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MILANO 1863

MSC. MUSIC AND ACOUSTIC ENGINEERING

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## **H6 - Design of a Recorder Flute**

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# 1 Resonator

## 1.1 Fixing the shape of the bore

The resonances of the instrument occur at the zeros of the series  $Z_m + Z_p$  of the impedances of the mouth and the resonator pipe. The pipe is a finite cone tapering towards the foot; we can write its input impedance neglecting the radiation load:

$$Z_p = \frac{j\rho c}{S_1} \frac{\sin kL \sin k\theta_1}{\sin k(L + \theta_1)}$$

where  $L = 0.45$  m is the length of the bore,  $S_1$  is the bore cross section at the mouth,  $x_1$  is the distance between the apex of the cone and  $S_1$ , and  $k\theta_1 = \arctan kx_1$ . The mouth impedance can be written as a pure inertance term  $j\omega M$ , where  $M = \rho l/S_m$  is determined by fixing arbitrarily the surface area and the depth of the mouth opening. Here we choose  $S = 0.65$  cm<sup>2</sup> and  $l = 3$  mm. The resonance condition becomes:

$$\frac{\rho}{S_1} \sin kL \sin k\theta_1 + kM \sin k(L + \theta_1) = 0.$$

Fixing  $k = 2\pi f_0/c$ , with  $f_0 = 392$  Hz, and solving this for  $x_1$ , we find  $x_1 = 99.9$  cm, which corresponds to a diameter at the resonator head  $d_1 = 26.2$  mm. The corresponding values at the foot are  $x_2 = 54.9$  cm and  $d_2 = 14.4$  mm.

## 1.2 Finger holes

If we ignore the tapering of the bore, we can model it as a cylinder with impedance  $Z = jZ_0 \tan kL'$ , where  $L'$  is the acoustic length of the pipe, i. e. its geometrical length plus the end correction at the foot. When adding a finger hole, it can be shown that the acoustic length of the bore is reduced by an amount:

$$\delta \simeq D + \Delta - \Delta'$$

where  $D$  is the distance from the foot to the finger hole,  $\Delta = 0.3d_2$  is the end correction at the foot opening and  $\Delta'$  is the end correction relative to the parallel between the finger hole and the terminating segment of length  $D$ . The latter correction can be written as:

$$\Delta' \simeq \frac{S_2 l (D + \Delta)}{S_h (D + \Delta) + S_2 l}$$

where  $S_2$  is the cross section of the bore at the foot and  $S_h$  is the cross section of the finger hole.

In this picture, we can also include an end correction at the mouth in the total acoustic length of the instrument: in this way, the resonances will occur at the zeros of  $Z_1 = jZ_0 \tan k(L' + \Delta L - \delta)$ . The end correction at the mouth can be recovered by imposing

$$j \frac{\rho c}{S_1} \tan k\Delta L \approx j \frac{\rho c}{S_1} k\Delta L = jkcM.$$

By choosing  $k = 2\pi f_1/c$ , with  $f_1 = 392$  Hz, we can solve the equation  $k(L' + \Delta L - \delta) = \pi$ , recovering  $D = 6.95$  cm.

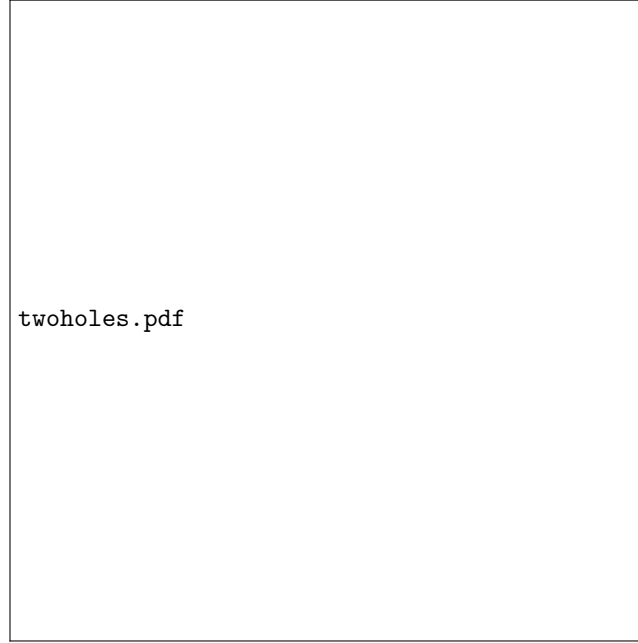


Figure 1: Electrical analog of the load impedance with two open holes. The subscripts  $h_1$  and  $p_1$  refer to the first hole and the section of the bore between the open holes, while  $h_2$  and  $p_2$  refer to the second hole and the terminating section of the pipe. The reference is the atmospheric pressure, while  $P_{in}$  is the pressure at the end of the closed section of the tube.

An alternative and simpler approach could be choosing  $S_h = S_2$ . This yields

$$\delta \approx D + \frac{\Delta^2}{D + 2\Delta}$$

. The resulting position of the hole is  $D = 4.12$  cm. This decrease in the distance from the foot reflects the increase in impedance of the finger hole.

The situation is analogous when there are multiple finger holes, albeit a little more complex. The radiation load will at the end of the closed section of the tube will now include four components, arranged as in Fig. ???. The resulting load impedance will be  $Z_r = Z_{h_1} \parallel (Z_{p_1} + Z_{h_2} \parallel Z_{p_2})$ .

## 2 Flue channel and mouth

### 2.1 Channel thickness

Given the pressure difference  $\Delta p$  between the player's mouth and the flue channel entrance, the flow velocity  $U_j$  in the flue channel can be easily recovered from the Bernoulli equation:

$$\frac{1}{2}\rho_0(v_2^2 - v_1^2) = p_2 - p_1 \quad \Rightarrow \quad U_j = \sqrt{\frac{2\Delta p}{\rho_0}}$$

where  $v_2 = U_j$  is the jet velocity in the channel,  $v_1 = 0$  is the air velocity in the player's mouth (assumed negligible) and  $p_1$  and  $p_2$  are the corresponding pressures.

Now, we know that the amplification of the jet perturbation is strongly dependent on the frequency of the acoustic field, and that it is strongest for frequencies around  $0.3U_j/h$ . This allows us to choose  $h$  according to the desired spectral characteristics of the instrument: we can “tune” the channel thickness to maximize the perturbation amplification at a target spectral centroid  $f_c$ , choosing  $h = 0.3U_j/f_c$ .

In our particular case we can compute:

$$\Delta p = 62 \text{ Pa}, \quad f_c = 2 \text{ kHz}, \quad \rho_0 = 1.2 \text{ kg m}^{-3} \quad \longrightarrow \quad \boxed{h = 1.5 \text{ mm}}$$

The structure of the jet can be characterized using the Reynolds number:

$$\text{Re} = \frac{U_j h}{\nu} = 1033.3$$

where  $\nu = 1.5 \cdot 10^{-5} \text{ m}^2 \text{ s}^{-1}$  is the kinematic viscosity of the air.

## 2.2 Boundary layer effects

Of course the result for the flow velocity we found above is only valid at the center of the flow. Indeed, even if the fluid can be safely assumed to be frictionless in open space, the boundary conditions in a channel make the viscosity effects significant near the walls. There will be a thin layer extending from the boundary into the fluid where the velocity increases rapidly from zero to the value  $U_j$ : this is the so-called *boundary layer*. The thickness of the boundary layer increases along the length of the channel: at position  $x$  along the channel it is:

$$\delta(x) \approx \sqrt{\frac{\nu x}{U_j}}.$$

This means that for a channel length of 20 mm we get  $\delta = 0.54 \text{ mm}$ . Notice that  $\delta/h = 0.36 < \frac{1}{2}$ , so we can expect the Bernoulli approximation to hold reasonably well in the center of the channel.