

Lezione 12 - Problema Elastico

$$\begin{cases} \left(\frac{\partial \sigma_x}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} \right) = \gamma_x \\ - \left(\frac{\partial \sigma_y}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} \right) = \gamma_y \end{cases}$$

$$\begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{bmatrix} = \begin{bmatrix} 1/E & -\nu/E & 0 \\ -\nu/E & 1/E & 0 \\ 0 & 0 & G \end{bmatrix} \begin{bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{bmatrix} \quad \left. \vphantom{\begin{bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{bmatrix}} \right\} \begin{array}{l} \text{Legge di} \\ \text{Hooke} \end{array}$$

$$\epsilon_x = \frac{\partial u}{\partial x}$$

$$\epsilon_y = \frac{\partial v}{\partial y}$$

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$

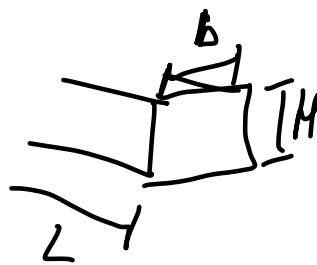
$$\frac{\partial^2 \epsilon_x}{\partial y^2} + \frac{\partial^2 \epsilon_y}{\partial x^2} = \frac{\partial^2 \gamma_{xy}}{\partial x \partial y}$$

$$\frac{\partial^2 \epsilon_x}{\partial y \partial z} = \frac{1}{2} \frac{\partial}{\partial x} \left(\frac{\partial \gamma_{xy}}{\partial z} + \frac{\partial \gamma_{xz}}{\partial y} + \frac{\partial \gamma_{zy}}{\partial x} \right)$$

Come si risolve problema elastico?

Ipotesi:

Trovi snelle : $l \gg b$ e $l \gg h$

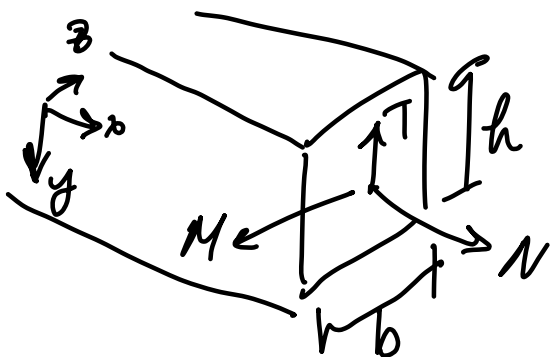


- Materiale Omogeneo, Isotropo, Elastico lineare

→ In campo lineare:
 $\sigma = E \epsilon$

- Forze di Volume nulle
- Forze di Superficie sulle basi
- Aggiungiamo lontananza dei carichi sufficientemente
- Corpo non vincolato, ma in equilibrio
- Sistema di riferimento principale di inerzia

Consideriamo una trave caricata solo su x e y :



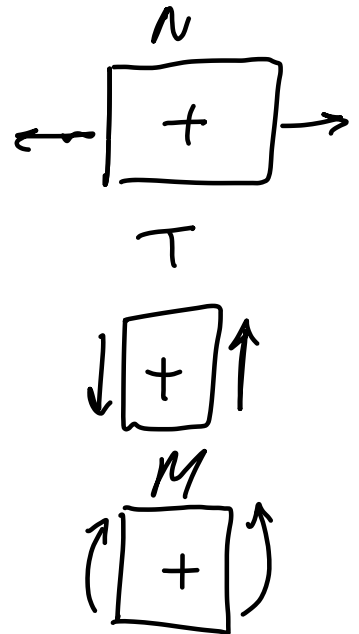
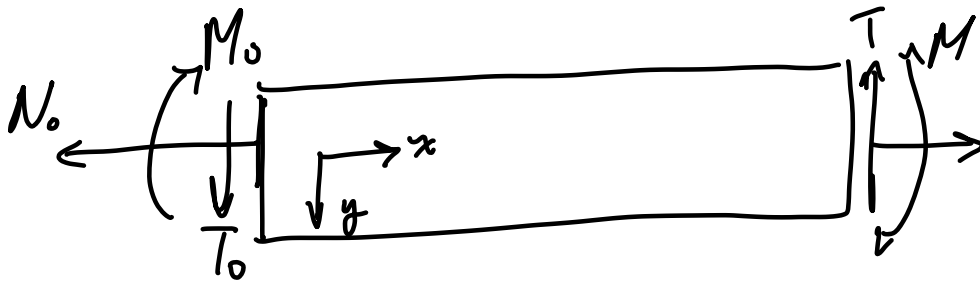
Usando il principio di
 San BERNARD:

$$N = b \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_x dy \quad \left. \begin{array}{l} \text{Integrale} \\ \text{della} \\ \text{stress su} \end{array} \right\}$$

superficie

$$T = b \int_{-h/2}^{h/2} \tau_{xy} dy$$

$$M = b \int_{-h/2}^{h/2} y \sigma_x dy =$$



$$(1) \begin{cases} N = N_0 \\ T = T_0 \\ M = M_0 - T_0 x \end{cases}$$

Saltiamo di passaggio
con San Venant:

Troniamo queste equazioni importanti

$$\sigma_x = a + a_2 y - x b_2 y$$

$$(2) \begin{cases} N = a A \\ M = (a_2 - x b_2) J_{zz} \end{cases}$$

Momento d'inerzia
↓
Inerzia per la rotazione

$$a = \frac{N}{A} = \frac{N_0}{A}$$

$$a_2 = \frac{M_0}{J_{zz}}$$

$$b_2 = \frac{T_0}{J_{zz}}$$

$$\rightarrow \sigma_x = \frac{N}{A} + \frac{M_0 y}{J_{zz}} - x \frac{T_0}{J_{zz}} \cdot y =$$

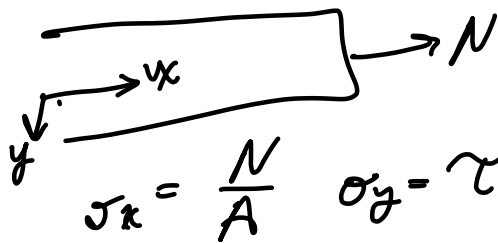
$$M_0 = M + T_0 \cdot x$$

$$\rightarrow \sigma_x = \frac{N}{A} + \frac{M}{J_{zz}} \cdot y + x \frac{T_0}{J_{zz}} y - x \frac{T_0}{J_{zz}} y$$

$$\sigma_x = \frac{N}{A} + \frac{M}{J_{zz}} y$$

linea elastica Assiale

Prendiamo una trave e sotto poniamo a N



$$\left\{ \begin{array}{l} \epsilon_x = \frac{\sigma_x}{E} \\ \epsilon_y = -\nu \frac{\sigma_x}{E} \\ \gamma_{xy} = 0 \end{array} \right. \quad \begin{array}{l} \text{Coefficiente} \\ \text{di Poisson} \end{array}$$

$$\epsilon_x = \frac{\partial u}{\partial x} \Rightarrow u(x, y) = \frac{N}{EA} x + C_1 \quad \leftarrow \text{costante di integrazione}$$

$$\frac{\partial u}{\partial x} = \frac{N}{EA}$$

$$\epsilon_y = \frac{\partial v}{\partial y} \Rightarrow v(x, y) = -\nu \frac{N}{EA} y + C_2$$

$$\left. \begin{aligned} u(x,0) &= \frac{N}{EA} + C_1 \\ v(x,0) &= C_2 \end{aligned} \right\} \text{In questo caso}$$

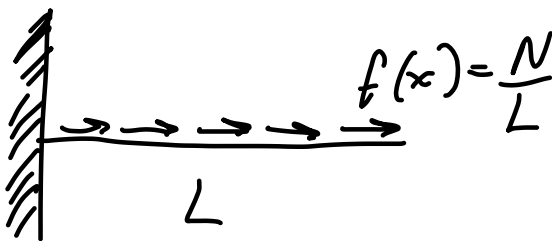
σ_x e E_x uniformi sulla sezione

→ Consideriamo un carico distribuito

$$-\frac{dN}{dx} = f(x)$$

$$N = \sigma_x A = E \cdot A \cdot E_x = E \cdot A \cdot \frac{\partial u}{\partial x} \Rightarrow EA \frac{\partial^2 u}{\partial x^2} = f(x)$$

Esempio



$$\frac{\partial^2 u}{\partial x^2} \cdot EA = f(x) = \frac{N}{L} \quad \left(\frac{\partial u}{\partial x} = -\frac{N}{EAL} x + C_1 \right)$$

$$\rightarrow u = -\frac{N}{EAL} \frac{x^2}{2} + C_1 x + C_2$$

Condizioni al contorno

$$\begin{cases} u(0) = 0 \\ N(L) = 0 \end{cases}$$

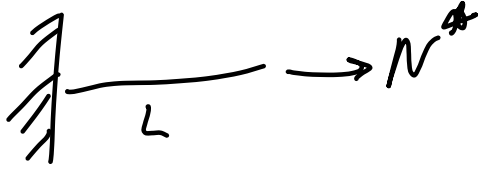
$$u(0) = C_2 = 0$$

$$N(x) = EA E_x = EA \frac{\partial u}{\partial x}$$

$$EA \left(-\frac{N}{EAL} x + C_1 \right) \Big|_{x=L} = 0$$

$$\Rightarrow -\frac{N}{EA} + C_1 = 0 \Rightarrow C_1 = \frac{N}{EA}$$

Altri Esempi



$$-EA \frac{\partial^2 u}{\partial x^2} = 0$$

$$\frac{\partial u}{\partial x} = C_1$$

$$u = C_1 x + C_2$$

Condizioni al Contorno

$$\begin{cases} u(0) = 0 \rightarrow C_2 = 0 \\ N(0) = N^* \end{cases}$$

$$EA \left. \frac{\partial u}{\partial x} \right|_{x=0} = N^*$$

$$EA \cdot C_1 = N^* \rightarrow C_1 = \frac{N^*}{EA}$$

$$u(x) = \frac{N^*}{EA} x$$

Da ricordare

Ipotesi di Saint Venant

$-EA \frac{\partial^2 u}{\partial x^2} = f(x)$ e procedeva per i due esempi
sugli unici
due esempi.

Saint Venant \rightarrow Saint-Venant