



Hamilton's principle and least action

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- ✓ D'Alembert principle was a tool to get rid of reaction forces if one describes kinematically admissible displacements.
- ✓ Lagrange equations provided a systematic tool for obtaining the equations without any need for performing vectorial calculations (Physics could be lost).
- ✓ Lagrange's equations is suitable if we are dealing with individual masses and stiffness elements. What if the relative position of 2 points within a body is not kinematically constraint or in other words the body's displacement is a function of location within the body?



Hamilton's principle



A general technique for
obtaining EoMs of discrete
and continuous media



**Based on the concept of
Energy**

D'Alembert principle for a system of N particles states that:

$$\sum_{k=1}^N \sum_{i=1}^3 (m_k \ddot{u}_{ik} - X_{ik}) \delta u_{ik} = 0$$

Let's assume

$$\sum_{k=1}^N \sum_{i=1}^3 X_{ik} \delta u_{ik} = \delta \bar{W}$$

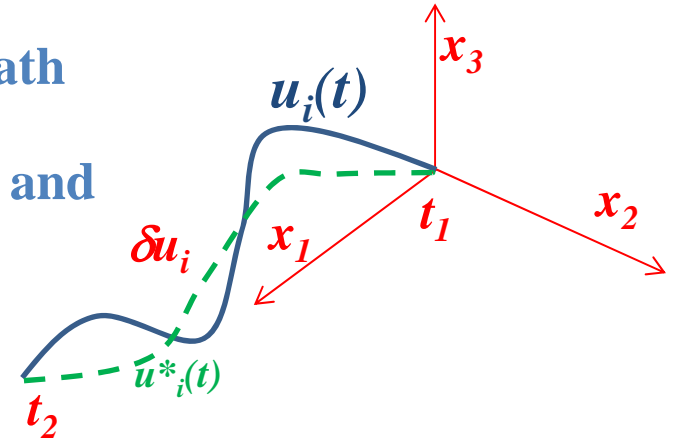
We can rewrite

$$\sum_{k=1}^N \sum_{i=1}^3 (m_k \ddot{u}_{ik}) \delta u_{ik} = \frac{d}{dt} \left(\sum_{k=1}^N \sum_{i=1}^3 (m_k \dot{u}_{ik}) \delta u_{ik} \right) - \sum_{k=1}^N \sum_{i=1}^3 (m_k \dot{u}_{ik}) \delta \dot{u}_{ik}$$

$$\delta \sum_{k=1}^N \sum_{i=1}^3 \left(\frac{1}{2} m_k \dot{u}_{ik}^2 \right) \Rightarrow \delta T$$

$$\sum_{k=1}^N \sum_{i=1}^3 (m_k \ddot{u}_{ik} - X_{ik}) \delta u_{ik} = 0 \quad \Rightarrow \quad \frac{d}{dt} \left(\sum_{k=1}^N \sum_{i=1}^3 (m_k \dot{u}_{ik}) \delta u_{ik} \right) = \delta T + \delta \bar{W}$$

As time unfolds the dynamic system traces the true path $u_i(t)$ in $3N$ space. Let's assume that a varied path also exists ($\bar{u}_i^*(t)$) that coincides with true path at times t_1 and t_2 . Integrating above equation between t_1 and t_2 :



$$\int_{t_1}^{t_2} (\delta T + \delta \bar{W}) dt = \int_{t_1}^{t_2} \frac{d}{dt} \left(\sum_{k=1}^N \sum_{i=1}^3 (m_k \dot{u}_{ik}) \delta u_{ik} \right) dt$$

$$\sum_{k=1}^N \sum_{i=1}^3 (m_k \dot{u}_{ik}) \delta u_{ik} \Big|_{t_1}^{t_2} = 0$$

Therefore:

$$\int_{t_1}^{t_2} (\delta T + \delta \bar{W}) dt = 0 \quad , \quad \delta u_{ik}(t_1) = \delta u_{ik}(t_2) = 0 \quad i = 1, 2, 3 \quad \& \quad k = 1, \dots, N$$

In terms of generalized coordinates

$$\int_{t_1}^{t_2} (\delta T + \delta \bar{W}) dt = 0 \quad , \quad \delta q_k(t_1) = \delta q_k(t_2) = 0 \quad k = 1, \dots, n$$

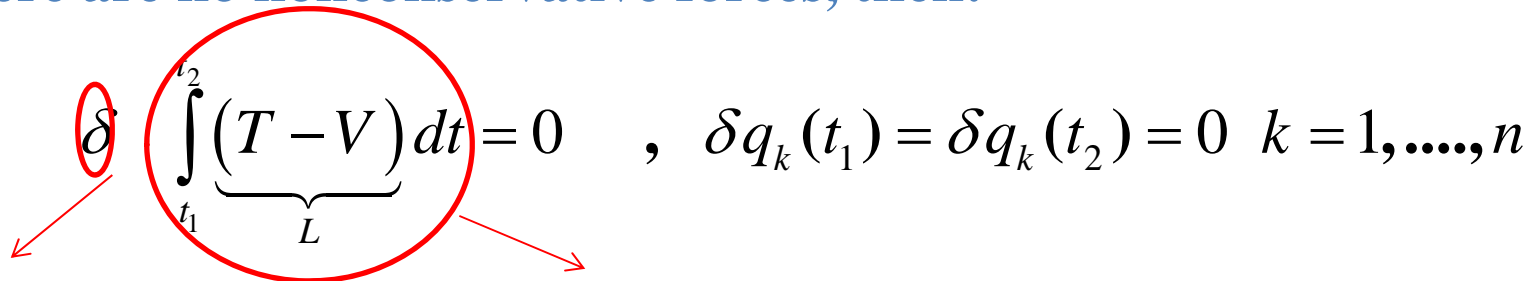
Also the virtual work includes contributions of conservative and non-conservative forces

$$\delta \bar{W} = \delta W^{cons} + \delta W^{ncons} = -\delta V + \delta W^{ncons}$$

Therefore:

$$\int_{t_1}^{t_2} (\delta T - \delta V + \delta W^{ncons}) dt = 0 \quad , \quad \delta q_k(t_1) = \delta q_k(t_2) = 0 \quad k = 1, \dots, n$$

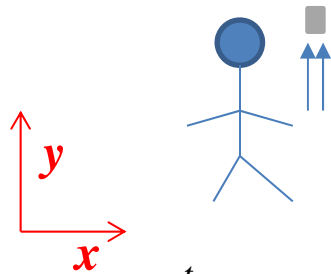
If there are no nonconservative forces, then:


$$\delta \int_{t_1}^{t_2} \underbrace{(T - V)}_L dt = 0 \quad , \quad \delta q_k(t_1) = \delta q_k(t_2) = 0 \quad k = 1, \dots, n$$

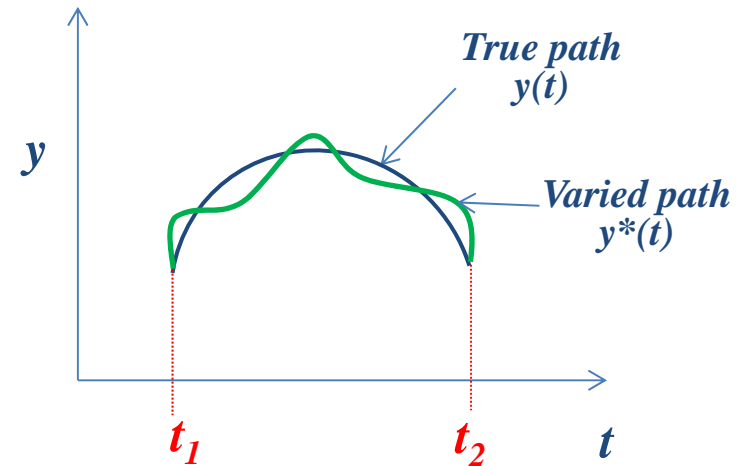
Variational operator

I=Action Integral

What does action integral mean and how Hamilton's principle could be interpreted?



$$y^*(t) - y(t) = \delta y$$



$$I_{true} = \int_{t_1}^{t_2} (T - V) dt = \int_{t_1}^{t_2} \left(\frac{1}{2} m \dot{y}^2 - mgy \right) dt$$

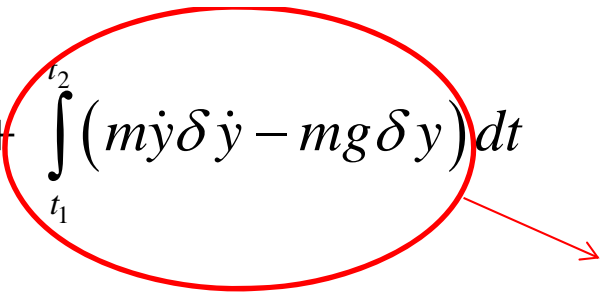
$$I_{var} = \int_{t_1}^{t_2} \left(\frac{1}{2} m (\dot{y} + \delta \dot{y})^2 - mg(y + \delta y) \right) dt$$

If we expand I_{var}

$$I_{var} = \int_{t_1}^{t_2} \left(\frac{1}{2} m \dot{y}^2 + m \dot{y} \delta \dot{y} + \frac{1}{2} m (\delta \dot{y})^2 - mgy - mg \delta y \right) dt$$

Very small=0

$$I_{var} = I_{true} + \int_{t_1}^{t_2} (m \dot{y} \delta \dot{y} - mg \delta y) dt$$

$$I_{\text{var}} = I_{\text{true}} + \int_{t_1}^{t_2} (m\dot{y}\delta\dot{y} - mg\delta y) dt$$


$$I_{\text{var}} = I_{\text{true}} + \underbrace{\delta \int_{t_1}^{t_2} \left(\frac{1}{2} m\dot{y}^2 - mgy \right) dt}_I$$

If $\delta I=0$ the true path renders the value of the action integral stationary with respect to all arbitrary variations of the path between two instants t_1 and t_2 provided that the path variations vanish at these two instants.

Therefore:

$$\int_{t_1}^{t_2} (m\dot{y}\delta\dot{y} - mg\delta y) dt = \underbrace{m\dot{y}\delta y \Big|_{t_1}^{t_2}}_{=0} - \int_{t_1}^{t_2} (m\ddot{y} + mg) \delta y dt$$

So the only way that $\delta I=0$ is:

$$m\ddot{y} + mg = 0$$

The double pendulum

$$\delta \int_{t_1}^{t_2} (T - V + W^{ncons}) dt = 0 \quad , \quad \delta q_k(t_1) = \delta q_k(t_2) = 0 \quad k = 1, \dots, n$$

$$k = 2 \quad , \quad q_1 = \theta_1 \quad , \quad q_2 = \theta_2$$

$$Q_1^{ncons} = Q_1^{ncons} = 0$$

$$\xi_{11} = u_{11} = l_1 \cos \theta_1$$

$$\xi_{21} = u_{21} = l_1 \sin \theta_1$$

$$\xi_{12} = u_{12} = \xi_{11} + l_2 \cos(\theta_1 + \theta_2) = l_1 \cos \theta_1 + l_2 \cos(\theta_1 + \theta_2)$$

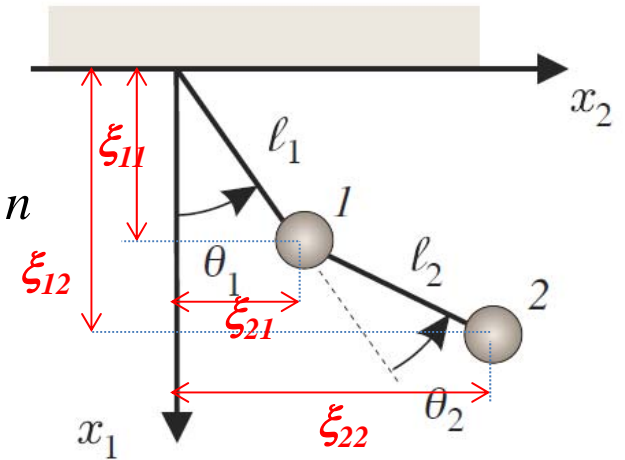
$$\xi_{22} = u_{22} = \xi_{21} + l_2 \sin(\theta_1 + \theta_2) = l_1 \sin \theta_1 + l_2 \sin(\theta_1 + \theta_2)$$

$$\dot{u}_{11} = -l_1 \dot{\theta}_1 \sin \theta_1$$

$$\dot{u}_{21} = l_1 \dot{\theta}_1 \cos \theta_1$$

$$\dot{u}_{12} = -l_1 \dot{\theta}_1 \sin \theta_1 - l_2 \dot{\theta}_1 \sin(\theta_1 + \theta_2) - l_2 \dot{\theta}_2 \sin(\theta_1 + \theta_2)$$

$$\dot{u}_{22} = l_1 \dot{\theta}_1 \cos \theta_1 + l_2 \dot{\theta}_1 \cos(\theta_1 + \theta_2) + l_2 \dot{\theta}_2 \cos(\theta_1 + \theta_2)$$



$$\begin{aligned}
T &= \frac{1}{2}m\dot{u}_{11}^2 + \frac{1}{2}m\dot{u}_{21}^2 + \frac{1}{2}m\dot{u}_{12}^2 + \frac{1}{2}m\dot{u}_{22}^2 \\
&= \frac{1}{2}ml_1^2\dot{\theta}_1^2 + \frac{1}{2}ml_1^2\dot{\theta}_1^2 + \frac{1}{2}ml_2^2(\dot{\theta}_1 + \dot{\theta}_2)^2 + ml_1l_2\dot{\theta}_1^2 \cos \theta_2 + ml_1l_2\dot{\theta}_1\dot{\theta}_2 \cos \theta_2
\end{aligned}$$

$$\begin{aligned}
V &= -mgu_{11} - mgu_{12} \\
&= mgl_1(1 - \cos \theta_1) + mgl_1(1 - \cos \theta_1) + mgl_2(1 - \cos(\theta_1 + \theta_2))
\end{aligned}$$

$$\begin{aligned}
\delta \int_{t_1}^{t_2} T dt &= \int_{t_1}^{t_2} \left[ml_1^2\dot{\theta}_1\delta\dot{\theta}_1 + ml_1^2\dot{\theta}_1\delta\dot{\theta}_1 + ml_2^2(\dot{\theta}_1 + \dot{\theta}_2)(\delta\dot{\theta}_1 + \delta\dot{\theta}_2) \right. \\
&\quad + 2ml_1l_2\dot{\theta}_1 \cos \theta_2\delta\dot{\theta}_1 - ml_1l_2\dot{\theta}_1^2 \sin \theta_2\delta\theta_2 + ml_1l_2\dot{\theta}_2 \cos \theta_2\delta\dot{\theta}_1 \\
&\quad \left. + ml_1l_2\dot{\theta}_1 \cos \theta_2\delta\dot{\theta}_2 - ml_1l_2\dot{\theta}_1\dot{\theta}_2 \sin \theta_2\delta\theta_2 \right] dt
\end{aligned}$$

$$\begin{aligned}
\delta \int_{t_1}^{t_2} V dt &= \int_{t_1}^{t_2} \left[mgl_1 \sin \theta_1\delta\theta_1 + mgl_1 \sin \theta_1\delta\theta_1 \right. \\
&\quad \left. + mgl_2 \sin(\theta_1 + \theta_2)\delta\theta_1 + mgl_2 \sin(\theta_1 + \theta_2)\delta\theta_2 \right] dt
\end{aligned}$$

Integration by parts

$$\delta \int_{t_1}^{t_2} T dt =$$

$$\left[ml_1^2 \dot{\theta}_1 \delta \theta_1 + ml_1^2 \dot{\theta}_1 \delta \theta_1 + ml_2^2 (\dot{\theta}_1 + \dot{\theta}_2) \delta \theta_1 + ml_2^2 (\dot{\theta}_1 + \dot{\theta}_2) \delta \theta_2 \right. \\ \left. + 2ml_1 l_2 \dot{\theta}_1 \cos \theta_2 \delta \theta_1 + ml_1 l_2 \dot{\theta}_2 \cos \theta_2 \delta \theta_1 + ml_1 l_2 \dot{\theta}_1 \cos \theta_2 \delta \theta_2 \right]_{t_1}^{t_2} -$$

Zero

$$\int_{t_1}^{t_2} \left[ml_1^2 \ddot{\theta}_1 \delta \theta_1 + ml_1^2 \ddot{\theta}_1 \delta \theta_1 + ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) (\delta \theta_1 + \delta \theta_2) \right.$$

$$+ \frac{d}{dt} (2ml_1 l_2 \dot{\theta}_1 \cos \theta_2) \delta \theta_1 - ml_1 l_2 \dot{\theta}_1^2 \sin \theta_2 \delta \theta_2 + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) \delta \theta_1$$

$$\left. + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) \delta \theta_2 - ml_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \sin \theta_2 \delta \theta_2 \right] dt$$

$$\delta \int_{t_1}^{t_2} T dt =$$

$$- \int_{t_1}^{t_2} \left[ml_1^2 \ddot{\theta}_1 + ml_1^2 \ddot{\theta}_1 + ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) + \frac{d}{dt} (2ml_1 l_2 \dot{\theta}_1 \cos \theta_2) + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) \right] \delta \theta_1 dt$$

$$- \int_{t_1}^{t_2} \left[ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) - ml_1 l_2 \dot{\theta}_1^2 \sin \theta_2 + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) - ml_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \sin \theta_2 \right] \delta \theta_2 dt$$

Finally

$$\begin{aligned} \delta \int_{t_1}^{t_2} (T - V) dt = \\ - \int_{t_1}^{t_2} \left[ml_1^2 \ddot{\theta}_1 + ml_1^2 \ddot{\theta}_1 + ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) + \frac{d}{dt} (2ml_1 l_2 \dot{\theta}_1 \cos \theta_2) + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) + 2mgl_1 \sin \theta_1 + mgl_2 \sin (\theta_1 + \theta_2) \right] \delta \theta_1 dt \\ - \int_{t_1}^{t_2} \left[ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) - ml_1 l_2 \dot{\theta}_1^2 \sin \theta_2 + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) - ml_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \sin \theta_2 + mgl_2 \sin (\theta_1 + \theta_2) \right] \delta \theta_2 dt \end{aligned}$$



First EoM

$$2ml_1^2 \ddot{\theta}_1 + ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) + \frac{d}{dt} (2ml_1 l_2 \dot{\theta}_1 \cos \theta_2) + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) + 2mgl_1 \sin \theta_1 + mgl_2 \sin (\theta_1 + \theta_2) = 0$$

Second EoM

$$ml_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) - ml_1 l_2 \dot{\theta}_1^2 \sin \theta_2 + \frac{d}{dt} (ml_1 l_2 \dot{\theta}_2 \cos \theta_2) - ml_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \sin \theta_2 + mgl_2 \sin (\theta_1 + \theta_2) = 0$$

Important property of the variational operator

$$\delta \left[\frac{\partial w(x,t)}{\partial t} \right] = \frac{\partial}{\partial t} [\delta w(x,t)] = \delta \dot{w} \quad , \quad \delta \left[\frac{\partial w(x,t)}{\partial x} \right] = \frac{\partial}{\partial x} [\delta w(x,t)] = \delta w'$$

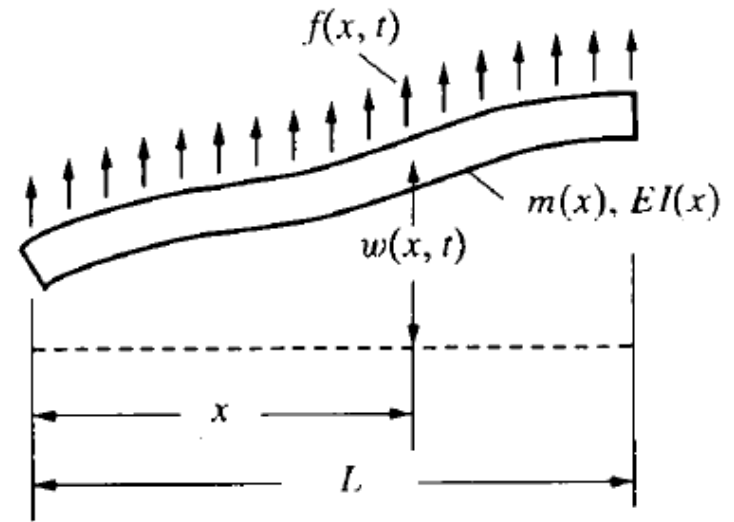
Beam in bending

$$\delta \int_{t_1}^{t_2} (T - V + W^{ncons}) dt = 0 \quad ,$$

$$T = \frac{1}{2} \int_0^L m(x) \left[\frac{\partial w(x,t)}{\partial t} \right]^2 dx$$

$$V = \frac{1}{2} \int_0^L EI(x) \left[\frac{\partial^2 w(x,t)}{\partial x^2} \right]^2 dx$$

$$W^{ncons} = \int_0^L f(x,t) w(x,t) dx$$



$$\delta \int_{t_1}^{t_2} \int_0^L \left(\frac{1}{2} m(x) \left[\frac{\partial w(x,t)}{\partial t} \right]^2 - \frac{1}{2} EI(x) \left[\frac{\partial^2 w(x,t)}{\partial x^2} \right]^2 + f(x,t) w(x,t) \right) dx dt$$

$$\int_{t_1}^{t_2} \int_0^L \left(\underbrace{m(x) \left[\frac{\partial w(x,t)}{\partial t} \right] \delta \left[\frac{\partial w(x,t)}{\partial t} \right]}_A - \underbrace{E(x) I(x) \left[\frac{\partial^2 w(x,t)}{\partial x^2} \right] \delta \left[\frac{\partial w^2(x,t)}{\partial x^2} \right]}_B + \underbrace{f(x,t) \delta w(x,t)}_C \right) dx dt$$

$$\text{A: } \int_{t_1}^{t_2} \int_0^L (m(x) \dot{w} \delta \dot{w}) dx dt = \int_0^L \underbrace{m(x) \dot{w} \delta w \Big|_{t_1}^{t_2}}_{=0} dx - \int_{t_1}^{t_2} \int_0^L (m(x) \ddot{w} \delta w) dx dt = - \int_{t_1}^{t_2} \int_0^L (m(x) \ddot{w} \delta w) dx dt$$

$$\begin{aligned} \text{B: } \int_{t_1}^{t_2} \int_0^L (E(x) I(x) w'' \delta w'') dx dt &= \int_{t_1}^{t_2} (E(x) I(x) w'') \delta w' \Big|_0^L dt - \int_{t_1}^{t_2} \int_0^L \left((E(x) I(x) w'')' \delta w' \right) dx dt \\ &= \int_{t_1}^{t_2} (E(x) I(x) w'') \delta w' \Big|_0^L dt - \int_{t_1}^{t_2} (E(x) I(x) w'')' \delta w \Big|_0^L dt + \int_{t_1}^{t_2} \int_0^L \left((E(x) I(x) w'')'' \delta w \right) dx dt \end{aligned}$$

$$\text{C: } \int_{t_1}^{t_2} \int_0^L f \delta w dx dt$$

Combining previous expressions

$$\int_{t_1}^{t_2} (\delta T - \delta V + \delta W^{ncons}) dt$$

$$= \int_{t_1}^{t_2} \int_0^L \left(m(x) \ddot{w} + \left((E(x) I(x) w'')'' \right) - f \right) \delta w dx dt + \int_{t_1}^{t_2} \left[(E(x) I(x) w'') \delta w' \Big|_0^L - (E(x) I(x) w'')' \delta w \Big|_0^L \right] dt$$

EOM

$$m(x) \ddot{w} + \left((E(x) I(x) w'')'' \right) - f = 0$$

Boundary conditions

$$\text{Either } \underbrace{E(x) I(x) w''}_{\text{bending moment}} = 0 \quad \text{or} \quad \underbrace{w'}_{\text{slope}} = 0 \quad \text{at } x = 0, L$$

$$\text{Either } \underbrace{(E(x) I(x) w'')'}_{\text{shearing force}} = 0 \quad \text{or} \quad \underbrace{w}_{\text{deflection}} = 0 \quad \text{at } x = 0, L$$