

Three-Dimensional EMHD Simulation Studies of Non-linear Magnetic Structures in Magnetically Confined Plasmas

* Mathematical Model

The numerical simulation is based on the nonlinear EMHD equations, which is derived from Faraday's & Ampere's Laws

$$\frac{1}{c} \frac{\partial \vec{B}}{\partial t} = -\nabla \times \vec{E} \quad [1]$$

$$\text{and} \quad \nabla \times \vec{B} = -\frac{4\pi e n_e \vec{v}_e}{c} \quad [2]$$

These together with the electron momentum equation

$$m_e \left(\frac{\partial}{\partial t} + \vec{v}_e \cdot \nabla \right) \vec{v}_e = e \left(\vec{E} + \frac{\vec{v}_e \times \vec{B}}{c} \right) \quad [3]$$

where \vec{B} & \vec{E} are the magnetic and electric fields, \vec{v}_e is the electron fluid velocity and n_e is the electron number density, and m_e is the electron mass and e is the magnitude of the electron charge and c is the speed of light in vacuum.

Equation [3] is true whenever $\partial/\partial t \gg$ the electron-ion collision frequency. Furthermore, since the divergence of

the electron flux is zero, the electron density fluctuations associated with the whistlers. The wave-equation interaction is absent in our equations. By eliminating \vec{E} and \vec{v}_e and noting that

$$(\vec{v}_e \cdot \nabla) \vec{v}_e = -\vec{v}_e \times (\nabla \times \vec{v}_e) + \nabla v_e^2/2, \text{ we obtain the nonlinear EMHD}$$

$$\frac{\partial}{\partial t} (\vec{B} - \lambda_e^2 \nabla^2 \vec{B}) = \frac{c^2}{4\pi e n_0} \nabla \times ((\vec{B} - \lambda_e^2 \nabla^2 \vec{B}) \times (\nabla \times \vec{B})), \quad [4]$$

which is the desired equation for computer simulations

Set up for simulation

The initial conditions for equation [4] are

$\vec{B} = B_0 \hat{z} + \nabla \times \vec{A} + B_{tor} \hat{\phi}$. In cylindrical coordinates the vector potential \vec{A} is taken to be of the form

$$\vec{A} = [A_{forward}(r, z) + A_{reverse}(r, z)] \hat{\phi}$$

Here the forward field is

$$A_{forward}(r, z) = A_0 \frac{r}{D} \exp \left[\frac{-(r-r_0)^2 + z^2}{D^2} \right]$$

is centered at zero ($z=0$) and the reverse vector potential is

$$A_{reverse}(r, z) = -A_0 \frac{r}{D} \exp \left[\frac{-(r-r_0)^2 + (z-S)^2}{D^2} \right]$$

is centered at $z=S$. The toroidal magnetic field is chosen to be of the form

$$B_{tor}(r, z) = -B_{0,tor} \frac{r}{D} \exp \left[\frac{-(r-r_0)^2 + (z-S)^2}{D^2} \right]$$