

# Numerical Solution with Spectral Method for Smoluchowski Equation on $S^2$

Bella My Phuong Quynh Duong

Heinrich-Heine University Düsseldorf

September 26, 2023



# Table of Contents

## 1 Introduction

- Motivation

## 2 Spectral method

## 3 Numerical Results

- Shear flow on  $S^2$ 
  - Stability analysis
- Elongational flow
  - Stability analysis

## 1 Introduction

- Motivation

## 2 Spectral method

## 3 Numerical Results

- Shear flow on  $S^2$ 
  - Stability analysis
- Elongational flow
  - Stability analysis

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.

## General about the project

## First step

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$

## General about the project

## First step

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$
- Construction a lower-dimensional model by using a hierarchy of moment equations

## General about the project

## First step

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$
- Construction a lower-dimensional model by using a hierarchy of moment equations
- The moment equations are a hyperbolic system with a source term

$$\partial_t Q(x, t) + A \partial_x Q(x, t) = \phi(Q(x, t))$$

The update of the source term is equivalent to a spectral method ([1]).

## General about the project

## First step

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$
- Construction a lower-dimensional model by using a hierarchy of moment equations
- The moment equations are a hyperbolic system with a source term

$$\partial_t Q(x, t) + A \partial_x Q(x, t) = \phi(Q(x, t))$$

The update of the source term is equivalent to a spectral method ([1]).

## General about the project

- Study the coupled system on  $S^2$ , which is a five dimensional problem (plus time)

## First step



## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$
- Construction a lower-dimensional model by using a hierarchy of moment equations
- The moment equations are a hyperbolic system with a source term

$$\partial_t Q(x, t) + A \partial_x Q(x, t) = \phi(Q(x, t))$$

The update of the source term is equivalent to a spectral method ([1]).

## General about the project

- Study the coupled system on  $S^2$ , which is a five dimensional problem (plus time)
- Reduction of the higher-dimensional model by a three-dimensional system of partial differential equations

## First step

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$
- Construction a lower-dimensional model by using a hierarchy of moment equations
- The moment equations are a hyperbolic system with a source term

$$\partial_t Q(x, t) + A \partial_x Q(x, t) = \phi(Q(x, t))$$

The update of the source term is equivalent to a spectral method ([1]).

## General about the project

- Study the coupled system on  $S^2$ , which is a five dimensional problem (plus time)
- Reduction of the higher-dimensional model by a three-dimensional system of partial differential equations

## First step

- Derive a spectral Method for the Smoluchowski Equation on  $S^2$

## In the earlier work

- Focus on a system of partial differential equations which couples a kinetic equation to a macroscopic flow equation. This system model (given by Helzel & Tzavaras) a sedimentation process in suspension of rigid rod-like particles.
- Study the coupled system on  $S^1$
- Construction a lower-dimensional model by using a hierarchy of moment equations
- The moment equations are a hyperbolic system with a source term

$$\partial_t Q(x, t) + A \partial_x Q(x, t) = \phi(Q(x, t))$$

The update of the source term is equivalent to a spectral method ([1]).

## General about the project

- Study the coupled system on  $S^2$ , which is a five dimensional problem (plus time)
- Reduction of the higher-dimensional model by a three-dimensional system of partial differential equations

## First step

- Derive a spectral Method for the Smoluchowski Equation on  $S^2$
-

## 1 Introduction

- Motivation

## 2 Spectral method

## 3 Numerical Results

- Shear flow on  $S^2$ 
  - Stability analysis
- Elongational flow
  - Stability analysis

# Mathematical Model for the Sedimentation of Rod-Like Particles [4]

Coupling of a Smoluchowski Equation and a Navier-Stokes Equation

$$\begin{aligned}
 \partial_t f + \nabla_x \cdot (\mathbf{u} f) + \nabla_n \cdot (P_n \perp \nabla_x \mathbf{u} n f) - \nabla_x \cdot ((I + \mathbf{n} \otimes \mathbf{n}) e_3 f) \\
 = D_r \Delta_n f + \gamma \nabla_x \cdot (I + \mathbf{n} \otimes \mathbf{n}) \nabla_x f, \\
 \sigma = \int_{S^{d-1}} (d \mathbf{n} \otimes \mathbf{n} - I) f d\mathbf{n}, \\
 \operatorname{Re} (\partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla_x) \mathbf{u}) = \Delta_x \mathbf{u} - \nabla_x p + \delta \gamma \nabla_x \cdot \sigma - \delta \int_{S^{d-1}} f d\mathbf{n} e_3, \\
 \nabla_x \cdot \mathbf{u} = 0,
 \end{aligned}$$

where  $f = f(t; x; n)$ ,  $x \in \mathbb{R}^d$ ,  $n \in S^{d-1}$ ,  $t \in \mathbb{R}$  is a density distribution function of particle orientation.  $D_r, \gamma, \delta$  and  $\operatorname{Re}$  are non-dimensional values.

Consider

$$\begin{aligned} \partial_t f + \nabla_x \cdot (uf) + \nabla_n \cdot (P_{n\perp} \nabla_x u f) - \nabla_x \cdot ((I + n \otimes n) e_3 f) \\ = D_r \Delta_n f + \gamma \nabla_x \cdot (I + n \otimes n) \nabla_x f. \end{aligned} \quad (1)$$

Rewrite the equation (1) in spherical coordinates ([2], [3]) and it follows

$$\sin \theta \partial_t f + \partial_\phi (a(\phi, \theta) f) + \partial_\theta (b(\phi, \theta) f) = D_r \left( \partial_\phi \left( \frac{1}{\sin \theta} \partial_\phi f \right) + \partial_\theta (\sin \theta \partial_\theta f) \right), \quad (2)$$

with  $\phi \in [0, 2 \cdot \pi]$  and  $\theta \in [0, \pi]$ . Solving the Smoluchowski equation (2) on  $S^2$  by using a [spectral method](#), which is based on the ansatz

$$f(\phi, \theta, t) = f_0(t) \cdot P_0^0 + \sum_{n=1}^{\infty} \sum_{i=-n}^n c_{2n}^i(t) \cdot P_{2n}^i(\phi, \theta), \quad (3)$$

where  $P_{2n}^i(\phi, \theta)$  are harmonic polynomial basis functions.

# Harmonic polynomial basis functions

Let  $P_n^i(\phi, \theta)$  with  $n = 0, \dots, \infty$  and  $i = n, \dots, -n$  be the basis function

*TODO*

The scalar product of any two basis functions over sphere is defined as follows

$$\langle P_n^i, P_m^l \rangle_{S^2} = \int_0^{2\pi} \int_0^\pi P_n^i(\phi, \theta) \cdot P_m^l(\phi, \theta) \cdot \sin(\theta) d\theta d\phi.$$

# Properties of harmonic polynomial basis functions

## Property 1

Let  $P_n^i(\phi, \theta)$  and  $P_m^l(\phi, \theta)$  are two different harmonic polynomial basis functions. Then

$$\langle P_n^i, P_m^l \rangle_{S^2} = 0,$$

for  $i \neq l$  or  $n \neq m$ .

## Property 2

Let  $P_n^i$  be the normalized harmonic polynomial basis functions. Then

$$\langle P_n^i, P_n^i \rangle_{S^2} = 1.$$

## Property 3

The spherical harmonic function are the eigenfunctions of Laplace-Beltrami operator with the eigenvalues  $(-n(n+1), n \in \mathbb{N}_0)$  (see [5])

$$\Delta_{S^2} P_n^i = -n(n+1)P_n^i.$$



Recall the Smochluchowski equation on  $S^2$

$$\underbrace{\sin \theta \partial_t f}_{(1)} + \underbrace{\partial_\phi (a(\phi, \theta) f) + \partial_\theta (b(\phi, \theta) f)}_{(2)} = D_r \underbrace{\left( \partial_\phi \left( \frac{1}{\sin \theta} \partial_\phi f \right) + \partial_\theta (\sin \theta \partial_\theta f) \right)}_{(3)} \quad (4)$$

and the ansatz

$$f(\phi, \theta, t) = f_0(t) \cdot P_0^0 + \sum_{n=1}^{\infty} \sum_{i=-n}^n c_{2n}^i(t) \cdot P_{2n}^i(\phi, \theta). \quad (5)$$

The Laplace Beltrami operator ([3]) on the unit sphere  $S^2$  is given by

$$\Delta_{S^2} f = \frac{1}{\sin^2 \theta} \partial_{\phi\phi} f + \frac{1}{\sin \theta} \partial_{\theta} (\sin \theta \partial_{\theta} f). \quad (6)$$

From property 3 and the equation (6), it follows for the term (3)

$$\Delta_{S^2} P_{2n}^i = -n(n+1)P_{2n}^i.$$

For the term (1) and (2), we inserting (5) in (4), multiplying with each basis function and integrate it over  $S^2$ . We derive a system of ODEs for the coefficients

$$\begin{pmatrix} f_0 \\ c_2^{-2} \\ \vdots \\ c_{2n}^i \end{pmatrix}' = A \begin{pmatrix} f_0 \\ c_2^{-2} \\ \vdots \\ c_{2n}^i \end{pmatrix}, \quad (7)$$

with  $A \in \mathbb{R}^{cn \times cn}$

$$cn = \begin{cases} \text{order} = 2 : & cn = 2 \cdot \text{order} + 2 \\ \text{order} = \text{even} : & c2 = 6 \\ & cn = c2 \\ & cn = c2 + \sum_{i=4}^{\text{order}} (2i + 1) \end{cases}$$

## Example: Shear flow

## Example: Shear flow

Consider the Smoluchowski equation (4) with the velocity gradient

$$\vec{u} = \begin{pmatrix} u(x, y, z) \\ v(x, y, z) \\ w(x, y, z) \end{pmatrix}, \nabla_x \vec{u}_{\text{ext}} = \begin{pmatrix} u_x & u_y & u_z \\ v_x & v_y & v_z \\ w_x & w_y & w_z \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

With the given velocity gradient it follows

$$\begin{aligned} \partial_t (\sin \theta f) + \partial_\theta (\sin \phi \cos \phi \sin^2 \theta \cos \theta f) + \partial_\phi (-\sin \theta \sin^2 \phi f) \\ = D_r \left( \partial_\theta (\sin \theta \partial_\theta f) + \partial_\phi \left( \frac{1}{\sin \theta} \partial_\phi f \right) \right). \end{aligned} \quad (8)$$

Consider the ansatz with the zeroth order

$$f(\phi, \theta, t) = f_0(t) \cdot P_0^0. \quad (9)$$

Insert the ansatz (9) in (8)

$$\partial_t (f_0(t) \cdot P_0^0) + \frac{1}{\sin \theta} (\partial_\theta (\dots) + \partial_\phi (\dots)) = \frac{1}{\sin \theta} D_r (\dots). \quad (10)$$

We know

$$\Delta_{S^2} P_0^0(\phi, \theta) = \frac{1}{\sin \theta} D_r(\dots) = \lambda_{2n,i} \cdot P_0^0(\phi, \theta),$$

where  $\lambda_{2n,i}$  is the corresponding eigenvalue.

Since  $P_0^0(\phi, \theta) = 1$  does not depend on  $\phi$  and  $\theta$ , the partial derivatives will be zero

$$\Delta_{S^2} P_0^0(\phi, \theta) = 0. \quad (11)$$

Consider the rest of the equation (10)

$$\underbrace{\partial_t(f_0(t) \cdot P_0^0)}_{(1)} + \underbrace{\frac{1}{\sin \theta} (\partial_\theta(\sin \phi \cos \phi \sin^2 \theta \cos \theta \cdot f_0(t) \cdot P_0^0) + \partial_\phi(-\sin \theta \sin^2 \phi \cdot f_0(t) \cdot P_0^0))}_{(2)}.$$

It is

$$\partial_t(f_0(t) \cdot P_0^0) = f_0'(t).$$

Let  $z(\phi, \theta) := (2)$

Project the solution  $z(\phi, \theta)$  onto all polynomials to find out which polynomial are needed

$$\begin{aligned}
 \int_0^{2\pi} \int_0^\pi z(\phi, \theta) \cdot P_2^{-2}(\phi, \theta) \sin \theta d\theta d\phi &\stackrel{Maple}{=} 0 \\
 \int_0^{2\pi} \int_0^\pi z(\phi, \theta) \cdot P_2^{-1}(\phi, \theta) \sin \theta d\theta d\phi &\stackrel{Maple}{=} 0 \\
 \int_0^{2\pi} \int_0^\pi z(\phi, \theta) \cdot P_2^0(\phi, \theta) \sin \theta d\theta d\phi &\stackrel{Maple}{=} 0 \\
 \int_0^{2\pi} \int_0^\pi z(\phi, \theta) \cdot P_2^1(\phi, \theta) \sin \theta d\theta d\phi &\stackrel{Maple}{=} 0 \\
 \int_0^{2\pi} \int_0^\pi z(\phi, \theta) \cdot P_2^2(\phi, \theta) \sin \theta d\theta d\phi &\stackrel{Maple}{=} -\frac{\sqrt{15}}{5}
 \end{aligned}$$

It follows

$$f_0(t) \cdot P_0^0 \cdot \frac{1}{\sin \theta} (\partial_\theta (\sin \phi \cos \phi \sin^2 \theta \cos \theta) + \partial_\phi (-\sin \theta \sin^2 \phi)) = f_0(t) \left[ -\frac{\sqrt{15}}{2} P_2^2 \right]. \quad (12)$$

Together we have

$$f_0'(t) - \frac{\sqrt{15}}{2} f_0(t) P_2^2(\phi, \theta) = 0 \cdot P_0^0(\phi, \theta) D_r.$$

For the ansatzfunction with higher order, the calculation is done in the same way.

As an example we obtain an ODE system with ansatzfunction of the 2<sup>nd</sup>. order

$$\begin{pmatrix} f_0'(t) \\ c_2^{-2} \\ c_2^{-1} \\ c_2^0 \\ c_2^1 \\ c_2^2 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -6D_r & 0 & 0 & 0 & 1 \\ 0 & 0 & -6D_r & 0 & 5/7 & 0 \\ 0 & 0 & 0 & -6D_r & 0 & -\frac{\sqrt{3}}{7} \\ 0 & 0 & -2/7 & 0 & -6D_r & 0 \\ \frac{\sqrt{15}}{5} & 1 & 0 & -\frac{\sqrt{3}}{7} & 0 & -6D_r \end{pmatrix} \cdot \begin{pmatrix} f_0 \\ c_2^{-2} \\ c_2^{-1} \\ c_2^0 \\ c_2^1 \\ c_2^2 \end{pmatrix} \quad (13)$$



## 1 Introduction

- Motivation

## 2 Spectral method

## 3 Numerical Results

- Shear flow on  $S^2$ 
  - Stability analysis
- Elongational flow
  - Stability analysis

Shear flow on  $S^2$

# Shear flow on $S^2$

Consider the externally imposed shear flow

$$\mathbf{u} = \begin{pmatrix} u(y) \\ 0 \\ 0 \end{pmatrix}, \quad \nabla_{\mathbf{x}} \mathbf{u} = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad (14)$$

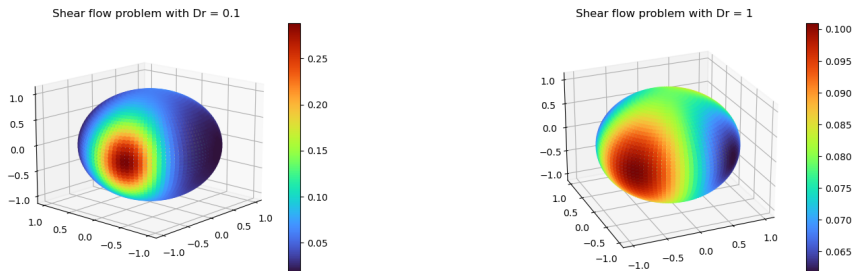


Figure 1: Steady state solution of the Smoluchowski equation approximated at  $T = 10$  using different values of  $D_r$ .

# Shear flow: Stability analysis

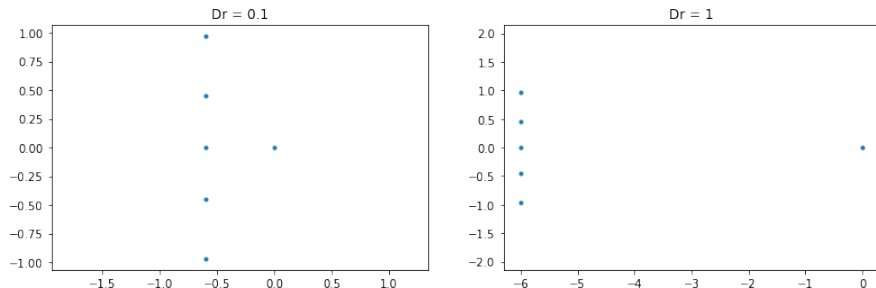


Figure 2: Eigenvalue of matrix  $A$  with basis functions of 2nd. order

## Proposition 1

Spectral method for 2nd. order is stable with for both small and large  $D_r$ .

# Shear flow: Stability analysis

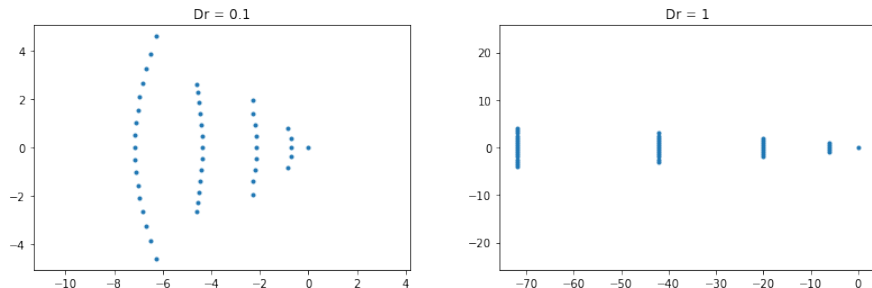


Figure 3: Eigenvalue of matrix  $A$  with basis functions of 8th. order

## Proposition 2

Spectral method for 8th. order is stable with for both small and large  $Dr$ .

## Elongational flow

# Elongational flow

Consider the externally imposed velocity gradient

$$\nabla_{\vec{x}} \vec{u}_{\text{ext}} = \begin{pmatrix} 2 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

The exact steady-state solution has the form

$$f_{\text{exact}}(\phi, \theta) = C_1 \exp\left(-\frac{3}{2D_r} (1 - \cos^2(\phi) \sin^2(\theta))\right), \quad (15)$$

with the constants  $C_1 = 2.30121384511755303190$  for  $D_r = 0.1$ .



Figure 4: Exact steady state solution with different  $D_r$

# Elongational flow

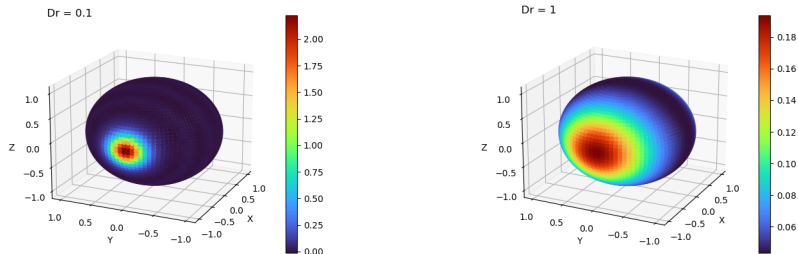


Figure 5: Numerical solution on  $S^2$  with basis functions of 14<sup>th</sup>. order



# Convergence Study

For the convergence study, the maximum norm error is used

$$E_{max} = \max |U_{exact} - U_{approx}|.$$

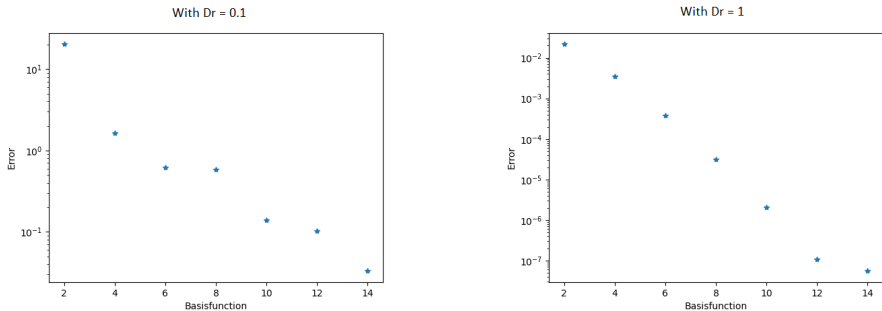
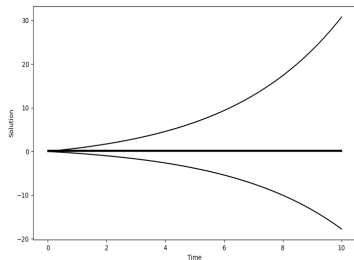
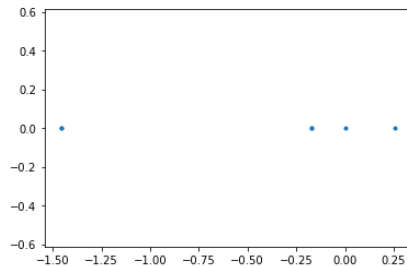


Figure 6: Error with respect to basisfunction with different  $D_r$

# Stability analysis



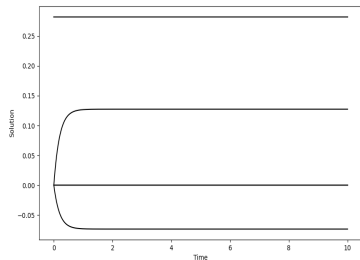
(a) Coefficients as function order



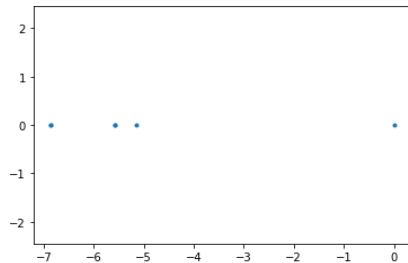
(b) Eigenwert of matrix  $A$  order

Figure 7: With 2nd. order and  $D_r = 0.1$

# Stability analysis



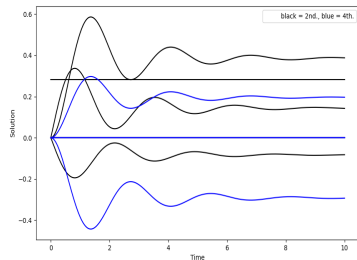
(a) Coefficients as function order



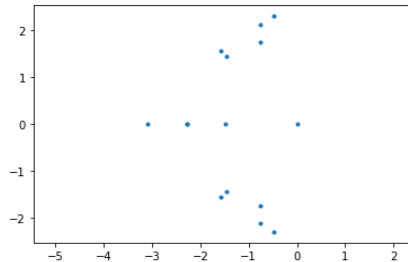
(b) Eigenwert of matrix  $A$  order

Figure 8: With 2nd. order and  $D_r = 1$

# Stability analysis



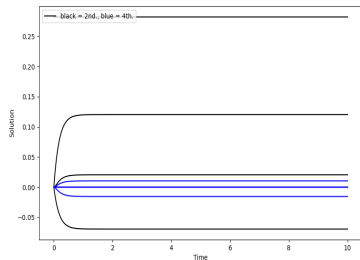
(a) Coefficients as function order



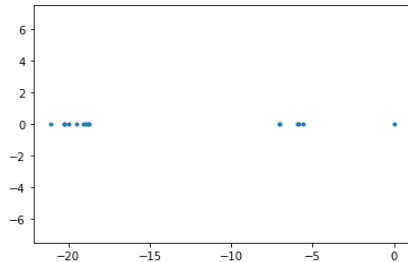
(b) Eigenwert of matrix  $A$  order

Figure 9: With 4th. order and  $D_r = 0.1$

# Stability analysis



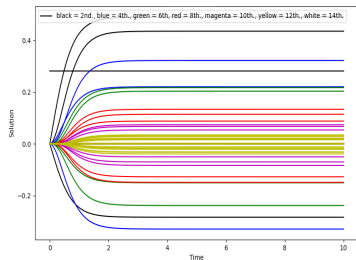
(a) Coefficients as function order



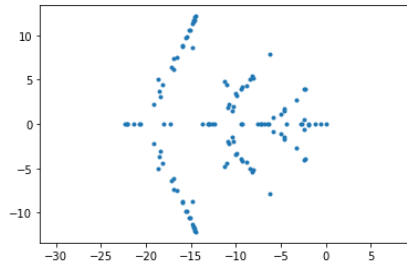
(b) Eigenwert of matrix  $A$  order

Figure 10: With 4th. order and  $D_r = 1$

# Stability analysis



(a) Coefficients as function order



(b) Eigenwert of matrix  $A$  order

Figure 11: With 14th. order and  $D_r = 0.1$

## Finding

- For small  $D_r$  we have seen that the error was relatively large.
- The eigenvalue of the matrix  $A$  of the 2nd. order ODE system for  $D_r = 0.1$  has a positive real part  $\rightarrow$  the ansatzfunction with 2nd. order is unsuitable.
- The spectral method becomes stable when taking more approach functions for small  $D_r$ .

Thank you for listening!





S. Dahm and C. Helzel.

Hyperbolic systems of moment equations describing sedimentation in suspensions of rod-like particles.

*Multiscale Modeling & Simulation*, 20(3):1002–1039, 2022.



C. Helzel and F. Otto.

Multiscale simulations for suspensions of rod-like molecules.

*J. Comput. Phys.*, 216(1):52–75, 2006.



C. Helzel and M. Schneiders.

Numerical approximation of the Smoluchowski equation using radial basis functions.

*J. Comput. Math.*, 38(1):176–194, 2020.



C. Helzel and A. E. Tzavaras.

A kinetic model for the sedimentation of rod-like particles.

*Multiscale Model. Simul.*, 15(1):500–536, 2017.



M. Morimoto.

*Analytic functionals on the sphere*, volume 178 of *Transl. Math. Monogr.*

Providence, RI: American Mathematical Society, 1998.