

Fields and Waves I

①

L 1

4 Parts

The Wave Equation

1) Transmission Lines

2) Electrostatics

3) Magnetostatics & Magneto dynamics

4) Plane Electromagnetic Waves
(Plane Waves)

extra

5) Antennas

6) Satellite Comm. Systems and Radars

②

Circuit Theory and Electromagnetics

→ uses simplified (Lumped) Models of components

There are limitations
Components are not simply R, L, C

→ we need details about how they work

Circuit Theory does not include details about

- Distributed properties of the Components
(Transmission lines)

- EM Waves ~ ultras, Radio Waves, Optics

- Capacitive Sensors - Noise

(3)

Linear Circuit Analysis fail in most
high-speed circuits. KCL, KVL fail

so we use EM Theory to analyze such
problems

Kayfabe

④

OK, Transmission Lines

What's a Field?

 \vec{E} electric field both vectors

\vec{B} magnetic field

-vector that vary w/ position (and time)

→ direction
→ magnitude

$I \rightarrow$ current (total current)

$V \rightarrow$ voltage

(5)

Phasor Analysis

From circuits

$$V = |V_b| \angle V_b = |V_o| \angle \Phi_v$$

remember "implied" $e^{j\omega t}$

→ from AC analysis (AC steady state)

$$V = V e^{j\omega t}$$

$$R_a \{ V \} = |V| \cos(\omega t + \Phi_v)$$

⑥

Electrostatics is a big deal in Fields and Waves

$$\oint \bar{E} \cdot d\bar{l} = - \frac{\partial}{\partial t} \phi = 0$$

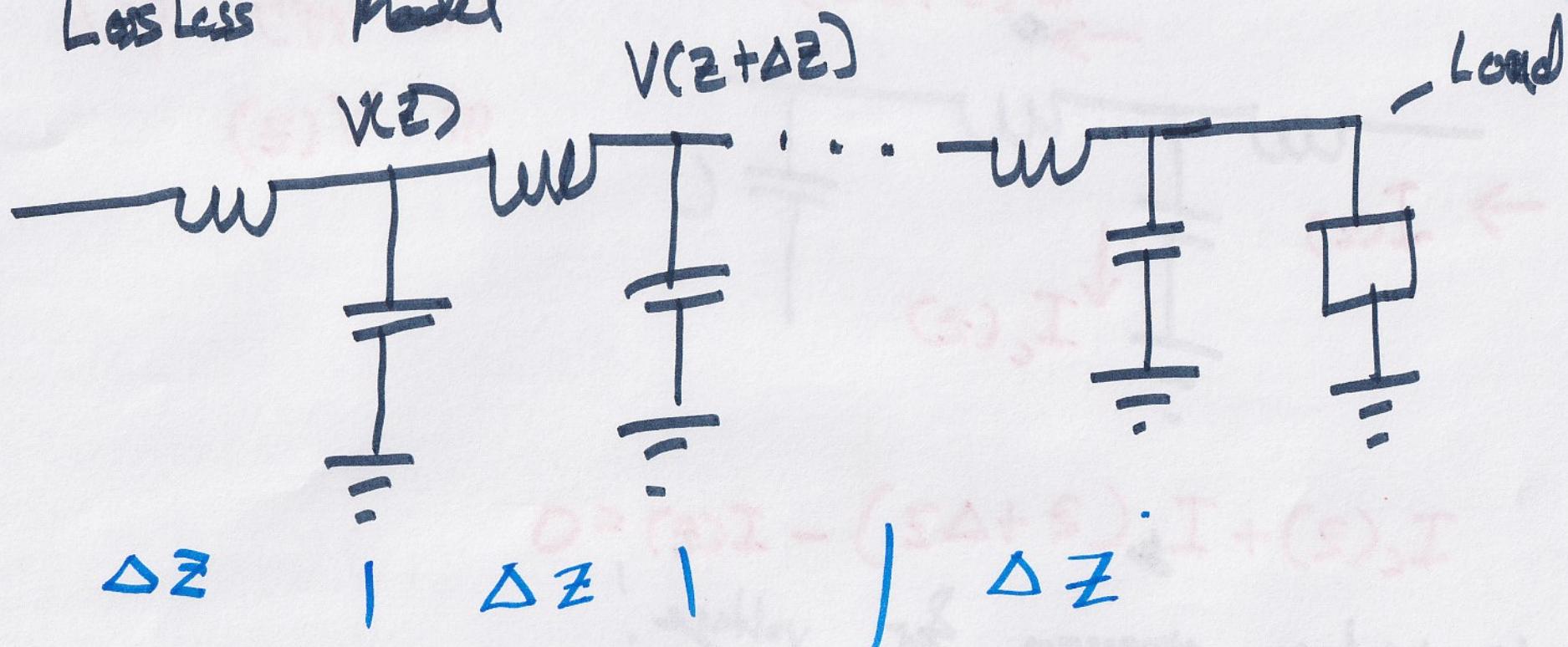
→ does not change with time!

Lumped parameter Model

- use capacitive and inductive concepts to model
a transmission line

Lossless Model

(7)

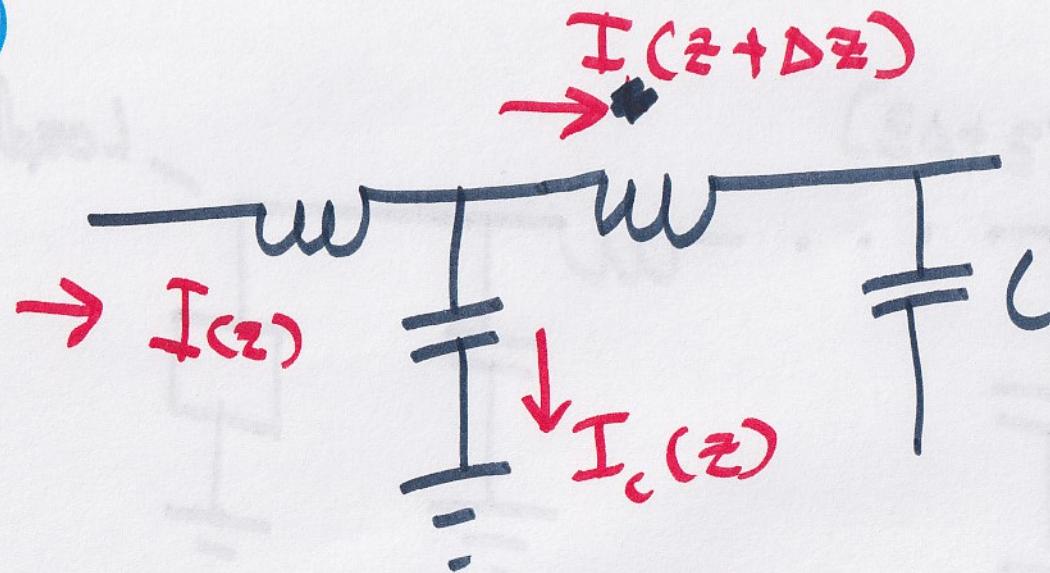


Transmission Line Model

Every \downarrow , section represents some length
of T-Line

— Now use Node Analysis!

(8)



We apply KCL
at $V(z)$

$$I_c(z) + I_{\downarrow}(z + \Delta z) - I(z) = 0$$

We want an expression for voltage!

Use multimeter

$$I_c(z, t) + I(z + \Delta z, t) - I(z, t) = 0$$

current, voltage are a function of both position and time

(9)

We know that

$$I_C = C \frac{d}{dt} V_C$$

$$V_L = L \frac{d}{dt} I_L$$

so $\Delta V = L \frac{d}{dt} I$ changing voltage across the inductor
 So we may write (we are going to need this)

$$C \frac{\partial}{\partial t} V_C(z, t) + \Delta z \left(\frac{I(z+\Delta z, t) - I(z, t)}{\Delta z} \right) = 0$$

$$\rightarrow C \cancel{V_C(z, t)} + \Delta z \frac{\partial}{\partial z} I(z, t) = 0$$

$$\frac{\partial}{\partial z}$$

Now take $\frac{\partial}{\partial t}$

$$⑩ \frac{\partial}{\partial t} \left[c \frac{\partial}{\partial t} V_c(z, t) + \Delta z \frac{\partial}{\partial z} I(z + \Delta z, t) \right] = 0$$

we can also write it like this

if a dc term exists (a constant) it goes away
(could lead to a problem)

$$0 = c \frac{\partial^2}{\partial t^2} V_c(z, t) + (\Delta z) \frac{\partial}{\partial t} \left[\frac{\partial}{\partial z} I(z + \Delta z, t) \right]$$

for some reason, we can do this

$$0 = c \frac{\partial^2}{\partial t^2} V_c(z, t) + (\Delta z) \frac{\partial}{\partial z} \left[\frac{\partial}{\partial t} I(z + \Delta z, t) \right]$$

$$0 = c \frac{\partial^2}{\partial t^2} V_c(z, t) + (\Delta z) \frac{c}{\Delta z} \left[-\frac{V(z+\Delta z, t) + V(z, t)}{L} \right] \quad 11$$

$$LC \frac{\partial^2}{\partial t^2} V_c(z, t) - (\Delta z)^2 \frac{\partial^2}{\partial z^2} V(z, t) = 0$$

$$\left(\frac{L}{\Delta z}\right) \left(\frac{c}{\Delta z}\right) \frac{\partial^2}{\partial t^2} V_c(z, t) - \frac{\partial^2}{\partial z^2} V(z, t) = 0$$

$$LC \frac{\partial^2}{\partial t^2} V_c(z, t) - \frac{\partial^2}{\partial z^2} V(z, t) = 0$$

$$L = \frac{m}{\omega} \quad C = \frac{\epsilon_0}{m}$$

wave equation!

going back to page 10 [From page 10)

(12)

$$\frac{\partial}{\partial t} \left[C \frac{\partial}{\partial t} V_c(z, t) + \Delta z \frac{\partial}{\partial z} I(z, t) \right] = 0$$

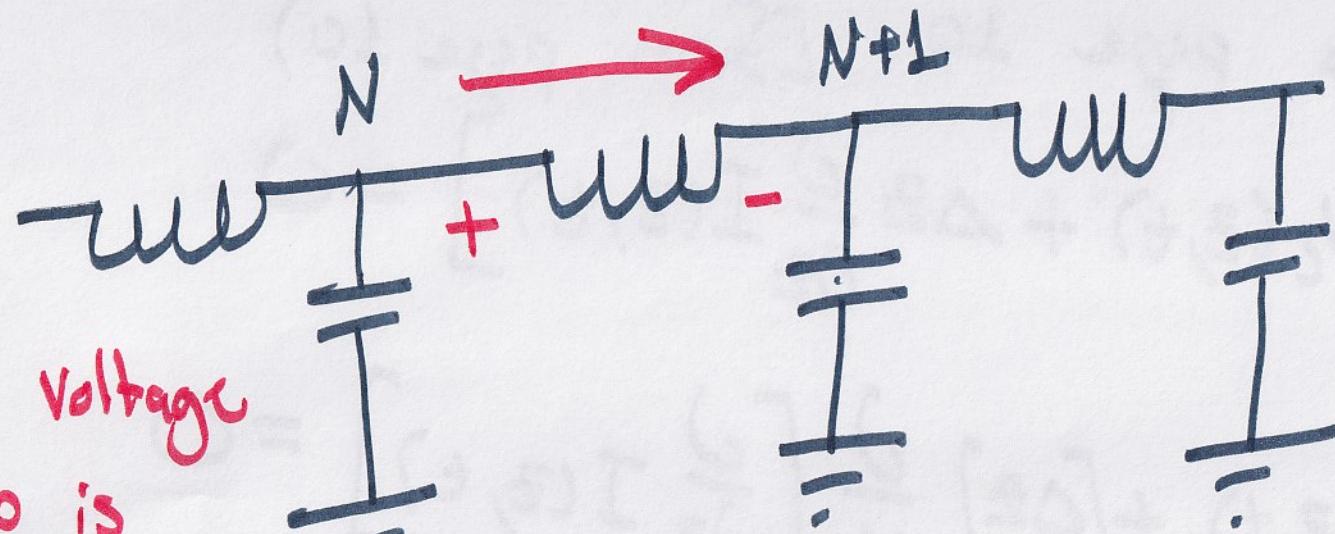
$$\rightarrow C \frac{\partial^2}{\partial t^2} V_c(z, t) + [\Delta z] \frac{\partial}{\partial t} \left[\frac{\partial}{\partial z} I(z, t) \right] = 0$$

$$C \frac{\partial^2}{\partial t^2} V_c(z, t) + [\Delta z] \frac{\partial}{\partial z} \left[\frac{\partial}{\partial t} I(z, t) \right] = 0$$

recall $V_L = L \frac{d}{dt} I = \Delta V$ charge in voltage
across the inductor

So what's the voltage across the inductor?

13



'+' Voltage

drop is
from N to N+1

[Passive

Sign

convention]

$$V_{N+1} - V_N = -V_L$$

$$V_{N+1} = V(z + \Delta z, t)$$

$$V_N = V(z, t)$$

$$\text{So } -\frac{\partial}{\partial t} I(z, t) = \frac{V(z + \Delta z, t) - V(z, t)}{L}$$

$$\rightarrow C \frac{\partial^2}{\partial t^2} V_c(z, t) + (\Delta z) \frac{\partial}{\partial z} \left[\frac{-V(z + \Delta z, t) + V(z, t)}{L} \right] = 0$$

$$\rightarrow LC \frac{\partial^2}{\partial t^2} V_c(z, t) - (\Delta z)^2 \frac{\partial^2}{\partial z^2} \left[\frac{V(z + \Delta z, t) + V(z, t)}{\Delta z} \right] = 0$$

$$LC \frac{\partial^2}{\partial t^2} V_c(z, t) - (\Delta z)^2 \frac{\partial^2}{\partial z^2} V(z, t) = 0$$

$$L = \frac{L}{m} \cdot \frac{H}{m} \quad C = \frac{C}{m} = \frac{F}{m}$$

$$\left(\frac{L}{\Delta z}\right)\left(\frac{C}{\Delta z}\right) \frac{\partial^2}{\partial t^2} V_c(z, t) - \frac{\partial^2}{\partial z^2} V(z, t) = 0$$

$$LC \frac{\partial^2}{\partial t^2} V_c(z, t) - \frac{\partial^2}{\partial z^2} V(z, t) = 0$$

(15)

Note!

Concerning $\frac{\partial}{\partial t} \left[\frac{\partial}{\partial z} I(z,t) \right] = \frac{\partial}{\partial z} \left[\frac{\partial}{\partial t} I(z,t) \right]$

We can fairly assume (and it is) $I(z,t)$ is of class C^2 (the second partials are continuous)

$$I_{zt} = I_{tz}$$

Basically if the partial derivatives exist in some neighborhood of (z,t) and are continuous you can do this