Tracing Magnetic Field Lines in BOUT++

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1 Field Line Equations

The equation of a magnetic field line is

$$\frac{d\mathbf{R}}{ds} = \frac{\partial \mathbf{R}}{\partial x} \frac{dx}{ds} + \frac{\partial \mathbf{R}}{\partial y} \frac{dy}{ds} + \frac{\partial \mathbf{R}}{\partial z} \frac{dz}{ds} = \frac{\mathbf{B}}{|\mathbf{B}|}.$$

The magnetic field can be expressed in contravariant form as

$$\boldsymbol{B} = (\boldsymbol{B} \cdot \boldsymbol{\nabla} x) \frac{\partial \boldsymbol{R}}{\partial x} + (\boldsymbol{B} \cdot \boldsymbol{\nabla} y) \frac{\partial \boldsymbol{R}}{\partial y} + (\boldsymbol{B} \cdot \boldsymbol{\nabla} z) \frac{\partial \boldsymbol{R}}{\partial z}.$$

Equating components on both sides of the equation one obtains

$$\frac{dx}{ds} = \frac{\boldsymbol{B} \cdot \boldsymbol{\nabla} x}{|\boldsymbol{B}|}, \quad \frac{dy}{ds} = \frac{\boldsymbol{B} \cdot \boldsymbol{\nabla} y}{|\boldsymbol{B}|}, \quad \frac{dz}{ds} = \frac{\boldsymbol{B} \cdot \boldsymbol{\nabla} z}{|\boldsymbol{B}|}.$$

In BOUT++ the equilibrium field is just $\boldsymbol{B}_0 = \nabla z \times \nabla x$ so that the equations become

$$\frac{dx}{ds} = \frac{\tilde{\boldsymbol{B}} \cdot \nabla x}{\left| \boldsymbol{B}_0 + \tilde{\boldsymbol{B}} \right|}, \quad \frac{dy}{ds} = \frac{\left(\boldsymbol{B}_0 + \tilde{\boldsymbol{B}} \right) \cdot \nabla y}{\left| \boldsymbol{B}_0 + \tilde{\boldsymbol{B}} \right|}, \quad \frac{dz}{ds} = \frac{\tilde{\boldsymbol{B}} \cdot \nabla z}{\left| \boldsymbol{B}_0 + \tilde{\boldsymbol{B}} \right|}.$$

To lowest order in $\frac{|\tilde{\boldsymbol{B}}|}{|\boldsymbol{B}_0|}$ these become

$$\frac{dx}{ds} \approx \frac{\dot{\mathbf{B}} \cdot \nabla x}{|\mathbf{B}_0|}, \quad \frac{dy}{ds} \approx \frac{\mathbf{B}_0 \cdot \nabla y}{|\mathbf{B}_0|}, \quad \frac{dz}{ds} \approx \frac{\ddot{\mathbf{B}} \cdot \nabla z}{|\mathbf{B}_0|}.$$

Using the field-aligned coordinate y in lieu of s for integration gives

$$\frac{dx}{dy} = \frac{dx}{ds}\frac{ds}{dy} \approx \frac{\tilde{\boldsymbol{B}} \cdot \nabla x}{\boldsymbol{B}_0 \cdot \nabla y}, \quad \frac{dz}{dy} = \frac{dz}{ds}\frac{ds}{dy} \approx \frac{\tilde{\boldsymbol{B}} \cdot \nabla z}{\boldsymbol{B}_0 \cdot \nabla y}.$$

The term in the denominator is simply $\boldsymbol{B}_0 \cdot \nabla y = (\nabla z \times \nabla x) \cdot \nabla y = \mathcal{J}^{-1} = \frac{B_{\theta}}{h_{\theta}}$ so to lowest order in the perturbation amplitude, the field-line equations are:

$$\frac{dx}{dy} \approx \left(\tilde{\boldsymbol{B}} \cdot \boldsymbol{\nabla} x\right) \frac{h_{\theta}}{B_{\theta}}, \quad \frac{dz}{dy} \approx \left(\tilde{\boldsymbol{B}} \cdot \boldsymbol{\nabla} z\right) \frac{h_{\theta}}{B_{\theta}}.$$

2 Deriving $\tilde{m{B}}$ From $m{ abla} imes \left(ilde{A}_{\parallel} m{b}_0 ight)$

First we use a vector identity to express the perturbed magnetic field as

$$\tilde{\boldsymbol{B}} = \boldsymbol{\nabla} \times \left(\tilde{A}_{\parallel} \boldsymbol{b}_{0} \right) = \left(\boldsymbol{\nabla} \tilde{A}_{\parallel} \right) \times \boldsymbol{b}_{0} + \tilde{A}_{\parallel} \boldsymbol{\nabla} \times \boldsymbol{b}_{0}. \tag{2.1}$$

Note that in order to remain divergence free, BOTH terms above must be kept:

$$\nabla \cdot \tilde{\boldsymbol{B}} = \nabla \cdot \left[\left(\nabla \tilde{A}_{\parallel} \right) \times \boldsymbol{b}_{0} \right] + \nabla \cdot \left[\tilde{A}_{\parallel} \nabla \times \boldsymbol{b}_{0} \right]$$

$$= \left\{ \boldsymbol{b}_{0} \cdot \left[\nabla \times \left(\nabla \tilde{A}_{\parallel} \right) \right] - \left(\nabla \tilde{A}_{\parallel} \right) \cdot \nabla \times \boldsymbol{b}_{0} \right\} + \left\{ \tilde{A}_{\parallel} \nabla \cdot \left(\nabla \times \boldsymbol{b}_{0} \right) + \left(\nabla \times \boldsymbol{b}_{0} \right) \cdot \left(\nabla \tilde{A}_{\parallel} \right) \right\}$$

$$= 0$$

The first term in Eq. 1.1 above becomes in covariant form

$$\begin{split} \left(\nabla\tilde{A}_{\parallel}\right)\times\boldsymbol{b}_{0} &= B_{0}^{-1}\left(\nabla\tilde{A}_{\parallel}\right)\times\boldsymbol{B}_{0} \\ &= B_{0}^{-1}\left(\frac{\partial\tilde{A}_{\parallel}}{\partial x}\nabla x+\frac{\partial\tilde{A}_{\parallel}}{\partial y}\nabla y+\frac{\partial\tilde{A}_{\parallel}}{\partial z}\nabla z\right)\times\left(\nabla z\times\nabla x\right) \\ &= B_{0}^{-1}\left\{\frac{\partial\tilde{A}_{\parallel}}{\partial x}\nabla x\times\left(\nabla z\times\nabla x\right)+\frac{\partial\tilde{A}_{\parallel}}{\partial y}\nabla y\times\left(\nabla z\times\nabla x\right)+\frac{\partial\tilde{A}_{\parallel}}{\partial z}\nabla z\times\left(\nabla z\times\nabla x\right)\right\} \\ &= B_{0}^{-1}\left\{\begin{array}{c} \frac{\partial\tilde{A}_{\parallel}}{\partial x}\left[\left(\nabla x\cdot\nabla x\right)\nabla z-\left(\nabla x\cdot\nabla z\right)\nabla x\right]+\frac{\partial\tilde{A}_{\parallel}}{\partial y}\left[\left(\nabla y\cdot\nabla x\right)\nabla z-\left(\nabla y\cdot\nabla z\right)\nabla x\right]\right. \\ &\left.+\frac{\partial\tilde{A}_{\parallel}}{\partial z}\left[\left(\nabla z\cdot\nabla x\right)\nabla z-\left(\nabla z\cdot\nabla z\right)\nabla x\right] \\ &= B_{0}^{-1}\left\{\left[g^{11}\frac{\partial\tilde{A}_{\parallel}}{\partial x}+g^{21}\frac{\partial\tilde{A}_{\parallel}}{\partial y}+g^{31}\frac{\partial\tilde{A}_{\parallel}}{\partial z}\right]\nabla z-\left[g^{13}\frac{\partial\tilde{A}_{\parallel}}{\partial x}+g^{23}\frac{\partial\tilde{A}_{\parallel}}{\partial y}+g^{33}\frac{\partial\tilde{A}_{\parallel}}{\partial z}\right]\nabla x\right\} \\ &\equiv f_{1}\nabla x+f_{3}\nabla z. \end{split}$$

The contravariant components are related via $B^i = g^{ij}B_j$ so that in contravariant form this becomes

$$\left(\nabla \tilde{A}_{\parallel}\right) \times \boldsymbol{b}_{0} = \left(g^{11}f_{1} + g^{13}f_{3}\right) \frac{\partial \boldsymbol{R}}{\partial x} + \left(g^{21}f_{1} + g^{23}f_{3}\right) \frac{\partial \boldsymbol{R}}{\partial y} + \left(g^{31}f_{1} + g^{33}f_{3}\right) \frac{\partial \boldsymbol{R}}{\partial z}.$$
 (2.2)

The second term in Eq. 1.1 can be manipulated using

$$\hat{\boldsymbol{b}}_{0} \times \boldsymbol{\kappa} = \hat{\boldsymbol{b}}_{0} \times \left[-\hat{\boldsymbol{b}}_{0} \times (\boldsymbol{\nabla} \times \boldsymbol{b}_{0}) \right]
= \hat{\boldsymbol{b}}_{0} \cdot \hat{\boldsymbol{b}}_{0} (\boldsymbol{\nabla} \times \boldsymbol{b}_{0}) - \left(\hat{\boldsymbol{b}}_{0} \cdot (\boldsymbol{\nabla} \times \boldsymbol{b}_{0}) \right) \hat{\boldsymbol{b}}_{0}$$

so that

$$\mathbf{\nabla} \times \mathbf{b}_0 = \hat{\mathbf{b}}_0 \times \mathbf{\kappa} + (\hat{\mathbf{b}}_0 \cdot (\mathbf{\nabla} \times \mathbf{b}_0)) \hat{\mathbf{b}}_0.$$

With this the second term is simply

$$\tilde{A}_{\parallel} \nabla \times \boldsymbol{b}_{0} = \tilde{A}_{\parallel} \hat{\boldsymbol{b}}_{0} \times \boldsymbol{\kappa} + \tilde{A}_{\parallel} \left(\hat{\boldsymbol{b}}_{0} \cdot (\nabla \times \boldsymbol{b}_{0}) \right) \hat{\boldsymbol{b}}_{0}.$$

Expressing this in contravariant form with $\hat{\boldsymbol{b}}_0 = \frac{(\boldsymbol{\nabla} z \times \boldsymbol{\nabla} x) \cdot \boldsymbol{\nabla} y}{B_0} \frac{\partial \boldsymbol{R}}{\partial y} = \frac{1}{\mathcal{J}B_0} \frac{\partial \boldsymbol{R}}{\partial y}$ we have

$$\tilde{A}_{\parallel} \nabla \times \boldsymbol{b}_{0} = \left[\tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_{0} \times \boldsymbol{\kappa} \right) \cdot \nabla x \right) \right] \frac{\partial \boldsymbol{R}}{\partial x}
+ \left[\tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_{0} \times \boldsymbol{\kappa} \right) \cdot \nabla y \right) + \frac{\tilde{A}_{\parallel}}{\mathcal{J}B_{0}} \left(\hat{\boldsymbol{b}}_{0} \cdot (\nabla \times \boldsymbol{b}_{0}) \right) \right] \frac{\partial \boldsymbol{R}}{\partial y}
+ \left[\tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_{0} \times \boldsymbol{\kappa} \right) \cdot \nabla z \right) \right] \frac{\partial \boldsymbol{R}}{\partial z}.$$
(2.3)

Combining Eqs. 1.2 and 1.3, we arrive at an expression for \hat{B} in contravariant form:

$$\tilde{\boldsymbol{B}} = \left\{ g^{11} f_1 + g^{13} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \nabla x \right) \right\} \frac{\partial \boldsymbol{R}}{\partial x} \\
+ \left\{ g^{21} f_1 + g^{23} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \nabla y \right) + \frac{\tilde{A}_{\parallel}}{\mathcal{J} B_0} \left(\hat{\boldsymbol{b}}_0 \cdot (\nabla \times \boldsymbol{b}_0) \right) \right\} \frac{\partial \boldsymbol{R}}{\partial y} \\
+ \left\{ g^{31} f_1 + g^{33} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \nabla z \right) \right\} \frac{\partial \boldsymbol{R}}{\partial z}. \tag{2.4}$$

3 Lowest Order Line Tracing Equations

The line-tracing equation depends upon $\tilde{\boldsymbol{B}} \cdot \nabla x$, $\tilde{\boldsymbol{B}} \cdot \nabla y$, and $\tilde{\boldsymbol{B}} \cdot \nabla z$. These can be read off from above:

$$\tilde{\boldsymbol{B}} \cdot \boldsymbol{\nabla} x = g^{11} f_1 + g^{13} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \boldsymbol{\nabla} x \right)$$
(3.1)

$$\tilde{\boldsymbol{B}} \cdot \boldsymbol{\nabla} y = g^{21} f_1 + g^{23} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \boldsymbol{\nabla} y \right) + \frac{\tilde{A}_{\parallel}}{\mathcal{I} B_0} \left(\hat{\boldsymbol{b}}_0 \cdot (\boldsymbol{\nabla} \times \boldsymbol{b}_0) \right)$$
(3.2)

$$\tilde{\boldsymbol{B}} \cdot \nabla z = g^{31} f_1 + g^{33} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \nabla z \right)$$
(3.3)

where

$$g^{11} = (RB_{\theta})^{2}$$

$$g^{12} = 0$$

$$g^{13} = -I(RB_{\theta})^{2}$$

$$g^{22} = 1/h_{\theta}^{2}$$

$$g^{23} = \nu/h_{\theta}^{2}$$

$$g^{33} = I^{2}(RB_{\theta})^{2} + B_{0}^{2}/(RB_{\theta})^{2}$$

and

$$f_{1} = -B_{0}^{-1} \left[g^{13} \frac{\partial \tilde{A}_{\parallel}}{\partial x} + g^{23} \frac{\partial \tilde{A}_{\parallel}}{\partial y} + g^{33} \frac{\partial \tilde{A}_{\parallel}}{\partial z} \right]$$

$$f_{3} = B_{0}^{-1} \left[g^{11} \frac{\partial \tilde{A}_{\parallel}}{\partial x} + g^{21} \frac{\partial \tilde{A}_{\parallel}}{\partial y} + g^{31} \frac{\partial \tilde{A}_{\parallel}}{\partial z} \right].$$

Line Tracing Equations The field line equations, to lowest order in $\frac{|\tilde{B}|}{|B_0|}$, are given by

$$\frac{dx}{dy} \cong \left[g^{11} f_1 + g^{13} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \boldsymbol{\nabla} x \right) \right] \frac{h_{\theta}}{B_{\theta}}$$
(3.4)

$$\frac{dz}{dy} \cong \left[g^{31} f_1 + g^{33} f_3 + \tilde{A}_{\parallel} \left(\left(\hat{\boldsymbol{b}}_0 \times \boldsymbol{\kappa} \right) \cdot \boldsymbol{\nabla} z \right) \right] \frac{h_{\theta}}{B_{\theta}}. \tag{3.5}$$