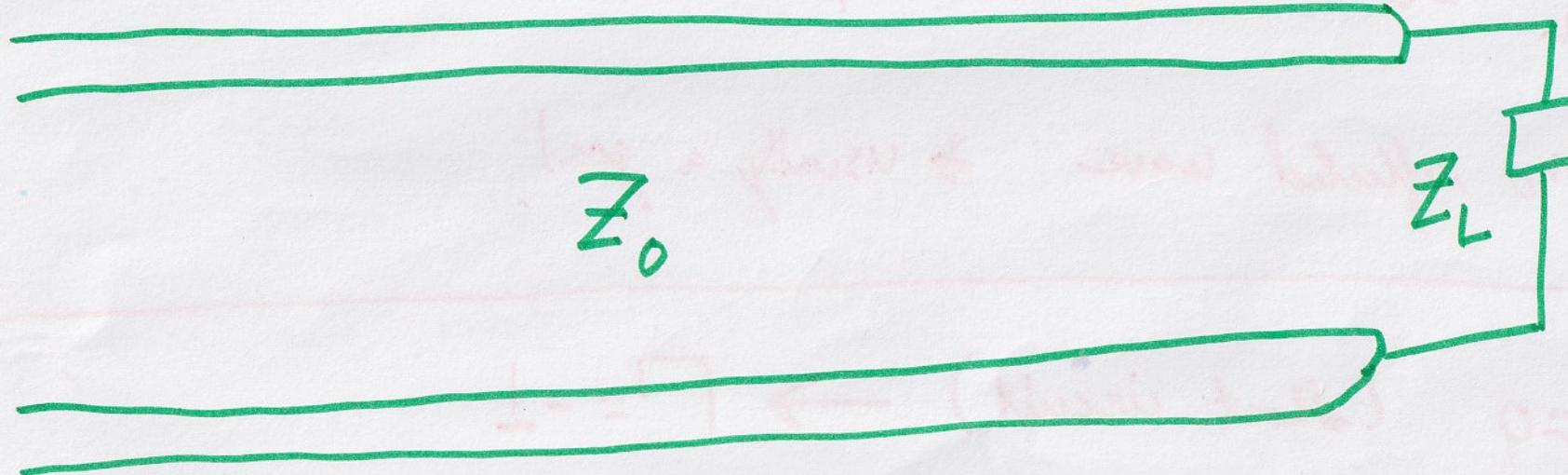


Fields and Waves I

L4

①

Input Impedance of a Transmission Line



Γ = reflection coefficient

$$\Gamma = \frac{|V_{back}|}{|V_{forward}|}$$

Amplitude of the forward and backward waves

② $\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{|V-|}{|V+|}$ + forward going
- backward going

$$Z_L = Z_0$$

$\Gamma = 0$, Matched Load

No reflected wave * usually a goal

$$Z_L = 0 \text{ (short circuit)} \rightarrow \Gamma = -1$$

$$Z_L = \infty \text{ (open circuit)} \rightarrow \Gamma = +1$$

Standing wave pattern

(3)

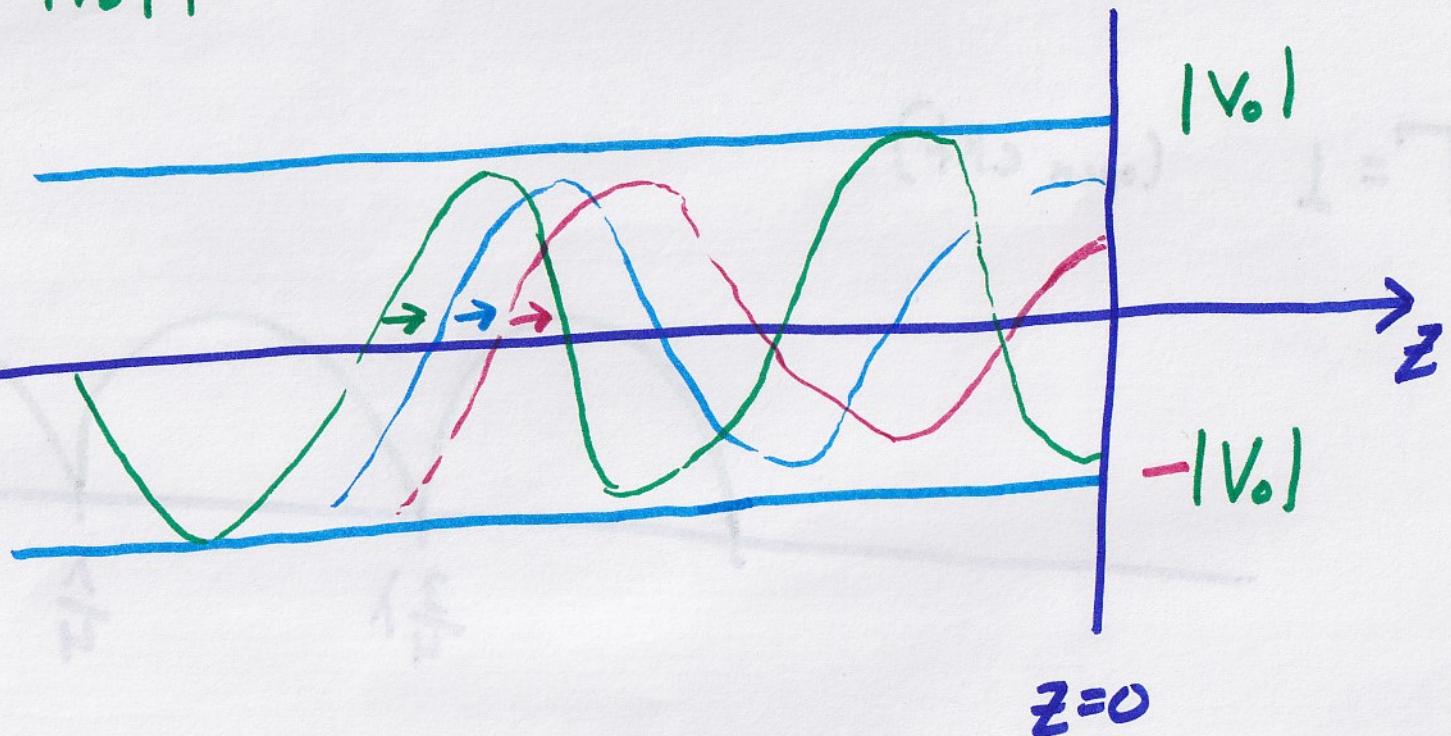
Max amplitude as a function of position

We looked at

(Voltage)

$$|V_{\text{total}}(z)| = |V_0| \left| \frac{1}{1 + \Gamma e^{+2j\beta z}} \right|$$

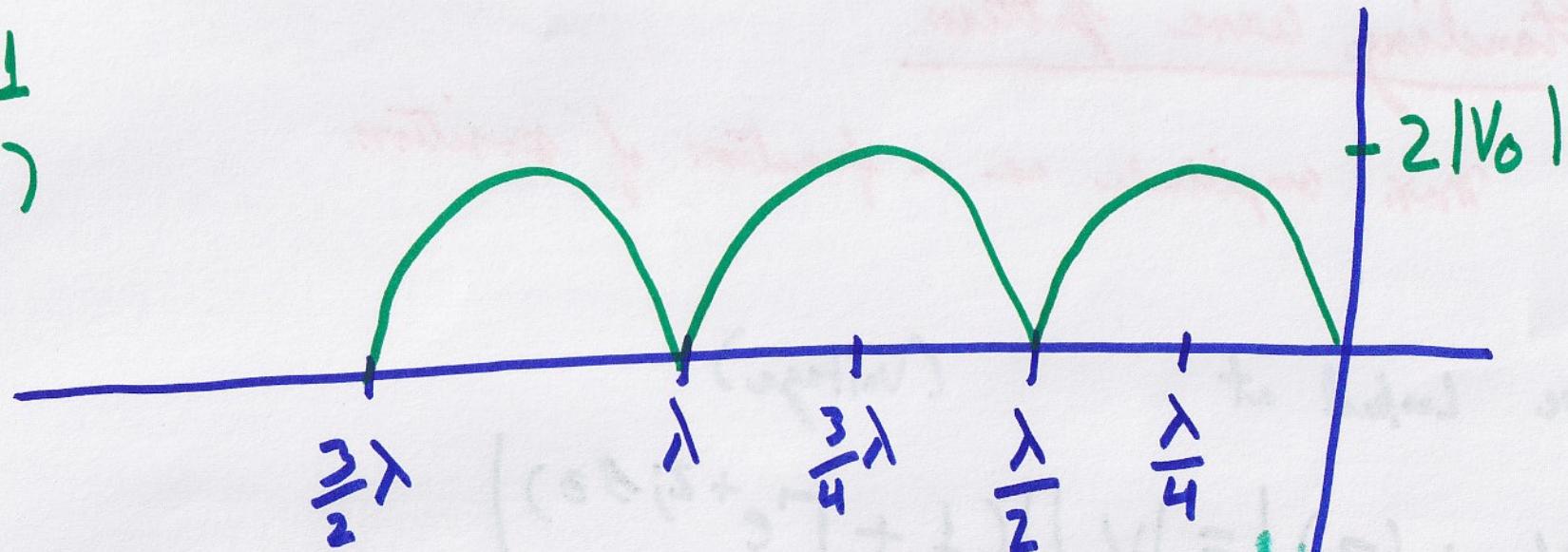
$\Gamma = 0$
(matched)



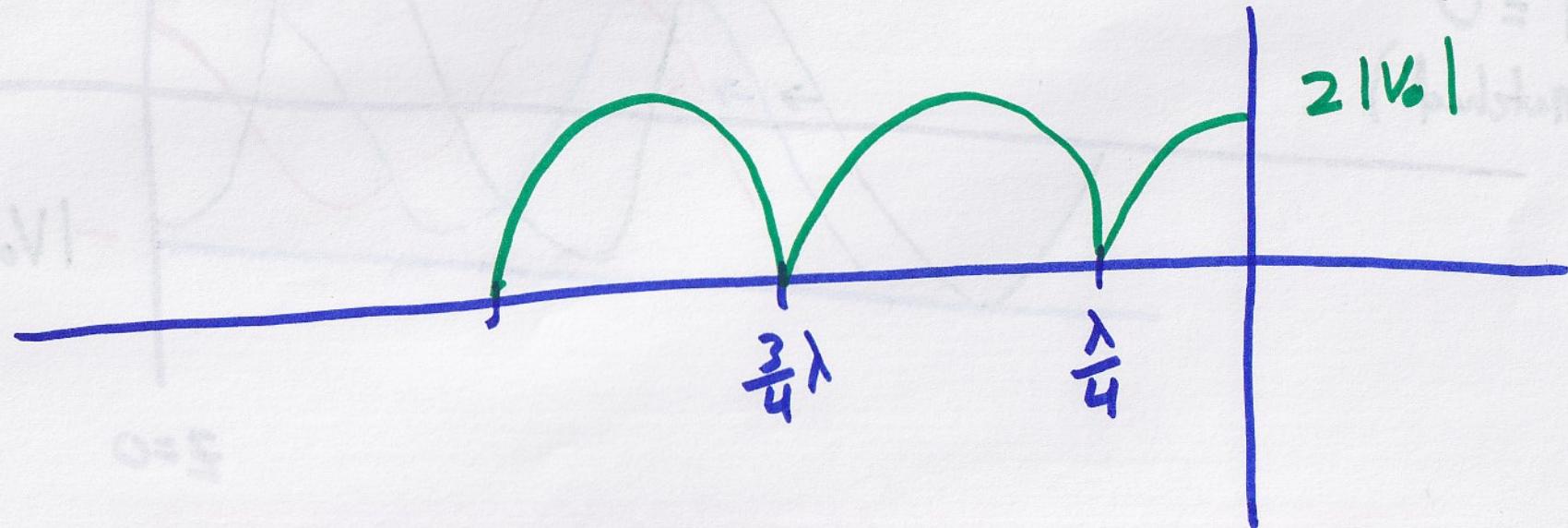
④

$$\Gamma = -1$$

(short)



$$\Gamma = 1 \quad (\text{open ckt})$$

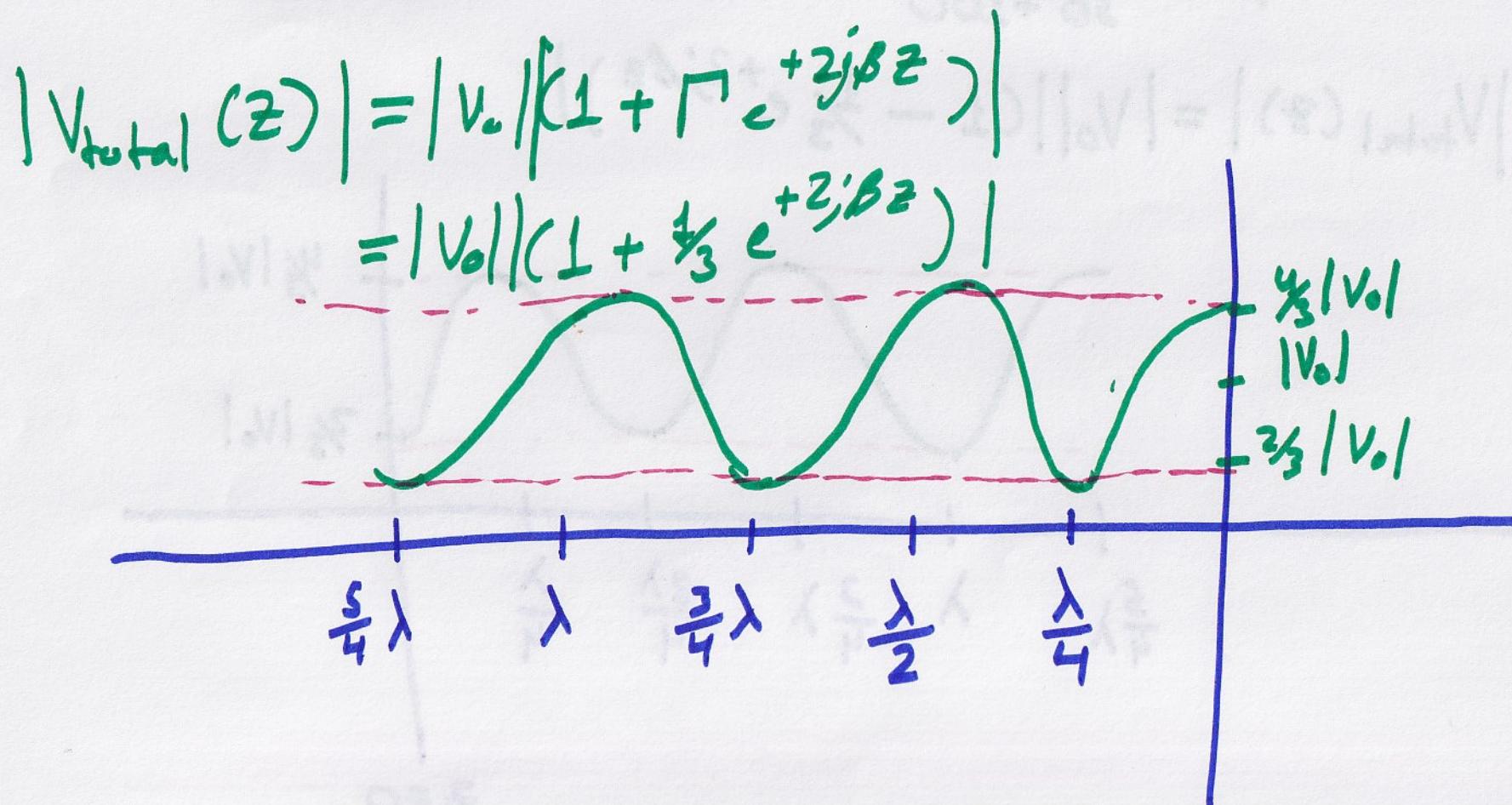


(5)

Example)

$$Z_0 = 100 \Omega \quad Z_L = 200 \Omega$$

$$\Gamma_c = \frac{200 - 100}{200 + 100} = \frac{1}{3}$$



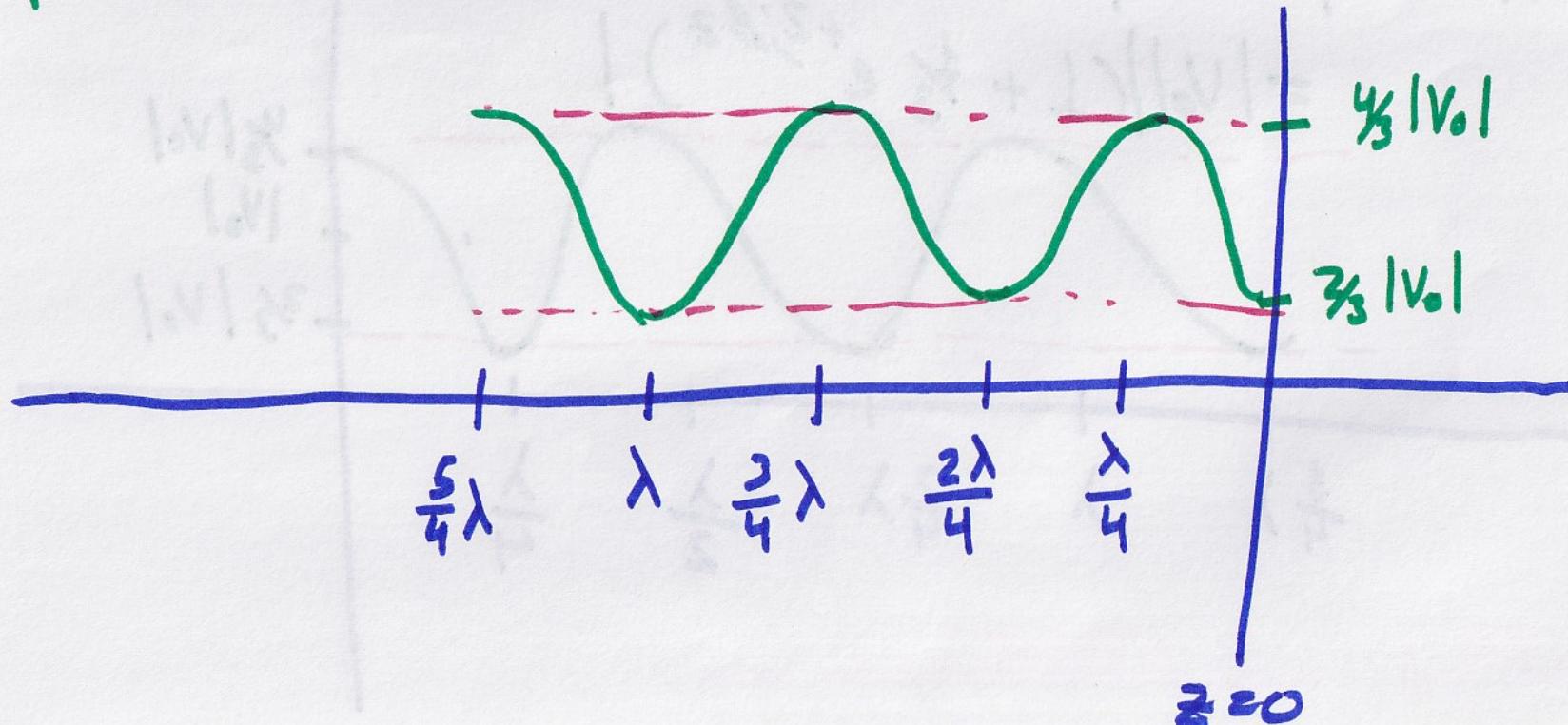
(6)

$V_{\text{total}}(z)$ is periodic with $n \frac{\lambda}{2}$ $n \in \mathbb{N}$

Example) $Z_0 = 100 \Omega$ $Z_L = 50 \Omega$

$$\Gamma = \frac{50 - 100}{50 + 100} = -\frac{1}{3}$$

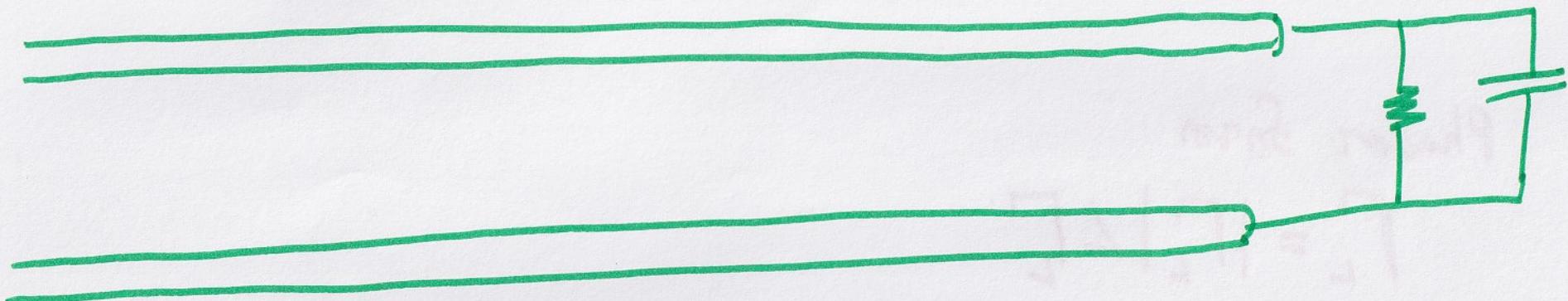
$$|V_{\text{total}}(z)| = |V_0| \left| \left(1 - \frac{1}{3} e^{+2jBz} \right) \right|$$



Position of Maxima / Minima

(7)

depends on the phase of the reflection coefficient



So what's the impedance?

$$\rightarrow R \parallel \frac{1}{j\omega C}$$

$$\rightarrow \left(\frac{1}{R} + j\omega L \right)^{-1} = \frac{1}{\frac{1}{R} + j\omega C} = \frac{R}{1 + j\omega RC} = \frac{R - j\omega R^2 C}{1 + \omega^2 R^2 C^2}$$

$$\frac{1}{Z} \frac{Z^*}{Z^*} = \frac{Z^*}{|Z|^2}$$

(8) Example, $Z_L = 100 - j50$ $Z_0 = 100 \Omega$

$$\Gamma_L = \frac{50 - j50 - 100}{50 - j50 + 100} = -0.2 - j0.4$$

Phasor form

$$\Gamma_L = |\Gamma_L| \angle \Gamma_L$$

$$|\Gamma_L| = (\Gamma_{L\text{ real}}^2 + \Gamma_{L\text{ imaginary}}^2)^{\frac{1}{2}}$$

$$\angle \Gamma_L = \tan^{-1}\left(\frac{\Gamma_{L\text{ imaginary}}}{\Gamma_{L\text{ real}}}\right) \quad (\text{know what quadrant you're in})$$

$$\text{so } \Gamma_L = 0.447 \angle -117^\circ$$

$$= 0.447 \angle -0.65\pi = 0.447 e^{-j0.65\pi}$$

$$|V_{\text{total}}(z)| = |V_0| \left| (1 + 0.447 e^{-j0.65\pi} e^{j\beta z}) \right| \\ = |V_0| \left| (1 + 0.447 e^{j(2\beta z - 0.65\pi)}) \right| \quad (9)$$

Find the Location on where is the maximum
so to start,

$$2\beta z - 0.65\pi = 0, 2\pi, 4\pi, \dots \\ = 2\pi n \quad n \in \mathbb{N}$$

→ Location of Voltage maxima on the T-Line
recall $\beta = \frac{2\pi}{\lambda}$

$$\text{so } 2\left(\frac{2\pi}{\lambda}\right)z - 0.65\pi = 0$$

$$\rightarrow z = 0.65\pi \frac{\lambda}{4\pi} + n \frac{\lambda}{2}$$

(10)

so now we have something called

$\text{SWR} \rightarrow \text{standing wave ratio}$

$$\text{SWR} = \frac{|V_{\max}|}{|V_{\min}|} = \frac{|V_o|(1 + |\Gamma|)}{|V_o|(1 - |\Gamma|)}$$

when the line is matched, $\text{SWR} = 1$

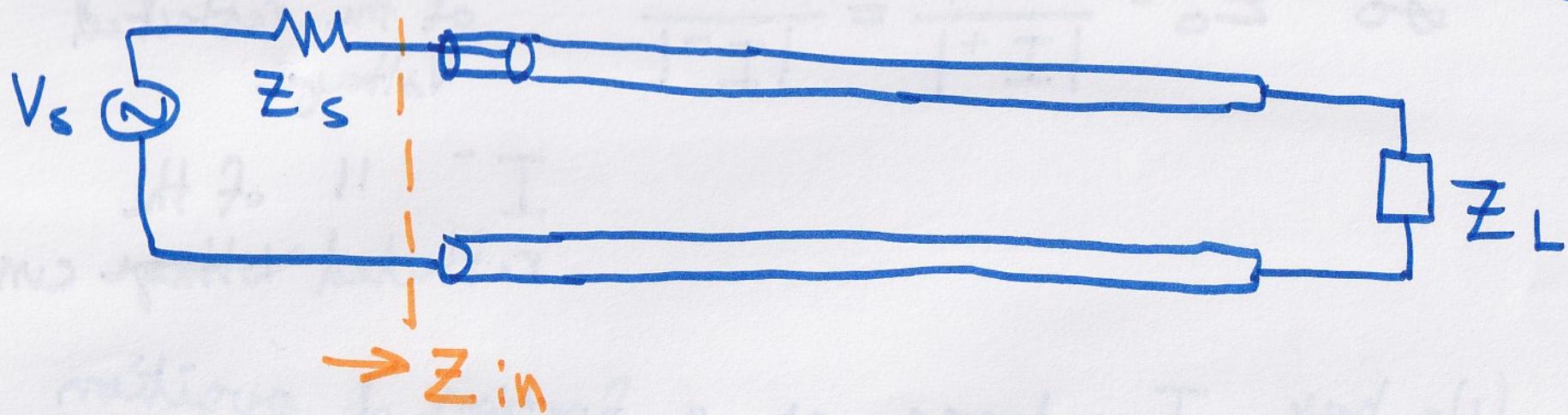
consequently, if we have open/short circuit Load
~~different load~~

$$\text{SWR} = \frac{2}{0} \rightarrow \infty \quad \text{not good}$$

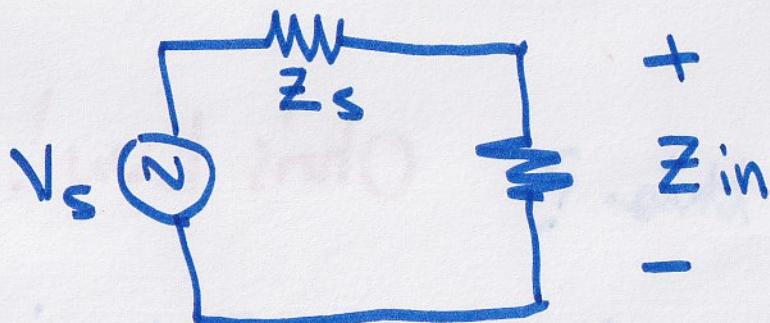
$$1 \leq \text{SWR} < \infty \quad |\Gamma| \leq 1$$

Look at the Circuit

11



Z_{in} is the Input impedance of the T-Line



$$V_{in} = \frac{Z_{in}}{Z_{in} + Z_s} V_s$$

V_{in} is the Total voltage at the input of the T-line

12

$$\text{So } Z_0 = \frac{|V^+|}{|I^+|} = \frac{|V^-|}{|I^-|}$$

V^- is the Amplitude
of the reflected
Voltage

I^- " of the
reflected ~~voltage~~ current

We have Impedance as a function of position

$Z(z)$
↑ position
Impedance

But how do we get this? Ohm's Law!

$$Z(z) = \frac{V_{\text{total}}(z)}{I_{\text{total}}(z)} = \frac{V_0 (1 + \Gamma e^{+2jBz}) e^{-jBz}}{I_0 (1 - \Gamma e^{+2jBz}) e^{-jBz}}$$

Note!

$Z(z)$ is the ratio of the total voltage (incident and reflected wave voltages) to the total current at any point z on the line.

This is in contrast with the characteristic impedance of the line Z_0 , which relates the voltage and current of each of the two waves individually

$$Z_0 = \frac{V_o^+}{I_o^+} = \frac{-V_o^-}{I_o^-}$$

So then, $Z(z) = \frac{V_o(1 + \Gamma e^{+2j\beta z})}{\frac{V_o}{Z_0}(1 - \Gamma e^{+2j\beta z})}$

(14)

$$Z(z) = Z_0 \frac{(1 + \Gamma e^{+2jBz})}{(1 - \Gamma e^{+2jBz})}$$

Interestingly, we have no voltage dependency
but we have frequency dependency

$$\beta = \frac{Z_0}{\lambda} = \frac{\omega}{V_p}$$

So we may write,

$$Z(z) = Z_0 \frac{Z_L + jZ_0 \tan(Bz)}{Z_0 + jZ_L \tan(Bz)}$$

Also, for a known Line length, l

$$Z_{in} = Z_0 \frac{Z_L + jZ_0 \tan(\beta l)}{Z_0 + jZ_L \tan(\beta l)}$$

$l \rightarrow$ length of T-line

$Z_0 \rightarrow$ characteristic Impedance

$Z_L \rightarrow$ Load impedance

$$\beta = \frac{2\pi}{\lambda} = \frac{\omega_p}{V_p} = \frac{2\pi f}{V_p}$$

Frequency controls
wavelength

⑯ The Input Impedance of an Oscilloscope

- typically we want a large input impedance for the oscilloscope

Ex) $C = V_p = 3 \times 10^3 \text{ nF}$ $f_{\text{sig}} = 1 \text{ MHz}$

$$\lambda f = c \quad \lambda = 300 \text{ [m]}$$

Coaxial cables connectors for the oscilloscope

$$Z_{\text{in}} = Z_0 \frac{Z_L + j Z_0 \tan \left[\left(\frac{2\pi}{300} \right) (1) \right]}{Z_0 + j Z_L \tan \left[\left(\frac{2\pi}{300} \right) (1) \right]} \quad \sim 1 \text{ m}$$



(17)

Scope is in parallel with R_2 , Ideally we want the scope to be an open circuit

For $\lambda \ll \lambda$

$$Z_{in} = Z_0 \frac{Z_L + jZ_0 \tan\left(\frac{2\pi}{\lambda}\right)\left(\frac{1}{f}\right)}{Z_0 + jZ_L \tan\left[\frac{2\pi}{\lambda}\left(\frac{1}{f}\right)\right]}$$

$$= Z_0 \frac{Z_L}{Z_0} = Z_L$$

In this case, Z_L would be the impedance of the oscilloscope

(18) The Ideal input impedance for an oscilloscope
is $Z_{\text{scope}} \rightarrow \infty$

So if the Load is ∞ (open)

$$Z_{\text{in}} = Z_0 \frac{\infty + j Z_0 \tan(\beta l)}{Z_0 + j \infty \tan(\beta l)}$$

$$= Z_0 \frac{\infty}{j \infty \tan(\beta l)} = Z_0 \frac{1}{j \tan(\beta l)}$$

So open ckt load

$$Z_{\text{in}_{\text{oc}}} = - \frac{j Z_0}{\tan(\beta l)}$$

(19)

so if $\lambda = 1$ and $\lambda \ll \lambda$
 $Z_{inoc} \rightarrow \infty$ (to be expected)

Ex) $\lambda = \frac{\lambda}{4}$ in Length

$$\tan(\beta L) = \frac{\sin(\beta L)}{\cos(\beta L)} = \frac{\sin\left[\frac{2\pi}{\lambda} \frac{\lambda}{4}\right]}{\cos\left[\frac{2\pi}{\lambda} \frac{\lambda}{4}\right]} = \frac{\sin\left(\frac{\pi}{2}\right)}{\cos\left(\frac{\pi}{2}\right)}$$

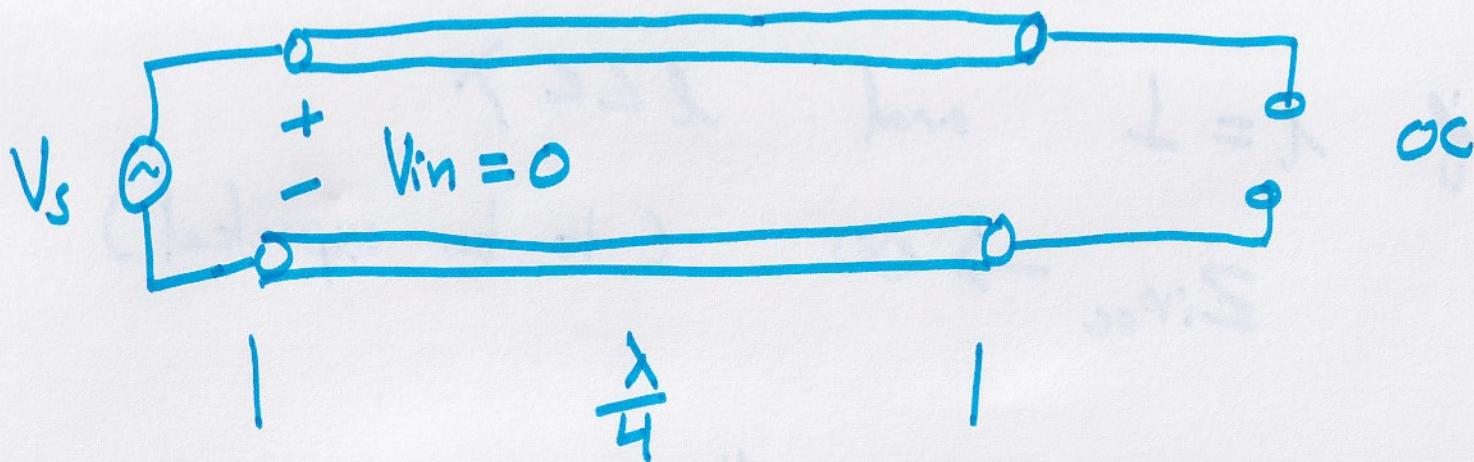
$$= \frac{1}{0} \rightarrow \infty$$

so then

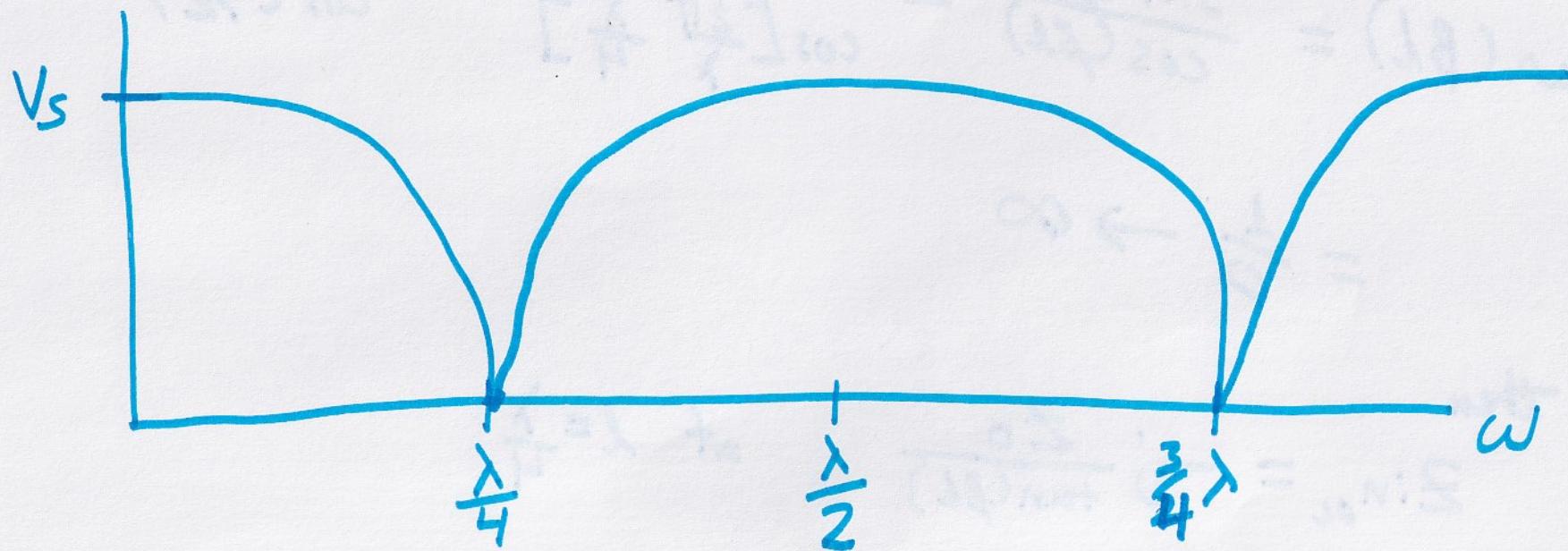
$$Z_{inoc} = -j \frac{Z_0}{\tan(\beta L)} \quad \text{at } \lambda = \frac{\lambda}{4}$$

$$\rightarrow Z_{inoc} = 0$$

(20)



If we plot V_{in}



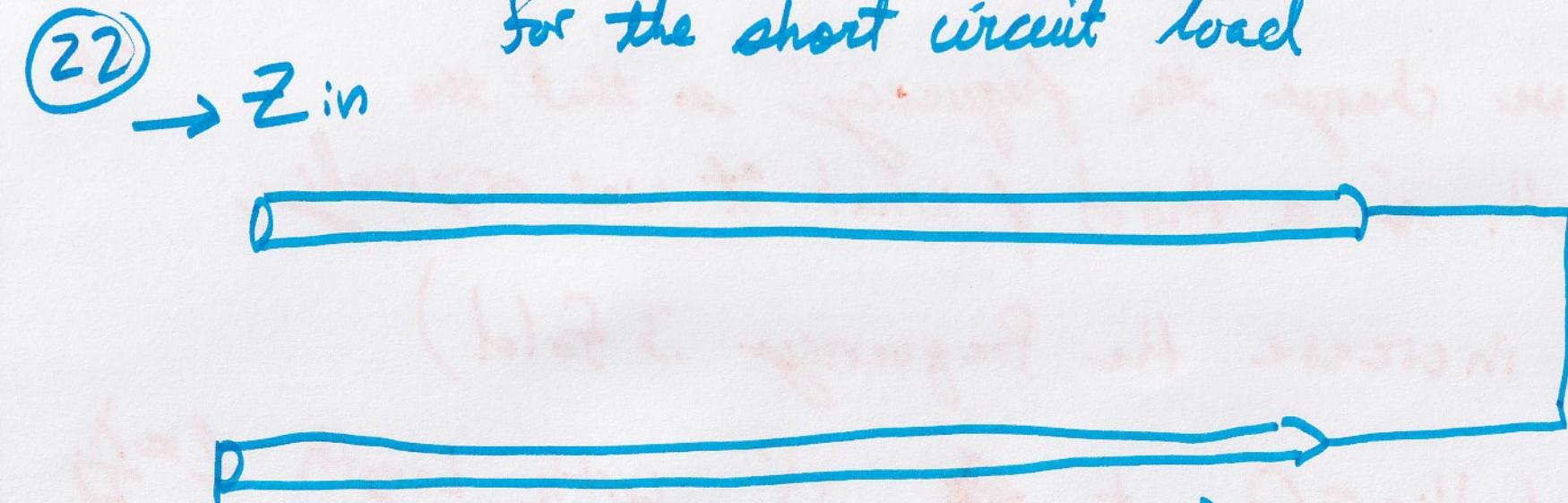
This is changing the frequency,

So if we change the frequency so that the wavelength is a third of what it was originally,

(increase the frequency 3 fold)

We get $V_{in} = 0$ at ~~at~~ $3f$ with Length $l = \frac{\lambda}{4}$

where λ is the original wave length



$Z_L = 0$ (short circuit load)

$$Z_{in} = Z_0 \frac{0 + j Z_0 \tan(\beta l)}{Z_0 + j 0 \tan(\beta l)}$$

$$= j Z_0 \tan(\beta l)$$

For lengths $l = 0, \frac{\lambda}{2}, \lambda, \frac{3}{2}\lambda, \dots$

$$\rightarrow Z_{in} = 0$$

For a matched Load ($Z_L = Z_0$)

(23)

$$Z_0 = Z_{in}$$

$$Z_{in} = Z_0 \frac{Z_L + jZ_0 \tan(\beta l)}{Z_0 + jZ_L \tan(\beta l)}$$

$$= Z_0 (1)$$

Whenever $\tan(\beta l) \rightarrow 0$ for $l = n\lambda$

$$Z_{in} = Z_0 \frac{Z_L + jZ_0 \cancel{\tan(\beta l)}}{Z_0 + jZ_L \cancel{\tan(\beta l)}}$$

$$= Z_L$$

(24)

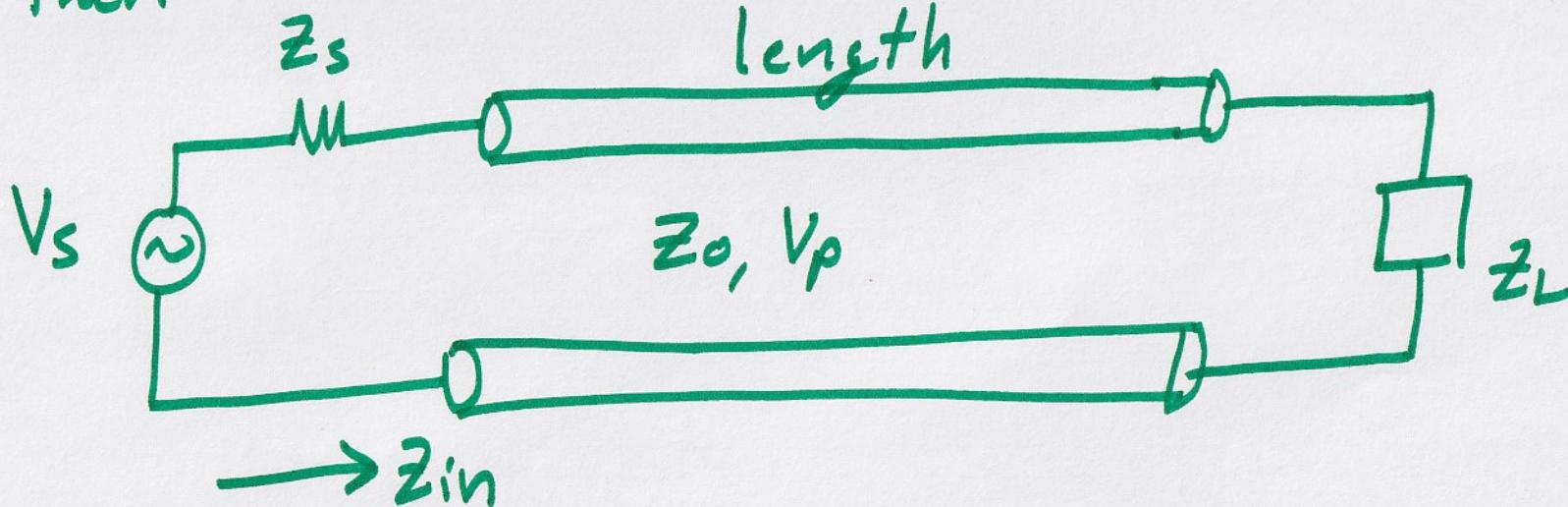
When ever $\tan(\beta l) \rightarrow \infty$ for $l = n \frac{\lambda}{2} + \frac{\lambda}{4}$

$$Z_{in} = Z_0 \frac{Z_L + j Z_0 \tan(\beta l)^\infty}{Z_0 + j Z_L \tan(\beta l)^\infty}$$

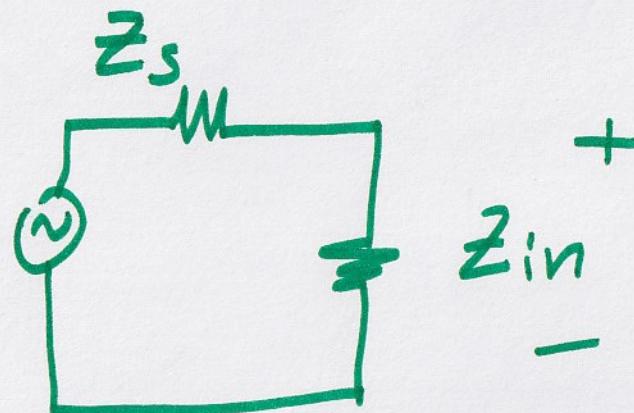
$$= Z_0 \frac{j Z_0(\infty)}{j Z_L(\infty)} = \frac{Z_0^2}{Z_L}$$

~~Ans~~

So then



(25)



$$V_{in} = \frac{z_{in}}{z_s + z_{in}} V_s$$

V_{in} is the voltage of z_{in} and that voltage is at the source end of the line

$$V_{in} |_{\text{source}} = V_{\text{total}} (-l) = V_o^+ (1 + \Gamma e^{+2j\beta l}) e^{-j\beta l}$$

we evaluate V_{total} at $z = -l$