

# Adjoint Computational Electromagnetics Eigenvalue Problems

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The sensitivities of eigenvalues and eigenvectors due to perturbations in the design variables are derived. The focus here is on linear, isotropic, nonmagnetic and transparent materials <sup>1</sup>.

<sup>1</sup> In transparent materials, both  $\epsilon$  and  $\mu$  are real and positive.

## Problem Formulation

Maxwell's equation for reads

$$\nabla \times \mathbf{E} = -j\mu_0\omega\mathbf{H} \quad (1a)$$

$$\nabla \times \mathbf{H} = j\omega\epsilon_0\epsilon(r)\mathbf{E} \quad (1b)$$

which can be combined into a single PDE <sup>2</sup>

$$\nabla \times \frac{1}{\epsilon} \nabla \times \mathbf{H} = \left(\frac{\omega}{c}\right)^2 \mathbf{H} \quad (2)$$

<sup>2</sup> We will refer to this equation as the master equation.

Equation. 2 is a standard eigenvalue problem of the form,

$$Ax = \lambda x \quad (3)$$

where we have identified the operator  $A = \nabla \times \frac{1}{\epsilon} \nabla \times$  and the eigenvector  $x = \mathbf{H}$  with associated eigenvalue  $\lambda = \omega^2$ . The operator  $A$  can be shown to be hermitian and positive-semidefinite. The eigenmodes can be normalized

$$x^T x = 1 \quad (4)$$

## Sensitivity Derivatives (Nondegenerate Case)

Our optimization problem can be formulated as

$$\begin{aligned} \min_p \quad & g(p, x, \lambda) \\ \text{s.t.} \quad & Ax = \lambda x \\ & x^T x = 1 \end{aligned} \quad (5)$$

where  $p$  is a vector of design parameters and  $g$  is the objective function we seek to minimize. To develop the eigen-adjoint variable method, we first define the Lagrangian as

$$L = g(p, x, \lambda) + \kappa^T (A - \lambda)x + \sigma(1 - x^T x) \quad (6)$$

where  $\sigma$  and  $\kappa$  are the Lagrange multipliers, for eigenvalues and eigenvectors, respectively <sup>3</sup>. Direct differentiation of equation.6 with respect to the design parameters

<sup>3</sup> We will choose the Lagrange multipliers in a clever way to cancel the annoying derivative terms later.

$$\frac{dL}{dp} = \frac{\partial g}{\partial p} + \frac{\partial g}{\partial x} \frac{dx}{dp} + \frac{\partial g}{\partial \lambda} \frac{d\lambda}{dp} + \kappa^T \left( \frac{dA}{dp} - \frac{d\lambda}{dp} \right) x + \kappa^T (A - \lambda) \frac{dx}{dp} - 2\sigma \left( x^T \frac{dx}{dp} \right) \quad (7)$$

Collecting terms in  $\frac{d\lambda}{dp}$  and  $\frac{dx}{dp}$

$$\frac{dL}{dp} = \frac{\partial g}{\partial p} + \left( \frac{\partial g}{\partial x} + \kappa^T (A - \lambda) - 2\sigma x^T \right) \frac{dx}{dp} + \left( \frac{\partial g}{\partial \lambda} - \kappa^T x \right) \frac{d\lambda}{dp} + \kappa^T \frac{dA}{dp} x \quad (8)$$

Requiring  $\kappa$  and  $\sigma$  to satisfy the following equations,

$$\kappa^T x = \frac{\partial g}{\partial \lambda} \quad (9)$$

$$(A - \lambda)\kappa - 2\sigma x = - \left( \frac{\partial g}{\partial x} \right)^T \quad (10)$$

causes the implicit design variation of the Lagrangian to vanish. The second equation is problematic:  $A - \lambda$  is singular, with a null space of  $x$ . So, we also require  $\sigma$  to make the solution orthogonal to the the

null space:  $x^T \left( \left( \frac{\partial g}{\partial x} \right)^T - 2\sigma x \right) = 0$ . Then,  $\sigma = \frac{x^T \left( \frac{\partial g}{\partial x} \right)^T}{2}$