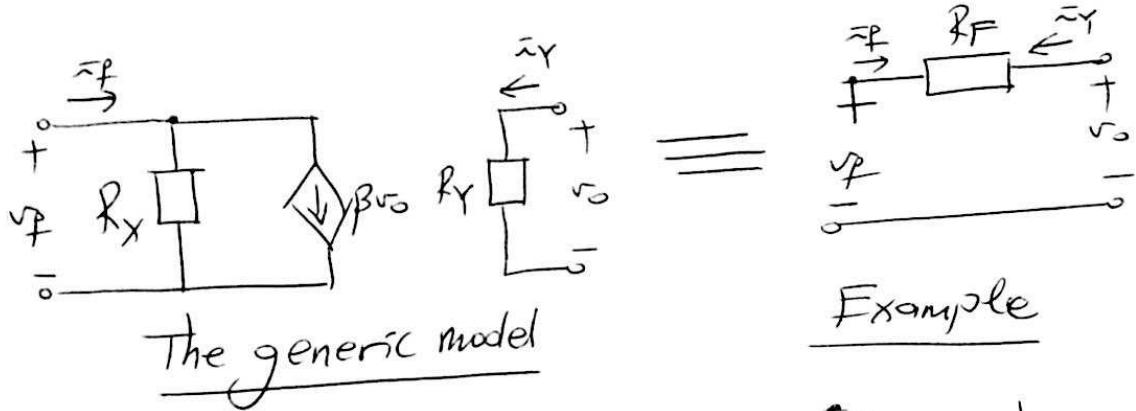
$$R_Y = \frac{v_Y}{v_o} \Big|_{i_f=0}$$

$$= r // (R_1 + R_2)$$

Voltage-Shunt Feedback Network Modelling

Worked on this on: Nov 19, 2014 / Wed

Analysis of the feedback networks
used in voltage-shunt feedback circuits



$$\beta = \frac{v_f}{v_o} \Big|_{v_f=0} = -\frac{\bar{A}_Y}{v_o} = -\frac{1}{R_2}$$

$$R_Y = \frac{v_o}{\bar{A}_Y} \Big|_{v_f=0} = R_2$$

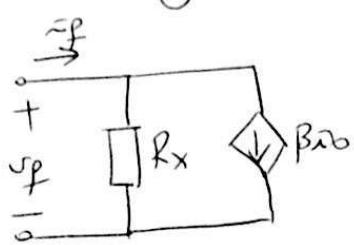
$$R_X = \frac{v_f}{\bar{A}_Y} \Big|_{v_o=0} = R_2$$

Modeling a current
hence
feedback network

used in current amplifiers

Worked on this on: Nov 19, 2014 / wed

Modeling a current-shunt feedback network

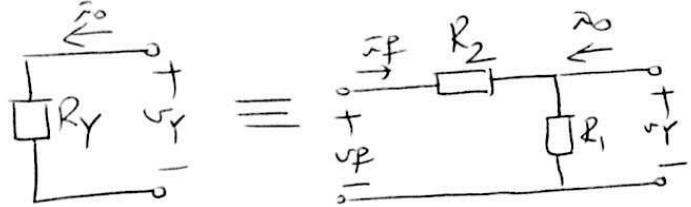


Model

$$R_x = \frac{v_p}{i_f} \Big|_{z_0=0}$$

$$R_Y = \frac{v_Y}{z_0} \Big|_{v_p=0}$$

$$\beta = \frac{i_f}{z_0} \Big|_{v_p=0}$$



Example

$$= R_2 + R_1$$

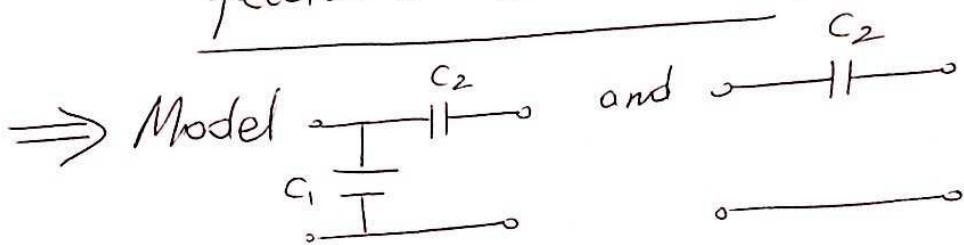
$$= R_2 // R_1$$

$$= - \frac{R_1}{R_1 + R_2}$$

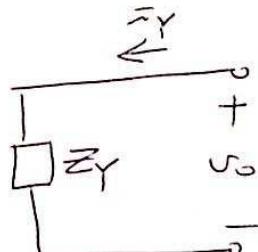
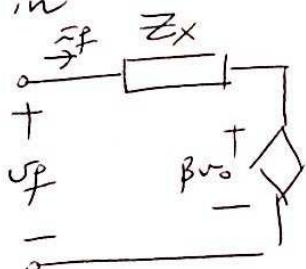
9.5 Auxiliaries

Analysis of capacitor feedback networks

typically used in
voltage-series (or series-shunt)
feedback ~~circuits~~ circuits



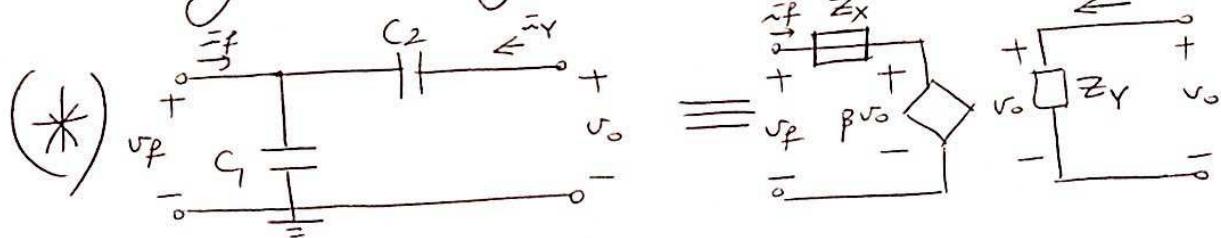
as in



Worked on this on: Nov 10, 2014 / Mon

Analysis of capacitor feedback networks

⇒ Analyze the following network (in feedback)



$$\beta = \frac{v_f}{v_o} \Big|_{\tilde{v}_f=0} = \frac{\frac{1}{C_1 s}}{\frac{1}{C_1 s} + \frac{1}{C_2 s}}$$

$$= \frac{\frac{1}{C_1}}{\frac{1}{C_1} + \frac{1}{C_2}} = \frac{\frac{1}{C_1}}{\frac{C_1 + C_2}{C_1 C_2}}$$

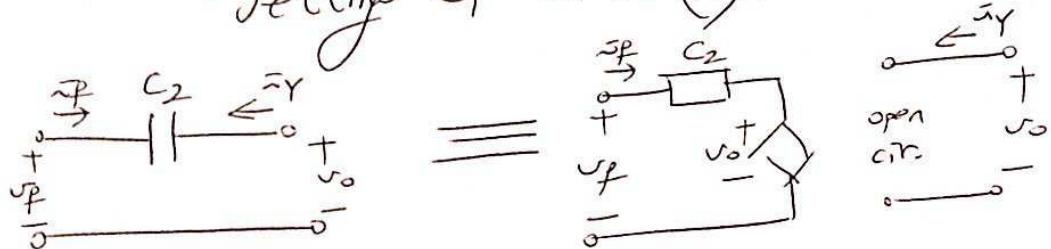
$$= \frac{C_2}{C_1 + C_2}$$

$$Z_Y = \frac{v_o}{\tilde{v}_Y} \Big|_{\tilde{v}_f=0} = \frac{1}{C_1 s} + \frac{1}{C_2 s} = \frac{1}{s} \frac{1}{\left(\frac{C_1 C_2}{C_1 + C_2}\right)}$$

$$Z_X = \frac{v_f}{\tilde{v}_f} \Big|_{v_o=0} = \frac{\frac{1}{C_1 s} \frac{1}{C_2 s}}{\frac{1}{C_1 s} + \frac{1}{C_2 s}} = \frac{1}{s} \frac{\frac{1}{C_1 C_2}}{\frac{C_1 + C_2}{C_1 C_2}}$$

$$= \frac{1}{s} \frac{1}{(C_1 + C_2)}$$

The following network is obtained by
setting $C_1 = 0$ in (*):



$$\beta = \frac{C_2}{C_1 + C_2} \Big|_{C_1=0} = 1$$

$$Z_Y = \frac{1}{s} \frac{1}{\frac{C_1 C_2}{C_1 + C_2}} \Big|_{C_1=0} = +\infty$$

$$Z_X = \frac{1}{s(C_1 + C_2)} \Big|_{C_1=0} = \frac{1}{sC_2}$$



9.6 Stability

Phase Margin as a Measure of Stability in Feedback Systems

8.73 An amplifier has a dc gain of 10^5 and poles at 10^5 Hz, 3.16×10^5 Hz, and 10^6 Hz. Find the value of β , and the corresponding closed-loop gain, for which a phase margin of 45° is obtained.

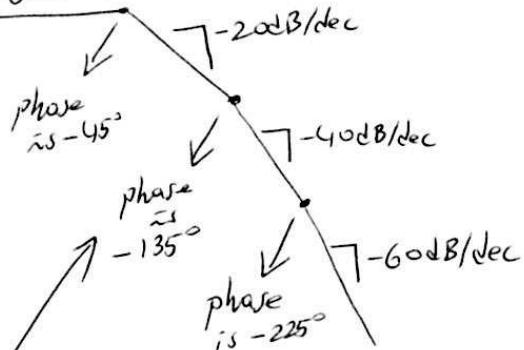
Notes: None.

Additional Tasks: None.

Necessary Knowledge and Skills: Open-loop and closed loop gains in feedback, feedback factor, stability, phase margin as a stability measure, transfer functions, magnitude and phase Bode plots, phases at the poles of a transfer function when the poles are far apart.

Magnitude Bode plot for $A(jf)$

$$20 \log_{10} |10^5| = 100 \text{ dB}$$



Sedra
8.73

The asymptotic Bode plot and the phase values are accurate if the poles of $A(jf)$ are far apart.

→ β is given to be indep of frequency, positive and real.

→ Note that

$$20 \log_{10} |A(jf)\beta| = 20 \log_{10} |A(jf)| - 20 \log_{10} \left| \frac{1}{\beta} \right|$$

$$\angle A(jf)\beta = \angle A(jf) \quad \text{since } \beta \text{ is real and positive.}$$

→ Phase margin to be 45°

$$\text{then } |A(jf_u)\beta| = 1 \quad \text{and} \quad \begin{aligned} \angle A(jf_u)\beta &= (180^\circ - 45^\circ) \\ &= -135^\circ \end{aligned}$$

since β is real and positive, the unity gain frequency f_u (for $A(jf)\beta$) needs to be at the second pole of $A(jf)$. See here-

→ Note that

$$A(jf) = \frac{10^5}{\left(1 + \frac{jf}{10^5}\right) \left(1 + \frac{jf}{3.16 \times 10^5}\right) \left(1 + \frac{jf}{10^6}\right)}$$

Now solve for β in

Sedra
8.73
contin-

$$|A(j\omega)\beta| = \left| \frac{10^5 \beta}{\left(1 + \frac{j\omega_u}{10^5}\right) \left(1 + \frac{j\omega_u}{3.16 \times 10^5}\right) \left(1 + \frac{j\omega_u}{10^6}\right)} \right| = 1$$

where $\omega_u = 3.16 \times 10^5 \text{ Hz}$ (the second pole)

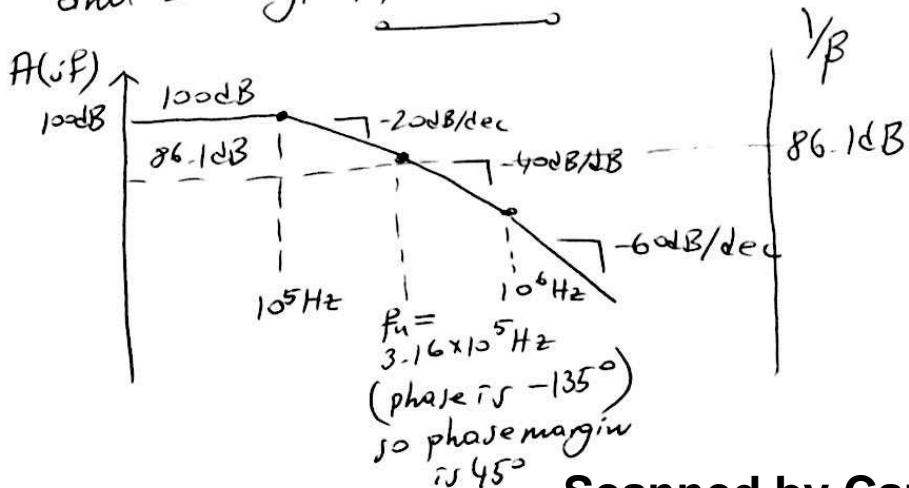
then

$$\begin{aligned} 1 &= \frac{10^5 \beta}{\sqrt{1 + \left(\frac{3.16 \times 10^5}{10^5}\right)^2} \sqrt{1 + \left(\frac{3.16 \times 10^5}{3.16 \times 10^5}\right)^2} \sqrt{1 + \left(\frac{3.16 \times 10^5}{10^6}\right)^2}} \\ &= \frac{10^5 \beta}{\sqrt{1 + (3.16)^2} \sqrt{2} \sqrt{1 + (0.316)^2}} \end{aligned}$$

$$\beta = 4.9 \times 10^{-5}$$

$$\text{then } \frac{1}{\beta} = 2.03 \times 10^4$$

$$\text{and } 20 \log_{10} \left| \frac{1}{\beta} \right| = 86.1 \text{ dB}$$



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midband gain with feedback

Sedra
8-73
contin -

$$A_f = \frac{10^5}{1 + 10^5(4.9 \times 10^{-5})} \approx 16.9 \times 10^3$$

$$= \frac{A}{1 + A\beta}$$

—————

Stability Analysis through the Phase Margin

10.56 If the phase margin of an amplifier is $PM = 40^\circ$ and the magnitude of the open-loop gain is $|A(j\omega)| = 50$, find the magnitude of the closed-loop gain $|A_f(j\omega)|$.

Notes: Assume a two-pole system for modeling the high-frequency response of the open-loop amplifier. Note that the ratio between these two poles is not given.

Additional Tasks: Observe and conclude that if the ratio between the two poles of the system had been given, there would be no missing information and you could have solved for the feedback factor.

Necessary Knowledge and Skills: Transfer function in terms of poles and midband gain, magnitude and phase computation through the transfer function, open-loop transfer function assembly, loop gain calculation for analyzing stability, phase margin calculation through the loop gain, Bode plot illustrations of the phase margin.

\Rightarrow assume that the open-loop amplifier is a two-pole system and β is real.

Pushed
10.56

Then

$$A(j\omega) = \frac{A \xrightarrow{\text{given as } 50}}{(1 + \frac{j\omega}{\omega_1})(1 + \frac{j\omega}{\omega_2})}$$

we are assuming also that $\omega_1 \ll \omega_2$.

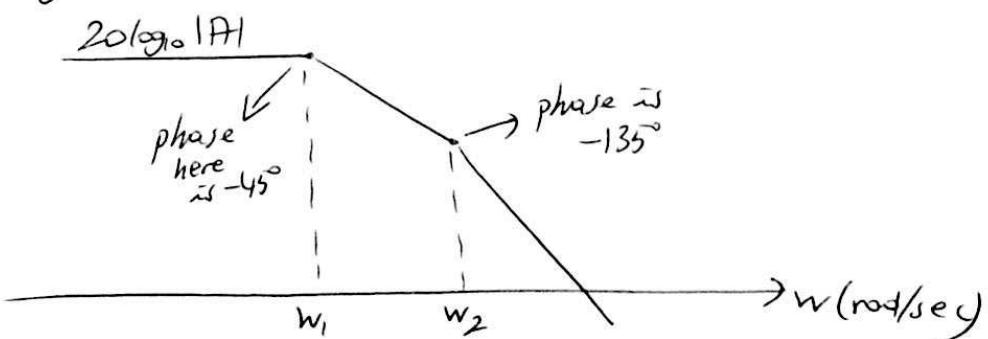
—————

\Rightarrow Note that

$$20 \log_{10} |A(j\omega)\beta| = 20 \log_{10} |A(j\omega)| - 20 \log_{10} \left| \frac{1}{\beta} \right|$$

\Rightarrow Bode plot of $A(j\omega)$ vs:

(asymptotic magnitude)



Note here that

$$A(j\omega_1) \approx \frac{A}{(1 + \frac{j\omega_1}{\omega_1})(1)} \Rightarrow \text{phase is } -45^\circ$$

$$A(j\omega_2) \approx \frac{A}{(\frac{j\omega_2}{\omega_1})(1 + \frac{j\omega_2}{\omega_2})} \Rightarrow \text{phase is } -135^\circ$$

We would like to set the phase margin to 45° ($PM = 45^\circ$) by designing β .

Part b
1056
contin-

$$\text{Phase margin is } \Rightarrow PM = 180^\circ - \left| \underbrace{\angle A(j\omega)}_{\text{phase of } A(j\omega)\beta} \right|$$

at ω_u : unity gain frequency for $T_L(j\omega) = A(j\omega)\beta$
loop gain

Note that there is no contribution to $T_L(j\omega) = A(j\omega)\beta$ in terms of phase by β

since β is positive and real.

Therefore, we take into consideration the phase of $A(j\omega)$ only. But β contributes to $T_L(j\omega)$ in terms of magnitude.

Remember

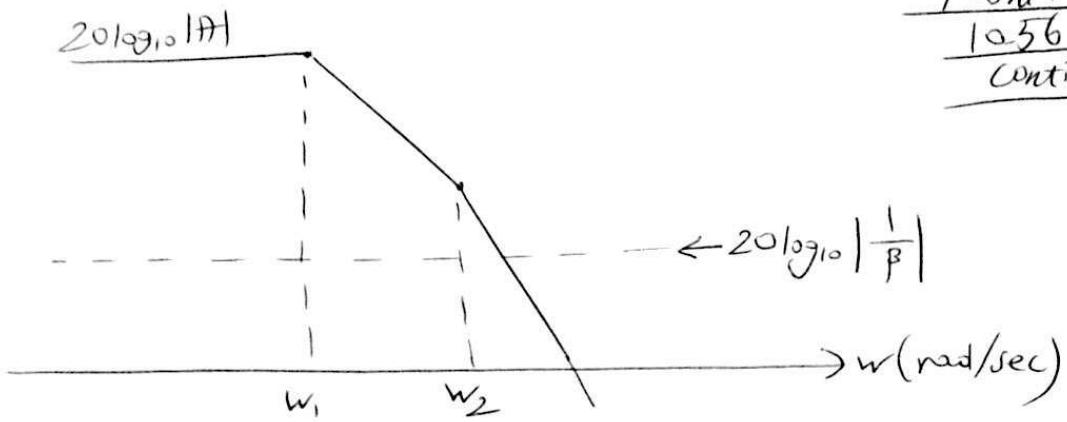
$$20 \log_{10} |A(j\omega)\beta| = 20 \log_{10} |A(j\omega)| - 20 \log_{10} \left| \frac{1}{\beta} \right|$$

Therefore the difference of the two magnitude

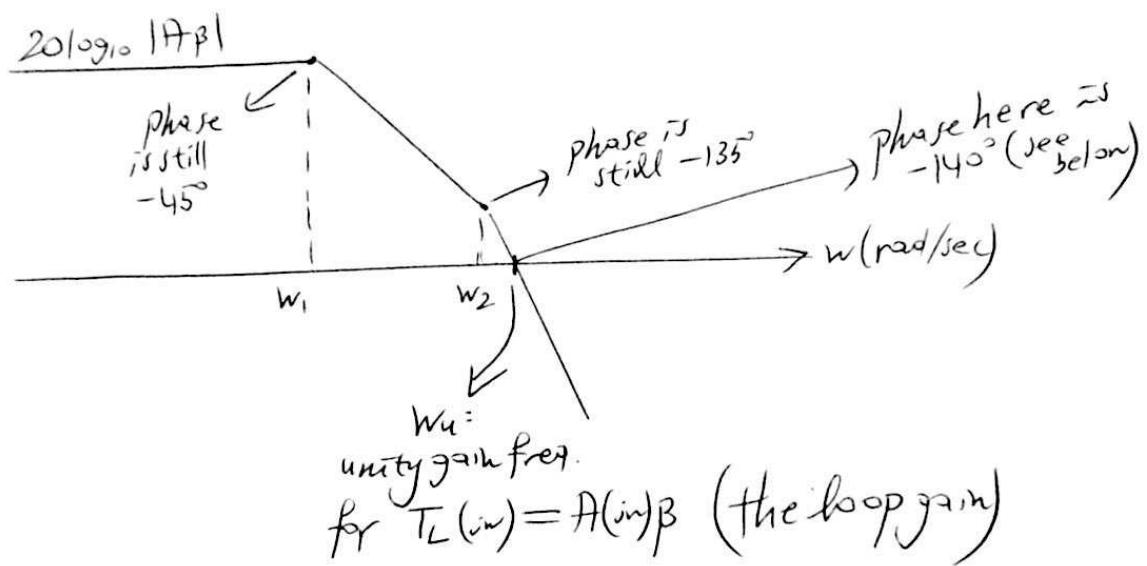
$$\begin{aligned} \text{Bode plots is equal to } & 20 \log_{10} |T_L(j\omega)| \\ & = 20 \log_{10} |A(j\omega)\beta| \end{aligned}$$

See the next page for the Bode plots.

Part b
10.56
contin-



Then the difference will be



Since the phase margin PM is

$$\begin{aligned}
 PM &= 180^\circ - |\angle A(jw)\beta| = 180^\circ - |\angle T_L(jw)| \\
 &= 40^\circ \Rightarrow \text{given in the question}
 \end{aligned}$$

Then $|\angle T_L(jw)| = 140^\circ$

$$\text{and } \angle T_L(jw) = -140^\circ$$

Note that we have

Rashed
10.56
contn.

$$T_L(j\omega_n) = \frac{A\beta}{\left(1 + \frac{j\omega_n}{\omega_1}\right)\left(1 + \frac{j\omega_n}{\omega_2}\right)}$$

$$\approx \frac{A\beta}{\left(\frac{j\omega_n}{\omega_1}\right)\left(1 + \frac{j\omega_n}{\omega_2}\right)}$$

since $\omega_n \gg \omega_1$
but $\omega_n > \omega_2$ only.

Then

$$\angle T_L(j\omega_n) = -90^\circ - \arctan\left(\frac{\omega_n}{\omega_2}\right)$$

$$= -140^\circ$$

and $\arctan\left(\frac{\omega_n}{\omega_2}\right) = 50^\circ$

$$\frac{\omega_n}{\omega_2} \approx \cancel{1.19} \implies \omega_n = \cancel{1.19} \omega_2$$

1.19

And by definition we have

$$\boxed{\boxed{|T_L(j\omega_n)| = 1}}$$

given as 50°
to be designed

$$\approx \frac{A\beta}{\left(\frac{\omega_n}{\omega_1}\right) \sqrt{1 + \frac{\omega_n^2}{\omega_2^2}}}$$

unknown can be computed from here

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Then

$$\beta \approx \left(\frac{w_n}{w_1} \right) \sqrt{1 + \frac{w_n^2}{w_2^2}}$$

~~for 1.55~~
 \approx
 A
 \sim
50

Pashol
10.56
contin.

if we had been given $\frac{w_2}{w_1}$

we could compute

$$\frac{w_n}{w_1} = \frac{w_2}{w_1} \left(\frac{w_n}{w_2} \right) \rightarrow \cancel{\beta} \cancel{for 1.19}$$

Then we could obtain the design value for β .

And the closed-loop gain would be (at low frequencies)

$$A_f = \frac{A}{1 + \beta A}$$

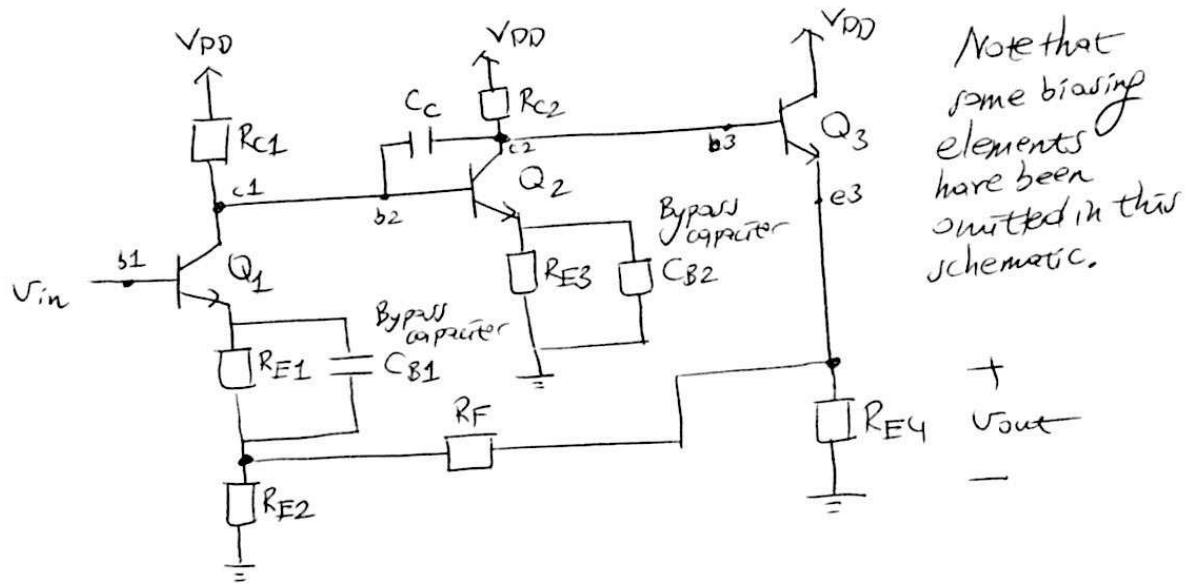


9.7 Compensation

Example of a bandwidth-limiting
compensation on a voltage amplifier
with voltage-series feedback

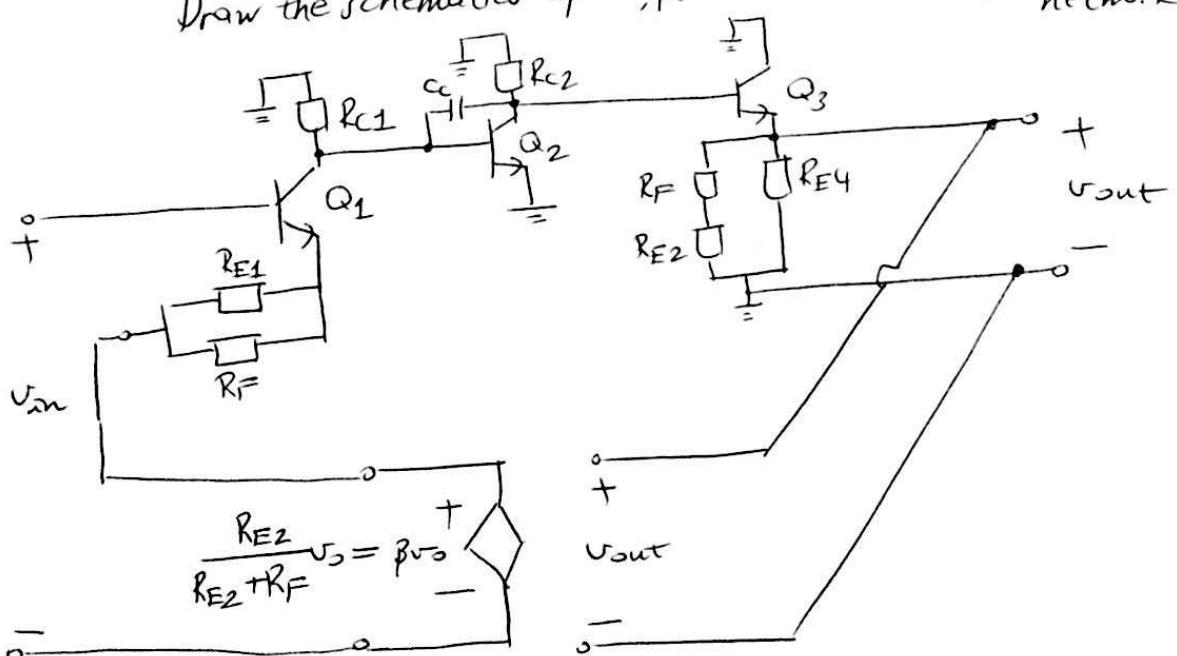
Worked on this on: Nov 26, 2014/Wed

Example of a bandwidth limiting compensation
on a voltage amplifier with voltage-series feedback



Above C_c is a small compensating capacitance.
There is negative feedback in the circuit above,
and it is a voltage amplifier with voltage series feedback.

Draw the schematics of V_{APe} and the idealized feedback network. (small-signal)



Analyze V_A, P_E by letting $\beta = 0$.

$$A_{VA, PE} = \left. \frac{V_{out}}{V_{in}} \right|_{\beta=0} = \frac{V_{C1}}{V_{B1}} \frac{V_{C2}}{V_{B2}} \frac{V_{E3}}{V_{B3}}$$

thus the
gain stage
(CE with
no source
degeneration
in small-signal)

Note \Rightarrow Therefore Miller multiplication
of the compensating cap.
is to be considered.

$$A_{VA, PE} \approx \left[\frac{-g_{m1}}{1 + g_{m1} [R_{E1} // R_F]} \right] \left[R_{C1} // \left[r_{o1} (1 + g_{m1} (R_{E1} // R_F)) \right] // r_{\pi2} \right]$$

$$\cdot \left[-g_{m2} \right] \left[R_{C2} // r_{o2} // \left[r_{\pi3} (1 + g_{m3} (R_{E4} // (R_F + R_{E2}))) \right] \right]$$

$$\cdot \frac{R_{E4} // [R_F + R_{E2}]}{g_{m3} + R_{E4} // [R_F + R_{E2}]}$$

C_c will be designed such that its time constant in OCTC analysis will dominate all the others
therefore

$$W_{H,VA,PE} \approx \frac{1}{\tau_{C_c}}$$

where $\tau_{C_c} = R_{C_c} C_c$

$$R_{C_c} = \left(1 - \frac{v_{C_2}}{v_{B_2}}\right) C_c$$

$$\cdot R_{C_1} \parallel \left[r_{O1} \left(1 + g_{m1} \left(R_{E1} \parallel R_F \right) \right) \right]$$

$$\parallel r_{T2}$$

recall that $\frac{v_{C_2}}{v_{B_2}} = \left[-g_{m2} \right] \left[R_{C_2} \parallel r_{O2} \parallel \left[r_{T3} \left(1 + g_{m3} \left(R_{E4} \parallel (R_F + R_{E2}) \right) \right) \right] \right]$

—————

Therefore we have

$$A_{VA,PE}(jw) \approx \frac{A_{VA,PE}}{1 + \frac{jw}{W_{H,VA,PE}}}$$

with $W_{H,VA,PE}$ having become the dominant pole.
 C_c limits the bandwidth but also is used to adjust the gain and phase margins of the loop gain

$$T_L(jw) = A_{VA,PE}(jw) \beta$$

Calculate the input and output impedances of V_A, P_L

$$r_{in, VA, PL} \approx r_{\pi 1} \left[1 + g_m (R_{E1} // R_F) \right]$$

$$r_{out, VA, PL} \approx R_{E4} // (R_F + R_{E2}) // \frac{r_{\pi 3} + (R_{C2} // r_{o2})}{1 + \beta_3}$$

Analysis of the feedback amplifier
 (V_A, P_L and the idealized
 feedback network)

$$A_f = \frac{A_{VA, PL}}{1 + \beta A_{VA, PL}} \Rightarrow A_f(jw) = \frac{\frac{A_{VA, PL}}{1 + \beta A_{VA, PL}}}{1 + \frac{jw}{w_{H, VA, PL}(1 + \beta A_{VA, PL})}}$$

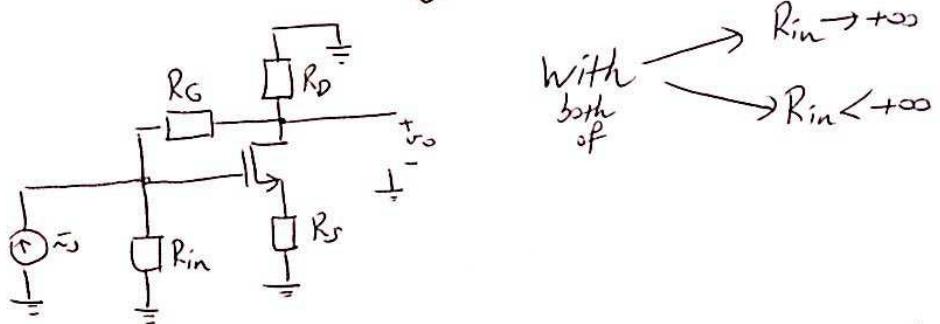
$$w_{H, f} = w_{H, VA, PL} (1 + \beta A_{VA, PL})$$

$$\begin{aligned} r_{in, f} &= r_{in, VA, PL} (1 + \beta A_{VA, PL}) \\ r_{out, f} &= \frac{r_{out, VA, PL}}{1 + \beta A_{VA, PL}} \end{aligned} \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{Due to the properties of voltage-series feedback}$$

9.8 Analysis in Practical Circuits

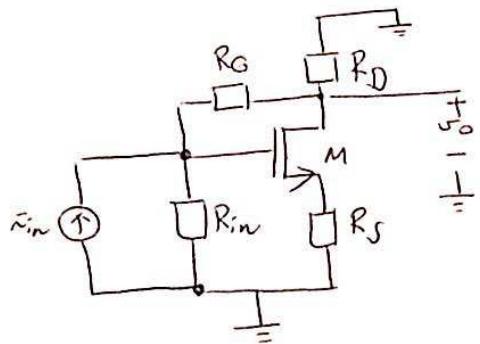
Application of
Voltage shunt feedback on
a transresistance amplifier
and related derivations

⇒ Applied the analysis techniques on

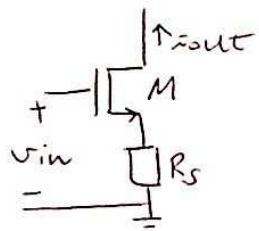


Worked on this on: Oct 27, 2014/Mon

⇒ Analysis of the following amplifier as a transresistance amplifier with voltage-shunt feedback



→ Note that the small signal model for the transistor M should be used in the schematic.

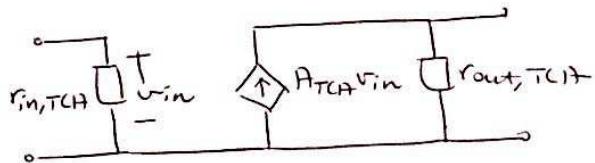


→ We are already acquainted with this circuit as a transconductance amplifier.

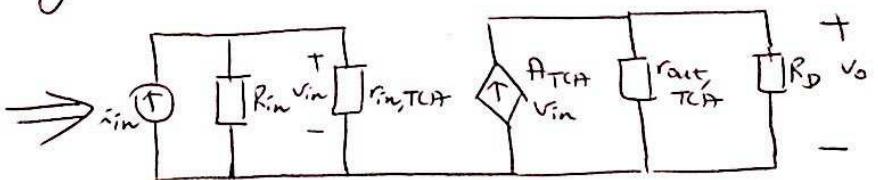
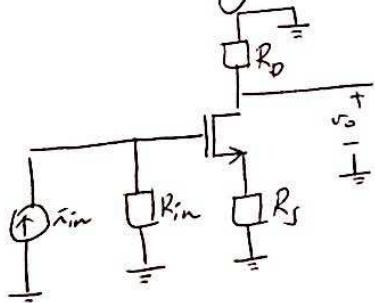
$$r_{in,TCA} = +\infty$$

$$A_{TCA} \approx -\frac{g_m}{1+g_m R_S}$$

$$r_{out,TCA} \approx (1+g_m R_S) r_0$$



To the TCA above, with the addition of R_{in} and R_D , we analyze the resulting amplifier as a transresistance amplifier.



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$$v_{in} = \bar{z}_{in} R_{in} \quad (\text{since } r_{in, TCA} = +\infty)$$

$$v_{out} = A_{TCA} \bar{z}_{in} R_{in} \left(r_{out, TCA} // R_D \right)$$

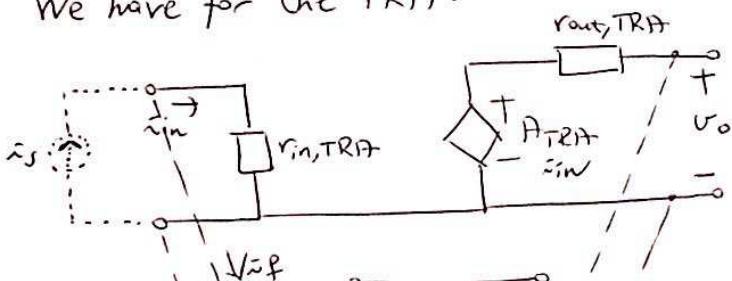
Then the following parameter values can easily be calculated for the transresistance amplifier (TRA).

$$A_{TRA} = - \frac{g_m}{1 + g_m R_s} R_{in} \underbrace{\left[r_{out, TCA} // R_D \right]}_{r_{out, TRA}} = \frac{v_{out}}{\bar{z}_{in}}$$

$$r_{in, TRA} = R_{in}$$

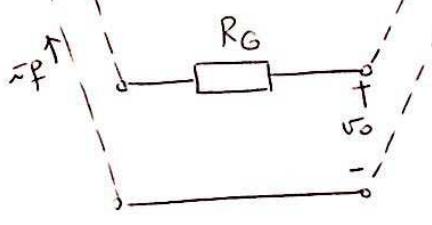
$$r_{out, TRA} = r_{out, TCA} // R_D \quad \xleftarrow{\text{see}}$$

We have for the TRA:



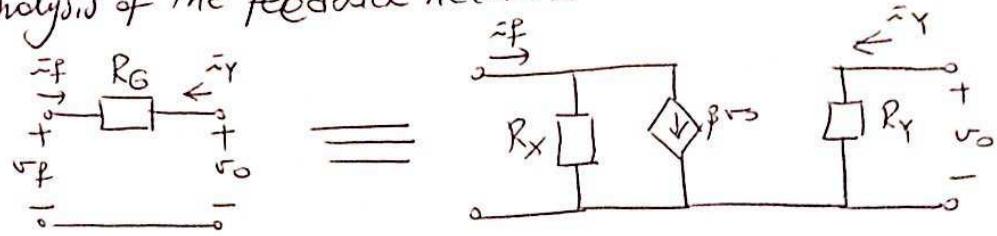
⇒ Dashed (--) lines are the connections between the TRA and the feedback network.

And we have for the feedback network:



⇒ The circuit is fully operational with the inclusion of the (dotted ...) source.

→ Analysis of the feedback network:

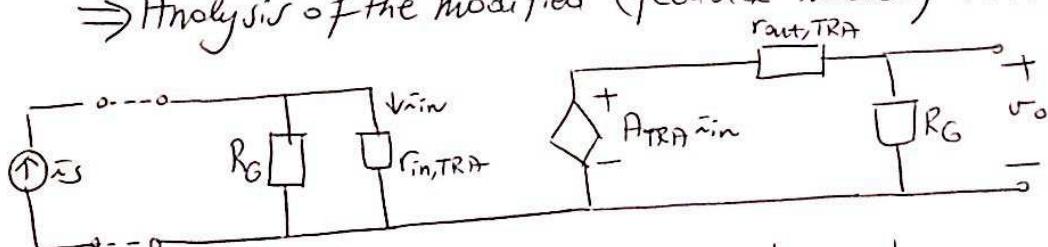


$$R_X = \left. \frac{v_f}{\bar{v}_f} \right|_{v_o=0} = R_G$$

$$R_Y = \left. \frac{v_o}{\bar{v}_Y} \right|_{v_f=0} = R_G$$

$$\beta = \left. \frac{\bar{v}_f}{v_o} \right|_{v_f=0} \Rightarrow -\bar{v}_f R_G = v_o \\ \text{therefore} \\ \beta = \frac{\bar{v}_f}{v_o} = -\frac{1}{R_G}$$

→ Analysis of the modified (feedback-loaded) TRA:



Note that above, the shunt resistances due to the feedback network now load the forward TRA amplifier network. Analyzing the new amplifier (TRA_{fl}) \Rightarrow feedback-loaded

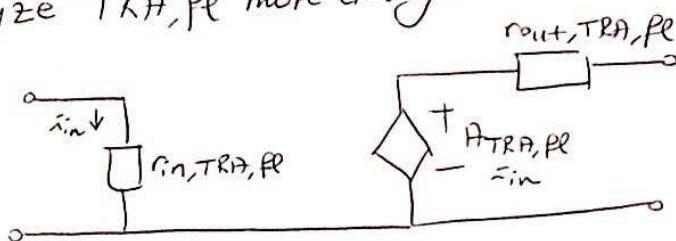
$$A_{\text{TRA, fl}} = \frac{R_G}{R_G + r_{in, \text{TRA}}} A_{\text{TRA}} \frac{R_G}{R_G + r_{out, \text{TRA}}}$$

$$r_{in, \text{TRA, fl}} = R_G // r_{in, \text{TRA}}$$

$$r_{out, \text{TRA, fl}} = R_G // r_{out, \text{TRA}}$$

⇒ Before analyzing the feedback amplifier
($TRA, pe + \text{feedback network}$),

analyze TRA, pe more closely:



- Let $R_{in} \rightarrow +\infty$. Note that $r_{in, TRA} = R_{in}$.

then

$$r_{in, TRA, pe} = R_G // r_{in, TRA} = R_G // +\infty = R_G$$

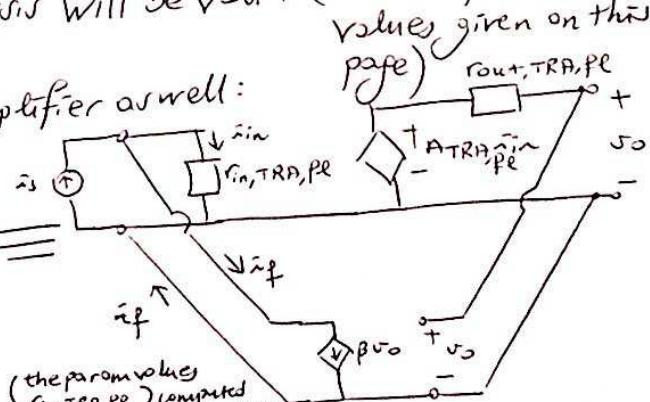
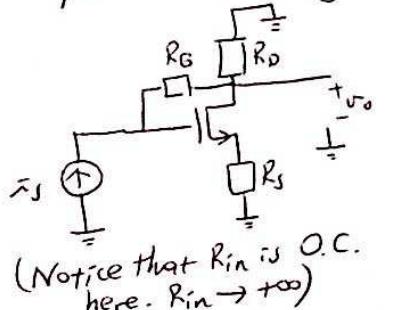
- Analyze $A_{TRA, pe}$ as well for $R_{in} \rightarrow +\infty$

$$\begin{aligned} A_{TRA, pe} &= \frac{R_G}{R_G + R_{in}} \left[\frac{-g_m}{1 + g_m R_S} \right] R_{in} r_{out, TRA} \frac{\frac{R_G}{R_G + r_{out, TRA}}}{R_G} \\ &= \left[\frac{R_G (-g_m)}{1 + g_m R_S} \right] \left[\frac{R_G r_{out, TRA}}{R_G + r_{out, TRA}} \right] \left[\frac{R_{in}}{R_G + R_{in}} \right] \end{aligned}$$

$$\underset{R_{in} \rightarrow +\infty}{A_{TRA, pe}} = \left[\frac{-g_m R_G}{1 + g_m R_S} \right] \left[R_G // r_{out, TRA} \right]$$

Therefore our analysis will be valid (with the parameter values given on this page)

for the following amplifier as well:



\Rightarrow In both cases (with $R_{in} \rightarrow +\infty$ or not)
 the analysis of the feedback amplifier
~~(TRApE + feedback network (not idealized))~~

proceeds as follows:

- Refer to the figure at the bottom of pg 4 (feedback ampli.)
- $\beta = -\frac{1}{R_G}$

	$R_{in} \rightarrow +\infty$	$R_{in} < +\infty$
$A_{TRA, PE}$ $r_{in, TRA, PE}$ $r_{out, TRA, PE}$	see pg 4	see pg 3

- See the note (a separate document) on voltage-shunt feedback on transresistance amplifiers. OR read Rashed pg 677-686.

We have of the feedback amplifier:

$$A_f = \frac{A_{TRA, PE}}{1 + \beta A_{TRA, PE}}$$

$$r_{in, f} = \frac{r_{in, TRA, PE}}{1 + \beta A_{TRA, PE}}$$

$$r_{out, f} = \frac{r_{out, TRA, PE}}{1 + \beta A_{TRA, PE}}$$

Note that
 $A_{TRA, PE} < 0$
 and
 $\beta < 0$,
 but $1 + \beta A_{TRA, PE} > 0$
 therefore negative feedback applies.

Voltage-Series Feedback Example

25. Consider the feedback circuit shown in Fig. 12.88, where $R_1 + R_2 \gg R_D$. Compute the

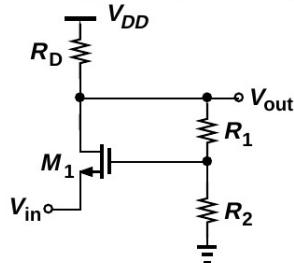


Figure 12.88

closed-loop gain and I/O impedances of the circuit. Assume $\lambda \neq 0$.

Notes: None.

Additional Tasks: None.

Necessary Knowledge and Skills: Voltage-Series Feedback, MOS small signal model, input/output impedance and voltage gain computation.

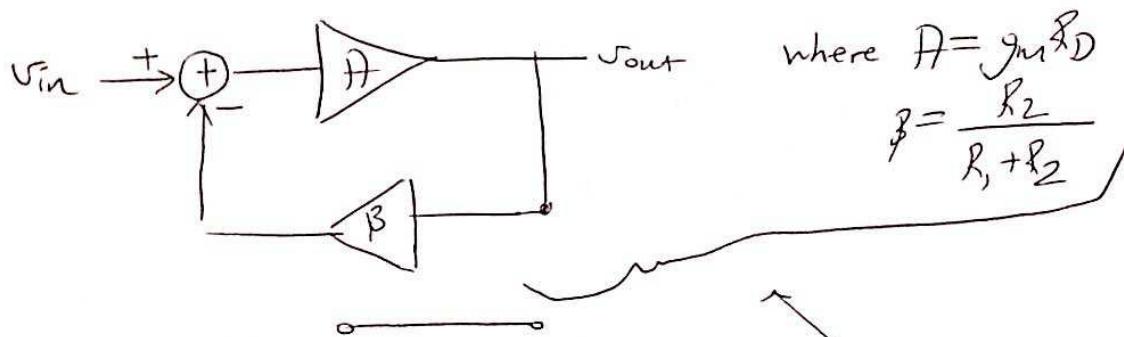
Hazari
12-25

With the condition

$$R_1 + R_2 \gg R_D$$

the current through R_1 and R_2 will be much less than that through R_D .

The circuit can then be accurately modeled by



$$\Delta I_{D1} = \Delta V \cdot g_{m1}$$

$$-R_D \Delta I_{D1} = \cancel{V_{out}}$$

$$\Delta V = V_{out} \frac{R_2}{R_1 + R_2} - V_{in}$$

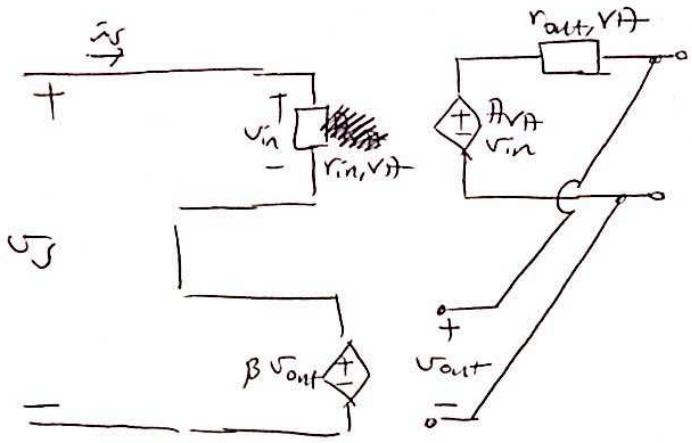
then

$$-R_D \left[V_{out} \frac{R_2}{R_1 + R_2} - V_{in} \right] g_{m1} = V_{out}$$

$$g_{m1} R_D V_{in} = V_{out} \left[1 + g_{m1} R_D \frac{R_2}{R_1 + R_2} \right]$$

$$\frac{V_{out}}{V_{in}} = \frac{g_{m1} R_D}{1 + g_{m1} R_D \frac{R_2}{R_1 + R_2}} = \frac{A}{1 + A\beta} \quad (\text{see the model above})$$

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 cont-



$$v_J - \beta v_{out} = v_{in}$$

$$A_{VA} v_{in} = v_{out} \Rightarrow A_{VA} (v_J - \beta v_{out}) = v_{out}$$

$$A_{VA} v_J = (A_{VA} \beta + 1) v_{out}$$

$$\frac{v_{out}}{v_J} = \frac{A_{VA}}{1 + A_{VA} \beta}$$

compute $\frac{v_J}{\bar{v}_J} = r_{in,f}$ (r_{in} with feedback)

$$\frac{v_J - \beta v_{out}}{r_{in,VA}} = \bar{v}_J = \frac{v_J - \beta \frac{A_{VA}}{1 + A_{VA} \beta} v_J}{r_{in,VA}} = \bar{v}_J$$

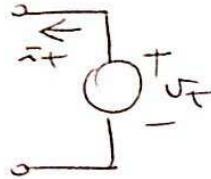
$$\Rightarrow \frac{v_J}{r_{in,VA}(1 + A_{VA} \beta)} = \bar{v}_J$$

$$\Rightarrow \frac{v_J}{\bar{v}_J} = r_{in,f} = r_{in,VA}(1 + A_{VA} \beta)$$

input imp multiplied
 by a large factor



Compute $\frac{v_t}{\bar{v}_t}$



(connect this source at the output and set $v_f = 0$)

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12.25
cont-

$$\beta v_{out} = \beta v_t = -v_{in}$$

$$v_t - \bar{v}_t r_{out,VA} = A_{VA} \underbrace{v_{in}}_{-\beta v_t}$$

$$v_t(1 + A_{VA}\beta) = \bar{v}_t r_{out,VA}$$

$$\frac{v_t}{\bar{v}_t} = \frac{r_{out,VA}}{1 + A_{VA}\beta} \Rightarrow \text{this is } r_{out,f} \quad (\text{output imp. with feedback decreased by the same factor with respect to } r_{out,VA})$$

This config. is called voltage series feedback

$$A_{VA} \rightarrow \frac{A_{VA}}{1 + A_{VA}\beta}$$

$$r_{out,VA} \rightarrow \frac{r_{out,VA}}{1 + A_{VA}\beta}$$

$$r_{in,VA} \rightarrow (1 + A_{VA}\beta) r_{in,VA}$$

$$\text{Note that } r_{in,VA} \approx \frac{1}{j_m 1} \quad r_{out,VA} \approx R_D \quad A_{VA} \approx j_m 1 R_D \\ \beta = \frac{R_2}{R_1 + R_2}$$

(3)

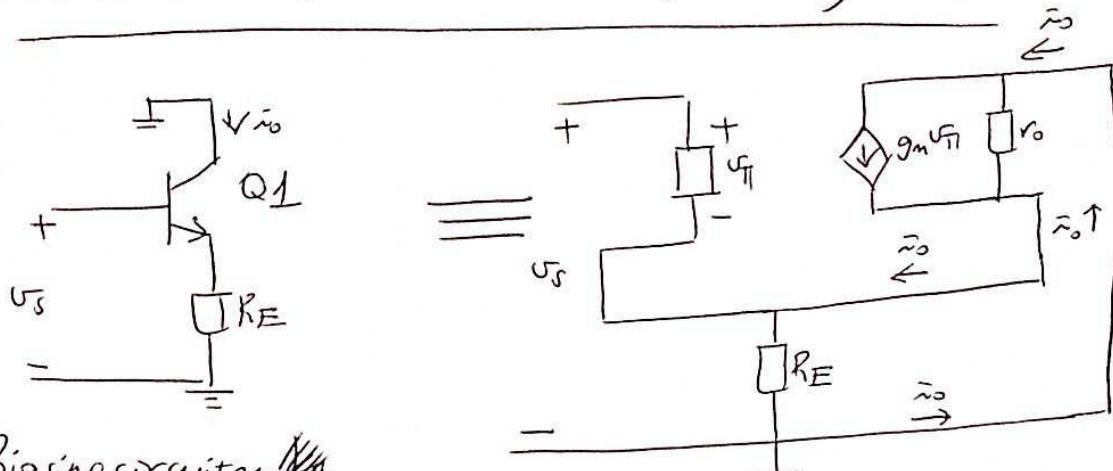
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Worked on this on: Nov 5, 2014/Wed

Analysis of the
Emitter-Degenerated
Common Emitter Configuration
as a Trans-Conductance Amplifier
with Current-Series (or Series-Series)
Feedback

- ⇒ Schematics
- ⇒ Identification of the feedback network
- ⇒ Modelling (computing the parameters) of the feedback network
- ⇒ Analysis of the feedback-loaded TCA ($TCA_{f,pe}$)
(This is a non-feedback or forward amplifier)
- ⇒ Applying the idealized feedback network to the $TCA_{f,pe}$ and computing the parameters of the feedback amplifier

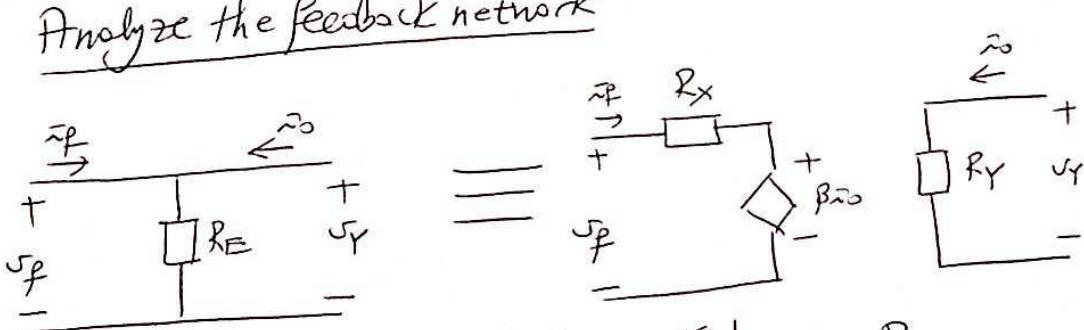
Analysis of the emitter-degenerated common-emitter configuration as a Trans-Conductance amplifier with current-series (or series-series feedback)



Biasing circuitry
not shown - already
in small-signal

TCA with the
feedback network

Analyze the feedback network



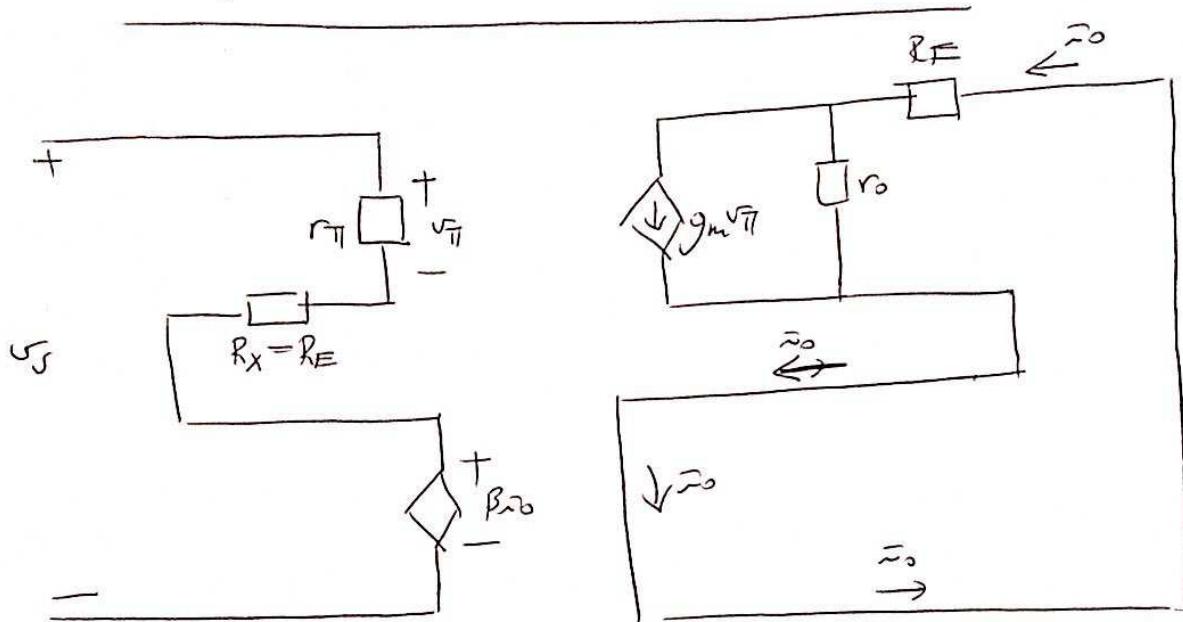
See
Rashid
section
10.7
for similar
derivations

$$\beta = \left. \frac{V_f}{\hat{z}_o} \right|_{\hat{z}_o=0} = R_E$$

$$R_X = \left. \frac{V_f}{\hat{z}_f} \right|_{\hat{z}_o=0} = R_E$$

$$R_Y = \left. \frac{V_o}{\hat{z}_o} \right|_{\hat{z}_o=0} = R_E$$

Draw the feedback-loaded TCA with the idealized network feedback



In the above schematic, set $\beta = 0$ and compute

$$\text{(gain of the feedback loaded TCA)} \quad A_{TCA, fe} = \frac{\tilde{z}_o}{v_s} = \frac{r_\pi}{r_\pi + R_E} g_m \frac{r_o}{r_o + R_E}$$

$$r_{in, TCA, fe} = r_\pi + R_E$$

$$r_{out, TCA, fe} = r_o + R_E$$

Note that above we assume $r_\pi \gg R_E$
 $r_o \gg R_E$

$$\text{then } A_{TCA, fe} \approx g_m$$

$$r_{in, TCA, fe} \approx r_\pi$$

$$r_{out, TCA, fe} \approx r_o$$

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The improvement factor of the current series feedback

$$1 + \beta \text{TF}_{\text{TCA, fe}} \approx 1 + g_m R_E$$

—————

Current-series feedback properties:

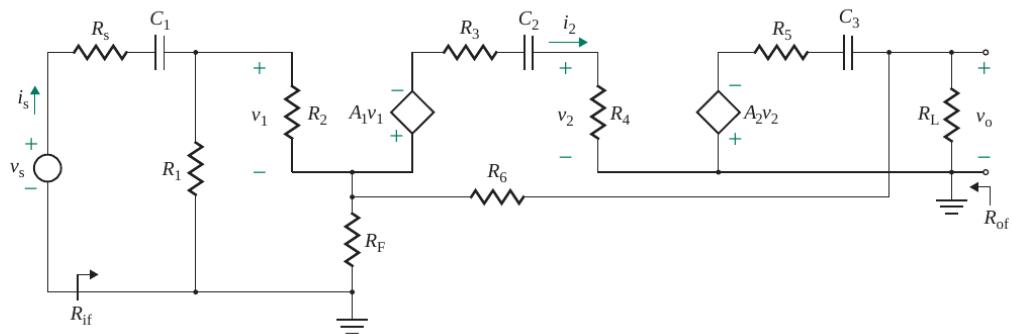
gain	$A_{\text{TCA, pe}} \approx g_m \implies A_f = \frac{A_{\text{TCA, pe}}}{1 + A_{\text{TCA, pe}} \beta} \approx \frac{g_m}{1 + g_m R_E}$
input impedance	$r_{in, \text{TCA, pe}} \approx r_\pi \implies r_{in, f} = r_{in, \text{TCA, pe}} (1 + A_{\text{TCA, pe}} \beta) \approx r_\pi (1 + g_m R_E)$
output impedance	$r_{out, \text{TCA, pe}} \approx r_o \implies r_{out, f} = r_{out, \text{TCA, pe}} (1 + A_{\text{TCA, pe}} \beta) \approx r_o (1 + g_m R_E)$

Gain is decreased by the improvement factor,
but both \approx impedances are multiplied by
the same factor. We have obtained a better
TCA at the expense of some gain.

Voltage Series Feedback on a Voltage Amplifier

- 10.13** The feedback amplifier in Fig. P10.13 has $A_1 = 50$, $A_2 = 60$, $R_s = 500 \Omega$, $R_1 = 15 \text{ k}\Omega$, $R_2 = 1.5 \text{ k}\Omega$, $R_3 = 250 \Omega$, $R_4 = 1.5 \text{ k}\Omega$, $R_5 = 250 \Omega$, $R_6 = 2 \text{ k}\Omega$, $R_L = 4.7 \text{ k}\Omega$, $R_F = 500 \Omega$, $C_1 = C_2 = C_3 = 0.1 \mu\text{F}$, and $v_s = 100 \text{ mV}$. Determine (a) the input resistance $R_{if} = v_s/i_s$, (b) the output resistance R_{of} , and (c) the overall voltage gain $A_f = v_o/v_s$. Assume C_1 , C_2 , and C_3 are shorted at the operating frequency.

FIGURE P10.13



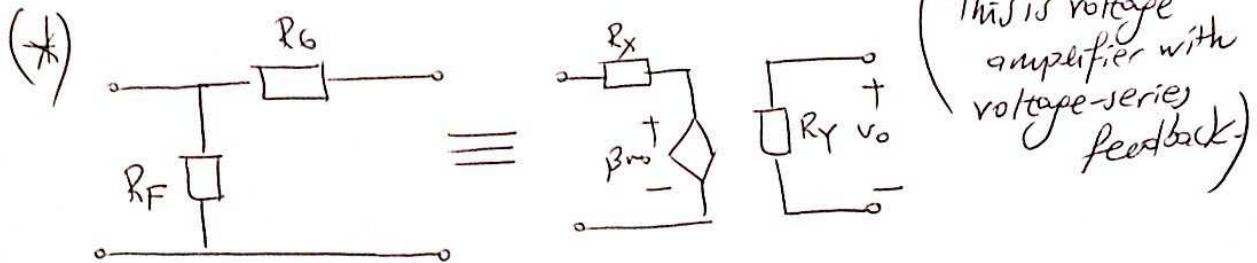
Notes: None.

Additional Tasks: Analyze this circuit as a voltage amplifier with voltage series feedback. Numerical values of the required quantities should be computed as the last task.

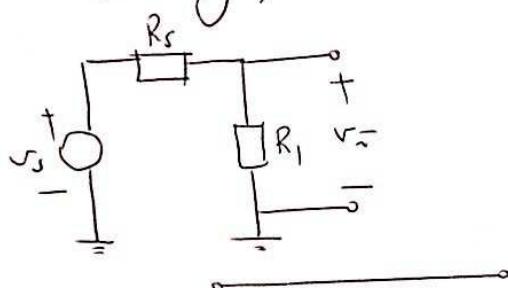
Necessary Knowledge and Skills: Modeling a given circuit as a voltage amplifier, input/output impedance and gain calculations, feedback network modeling, idealizing the feedback network by carrying its impedance loads into the non-feedback amplifier, remodeling the non-feedback amplifier, analyzing the feedback amplifier (non-feedback amplifier and the feedback network combined) to compute the input/output impedances and gain.

Model the feedback network

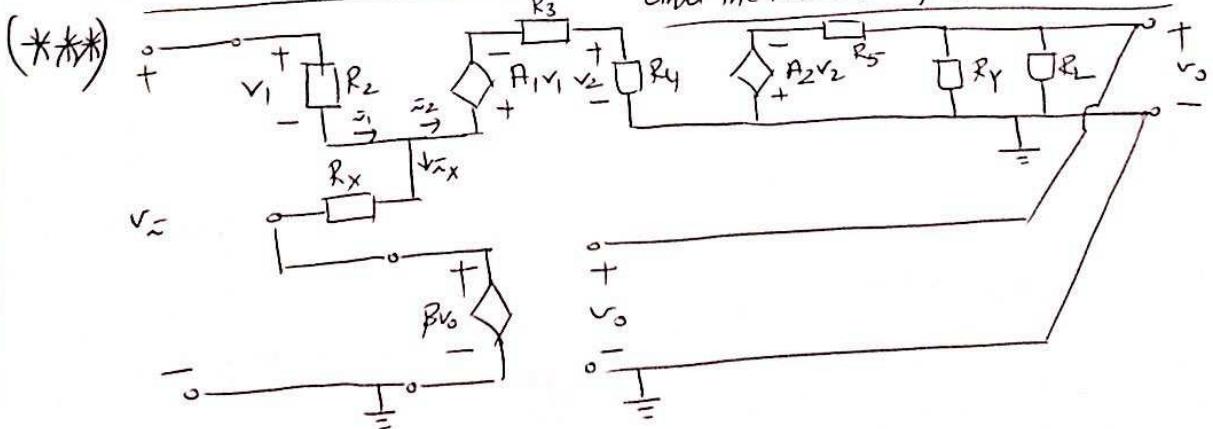
Rashid
10.13



(**) Leave the following part until later.



Schematics of the (VA, PL) Feedback-Loaded Voltage Amplifier
and the realized Feedback network



Kill β ($\beta=0$) and analyze VA, PL

Rashid
10.13
contin.

Compute $\frac{V_1}{V_{\bar{z}}}, \frac{V_2}{V_{\bar{z}}}, \frac{V_{\bar{z}}}{V_1}$ (Refer to (***)

$$(KVL) + V_{\bar{z}} - V_1 - A_1 V_1 - (R_3 + R_4) \bar{z}_2 = 0$$

$$(KVL) + V_{\bar{z}} - V_1 - \bar{z}_x R_x = 0$$

$$(KCL) \bar{z}_1 = \bar{z}_2 + \bar{z}_x$$

$$V_1 = \bar{z}_1 R_2$$

$$V_2 = \bar{z}_2 R_4$$

—————

$$V_{\bar{z}} = (1 + A_1) V_1 + (R_3 + R_4) \bar{z}_2$$

$$\frac{V_1}{R_2} = \bar{z}_2 + \bar{z}_x$$

$$\frac{V_{\bar{z}} - V_1}{R_x} = \bar{z}_x$$

$$V_{\bar{z}} = (1 + A_1) V_1 + (R_3 + R_4) \left[\frac{V_1}{R_2} - \frac{V_{\bar{z}} - V_1}{R_x} \right]$$

$$V_{\bar{z}} \left[1 + \frac{R_3 + R_4}{R_x} \right] = V_1 \left[1 + A_1 + \frac{R_3 + R_4}{R_2 // R_x} \right]$$

$$\frac{V_1}{V_{\bar{z}}} = \frac{\left[1 + \frac{R_3 + R_4}{R_x} \right]}{\left[(1 + A_1) + \left(\frac{R_3 + R_4}{R_2 // R_x} \right) \right]}$$

(2)

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Note that

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10.13
contd.

$$\frac{V_1}{V_{\infty}} = \frac{z_1 R_2}{V_{\infty}} = R_2 \left(\frac{V_{\infty}}{V_1} \right)^{-1}$$

$$r_{in, VA, PL} = \frac{V_{\infty}}{z_1} = R_2 \left(\frac{V_1}{V_{\infty}} \right)^{-1}$$

see the
bottom of pg 2



$$\frac{V_2}{V_{\infty}} = \frac{R_4 z_2}{V_{\infty}} = \frac{R_4}{V_{\infty}} \left[\frac{V_1}{R_2} - \frac{V_{\infty} - V_1}{R_X} \right]$$

$$= \frac{R_4}{V_{\infty}} \left[\frac{V_1}{R_2 // R_X} - \frac{V_{\infty}}{R_X} \right]$$

$$= \frac{R_4}{R_2 // R_X} \underbrace{\frac{V_1}{V_{\infty}}}_{< 0} - \frac{R_4}{R_X} < 0$$

see
the bottom
of pg 2

check this
numerically



Refer to (***)

$$\frac{V_o}{V_2} = (-A_2) \frac{R_Y // R_L}{R_Y // R_L + R_S} < 0 \text{ if } A_2 > 0$$

then

$$\frac{V_o}{V_{\infty}} = \underbrace{\frac{V_2}{V_{\infty}}}_{\text{see pg 3}} \underbrace{\frac{V_o}{V_2}}_{\text{see above}} = A_{VA, PL}$$

$$r_{out, VA, PL} = R_S // R_Y // R_L$$

(3)

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Now apply feedback theory to compute:

Rashed
10.B
contin.

$$A_f = \frac{A_{VA, PE}}{1 + \beta A_{VA, PE}}$$

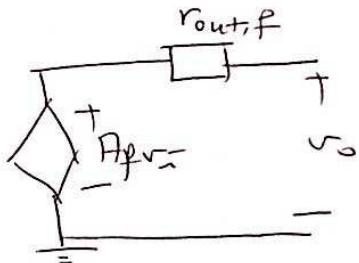
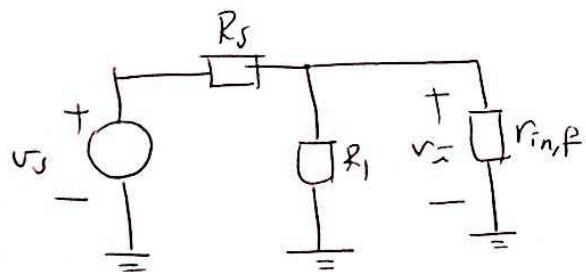
$$r_{in,f} = (1 + \beta A_{VA, PE}) r_{in, VA, PE}$$

$$r_{out,f} = \frac{r_{out, VA, PE}}{1 + \beta A_{VA, PE}}$$

Analysis of
the feedback
amplifier

Now analyze the whole circuit:

Refer to (**) and above



$$\frac{v_o}{v_s} = \frac{r_{in,f} // R_1}{r_{in,f} // R_1 + R_s} A_f$$

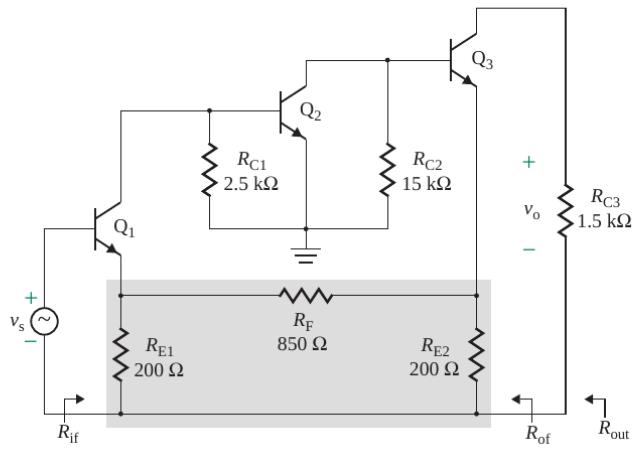
(4)

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Current-Series Feedback on a TCA

- 10.28** The AC equivalent circuit of a feedback amplifier is shown in Fig. P10.28. The circuit values are $R_{C1} = 2.5 \text{ k}\Omega$, $R_{C2} = 5 \text{ k}\Omega$, $R_{C3} = 1.5 \text{ k}\Omega$, $R_{E1} = 100 \Omega$, $R_{E2} = 100 \Omega$, $R_F = 750 \Omega$, and $R_s = 0$. The transistor parameters are $h_{fe} = 100$, $r_\pi = 2.5 \text{ k}\Omega$, $r_o = 25 \text{ k}\Omega$, and $r_\mu = \infty$. Use the techniques of feedback analysis to calculate (a) the input resistance R_{if} , (b) the output resistance R_{of} , and (c) the closed-loop voltage gain A_f .

FIGURE P10.28



Notes: None.

Additional Tasks: None.

Necessary Knowledge and Skills: Modeling the CE stage as a transconductance amplifier, current-series feedback network modeling in a practical amplifier, gain and i/o impedance computations, feedback-loaded TCA analysis, effect of feedback on gain and i/o impedances.

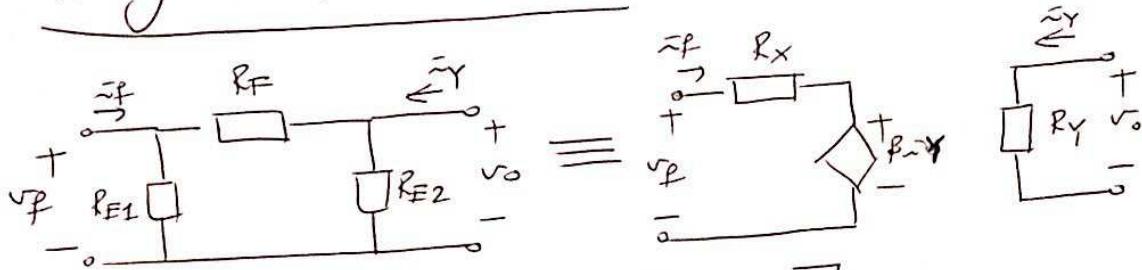
There is current-series feedback

Rashid
10.28

involved in this circuit.

The whole amplifier identifying the feedback network has already been drawn.

Analyze the feedback network (see pg 668 of Rashid)



$$R_x = \frac{v_f}{i_f} \Big|_{i_y=0} = R_{E1} \parallel [R_F + R_{E2}]$$

$$R_Y = \frac{v_o}{i_Y} \Big|_{i_f=0} = R_{E2} \parallel [R_{E1} + R_F]$$

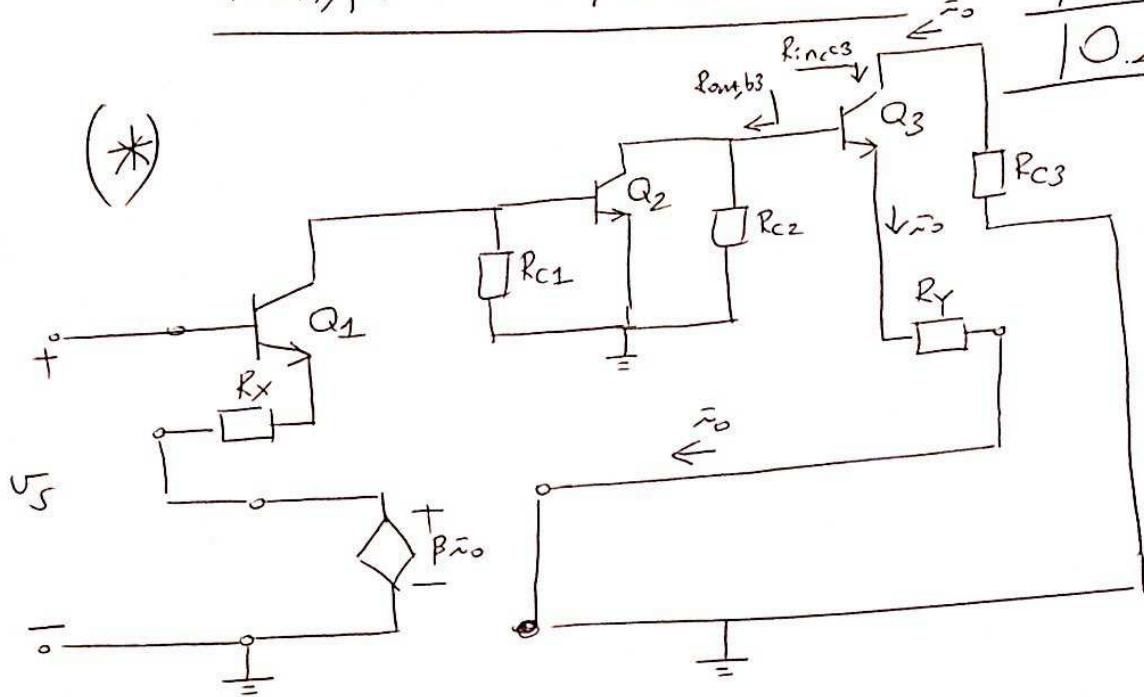
$$\beta = \frac{v_f}{i_Y} \Big|_{i_f=0} = \frac{R_{E2}}{R_{E2} + R_{E1} + R_F} \cdot R_{E1}$$

Next \Rightarrow Draw the schematic of TCA_{fl} and the current-series feedback network.

TCA, PL and the feedback network

Rashid
10.28

(*)



Above in (*), set $\beta = 0$, Analyze TCA, PL

$$A_{TCA, PL} \approx \left[\frac{-g_{m1}}{1 + g_{m1} R_x} \right] \left[\frac{[(1 + g_{m1} R_x) r_{o1}]}{r_{o2} // R_{C2} // [r_{T3}(1 + g_{m3} R_Y)]} \right]$$

$$\cdot \left[-g_{m2} \right] \cdot \left[r_{o2} // R_{C2} // [r_{T3}(1 + g_{m3} R_Y)] \right]$$

$$\cdot \left[\frac{g_{m3}}{1 + g_{m3} R_Y} \right]$$

$$r_{in, TCA, pe} \approx r_{\pi 1} (1 + g_{m1} R_x)$$

frashd
10.28
contin.

$$r_{out, TCA, pe} \approx R_{C3} + R_{in, c3}$$

$$\text{note that } R_{out, b3} = R_{C2} // r_{o2}$$

$$\text{then } R_{in, c3} = (1 + g_{m3} r_{o3}) \frac{R_Y r_{\pi 3}}{R_Y + r_{\pi 3} + R_{out, b3}} + r_o$$

(see related document on the derivation)

$$+ \frac{R_Y R_{out, b3}}{R_Y + r_{\pi 3} + R_{out, b3}}$$

if the trailing two terms can be neglected

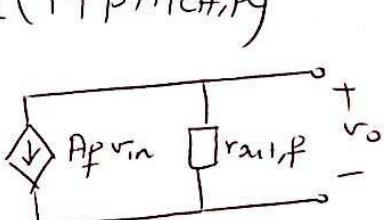
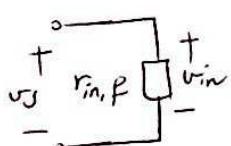
$$\approx (1 + g_{m3} r_{o3}) \left[\frac{R_Y r_{\pi 3}}{R_Y + r_{\pi 3} + R_{out, b3}} \right]$$

Now analyze the whole amplifier:

$$A_f = \frac{A_{TCA, pe}}{1 + \beta A_{TCA, pe}}$$

$$r_{in, f} = r_{in, TCA, pe} (1 + \beta A_{TCA, pe})$$

$$r_{out, f} = r_{out, TCA, pe} (1 + \beta A_{TCA, pe})$$



$$\frac{v_o}{v_s} = -A_f r_{out, f}$$

Analysis of Feedback in OpAmp Circuits

***8.26** For each of the op-amp circuits shown in Fig. P8.26, identify the feedback topology and indicate the output variable being sampled and the feedback signal. In each case, assuming the op amp to be ideal, find an expression for β , and hence find A_f .

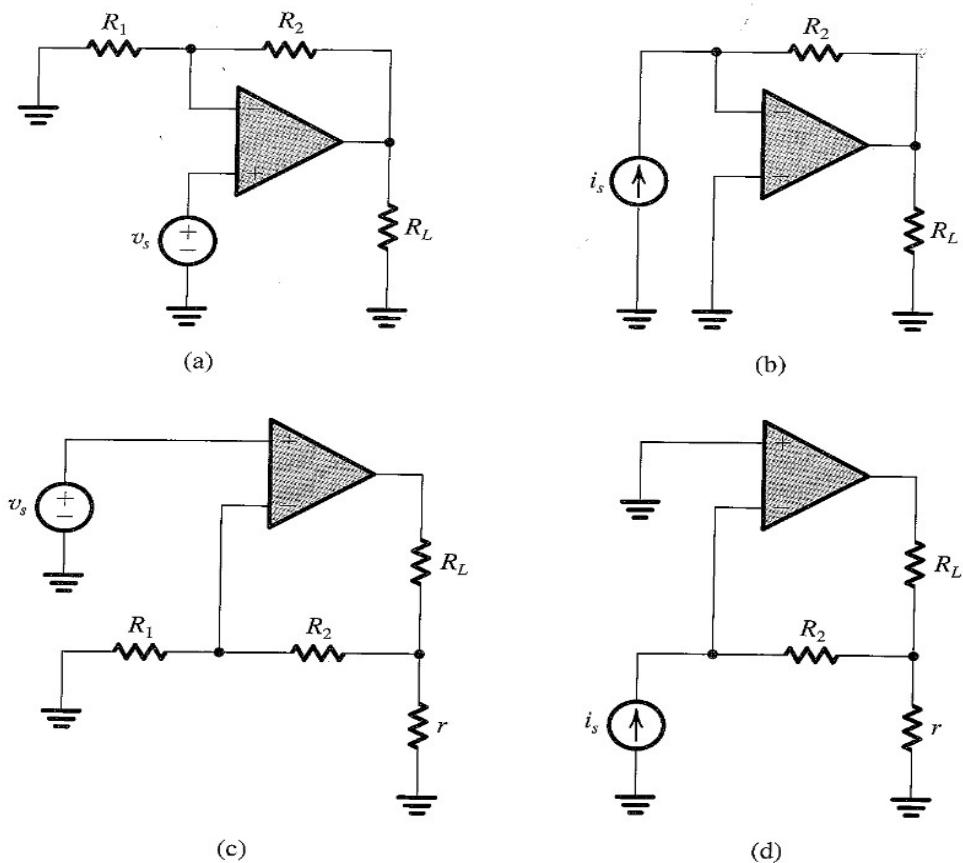


FIGURE P8.26

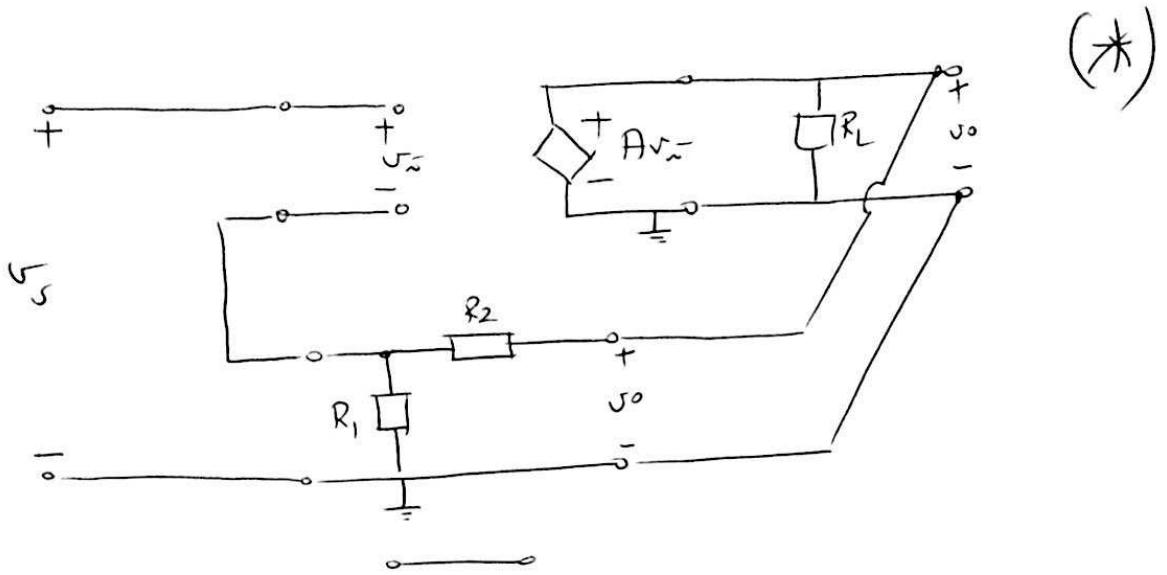
Notes: None.

Additional Tasks: None.

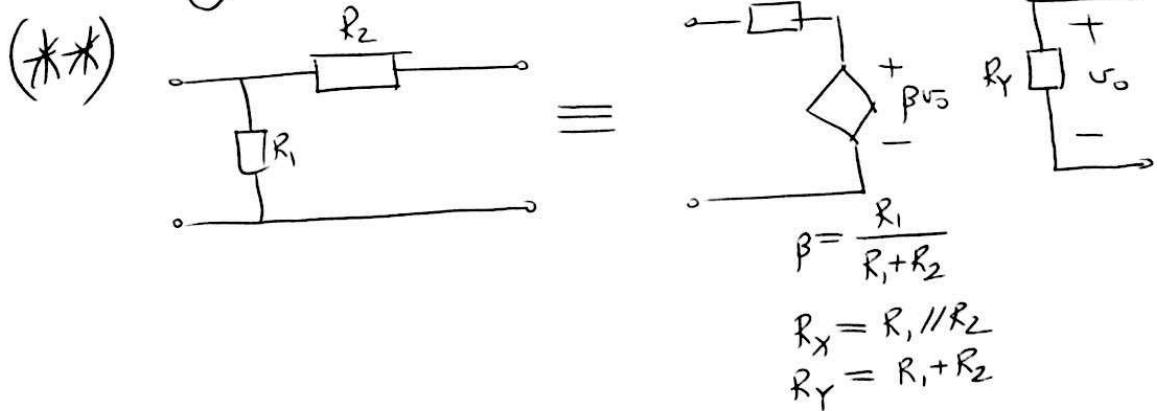
Necessary Knowledge and Skills: The four types of feedback, identification and modeling of feedback networks, feedback-loaded amplifier analysis, gain and input/output imp. computation.

This is a voltage amplifier with voltage series feedback.

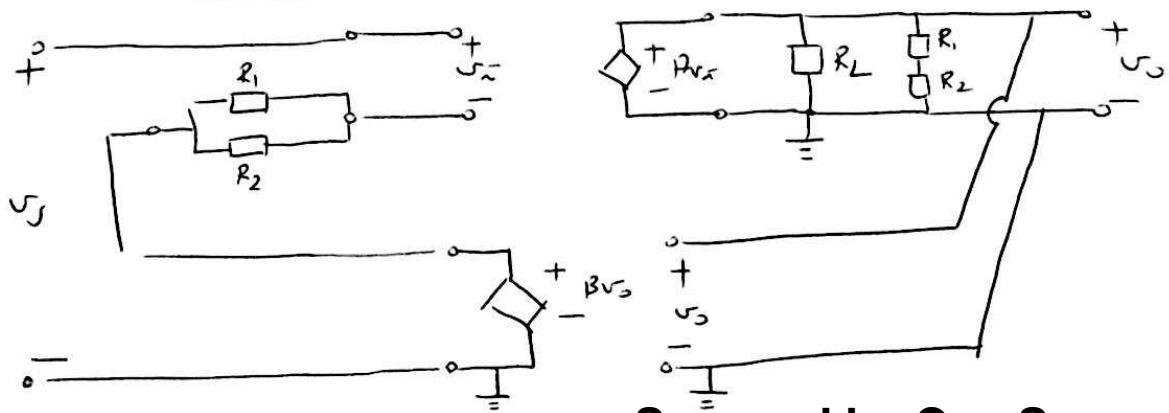
Sedra 8.260)



Analyze the feedback network:



(****) Schematics of VAF, FP and the feedback network (idealized)



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$$\Rightarrow k_i \parallel \beta \quad (\beta = 0)$$

\Rightarrow Analyze V_A, R_L

$$r_{in, VA, RL} = \infty \quad \begin{matrix} \text{(would be } r_{in} + R_1 \parallel R_2 \\ \text{if the opamp were nonideal)} \end{matrix}$$

$$r_{out, VA, RL} = 0 \quad \begin{matrix} \text{(would be } r_{out} \parallel R_L \parallel (R_1 + R_2) \\ \text{if the opamp were nonideal)} \end{matrix}$$

$$A_{VA, RL} = A \quad \begin{matrix} \text{(would be } \frac{r_{in}}{r_{in} + R_1 \parallel R_2} A \quad \frac{R_L \parallel (R_1 + R_2)}{R_L \parallel (R_1 + R_2) + r_{out} \\ \text{if the opamp were nonideal)} \end{matrix}$$

Now analyze the whole feedback amplifier

$$r_{in,f} = r_{in, VA, RL} \quad (1 + \beta A_{VA, RL}) = \infty \quad \begin{matrix} \text{(or phs in the} \\ \text{value above} \\ \text{and } \beta) \end{matrix}$$

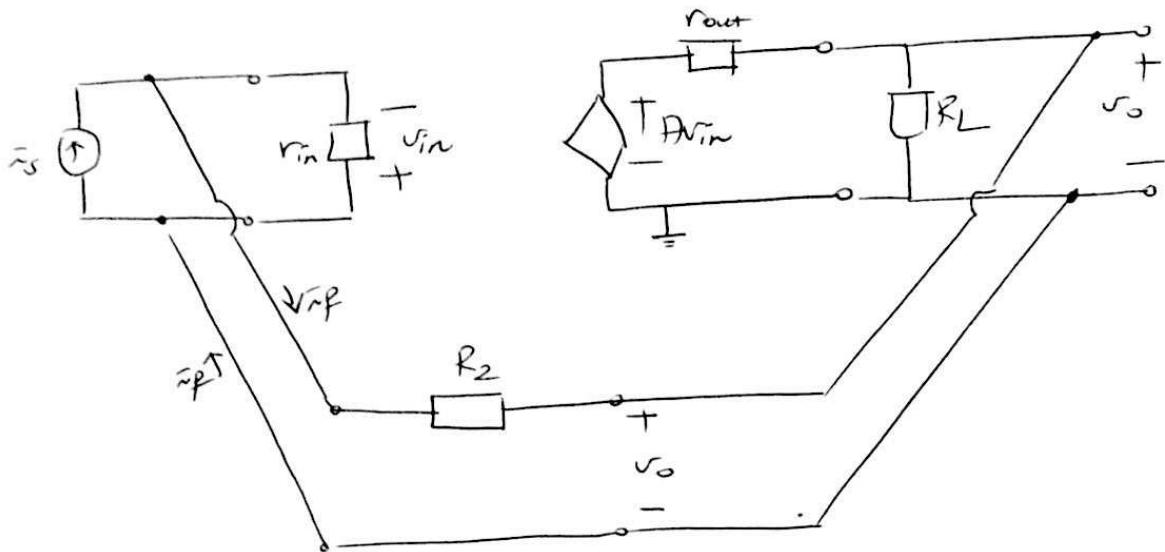
$$r_{out,f} = \frac{r_{out, VA, RL}}{1 + \beta A_{VA, RL}} = 0$$

$$A_f = \frac{A_{VA, RL}}{1 + \beta A_{VA, RL}} = \frac{A}{1 + \frac{R_1}{R_1 + R_2} A} \approx \frac{R_1 + R_2}{R_1} \\ = 1 + \frac{R_2}{R_1} \quad (\sim F \gg 1)$$

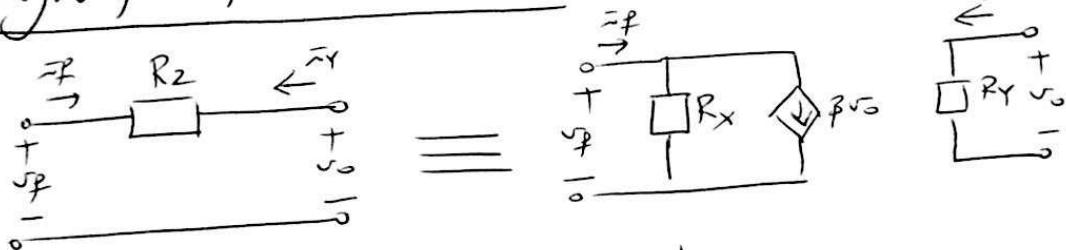
Sedra
8.26x
contin.

This is a transresistance amplifier with voltage-shunt feedback.

Sedra 8.26b)



Analysis of the feedback network



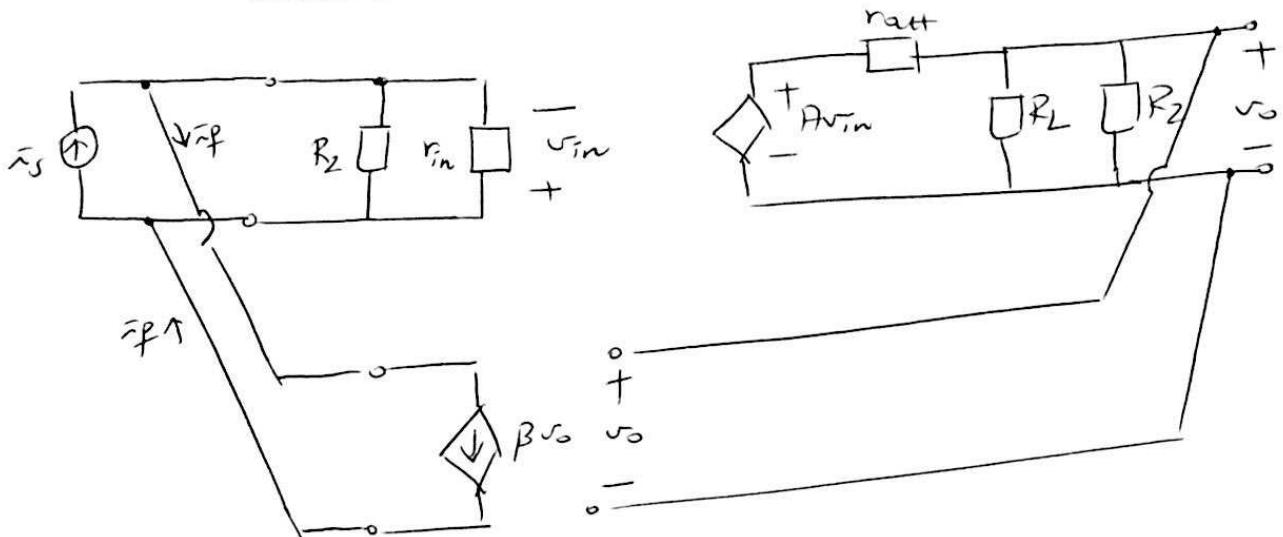
$$\beta = \frac{\bar{v}_f}{v_o} \Big|_{v_f=0} = \frac{-\bar{v}_Y}{v_o} = -\frac{1}{R_2}$$

$$R_Y = \frac{v_o}{\bar{v}_Y} \Big|_{v_f=0} = R_2$$

$$R_X = \frac{v_f}{\bar{v}_f} \Big|_{v_o=0} = R_2$$

TRA, PL and the idealized feedback network

Sedra 8 26b)



Kill β ($\beta=0$) and analyze ~~TRA, PL~~ TRA, PL

$$r_{in, \text{TRA, PL}} = R_2 // r_{in}$$

$$r_{out, \text{TRA, PL}} = r_{out} // R_L // R_2$$

$$A_{\text{TRA, PL}} = \left[-R_2 // r_{in} \right] \cdot \left[\frac{R_L // R_2}{R_L // R_2 + r_{out}} \right]$$

Note that with $r_{in} = +\infty$ and $r_{out} = 0$
the above would be

$$r_{in, \text{TRA, PL}} = R_2$$

$$r_{out, \text{TRA, PL}} = 0$$

$$A_{\text{TRA, PL}} = -R_2 \cdot A$$

Now analyze the whole feedback ampli.

Sedro 826
b)
contin

$$r_{in,f} = \frac{r_{in,TRA,PE}}{1 + \beta A_{TRA,PE}}$$

$$= \frac{R_2 // r_{in}}{1 + \left(-\frac{1}{R_2}\right)(-R_2 // r_{in}) A \left(\frac{R_L // R_2}{R_L // R_2 + r_{out}}\right)}$$

with $r_{in} = +\infty$ $r_{out} = 0$

$$\approx \frac{R_2}{1 + \left(-\frac{1}{R_2}\right)(-R_2) A}$$

$$= \frac{R_2}{1 + A}$$

with $A \gg 1$

$$\approx \textcircled{O}$$

$$r_{out,f} = \frac{r_{out,TRA,PE}}{1 + \beta A_{TRA,PE}}$$

with $r_{in} = +\infty$ $r_{out} = 0$

$$\approx \frac{\textcircled{O}}{1 + \left(-\frac{1}{R_2}\right)(-R_2 A)}$$

$$A_f = \frac{A_{TRA,PE}}{1 + \beta A_{TRA,PE}} = \frac{\left[-R_2 // r_{in}\right] A \left[\frac{R_L // R_2}{R_L // R_2 + r_{out}}\right]}{1 + \left(-\frac{1}{R_2}\right)\left[-R_2 // r_{in}\right] A \left[\frac{R_L // R_2}{R_L // R_2 + r_{out}}\right]}$$

$$\begin{aligned} A_f &\stackrel{\sim}{=} \frac{-R_2 A}{1 + \left(-\frac{1}{R_2}\right)(-R_2)A} \\ &= \frac{-R_2 A}{1 + A} \end{aligned}$$

with $r_{in} = +\infty$
 $r_{out} = 0$

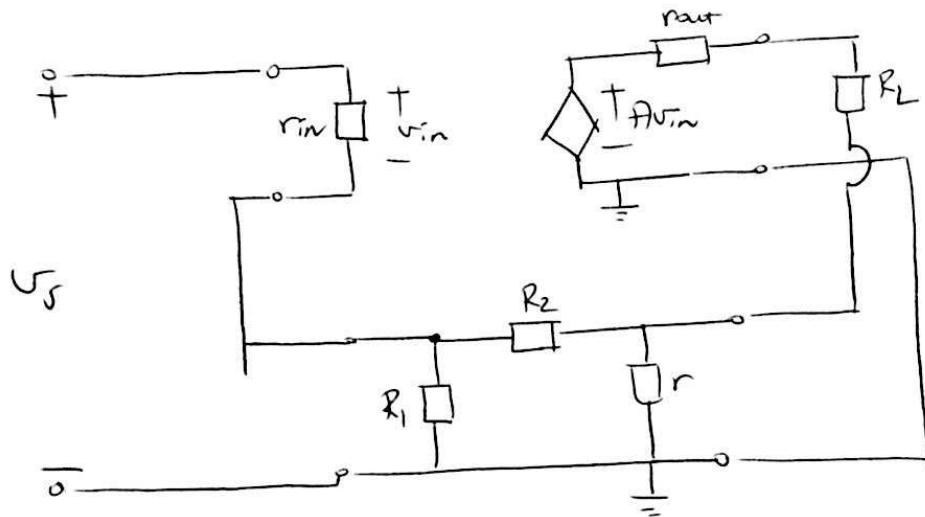
$$A \gg 1 \quad \approx -R_2$$

Sedro
8.26
cont-

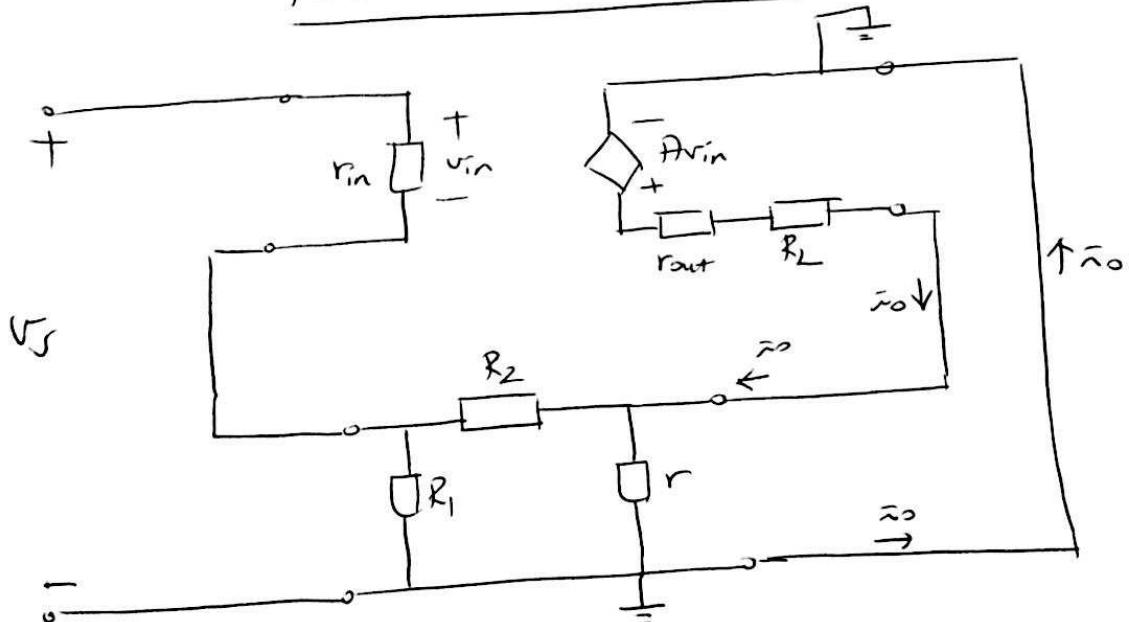


Transconductance amplifier with current-series feedback

Sedra
8.26c



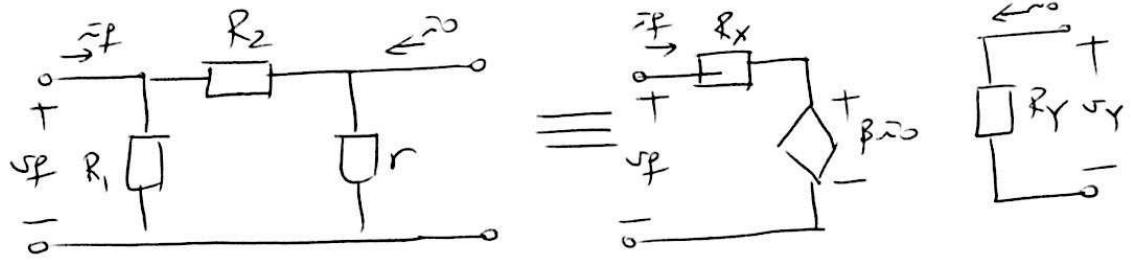
Redraw the above schematics:



Analyze the feedback network

Sedra
8.26c

cont -

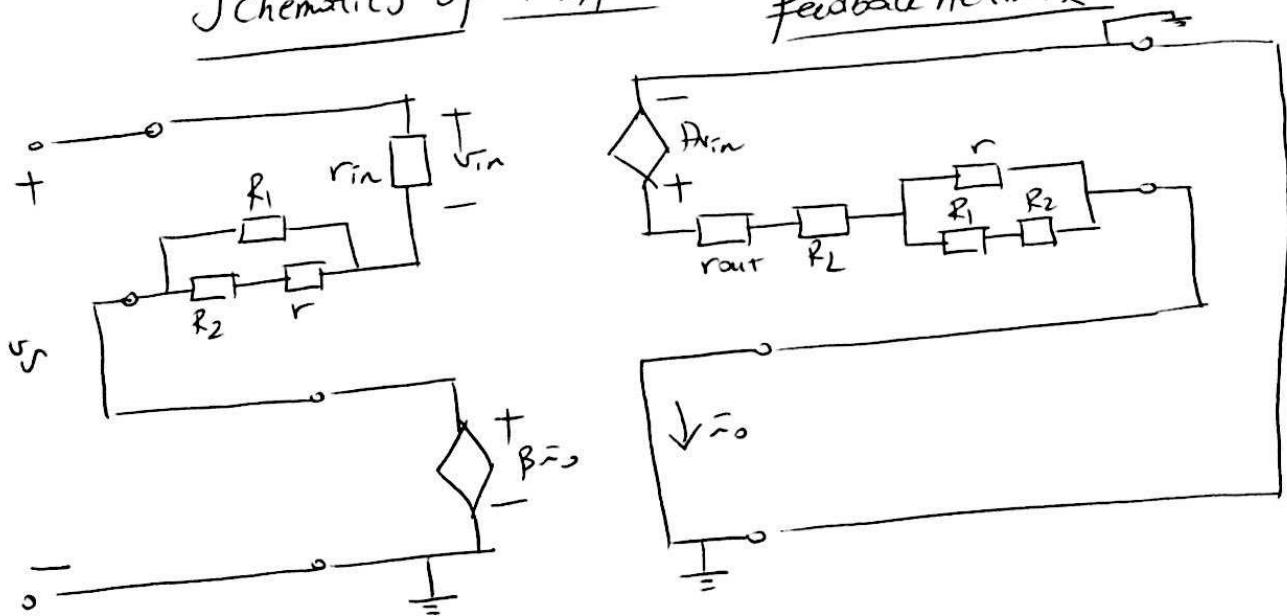


$$\beta = \frac{v_f}{v_o} \Big|_{v_f=0} = \frac{r}{r+R_1+R_2} \cdot R_1$$

$$R_X = \frac{v_f}{r} \Big|_{v_o=0} = R_1 \parallel (R_2 + r)$$

$$R_Y = \frac{v_Y}{v_o} \Big|_{v_f=0} = r \parallel (R_1 + R_2)$$

Schematics of $T(s)$, f_L and the idealized feedback network



$\Rightarrow \text{kill } \beta \ (\beta=0)$

$\Rightarrow \text{Analyze TCA, PL}$

Sedra
8.26c
contin.

$$r_{in, TCA, PL} = r_{in} + R_1 // (R_2 + r)$$

$$r_{out, TCA, PL} = r_{out} + R_L + r // (R_1 + R_2)$$

$$A_{TCA, PL} = \frac{r_{in}}{r_{in} + R_1 // (R_2 + r)} A \frac{1}{r_{out} + R_L + r // (R_1 + R_2)}$$



Note that with $r_{in} = +\infty$
 $r_{out} = 0$

$$r_{in, TCA, PL} \approx +\infty$$

$$r_{out, TCA, PL} \approx R_L + r // (R_1 + R_2)$$

$$A_{TCA, PL} \approx A \frac{1}{R_L + r // (R_1 + R_2)}$$



Now analyze the whole feedback amplifier

Sedra
8.26c
contin-

$$r_{in,f} = (1 + \beta A_{TCA,pe}) r_{in,TCA,pe}$$

$$r_{out,f} = (1 + \beta A_{TCA,pe}) r_{out,TCA,pe}$$

$$A_f = \frac{A_{TCA,pe}}{1 + \beta A_{TCA,pe}}$$

Note that with $\frac{r_{in}}{r_{out}} = \infty$

$$\rightarrow r_{in,f} = \infty$$
$$\rightarrow r_{out,f} = \left(1 + \frac{rR_1}{r+R_1+R_2} A \frac{1}{R_L + r/(R_1+R_2)} \right) \cdot \left[R_L + r/(R_1+R_2) \right]$$

$$\text{with } A \gg 1 \approx \frac{rR_1 A}{r+R_1+R_2}$$

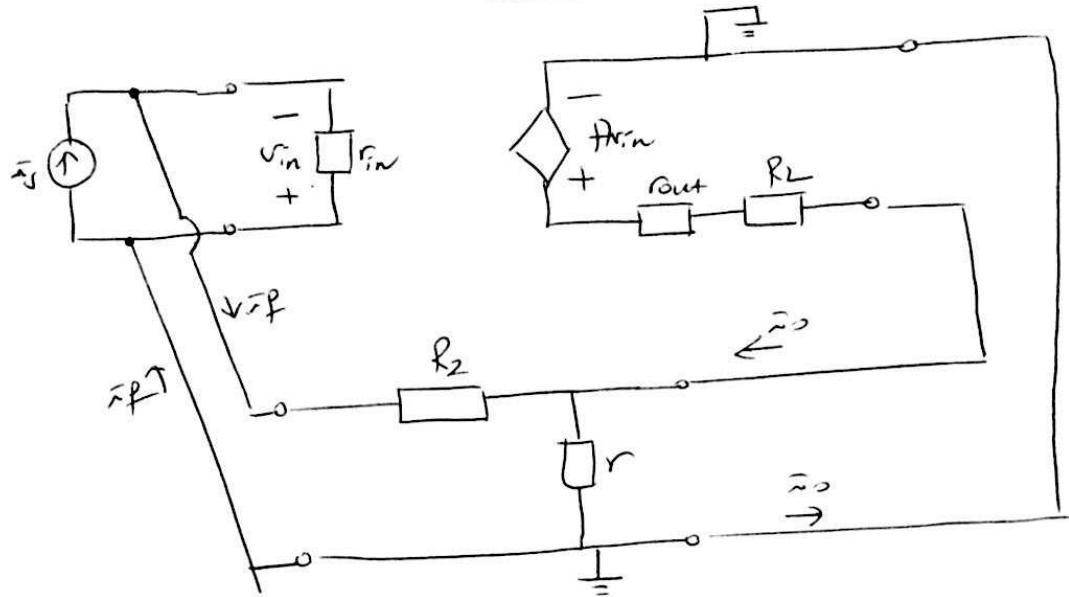
$$\rightarrow A_f \approx \frac{A \frac{1}{R_L + r/(R_1+R_2)}}{1 + \frac{rR_1}{r+R_1+R_2} A \frac{1}{R_L + r/(R_1+R_2)}}$$

$$\text{with } A \gg 1 \approx \frac{r+R_1+R_2}{rR_1} = \frac{1}{r/R_1} + \frac{R_2}{rR_1}$$

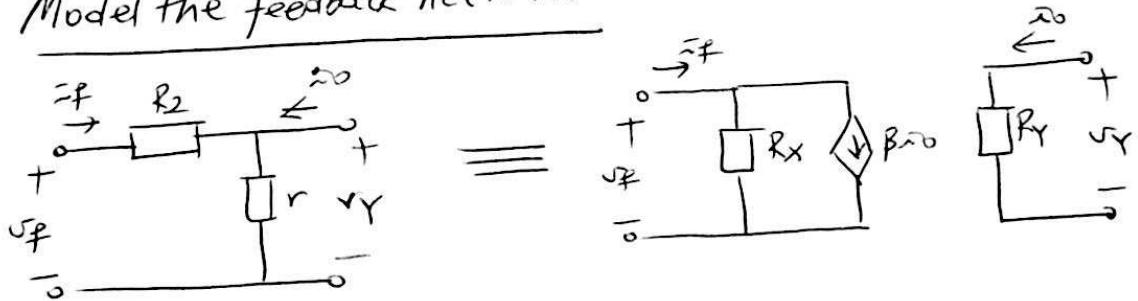
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This is \Rightarrow current amplifier with
current-shunt feedback

Sedra
8.26d



Model the feedback network



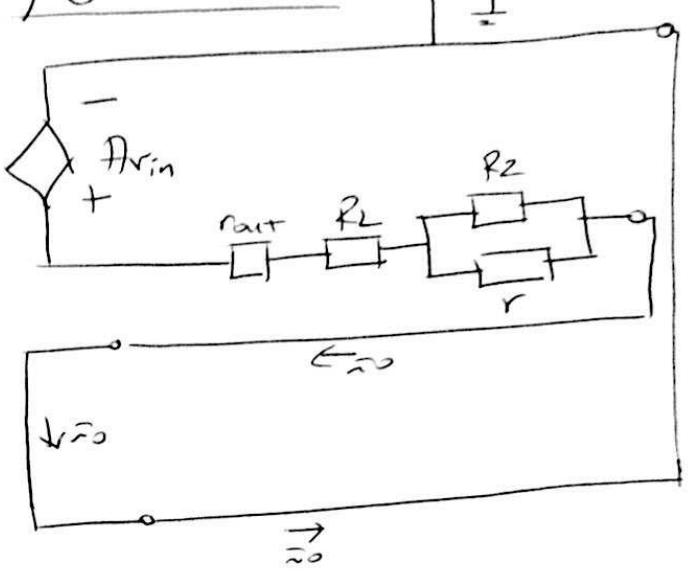
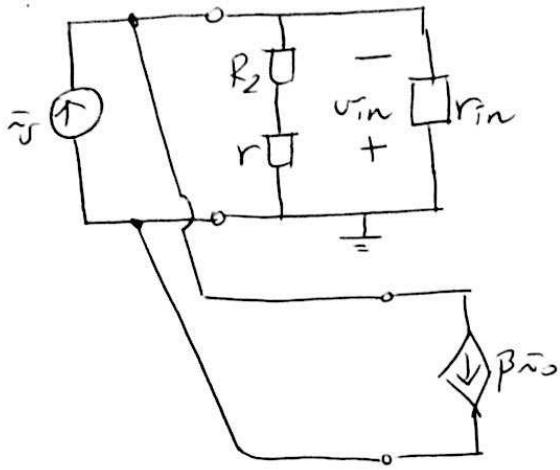
$$R_X = \frac{v_f}{i_f} \Big|_{i_o=0} = R_2 + r$$

$$R_Y = \frac{v_f}{i_o} \Big|_{v_f=0} = R_2 \parallel r$$

$$\beta = \frac{i_f}{i_o} \Big|_{v_f=0} = -\frac{r}{r+R_2}$$

Schematics of CA, PL and the idealized feedback network

Sedra 8.26d
cont.



\Rightarrow kill β ($\beta = 0$)

\Rightarrow Analyze CA, PL

$$r_{in, CA, PL} = r_{in} \parallel (R_2 + r)$$

$$r_{out, CA, PL} = r_{out} + R_L + R_2 \parallel r$$

$$A_{CA, PL} = -\frac{R_2 + r}{R_2 + r + r_{in}} \cdot \frac{1}{r_{out} + R_L + R_2 \parallel r}$$



With $r_{in} = +\infty$
 $r_{out} = 0$

Jedra
8.26 d)

$$r_{in, CA, PE} \approx R_2 + r$$

$$r_{out, CA, PE} \approx R_L + R_2 // r$$

$$A_{CA, PE} \approx -A \frac{(R_2 + r)}{R_L + (R_2 // r)}$$

→

Analyze the whole current amplifier:

$$A_f = \frac{A_{CA, PE}}{1 + \beta A_{CA, PE}}$$

$$\begin{aligned} \text{with } & r_{in} = +\infty \approx \frac{-A \frac{(R_2 + r)}{R_L + (R_2 // r)}}{1 + \left(\frac{-r}{r + R_2} \right) (-A) \left(\frac{R_2 + r}{R_L + R_2 // r} \right)} \\ \text{and } & r_{out} = 0 \end{aligned}$$

$$\begin{aligned} \text{with } & A \gg 1 \approx \frac{-A \frac{(R_2 + r)}{R_L + R_2 // r}}{rA \frac{1}{R_L + R_2 // r}} = -\frac{R_2 + r}{r} = -\left[1 + \frac{R_2}{r} \right] \end{aligned}$$

$$r_{in, f} = \frac{r_{in, CA, PE}}{1 + \beta A_{CA, PE}} \approx \frac{R_2 + r}{1 + \left(\frac{-r}{r + R_2} \right) (-A) \left(\frac{R_2 + r}{R_L + R_2 // r} \right)}$$

with
 $r_{in} = +\infty$
 $r_{out} = 0$

$$\begin{aligned} \text{with } & A \gg 1 \approx \frac{1}{rA} \left[(R_2 + r) \left[R_L + \frac{R_2 r}{R_2 + r} \right] \right] \\ & = \frac{(R_2 + r) R_L + R_2 r}{rA} \end{aligned}$$

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$$r_{out,f} = (1 + \beta A_{CA,PL}) r_{out,CA,PL}$$

Sectra
8.26 d)
contd.

with
 $r_{in} = \infty$
 $r_{out} = 0$

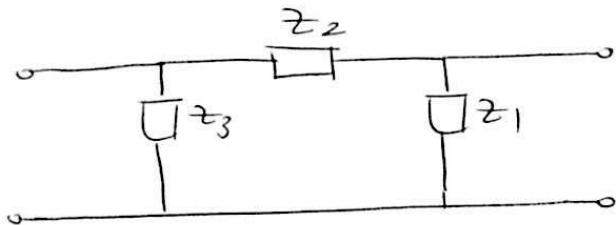
$$\approx \left[1 + \left(\frac{-r}{r+R_2} \right) (-A) \left[\frac{R_2+r}{R_L+R_2/r} \right] \right] (P_L + P_2/r)$$

$$\text{with } A \gg 1 \approx Ar \frac{\cancel{R_L + R_2/r}}{\cancel{R_L + R_2/r}} = Ar$$

10. Oscillators

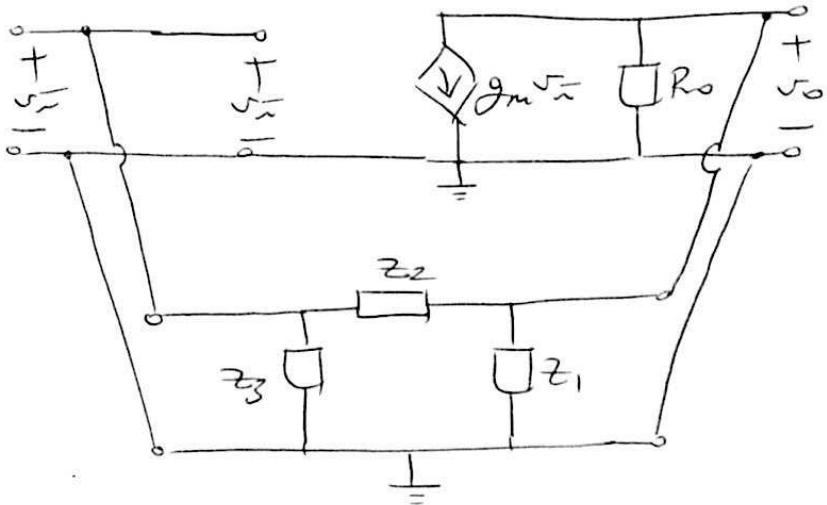
10.1 Auxiliaries

Analysis of LC networks
encountered in RF oscillators



Worked on this on: Nov 27, 2014/Th

Analysis of common LC networks
encountered in RF oscillators



$$A \cdot \beta = -g_m \left[R_o / \parallel z_1 / \parallel [z_3 + z_2] \right] \frac{z_3}{z_3 + z_2}$$

$$= -g_m \frac{1}{\frac{1}{R_o} + \frac{1}{z_1} + \frac{1}{z_3 + z_2}} \frac{z_3}{z_3 + z_2}$$

$$= -g_m \frac{\frac{z_3 + z_2}{z_3 + z_2}}{\frac{z_3 + z_2}{R_o} + \frac{z_3 + z_2}{z_1} + 1} \frac{z_3}{z_3 + z_2}$$

$$= \frac{-g_m R_o z_1 z_3}{z_1 z_3 + z_1 z_2 + z_2 z_3 + R_o (z_1 + z_2 + z_3)}$$

If each element Z_1, Z_2, Z_3 is an inductor or a capacitor:

$$Z_i = jX_i$$

Then we have

$$\alpha \cdot \beta = \frac{-g_m R_o (-X_1 X_3)}{-(X_1 X_3 + X_1 X_2) + j R_o (X_1 + X_2 + X_3)}$$

The phase condition requires

$$X_1 + X_2 + X_3 = 0$$

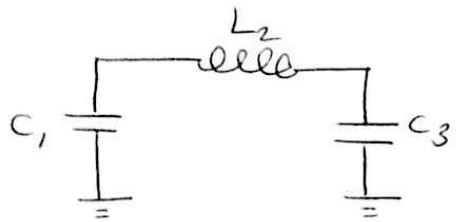
For self starting oscillations we have:

$$\frac{-g_m R_o (X_1 X_3)}{X_1 (X_2 + X_3)} \geq 1$$

$$\frac{-g_m R_o X_3}{X_3 + X_2} \geq 1$$

—————

Example = Colpitts Oscillator



We have for the phase condition:

$$\frac{1}{jwC_1} + jwL_2 + \frac{1}{jwC_3} = 0$$

$$\frac{-1}{wC_1} + wL_2 - \frac{1}{wC_3} = 0$$

$$wL_2 = \frac{1}{w} \frac{C_1 + C_3}{C_1 C_3}$$

then $w_o = \frac{1}{\sqrt{L_2 \frac{C_1 C_3}{C_1 + C_3}}}$

For self-starting oscillations:

$$\frac{\frac{g_m R_o \left(\frac{1}{w_o C_3} \right)}{-\frac{1}{w_o C_3} + w_o L_2}}{w_o^2 L_2 C_3 - 1} \geq 1$$

$$= \frac{\frac{g_m R_o}{\frac{1}{C_1 + C_3} L_2 C_3 - 1}}{w_o^2 L_2 C_3 - 1} \geq 1$$

$$g_m R_o \geq \frac{C_3}{C_1} \quad \text{for self-starting osc}$$

in a Colpitts
osc.

Audio Oscillator Analysis

13.13 For the circuit in Fig. P13.13 find $L(s)$, $L(j\omega)$, the frequency for zero loop phase, and R_2/R_1 for oscillation.

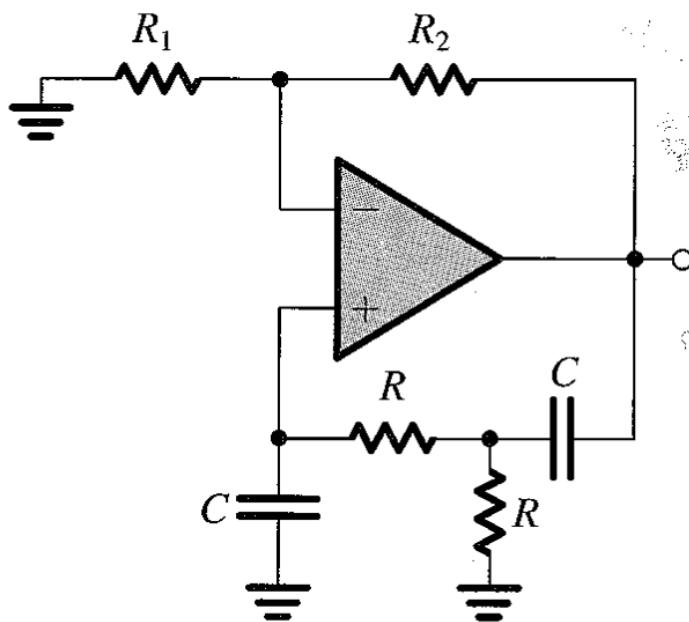


FIGURE P13.13

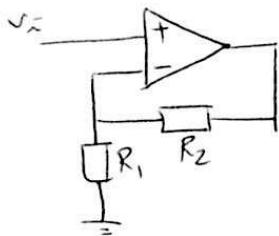
Notes: None.

Additional Tasks: None.

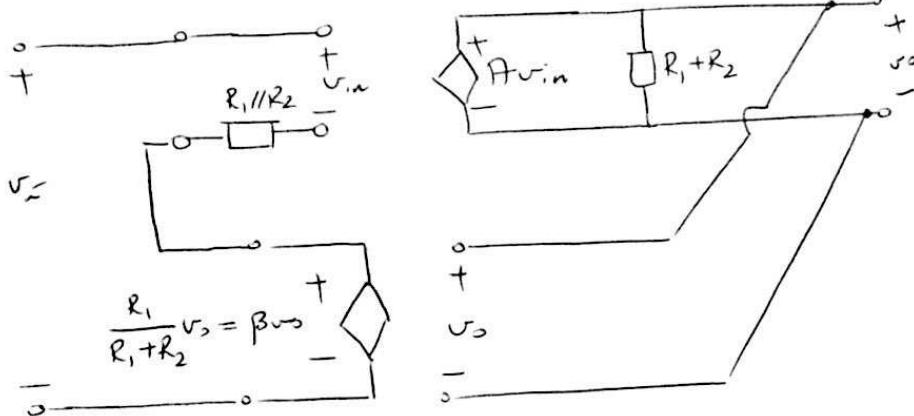
Necessary Knowledge and Skills: Oscillators, Barkhausen criteria, analysis of feedback networks, oscillation frequency, phase condition, amplitude condition.

Analyze the amplifier

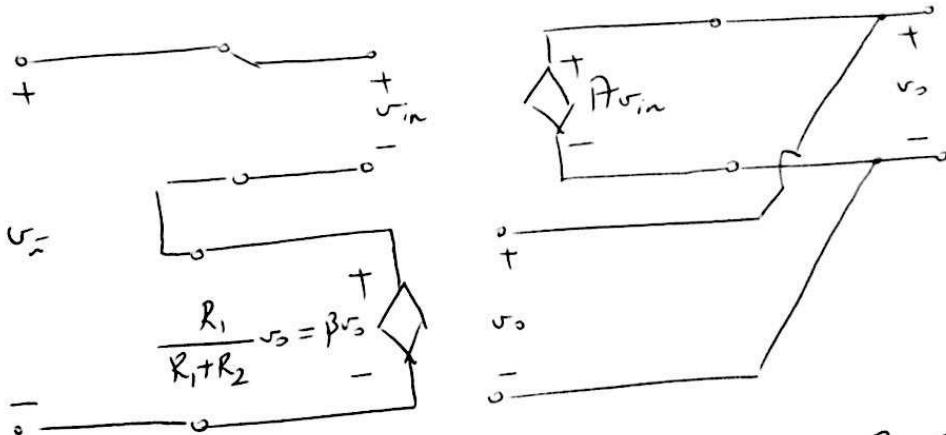
Seatra
13-13



This is a voltage amplifier
with voltage series feedback



This is equivalent to



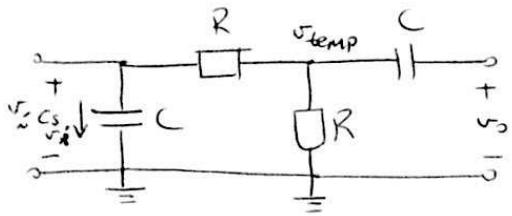
$$\text{Then } \frac{v_o}{v_{in}} = \frac{A}{1 + A \frac{R_1}{R_1 + R_2}} \underset{\text{with } A \rightarrow +\infty}{\approx} \frac{R_1 + R_2}{R_1} = 1 + \frac{R_2}{R_1}$$

and $r_{in,f} \rightarrow +\infty$
 $r_{out,f} \rightarrow 0$

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Analyze the feedback network

Sedra
13-13
contin-



$$\frac{v_i}{v_o} = \frac{\left(R + \frac{1}{C_s}\right) // R}{\left(R + \frac{1}{C_s}\right) // R + \frac{1}{C_s}} \cdot \frac{\frac{1}{C_s}}{\frac{1}{C_s} + R}$$

OR do it like this

$$\frac{v_{temp} - v_{\bar{o}}}{R} = \frac{v_{\bar{o}}}{\frac{1}{C_s}} = C_s v_{\bar{o}}$$

$$\begin{aligned} v_{temp} &= v_{\bar{o}} + R C_s v_{\bar{o}} \\ &= (1 + R C_s) v_{\bar{o}} \end{aligned}$$

and KCL at the node with v_{temp}

$$\frac{v_{temp}}{R} + C_s v_i + C_s (v_{temp} - v_{\bar{o}}) = 0$$

$$v_{temp} + R C_s v_{\bar{o}} + R C_s (v_{temp} - v_{\bar{o}}) = 0$$

$$v_{temp} (1 + R C_s) + R C_s v_{\bar{o}} = R C_s v_{\bar{o}}$$

$$v_{\bar{o}} (1 + R C_s)^2 + R C_s v_{\bar{o}} = R C_s v_{\bar{o}}$$

$$v_{\bar{o}} [R^2 C_s^2 + 3 R C_s + 1] = R C_s v_{\bar{o}}$$

$$\frac{v_{\bar{o}}}{v_{\bar{o}}} = \frac{R C_s}{R^2 C_s^2 + 3 R C_s + 1}$$

$$\frac{v_{\bar{o}}}{v_{\bar{o}}} = \frac{1}{3 + \frac{1}{R C_s} + R C_s}$$

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Therefore

$$\beta(j\omega) = \frac{v_o(j\omega)}{v_o(j\omega)} = \frac{1}{3 + j\left(wCR - \frac{1}{wCR}\right)}$$

Sedra
13.13
contin-

Recall that A_{VA} of the forward amplifier is

$$A_{VA} = 1 + \frac{R_2}{R_1}$$

$$A_{VA} \beta(j\omega) = \frac{1 + \frac{R_2}{R_1}}{3 + j\left(wCR + \frac{1}{wCR}\right)}$$

Now apply the Barkhausen criterion:
Zero phase condition requires

$$wCR - \frac{1}{wCR} = 0$$

$$\Rightarrow w_o = \frac{1}{RC} \quad (\text{osc frequency})$$

unity magnitude condition requires (at the osc freq. w_o)

$$|A_{VA} \beta(jw_o)| = 1$$

$$\frac{1 + \frac{R_2}{R_1}}{3} = 1 \quad \text{and for self-starting osc}$$

$$1 + \frac{R_2}{R_1} \geq 3$$

$$\frac{R_2}{R_1} \geq 2$$