

# Lecture 12

Wednesday, February 20, 2019

1:30 PM

## The ADM Equations

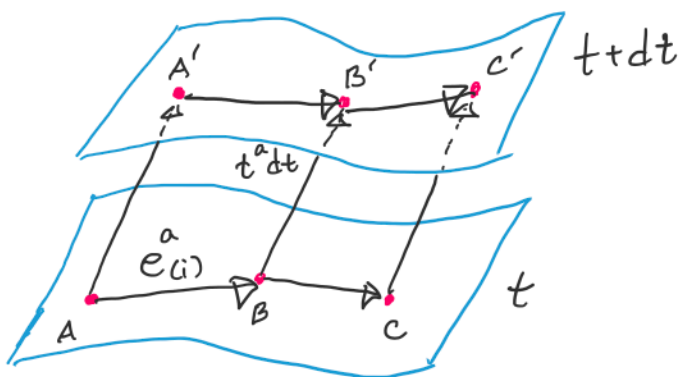
- Evolution eqns. above simplify if we choose our basis wisely.
- introduce spatial basis vectors  $e^a_{(i)}$ , where  $i = 1, 2, 3$  distinguish the 3 vectors, not their components (represented by "a") such that,

$$\Omega_a e^a_{(i)} = 0$$

↳ vectors reside entirely in time slice  $\Sigma$

- vectors are extended to other slices by "Lie dragging" them along  $t_a$

$$\mathcal{L}_{t_a} e^a_{(i)} = 0$$



- In ADM basis notation choose  $\gamma^a = t^a$

for given hypersurface, choose  $e_{(0)} = \dots$

- want this to connect events on neighbouring slices w/ same spatial coords.

• recall duality,  $t^a n_a = 1$

$$\hookrightarrow t^a = e^a_{(0)} = (1, 0, 0, 0)$$

• Lie derivative reduces to partial derivative along  $t$ :  $\mathcal{L}_{\vec{t}} = \partial_t$

• now,  $n_a e^a_{(i)} = -\frac{1}{\alpha} n_a e^a_{(i)} = 0$

• this choice of coordinates implies

$$n_i = 0$$

$\hookrightarrow$  all tensors with contravariant index  $= 0$  vanish, e.g.,

$$n_a \beta^a = n_0 \beta^0 = 0$$

$$\rightarrow \beta^a = (0, \beta^i)$$

• so,  $t^a = \alpha n^a + \beta^a = (1, 0, 0, 0)$

$$\rightarrow n^a = (\alpha^{-1}, -\alpha^{-1} \beta^i)$$

and

$$n_a n^a = -1 \text{ as it must}$$

$$\hookrightarrow n_a = (-\alpha, 0, 0, 0)$$

• from  $\gamma_{ab} = g_{ab} + n_a n_b$

$$\rightarrow \gamma_{ij} = g_{ij}$$

- spatial metric is just the spatial components of the full metric
  - since all  $\gamma^{a0} = 0$ ,
- $$g^{ab} = \gamma^{ab} - n^a n^b = \begin{pmatrix} -\alpha^{-2} & \alpha^{-2} \beta^i \\ \alpha^{-2} \beta^j & \gamma^{ij} - \alpha^{-2} \beta^i \beta^j \end{pmatrix}$$

- $\gamma^{ij}$  and  $\gamma_{ij}$  are 3D inverses,
- $$\gamma^{ik} \gamma_{kj} = \delta^i_j$$

→ can be used for raising & lowering indices of spatial tensors,

$$\beta_i = \gamma_{ij} \beta^j$$

- inverting the inverse 4-metric,

$$g_{ab} = \begin{pmatrix} -\alpha^2 + \beta_i \beta^i & \beta_i \\ \beta_j & \gamma_{ij} \end{pmatrix}$$

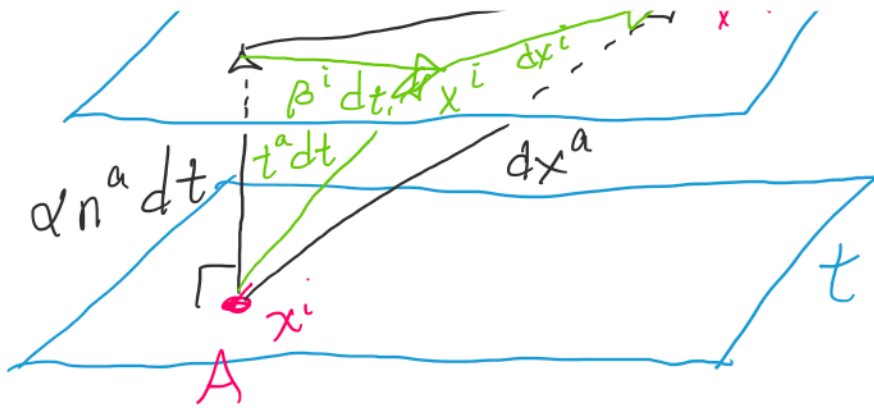
or,

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij} (dx^i + \beta^i dt)(dx^j + \beta^j dt)$$

.... 3+1 metric

- like Pythagorean theorem in 4D, determines invariant interval between events





- with this choice of 3+1 decomposition and basis vectors, entire content of Einstein eqns. is in their spatial components alone

∴ Hamiltonian constraint,

$$R + K^2 - K_{ij} K^{ij} = 16\pi \rho$$

• Momentum constraint,

$$D_j (K^{ij} - \gamma^{ij} K) = 8\pi S^i$$

• evolution equation for extrinsic curvature,

$$\begin{aligned} \partial_t K_{ij} = & -D_i D_j \alpha + \alpha (R_{ij} - 2K_{ik} K^k_j + K K_{ij}) \\ & - 8\pi \alpha (S_{ij} - \frac{1}{2} \gamma_{ij} (S - \rho)) \\ & + \beta^k D_k K_{ij} + K_{ik} D_j \beta^k + K_{kj} D_i \beta^k \end{aligned}$$

• evolution for spatial metric,

$$\partial_t \gamma_{ij} = -2\alpha K_{ij} + D_i \beta_j + D_j \beta_i$$

--- Arnowitt, Deser, Misner (ADM) eqns.

• determinant  $\gamma = \det(\gamma_{ij})$  satisfies

$$\partial_t \gamma^{1/2} = -\gamma^{1/2} K$$

$$\partial_t m^0 = -\alpha \Gamma + \omega_i p$$

trace  $K = K^i_i$  satisfies,

$$\partial_t K = -\Delta^2 \alpha + \alpha (K_{ij} K^{ij} + 4\pi (\rho + S)) + \beta^i \Delta_i K$$

where  $\Delta^2 = \gamma^{ij} \Delta_i \Delta_j$  is Laplacian

• matter source terms:

$$\rho = n_a n_b T^{ab} \quad S^i = -\gamma^{ij} n^a T_{aj}$$

$$S_{ij} = \gamma_{ia} \gamma_{jb} T^{ab} \quad S = \gamma^{ij} S_{ij}$$

• Evaluating RHS of ADM eqns, can disregard all "0" components, eg,  $\Delta_a \Delta_b \alpha = \partial_a \partial_b \alpha - \partial_c \alpha \Gamma^c_{ba}$   
 - from  $\gamma^{a0} = 0$ ,  $\Gamma^0_{ba} = 0$

∴ connection coeffs can be computed from spatial components alone:

$$\Gamma^i_{jk} = \frac{1}{2} \gamma^{il} (\partial_k \gamma_{lj} + \partial_j \gamma_{lk} - \partial_l \gamma_{jk})$$

• Ricci tensor,

$$R_{ij} = \partial_k \Gamma^k_{ij} - \partial_j \Gamma^k_{ik} + \Gamma^k_{ij} \Gamma^l_{kl} - \Gamma^k_{il} \Gamma^l_{jk}$$

$$\text{or } R_{ij} = \frac{1}{2} \gamma^{kl} (\partial_i \partial_l \gamma_{kj} + \partial_k \partial_j \gamma_{il} - \partial_i \partial_j \gamma_{kl} - \partial_k \partial_l \gamma_{ij}) + \gamma^{kl} (\Gamma^m_{il} \Gamma^m_{kj} - \Gamma^m_{ij} \Gamma^m_{kl})$$

• for Schwarzschild:

$$ds^2 = -\left(\frac{1-M/2r}{1+M/2r}\right)^2 dt^2 + \left(1 + \frac{M}{2r}\right)^4 (dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2)$$

→ lapse:

$$\alpha = \frac{1-M/2r}{1+M/2r}$$

$$\alpha = 1 + M/2r$$

shift:  $\beta^i = 0$

spatial metric:

$$\gamma_{ij} = \left(1 + \frac{M}{2r}\right)^4 \text{diag}(1, r^2, r^2 \sin^2 \theta)$$

- takes the form,

$$\gamma_{ij} = \psi^4 \eta_{ij}$$

where  $\psi = 1 + M/2r$  is the *conformal factor*

• since  $\gamma_{ij}$  is independent of  $\beta^i$ ,

$$K_{ij} = 0$$

• Ricci tensor,

$$R_{rr} = -\frac{8rM}{(2r^2 + Mr)^2} \quad R_{\theta\theta} = \frac{4r^3 M}{(2r^2 + Mr)^2}$$

$$R_{\phi\phi} = \sin^2 \theta R_{\theta\theta}$$

[in vacuum,  ${}^{(4)}R_{ij} = 0$ , but  $R_{ij} \neq 0$ ]

• above relations satisfy constraints

- Hamiltonian:  $R + K^2 - K_{ij} K^{ij} = 16\pi \rho$

$$\rightarrow K_{ij} = \rho = 0 \rightarrow R = 0$$

$$\text{so, } \gamma^{ij} R_{ij} = 0$$

• momentum constraint,

$$\Delta_j (K^{ij} - \gamma^{ij} K) = 8\pi S^i$$

$$K^{ij} = K = S^i = 0 \dots \text{so easy}$$

• spatial metric evolution eqn.  
gives  $K_{ij} = 0$

• only non-trivial eqn. then is extrinsic  
curvature evolution equation

$\hookrightarrow \partial_t K_{ij} = 0 \quad \dots$  as expected