Introduction to Eigenvalue Solver for ITG and Gyrofluid Simulation using BOUT++ GLF code in Core Region

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Introduction to Eigenvalue Solver for ITG

Governing equations for ITG

(Ottaviani et. al., PoP '99)

Vorticity equation

$$\frac{\partial \Omega_{i}}{\partial t} + \mathbf{V}_{E} \cdot \nabla \Omega_{i} = -n \nabla_{\parallel} V_{\parallel} + n \left(\mathbf{V}_{E} + \mathbf{V}_{p_{i}} \right) \cdot \left(\mathbf{\kappa} + \nabla \ln B \right) + n \mathbf{V}_{p_{i}} \cdot \nabla \left(\frac{n_{1} - \Omega_{i}}{n} \right) - \mathbf{V}_{E} \cdot \nabla n_{eq} + D_{c} \Omega_{i}$$

where
$$\Omega_i = n_1 - n_{pol}$$
: generalized vorticity, $n_{pol} = \frac{ne}{m_i \omega_{ci}^2} \nabla_{\perp}^2 \varphi$, $n_1 = n_{eq} \frac{e \varphi_1}{T_e}$: adiabatic response

Ion temperature equation

$$\mathbf{V}_{E} = \frac{c}{B} \mathbf{b} \times \nabla \varphi, \quad \mathbf{V}_{p_{i}} = \frac{c}{neB} \mathbf{b} \times \nabla p_{i}$$

$$\frac{\partial T_i}{\partial t} + \mathbf{V}_E \cdot \nabla T_i = -\frac{2}{3} T_i \nabla_{\parallel} V_{\parallel} + D_c T_i + D_{glf} T_i$$

Ion parallel velocity equation

$$\frac{\partial V_{\parallel}}{\partial t} + \mathbf{V}_{E} \cdot \nabla V_{\parallel} = -\frac{e}{m_{i}} \nabla_{\parallel} \varphi - \frac{1}{m_{i} n} \nabla_{\parallel} p_{i} + D_{c} V_{\parallel}$$

where
$$D_c F = \mu_1 \nabla_{\perp}^2 F - \mu_2 \nabla_{\perp}^4 F$$
: viscous damping, $D_{glf} T_i = -\sqrt{\frac{8T_{i,eq}}{\pi m_i}} |\nabla_{\parallel}| T_{i1}$: Landau damping.

Linearized equations

$$\begin{split} &\frac{\partial \widetilde{\Omega}_{i}}{\partial \bar{t}} = -\hat{n}_{i} \hat{\nabla}_{\parallel} \widetilde{V}_{i\parallel} + i \Big(\hat{n}_{i} \omega_{di} \widetilde{\varphi} + \omega_{di} \widetilde{p}_{i} \Big) + \left[\frac{\hat{p}_{i}}{\rho_{*}}, \widetilde{\nabla}_{\perp}^{2} \widetilde{\varphi} \right] - \left[\widetilde{\Psi}, \frac{\hat{n}_{i}}{\rho_{*}} \right] + \overline{D}_{c} \, \widetilde{\Omega}_{i} \\ &\frac{\partial \widetilde{T}_{i}}{\partial \bar{t}} + \left[\widetilde{\Psi}, \frac{\hat{T}_{i}}{\rho_{*}} \right] = -\frac{2}{3} \, \hat{T}_{i} \nabla_{\parallel} \widetilde{V}_{i\parallel} + \overline{D}_{c} \, \widetilde{T}_{i} + \overline{D}_{glf} \, \widetilde{T}_{i} \\ &\frac{\partial \widetilde{V}_{i\parallel}}{\partial \bar{t}} = -\hat{\nabla}_{\parallel} \widetilde{\varphi} - \frac{1}{\hat{n}_{i}} \, \hat{\nabla}_{\parallel} \, \widetilde{p}_{i} + \overline{D}_{c} \, \widetilde{V}_{i\parallel} \end{split}$$

where

$$\widetilde{\Omega}_i = \widetilde{n}_i - \hat{n}_i \widetilde{\nabla}_{\perp}^2 \widetilde{\varphi}, \quad \widetilde{n}_i = rac{\hat{n}_i}{\hat{T}_a} \widetilde{\varphi},$$

$$\widetilde{\Psi} = \Gamma_0^{1/2} \widetilde{\varphi} = \left(1 - \frac{1}{2} \widetilde{\nabla}_{\perp}^2\right)^{-1} \widetilde{\varphi}$$

is gyroaveraged potential.

following normalizations are used,

$$F = F_{eq} + F_{1}, \quad F = (\Omega, V_{\parallel}, n, T, \varphi), \quad F_{0} = (n_{0}, c_{s0}, n_{0}, T_{0}, \frac{T_{0}}{e})$$

$$\hat{F} = \frac{F_{eq}}{F_{0}}, \quad \tilde{F} = \frac{F_{1}}{\rho_{*}F_{0}}, \quad \tilde{\nabla}_{\perp} = \rho_{*}a\nabla_{\perp}, \quad \hat{\nabla}_{\parallel} = a\nabla_{\parallel}, \quad \rho = \frac{r}{a}, \quad \tilde{t} = \frac{t}{t_{0}}$$

$$\rho_{*} = \frac{\rho_{s0}}{a}, \quad \rho_{s0} = \frac{c_{s0}}{\omega_{c0}}, \quad c_{s0} = \sqrt{\frac{T_{0}}{m_{i}}}, \quad \omega_{c0} = \frac{e_{i}B_{0}}{m_{i}c}, \quad t_{0} = \frac{a}{c_{s0}}$$

$$[f, g] \equiv \rho_{s0}^{2}\mathbf{b} \times \nabla f \cdot \nabla g = \frac{\rho_{*}^{2}}{\rho} \left(\frac{\partial f}{\partial \rho} \frac{\partial g}{\partial \theta} - \frac{\partial g}{\partial \rho} \frac{\partial f}{\partial \theta}\right)$$

$$i\omega_{di}\tilde{F} \equiv 2\rho_{s0}a\mathbf{b} \times \nabla \tilde{F} \cdot \nabla \ln B = 2\rho_{*} \frac{a}{R} \left(\frac{1}{\rho}\cos\theta \frac{\partial \tilde{F}}{\partial \theta} + \sin\theta \frac{\partial \tilde{F}}{\partial \rho}\right)$$

$$Note Fig. 12.$$

Note: F_0 is constant (not the value on magnetic axis)

$$\frac{\partial \widetilde{\Omega}_{i}}{\partial \bar{t}} = -\hat{n}_{i} \hat{\nabla}_{\parallel} \widetilde{V}_{i\parallel} + i \left(\hat{n}_{i} \omega_{di} \widetilde{\varphi} + \omega_{di} \widetilde{p}_{i}\right) + \left[\frac{\hat{p}_{i}}{\rho_{*}}, \widetilde{\nabla}_{\perp}^{2} \widetilde{\varphi}\right] + \left[\frac{\hat{n}_{i}}{\rho_{*}}, \widetilde{\Psi}\right] + \overline{D}_{c} \widetilde{\Omega}_{i}$$

$$\omega M\widetilde{\varphi} = \hat{n}_{i}k_{\parallel}\widetilde{V}_{i\parallel} - a_{W}\widetilde{\varphi} - b_{W}\widetilde{T}_{i} + G_{\hat{p}_{i}}L\widetilde{\varphi} + G_{\hat{n}_{i}}J_{0}\widetilde{\varphi} + i\overline{D}_{c}M\widetilde{\varphi}$$

$$\frac{\partial \widetilde{T}_{i}}{\partial \bar{t}} = \left[\frac{\hat{T}_{i}}{\rho_{*}}, \widetilde{\Psi}\right] - \frac{2}{3}\hat{T}_{i}\nabla_{\parallel}\widetilde{V}_{i\parallel} + \overline{D}_{c}\,\widetilde{T}_{i} + \overline{D}_{glf}\widetilde{T}_{i}$$

$$\frac{\partial \widetilde{V}_{i\parallel}}{\partial \bar{t}} = -\hat{\nabla}_{\parallel} \widetilde{\varphi} - \frac{1}{\hat{n}_{i}} \hat{\nabla}_{\parallel} \widetilde{p}_{i} + \overline{D}_{c} \widetilde{V}_{i\parallel}$$

where we introduce

$$rac{\partial}{\partial ar{t}} \equiv -i\omega, \quad \ \, \widetilde{
abla}_{\perp}^2 \widetilde{arphi} \equiv L \widetilde{arphi}, \quad \, \hat{
abla}_{\parallel} \equiv -ik_{\parallel}$$

$$\widetilde{\Omega}_{i} = \left(\frac{\hat{n}_{i}}{\hat{T}_{e}} - \hat{n}_{i}\widetilde{\nabla}_{\perp}^{2}\right)\widetilde{\varphi} \equiv M\widetilde{\varphi}$$

$$\widetilde{\Psi} = \left(1 - \frac{1}{2}\widetilde{\nabla}_{\perp}^{2}\right)^{-1}\widetilde{\varphi} \equiv J_{0}\widetilde{\varphi}$$

$$\begin{bmatrix} \hat{F} \\ \rho_* \end{bmatrix} = \rho_{s0}^2 \mathbf{b} \times \nabla \frac{\hat{F}}{\rho_*} \cdot \nabla \equiv -iG_{\hat{F}}$$
$$i(\hat{n}_i \omega_{di} \widetilde{\varphi} + \omega_{di} \widetilde{p}_i) \equiv i(a_W \widetilde{\varphi} + b_W \widetilde{T}_i)$$

and use

$$\widetilde{p}_i = rac{\hat{p}_i}{\hat{T}_e}\widetilde{arphi}_i + \hat{n}_i\widetilde{T}_i.$$

$$\frac{\partial \widetilde{\Omega}_{i}}{\partial \bar{t}} = -\hat{n}_{i} \hat{\nabla}_{\parallel} \widetilde{V}_{i\parallel} + i \left(\hat{n}_{i} \omega_{di} \widetilde{\varphi} + \omega_{di} \widetilde{p}_{i}\right) + \left[\frac{\hat{p}_{i}}{\rho_{*}}, \widetilde{\nabla}_{\perp}^{2} \widetilde{\varphi}\right] + \left[\frac{\hat{n}_{i}}{\rho_{*}}, \widetilde{\Psi}\right] + \overline{D}_{c} \widetilde{\Omega}_{i}$$

$$\frac{\partial \widetilde{T}_{i}}{\partial \bar{t}} = \left[\frac{\hat{T}_{i}}{\rho_{*}}, \widetilde{\Psi}\right] - \frac{2}{3}\hat{T}_{i}\nabla_{\parallel}\widetilde{V}_{i\parallel} + \overline{D}_{c}\,\widetilde{T}_{i} + \overline{D}_{glf}\widetilde{T}_{i}$$

$$\frac{\partial \widetilde{V}_{i\parallel}}{\partial \bar{t}} = -\hat{\nabla}_{\parallel} \widetilde{\varphi} - \frac{1}{\hat{n}_{i}} \hat{\nabla}_{\parallel} \widetilde{p}_{i} + \overline{D}_{c} \, \widetilde{V}_{i\parallel}$$

where we introduce

$$rac{\partial}{\partial ar{t}} \equiv -i\omega, \quad \ \, \widetilde{
abla}_{\perp}^2 \widetilde{arphi} \equiv L \widetilde{arphi}, \quad \, \hat{
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$$\widetilde{\Omega}_{i} = \left(\frac{\hat{n}_{i}}{\hat{T}_{e}} - \hat{n}_{i}\widetilde{\nabla}_{\perp}^{2}\right)\widetilde{\varphi} \equiv M\widetilde{\varphi}$$

$$\widetilde{\Psi} = \left(1 - \frac{1}{2}\widetilde{\nabla}_{\perp}^{2}\right)^{-1}\widetilde{\varphi} \equiv J_{0}\widetilde{\varphi}$$

$$\begin{bmatrix} \hat{F} \\ \rho_* \end{bmatrix} = \rho_{s0}^2 \mathbf{b} \times \nabla \frac{\hat{F}}{\rho_*} \cdot \nabla \equiv -iG_{\hat{F}}$$
$$i(\hat{n}_i \omega_{di} \tilde{\varphi} + \omega_{di} \tilde{p}_i) \equiv i(a_W \tilde{\varphi} + b_W \tilde{T}_i)$$

and use

$$\widetilde{p}_i = rac{\hat{p}_i}{\hat{T}_e}\widetilde{arphi}_i + \hat{n}_i\widetilde{T}_i.$$

$$\frac{\partial \widetilde{\Omega}_{i}}{\partial \bar{t}} = -\hat{n}_{i} \hat{\nabla}_{\parallel} \widetilde{V}_{i\parallel} + i \left(\hat{n}_{i} \omega_{di} \widetilde{\varphi} + \omega_{di} \widetilde{p}_{i}\right) + \left[\frac{\hat{p}_{i}}{\rho_{*}}, \widetilde{\nabla}_{\perp}^{2} \widetilde{\varphi}\right] + \left[\frac{\hat{n}_{i}}{\rho_{*}}, \widetilde{\Psi}\right] + \overline{D}_{c} \widetilde{\Omega}_{i}$$

Those operators can be represented as matrices in a functional space.

$$\frac{\partial \widetilde{T}_{i}}{\partial \overline{t}} = \left[\frac{\hat{T}_{i}}{\rho_{*}}, \widetilde{\Psi}\right] - \frac{2}{3}\hat{T}_{i}\nabla_{\parallel}\widetilde{V}_{i\parallel} + \overline{D}_{c}\,\widetilde{T}_{i} + \overline{D}_{glf}\widetilde{T}_{i}$$

$$\widehat{\boldsymbol{\omega}} \widetilde{\boldsymbol{T}}_{i} = G_{\hat{T}_{i}} \boldsymbol{J}_{0} \widetilde{\boldsymbol{\varphi}}_{i} + \widehat{\boldsymbol{D}}_{c} + \widehat{\boldsymbol{D}}_{c} \widehat{\boldsymbol{J}}_{i} \widehat{\boldsymbol{T}}_{i} + \frac{2}{3} \hat{T}_{i} k_{\parallel} \widetilde{\boldsymbol{V}}_{i\parallel}$$

$$\frac{\partial \widetilde{V}_{i\parallel}}{\partial \bar{t}} = -\hat{\nabla}_{\parallel} \widetilde{\varphi} - \frac{1}{\hat{n}_{i}} \hat{\nabla}_{\parallel} \widetilde{p}_{i} + \overline{D}_{c} \, \widetilde{V}_{i\parallel}$$

where we introduce
$$\frac{\partial}{\partial \bar{t}} \equiv -i\omega, \quad \tilde{\nabla}_{\perp}^{2} \tilde{\varphi} \equiv L \tilde{\varphi}, \quad \hat{\nabla}_{\parallel} \equiv -k_{\parallel}$$

$$\tilde{\Omega}_{i} = \left(\frac{\hat{n}_{i}}{\hat{T}_{e}} - \hat{n}_{i} \tilde{\nabla}_{\perp}^{2}\right) \tilde{\varphi} \equiv M \tilde{\varphi}$$

$$\tilde{\Psi} = \left(1 - \frac{1}{2} \tilde{\nabla}_{\perp}^{2}\right)^{-1} \tilde{\varphi} \equiv J_{0} \tilde{\varphi}$$

$$\left[\frac{\hat{F}}{\rho_{*}}, \right] = \rho_{s0}^{2} \mathbf{b} \times \nabla \frac{\hat{F}}{\rho_{*}} \cdot \nabla \equiv -k G_{\hat{F}}$$

$$i(\hat{n}_{i} \omega_{di} \tilde{\varphi} + \omega_{di} \tilde{p}_{i}) \equiv i(a_{W} \tilde{\varphi} + b_{W} \tilde{T}_{i})$$

and use

$$\widetilde{p}_i = rac{\hat{p}_i}{\hat{T}_a}\widetilde{arphi}_i + \hat{n}_i\widetilde{T}_i.$$

Final form of eigenvalue equation

$$\omega \widetilde{\varphi} = A_{11} \widetilde{\varphi} + A_{12} \widetilde{T}_i + A_{13} \widetilde{V}_{i|}$$

$$\omega \widetilde{T}_{i} = A_{21} \widetilde{\varphi} + A_{22} \widetilde{T}_{i} + A_{23} \widetilde{V}_{i\parallel}$$



$$\omega \mathbf{F} = \mathbf{A} \cdot \mathbf{F}$$

$$\omega \widetilde{V}_{i\parallel} = A_{31} \widetilde{\varphi} + A_{32} \widetilde{T}_i + A_{33} \widetilde{V}_{i\parallel}$$

$$\mathbf{F} = \begin{pmatrix} \widetilde{\boldsymbol{\varphi}} \\ \widetilde{\mathbf{T}}_{\mathbf{i}} \\ \widetilde{\mathbf{V}}_{\mathbf{i}||} \end{pmatrix}, \quad \mathbf{A} = \begin{pmatrix} \mathbf{A}_{11} & \mathbf{A}_{12} & \mathbf{A}_{13} \\ \mathbf{A}_{21} & \mathbf{A}_{22} & \mathbf{A}_{23} \\ \mathbf{A}_{31} & \mathbf{A}_{32} & \mathbf{A}_{33} \end{pmatrix}$$

$$\omega \widetilde{\varphi} = A_{11} \widetilde{\varphi} + A_{12} \widetilde{T}_{i} + A_{13} \widetilde{V}_{i\parallel}$$

$$\omega \widetilde{T}_{i} = A_{21} \widetilde{\varphi} + A_{22} \widetilde{T}_{i} + A_{23} \widetilde{V}_{i\parallel}$$

$$\omega \widetilde{F} = A \cdot F$$

$$\widetilde{\varphi} = \sum_{k} \widetilde{\varphi}_{k} |k\rangle = \begin{pmatrix} \widetilde{\varphi}_{11} & A_{12} \\ A_{21} & A_{22} \\ A_{31} & A_{32} \end{pmatrix}$$

$$\widetilde{\varphi} = \sum_{k} \widetilde{\varphi}_{k} |k\rangle = \begin{pmatrix} \widetilde{\varphi}_{1} \\ \widetilde{\varphi}_{2} \\ \widetilde{\varphi}_{3} \\ \vdots \end{pmatrix}, \dots$$

$$|k\rangle \text{ is a basis function}$$

 $|k\rangle$ is a basis function

where

$$A_{11} = M^{-1} \left(-a_W + G_{\hat{p}_i} L + G_{\hat{n}_i} J_0 + i \overline{D}_c M \right), \qquad A_{12} = -M^{-1} b_w, \qquad A_{13} = M^{-1} \hat{n}_i k_\parallel$$

$$A_{12} = -M^{-1}b_w, \quad A_{13} = M^{-1}\hat{n}_i k$$

$$A_{21}=G_{\hat{T}_i}J_0,$$

$$A_{21} = G_{\hat{T}_i} J_0, \qquad \qquad A_{22} = i \overline{D}_c \ + i \overline{D}_{glf}, \qquad \qquad A_{23} = rac{2}{3} \hat{T}_i k_{\parallel}$$

$$A_{23} = \frac{2}{3}\hat{T}_i k$$

$$A_{31} = \left(1 + \frac{\hat{T}_i}{\hat{T}_e}\right) k_{\parallel}, \qquad A_{32} = k_{\parallel},$$

$$A_{33} = i\overline{D}_c$$

Basis functions

Field F can represented as $F(\rho, \theta, \zeta, t) = e^{-i\omega t} \sum_{k=mnp} F_k W_k(\rho) e^{i(m\theta-n\zeta)}$ in concentric - circular geometry:

$$\mathbf{F} = |F\rangle = e^{-i\omega t} \sum_{k} F_{k} |k\rangle, \quad |k\rangle = |mnp\rangle = W_{mnp}(\rho) e^{i(m\theta - n\zeta)},$$

inner product:

$$\langle k | F \rangle \equiv \frac{1}{4\pi^2} \int_0^1 \rho d\rho \int_0^{2\pi} d\theta \int_0^{2\pi} d\zeta W_{mnp}(\rho) e^{-i(m\theta - n\zeta)} F, \qquad \langle k | k' \rangle = \delta_{kk'}$$

Radial basis function

For Hermite basis,

$$W_{mnp}(\rho) = \frac{1}{\sqrt{2\rho w_{mn}} v_p} H_p(x) e^{-x^2/2}$$

where
$$x = \frac{\rho - \rho_{mn}}{w_{mn}} \left(\rho_{mn} \text{ is radial position of rational surface satisfying } q = \frac{m}{n} \right), \quad v_p = 2^{p/2} \Gamma(p+1)^{1/2} \pi^{1/4}$$

For Bessel basis,

$$W_{mnp}(\rho) = \frac{\sqrt{2}}{J_{m+1}(\alpha_{mp})} J_m(\alpha_{mp}\rho) \text{ where } J_m(\alpha_{mp}) = 0$$

Matrix (M, J0, Dc) involving Laplacian L

For Bessel basis,

$$L_{kk'} = \left\langle k \left| \widetilde{\nabla}_{\perp}^{2} \right| k' \right\rangle = \rho_{*}^{2} \left\langle k \left| \left(\frac{\partial^{2}}{\partial \rho^{2}} + \frac{1}{\rho} \frac{\partial}{\partial \rho} - \frac{m'^{2}}{\rho^{2}} \right) \right| k' \right\rangle = -k_{\perp}^{2} \delta_{kk'}$$

where $k_{\perp} = \rho_* \alpha_{mp}$

$$M_{kk'} pprox rac{\hat{n}_i}{\hat{T}_e} igg|_{
ho =
ho_{mn}} \delta_{kk'} - \hat{n}_i igg|_{
ho =
ho_{mn}} L_{kk'}$$

$$J_{0,kk'} = \left(1 - \frac{1}{2} L_{kk'}\right)^{-1}$$

$$\overline{D}_{c,kk'} = \overline{\mu}_1 L_{kk'} - \overline{\mu}_2 L_{kk'}^2$$

$$L\widetilde{\varphi} = \widetilde{\nabla}^2 \widetilde{\varphi}$$

$$egin{aligned} L\widetilde{arphi} &= \widetilde{
abla}_{\perp}^2 \widetilde{arphi} \ \\ M\widetilde{arphi} &= igg(rac{\hat{n}_i}{\hat{T}_e} - \hat{n}_i \widetilde{
abla}_{\perp}^2igg) \widetilde{arphi} &= \widetilde{\Omega}_i \end{aligned}$$

$$\overline{D}_{c}F = \overline{\mu}_{1}\widetilde{\nabla}_{\perp}^{2}F - \overline{\mu}_{2}\widetilde{\nabla}_{\perp}^{4}F$$

Matrix (k_{||}, D_{g|f}) involving parallel wavenumber

$$\hat{F}k_{\parallel,kk'} = i\langle k | \hat{F}\hat{\nabla}_{\parallel} | k' \rangle = \delta_{mm'}\delta_{nn'}\frac{a}{R}\int_{0}^{1} \hat{F}(\rho)\left(\frac{m}{q(\rho)} - n\right)W_{mnp}(\rho)W_{mnp'}(\rho)\rho d\rho$$

$$\overline{D}_{glf,kk'} = -\left\langle k \left| \sqrt{\frac{8\hat{T}_{i}}{\pi}} \right| \hat{\nabla}_{\parallel} \right| k' \right\rangle = -\delta_{mm'} \delta_{nn'} \sqrt{\frac{8}{\pi}} \frac{a}{R} \int_{0}^{1} \sqrt{\hat{T}_{i}(\rho)} \left| \frac{m}{q(\rho)} - n \right| W_{mnp}(\rho) W_{mnp'}(\rho) \rho d\rho$$

Matrix (G) for ExB and diamagnetic drifts

$$G_{\hat{F}}\widetilde{arphi}=iiggl[rac{\hat{F}}{
ho_*},\widetilde{arphi}iggr]:$$

$$G_{\hat{F},kk'} = i \langle k \left| \frac{\hat{F}}{\rho_*}, \right| k' \rangle = i \rho_* \langle k \left| \frac{1}{\rho} \frac{\partial \hat{F}}{\partial \rho} \frac{\partial}{\partial \theta} \right| k' \rangle = -\delta_{mm'} \delta_{nn'} \rho_* m \int_0^1 \frac{\partial \hat{F}}{\partial \rho} W_{mnp}(\rho) W_{mnp'}(\rho) d\rho$$

Matrix (aw, bw) involving curvature

$$\begin{split} \hat{\beta}\omega_{di}\hat{\alpha}\big|k'\big\rangle &= -2i\rho_*\,\frac{a}{R}\,\hat{\beta}\bigg[\frac{1}{\rho}\cos\theta\,\frac{\partial}{\partial\theta} + \sin\theta\,\frac{\partial}{\partial\rho}\bigg]\hat{\alpha}\big|k'\big\rangle \\ &= \rho_*\,\frac{a}{R}\bigg[\bigg(\frac{m'}{\rho}\,\hat{\beta}\hat{\alpha}W_{k'} + \hat{\beta}\,\frac{\partial\hat{\alpha}W_{k'}}{\partial\rho}\bigg)e^{i(m'-1)\theta} + \bigg(\frac{m'}{\rho}\,\hat{\beta}\hat{\alpha}W_{k'} - \hat{\beta}\,\frac{\partial\hat{\alpha}W_{k'}}{\partial\rho}\bigg)e^{i(m'+1)\theta}\bigg]e^{-in'\zeta} \\ &\Longrightarrow \langle k\,\big|\hat{\beta}\omega_{di}\hat{\alpha}\big|k'\big\rangle &= \rho_*\,\frac{a}{R}\Bigg[\bigg\{(m+1)\int_0^1\hat{\beta}\hat{\alpha}W_{mnp}W_{m+1,n,p'}d\rho + \int_0^1\hat{\beta}W_{mnp}\frac{\partial\hat{\alpha}W_{m+1,n,p'}}{\partial\rho}\,\rho d\rho\bigg\}\delta_{m+1,m'}\delta_{nn'} \\ &+ \bigg\{(m-1)\int_0^1\hat{\beta}\hat{\alpha}W_{mnp}W_{m-1,n,p'}d\rho - \int_0^1\hat{\beta}W_{mnp}\frac{\partial\hat{\alpha}W_{m-1,n,p'}}{\partial\rho}\,\rho d\rho\bigg\}\delta_{m-1,m'}\delta_{nn'}\bigg] \end{split}$$



$$a_{W,kk'} = \left\langle k \left| \hat{n}_{i} \omega_{di} + \omega_{di} \frac{\hat{p}_{i}}{\hat{T}_{e}} \right| k' \right\rangle = a_{W}^{+}(k, p') \delta_{m+1,m'} \delta_{nn'} + a_{W}^{-}(k, p') \delta_{m-1,m'} \delta_{nn'}$$

$$a_{W} \widetilde{\varphi} = \hat{n}_{i} \omega_{di} \widetilde{\varphi} + \omega_{di} \frac{\hat{p}_{i}}{\hat{T}_{e}} \widetilde{\varphi}$$

$$b_{W,kk'} = \left\langle k \left| \omega_{di} \hat{n}_{i} \right| k' \right\rangle = b_{W}^{+}(k, p') \delta_{m+1,m'} \delta_{nn'} + b_{W}^{-}(k, p') \delta_{m-1,m'} \delta_{nn'}$$

$$b_{W} \widetilde{T}_{i} = \omega_{di} \hat{n}_{i} \widetilde{T}_{i}$$

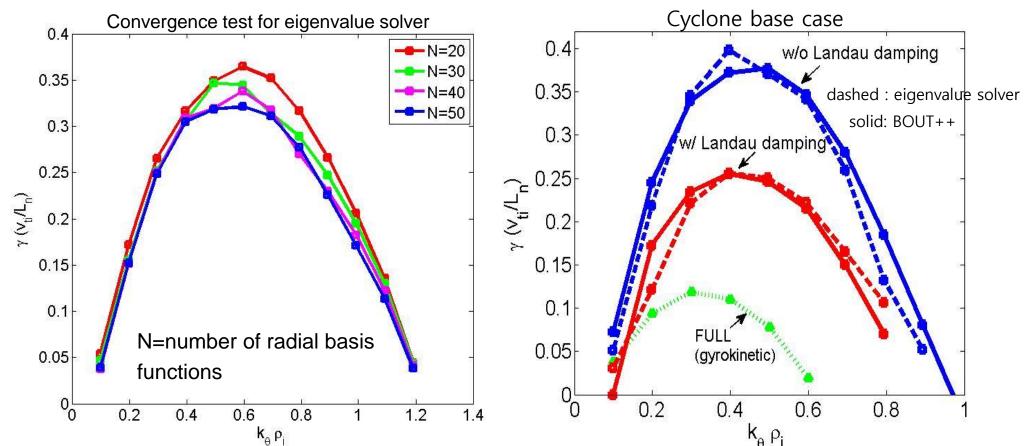
where

$$a_{W}^{\pm}(k,p') = \rho_{*} \frac{a}{R} \left[\int_{0}^{1} \hat{n}_{i} \left\{ (m\pm 1)(1+\tau) \pm \left(\frac{\partial \ln \hat{n}_{i}}{\partial \rho} + \frac{\partial \ln \tau}{\partial \rho} \right) \tau \right\} W_{mnp} W_{m\pm 1,n,p'} d\rho \pm \int_{0}^{1} \hat{n}_{i} (1+\tau) W_{mnp} \frac{\partial W_{m\pm 1,n,p'}}{\partial \rho} \rho d\rho \right]$$

$$b_{W}^{\pm}(k,p') = \rho_{*} \frac{a}{R} \left[\int_{0}^{1} \left(m\pm 1 + \frac{\partial \ln \hat{n}_{i}}{\partial \rho} \right) \hat{n}_{i} W_{mnp} W_{m\pm 1,n,p'} d\rho \pm \int_{0}^{1} \hat{n}_{i} W_{mnp} \frac{\partial W_{m\pm 1,n,p'}}{\partial \rho} \rho d\rho \right], \quad \tau = \frac{\hat{T}_{i}}{\hat{T}_{e}}$$

Comparison of BOUT++ with eigenvalue solver

- In eigenvalue solver, $W_{mnp} = \frac{\sqrt{2}}{J_{m+1}(\alpha_{mp})} J_m \left(\frac{\alpha_{mp}}{a}r\right) e^{i(m\theta n\zeta)} \left(p = 1 \sim N\right)$ is used as basis function.
- The fluid equations are then projected on to the set of basis functions.
- Eigenvalues are obtained by using matlab.
- BOUT++ results agree well with eigenvalues, which means BOUT++ correctly solves the given equations.



Gyrofluid Simulation using BOUT++ GLF Code in Core Region

Gyrofluid equation for ITG

(M. A. Beer and G. W. Hammett, PoP '96)

Guiding center density equation

$$\frac{\partial \widetilde{n}_{i}}{\partial t} + \mathbf{V}_{\widetilde{\Phi}} \cdot \nabla \widetilde{n}_{i} + \frac{n_{0}}{2T_{i0}} \left[\hat{\nabla}_{\perp}^{2} \mathbf{V}_{\widetilde{\Phi}} \right] \cdot \nabla \widetilde{T}_{i\perp} + B \nabla_{\parallel} \frac{n_{0} \widetilde{V}_{i\parallel}}{B} - \frac{e n_{0}}{T_{i0}} \left(1 + \frac{1}{2} \eta_{i} \hat{\nabla}_{\perp}^{2} \right) i \omega_{*} \widetilde{\Phi} + \frac{e n_{0}}{T_{i0}} \left(2 + \frac{1}{2} \hat{\nabla}_{\perp}^{2} \right) i \omega_{d} \widetilde{\Phi} + \frac{1}{T_{i0}} i \omega_{d} \left(\widetilde{p}_{i\parallel} + \widetilde{p}_{i\perp} \right) = 0$$

Guiding center parallel velocity equation

$$\frac{\partial \widetilde{V}_{i\parallel}}{\partial t} + \mathbf{V}_{\widetilde{\Phi}} \cdot \nabla \widetilde{V}_{i\parallel} + \frac{e}{m_i} \nabla_{_{\parallel}} \widetilde{\Phi} + \frac{B}{m_i n_0} \nabla_{_{\parallel}} \frac{\widetilde{p}_{i\parallel}}{B} + \left(\frac{\widetilde{T}_{i\perp}}{m_i} + \frac{e}{2m_i} \hat{\nabla}_{_{\perp}}^2 \widetilde{\Phi} \right) \nabla_{_{\parallel}} \ln B + \frac{1}{p_{i0}} i \omega_d \left(\widetilde{q}_{i\parallel} + \widetilde{q}_{i\perp} \right) + 4i \omega_d \widetilde{V}_{i\parallel} = 0$$

Guiding center parallel & perpendicular pressure equation

$$\begin{split} &\frac{\partial \widetilde{p}_{i\parallel}}{\partial t} + \mathbf{V}_{\widetilde{\Phi}} \cdot \nabla \widetilde{p}_{i\parallel} + \frac{n_0}{2} \left[\hat{\nabla}_{\perp}^2 \mathbf{V}_{\widetilde{\Phi}} \right] \cdot \nabla \widetilde{T}_{i\perp} + B \nabla_{\parallel} \frac{1}{B} \left(\widetilde{q}_{i\parallel} + p_{i0} \widetilde{V}_{i\parallel} \right) + 2 p_{i0} B \nabla_{\parallel} \frac{\widetilde{V}_{i\parallel}}{B} + 2 \widetilde{q}_{i\perp} \nabla_{\parallel} \ln B \\ &- e n_0 \left(1 + \eta_i + \frac{1}{2} \eta_i \hat{\nabla}_{\perp}^2 \right) i \omega_* \widetilde{\Phi} + e n_0 \left(4 + \frac{1}{2} \hat{\nabla}_{\perp}^2 \right) i \omega_d \widetilde{\Phi} + 4 i \omega_d \widetilde{p}_{i\parallel} + n_0 i \omega_d \left(3 \widetilde{T}_{i\parallel} + \widetilde{T}_{i\perp} \right) + 2 n_0 \left| \omega_d \right| \left(v_1 \widetilde{T}_{i\parallel} + v_2 \widetilde{T}_{i\perp} \right) = 0 \\ &\frac{\partial \widetilde{p}_{i\perp}}{\partial t} + \mathbf{V}_{\widetilde{\Phi}} \cdot \nabla \widetilde{p}_{i\perp} + \frac{1}{2} \left[\hat{\nabla}_{\perp}^2 \mathbf{V}_{\widetilde{\Phi}} \right] \cdot \nabla \widetilde{p}_{i\perp} + n_0 \left[\hat{\nabla}_{\perp}^2 \mathbf{V}_{\widetilde{\Phi}} \right] \cdot \nabla \widetilde{T}_{i\perp} + B^2 \nabla_{\parallel} \frac{1}{B^2} \left(\widetilde{q}_{i\perp} + p_{i0} \widetilde{V}_{i\parallel} \right) \\ &- e n_0 \left[1 + \frac{1}{2} \hat{\nabla}_{\perp}^2 + \eta_i \left(1 + \frac{1}{2} \hat{\nabla}_{\perp}^2 + \hat{\nabla}_{\perp}^2 \right) \right] i \omega_* \widetilde{\Phi} + e n_0 \left(3 + \frac{3}{2} \hat{\nabla}_{\perp}^2 + \hat{\nabla}_{\perp}^2 \right) i \omega_d \widetilde{\Phi} + 3 i \omega_d \widetilde{p}_{i\perp} + n_0 i \omega_d \left(\widetilde{T}_{i\parallel} + 2 \widetilde{T}_{i\perp} \right) + 2 n_0 \left| \omega_d \right| \left(v_3 \widetilde{T}_{i\parallel} + v_4 \widetilde{T}_{i\perp} \right) = 0 \end{split}$$

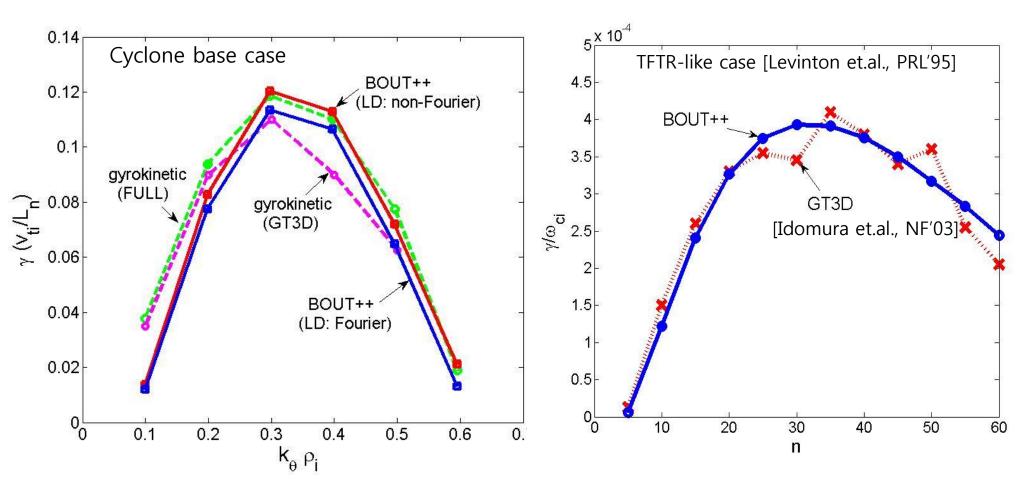
Quasi-neutrality

$$\frac{n_{0}}{T_{e0}}\widetilde{\varphi} + \frac{n_{0}}{T_{i0}}(1 - \Gamma_{0})\widetilde{\varphi} = \Gamma_{0}^{1/2}\widetilde{n}_{i} + \frac{n_{0}}{T_{i0}}b\frac{\partial\Gamma_{0}^{1/2}}{\partial b}\widetilde{T}_{i} \quad \text{where } \widetilde{\Phi} = \Gamma_{0}^{1/2}\widetilde{\varphi}, \quad \frac{1}{2}\widehat{\nabla}_{\perp}^{2}\widetilde{\Phi} = \left(b\frac{\partial\Gamma_{0}^{1/2}}{\partial b}\right)\widetilde{\varphi}, \quad \widehat{\nabla}_{\perp}^{2}\widetilde{\Phi} = b\frac{\partial^{2}}{\partial b^{2}}\left(b\Gamma_{0}^{1/2}\right)\widetilde{\varphi}, \quad \Gamma_{0}^{1/2} \cong \frac{1}{1 + b/2}$$

$$i\omega_{d} = \left(v_{t}^{2} / \omega_{c}\right)\mathbf{b} \times \nabla \ln B \cdot \nabla, \quad b = -\rho_{i}^{2}\nabla_{\perp}^{2}$$

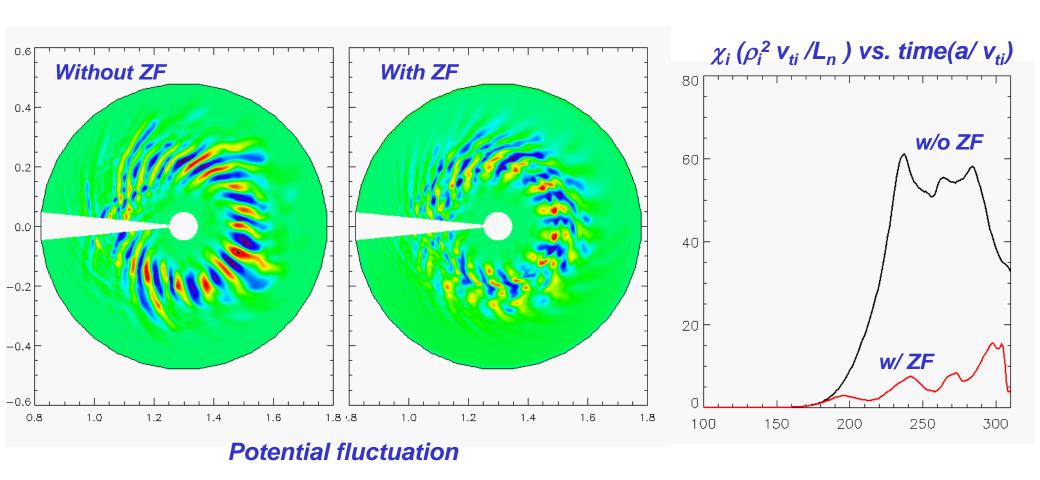
Linear benchmarks

- BOUT++ using Beer's "3+1" field model agrees well with gyrokinetic codes.
- Non-Fourier method for Landau damping shows good agreement with Fourier method.



Nonlinear simulations of ITG

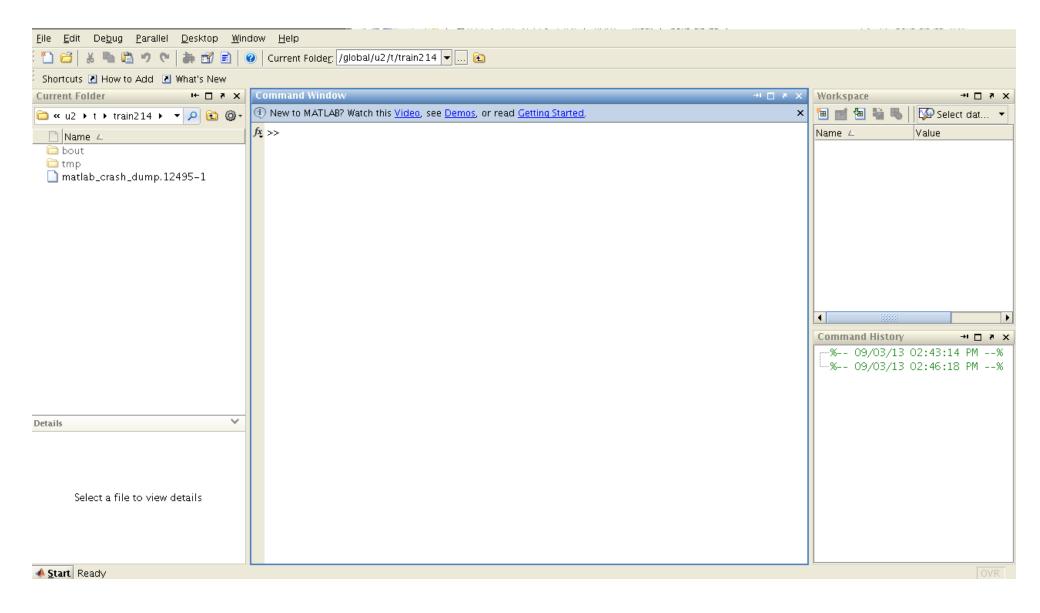
- Global nonlinear simulations using Beer model performed at fixed profile
- Turbulence suppression by zonal flow observed



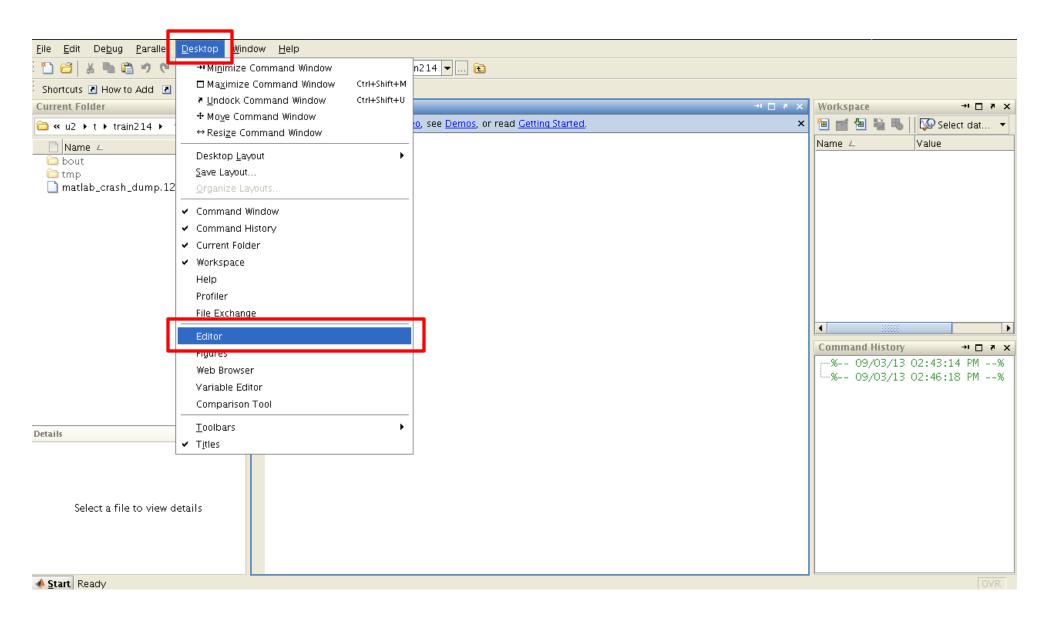
Hands-on Exercise for Eigenvalue Solver using MATLAB

- 0) Run matlab:
- > module load matlab-nofonts
- > matlab
- 1) Run eigensolver_init.m and eigensolver_ITG.m in matlab:
- >> eigensolver_init: assign equilibrium profiles, calculate radial basis functions, arrange mode index, evaluate matrix elements involving integrals
- >> eigensolver_ITG : set up matrix and solve the eigenvalue equation for ITG
- 2) Run plot_eigenvalue.m and plot_eigenmode.m in matlab:
- >> plot_eigenvalue : plot linear ITG growth rate and real frequency vs. k_theta*rho_i
- >> plot_eigenmode : choose toroidal mode number and plot eigenmode structure

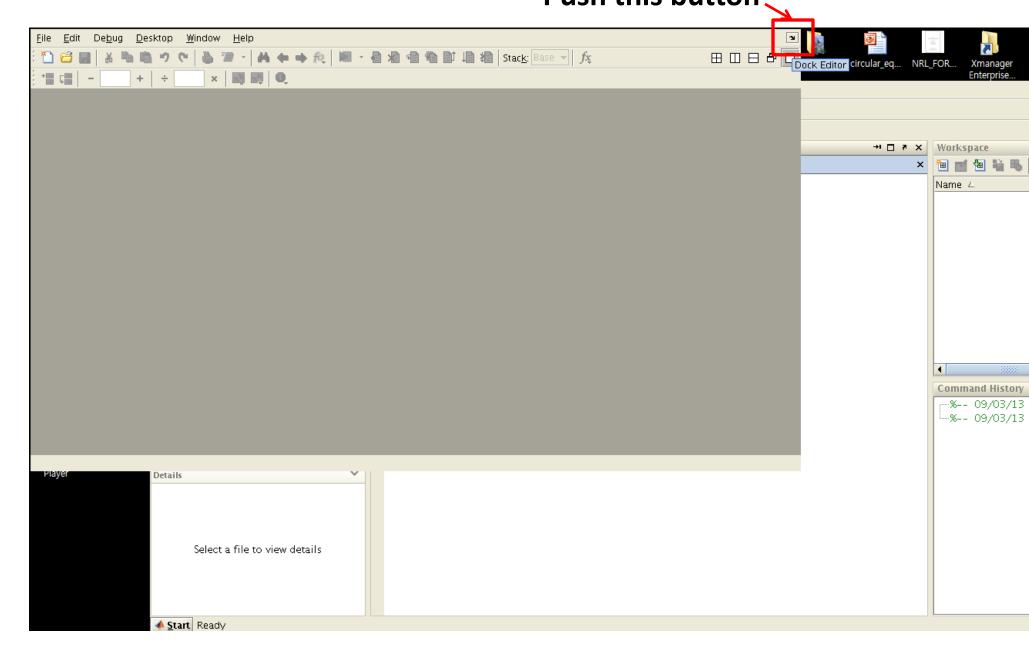
Run matlab: > module load matlab-nofonts > matlab

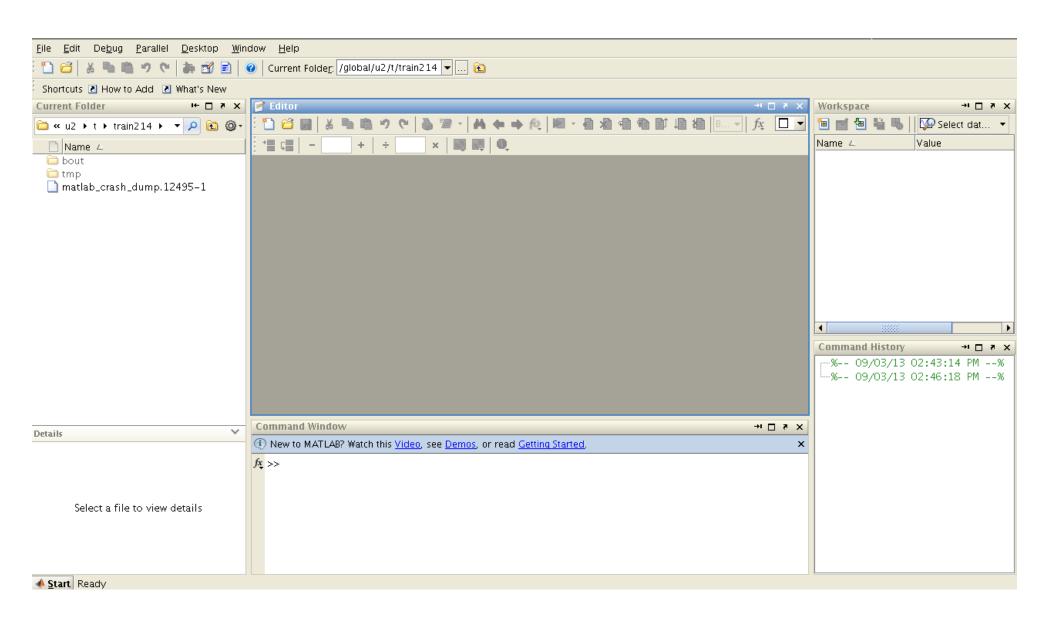


Run Editor:

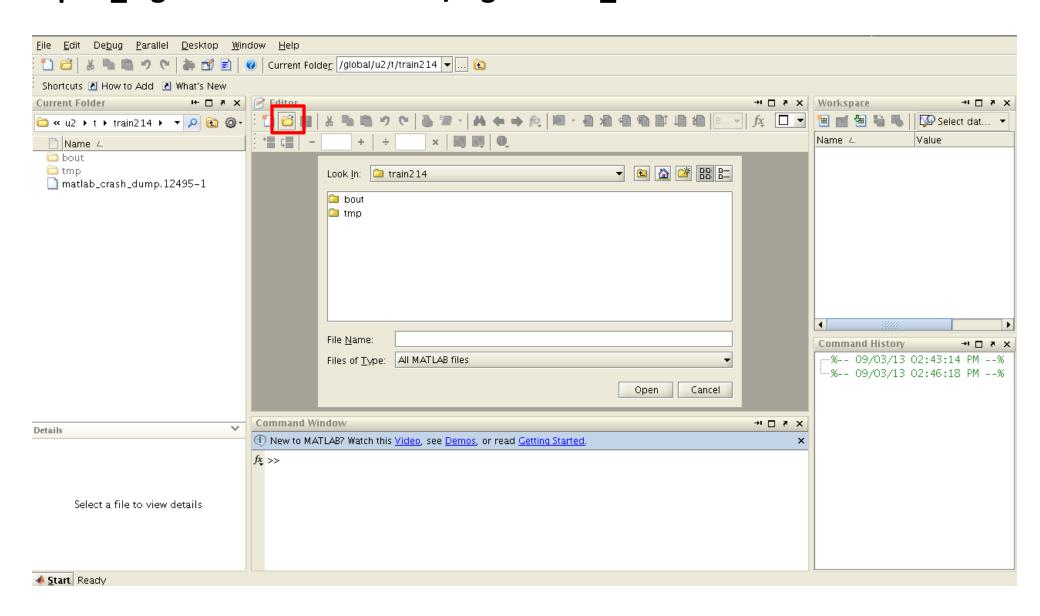


Dock Editor:
Push this button

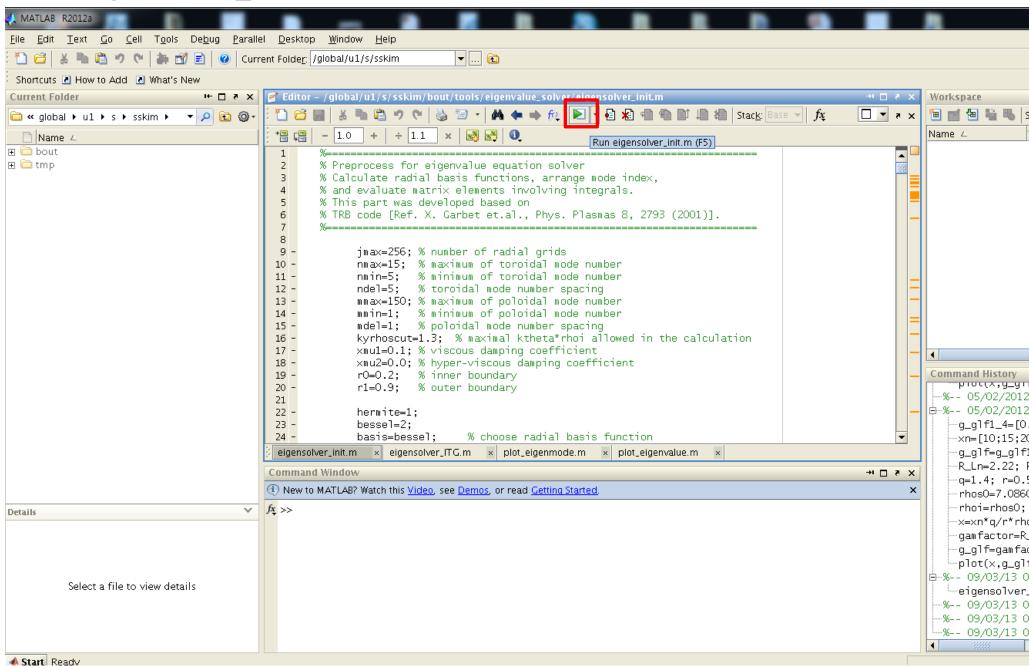


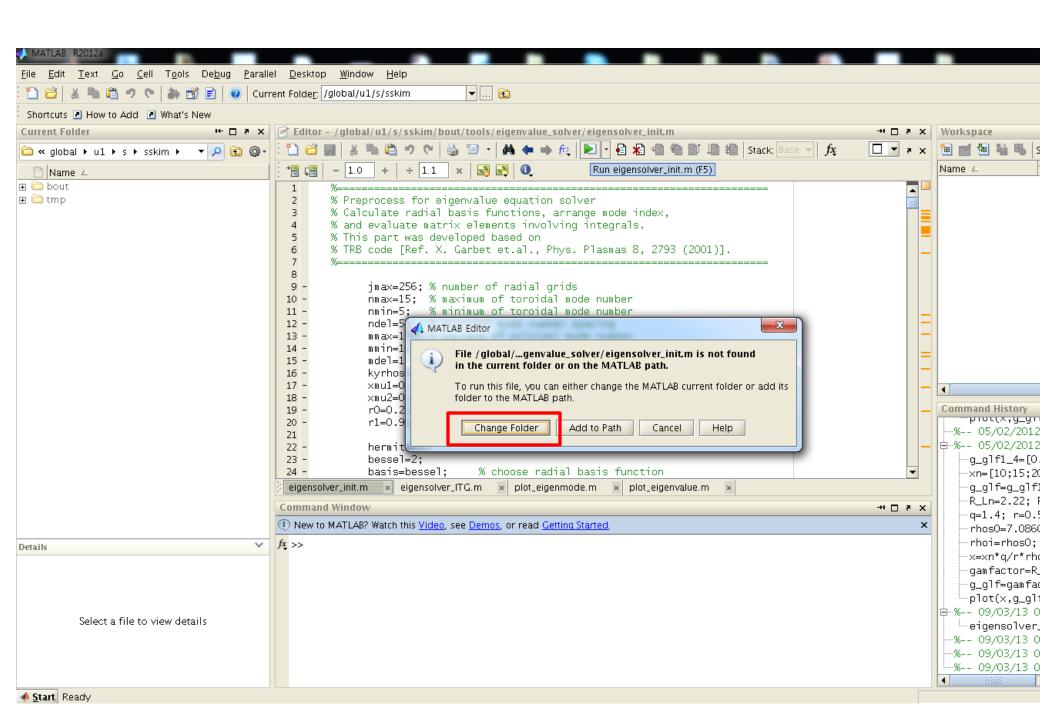


Open eigensolver_init.m, eigensolver_ITG.m, plot_eigenvalue.m, plot_eigenmode.m in tools/eigenvalue_solver.

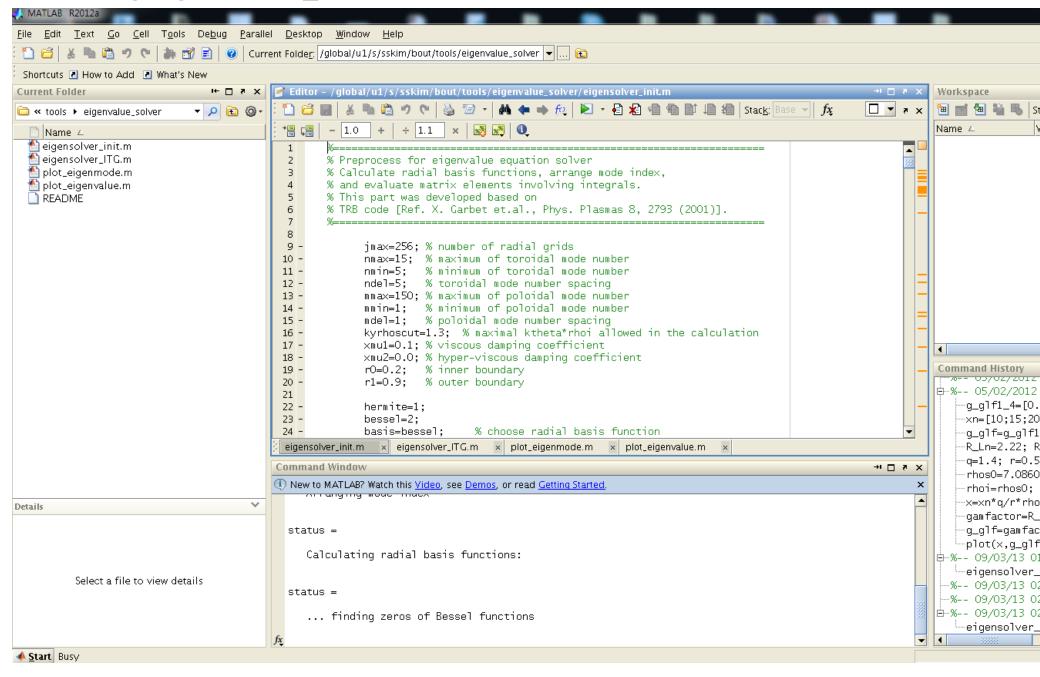


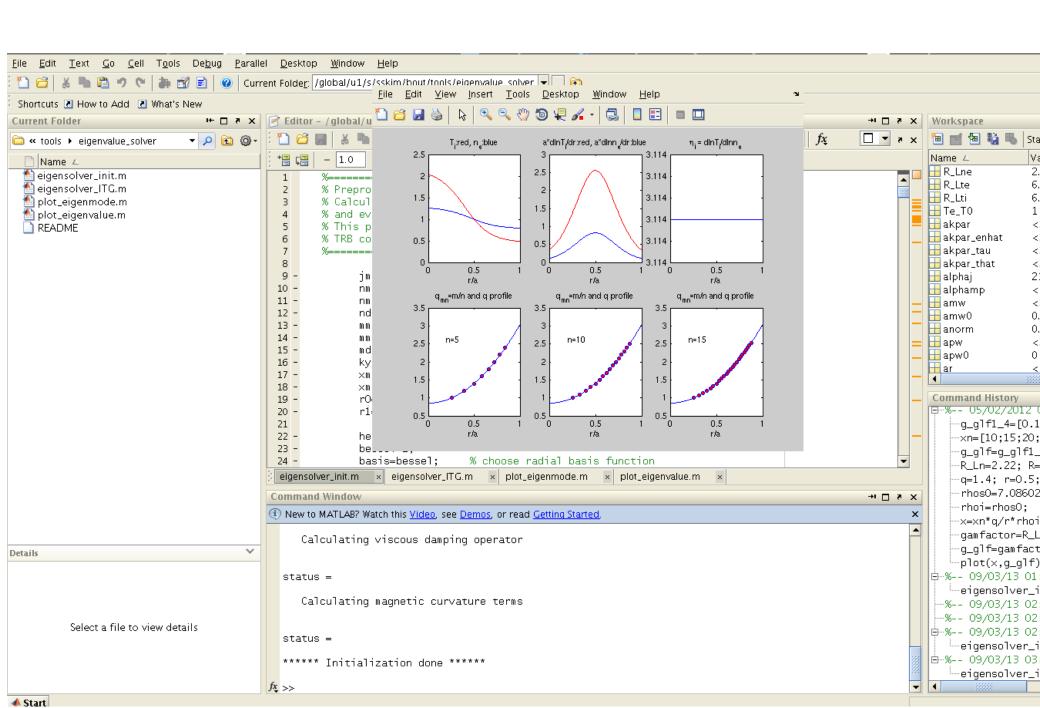
Run eigensolver_init.m;



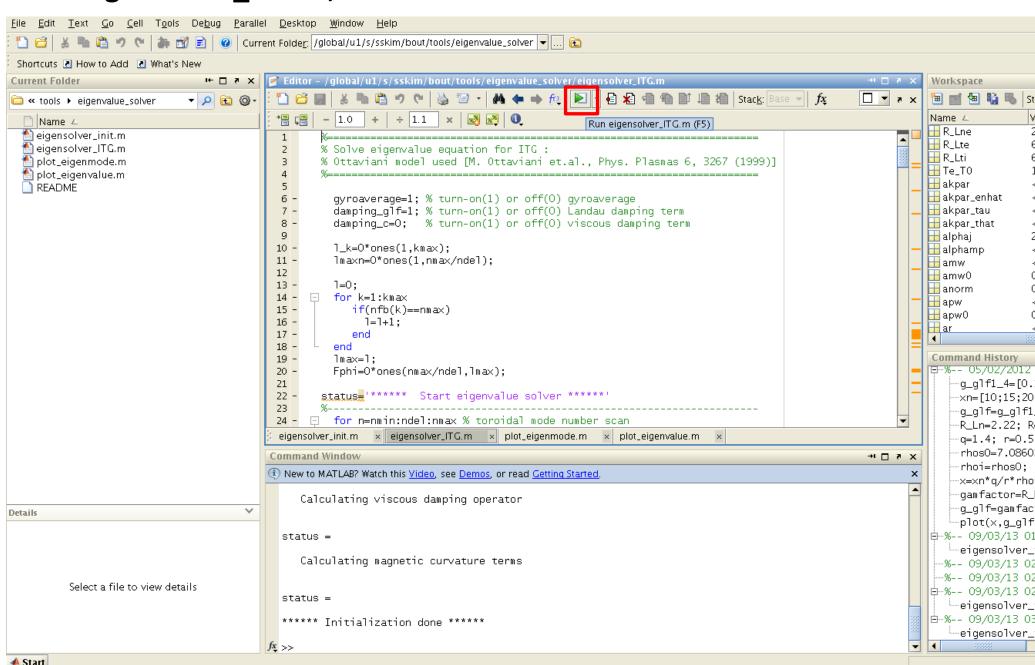


Running eigensolver_init.m ... : it takes about 5 min.

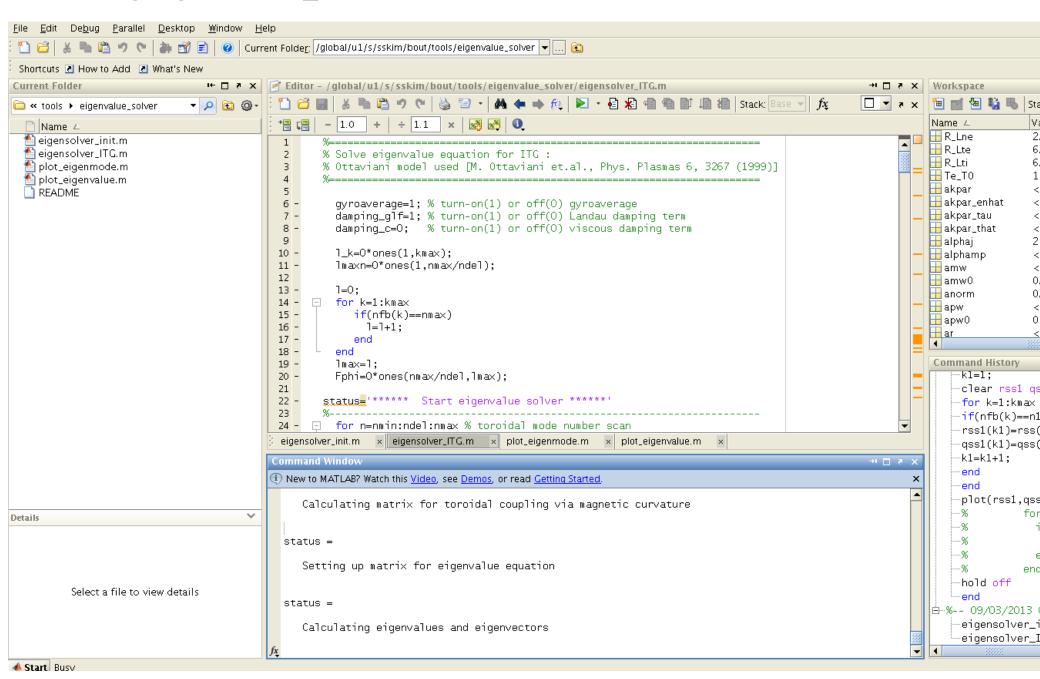




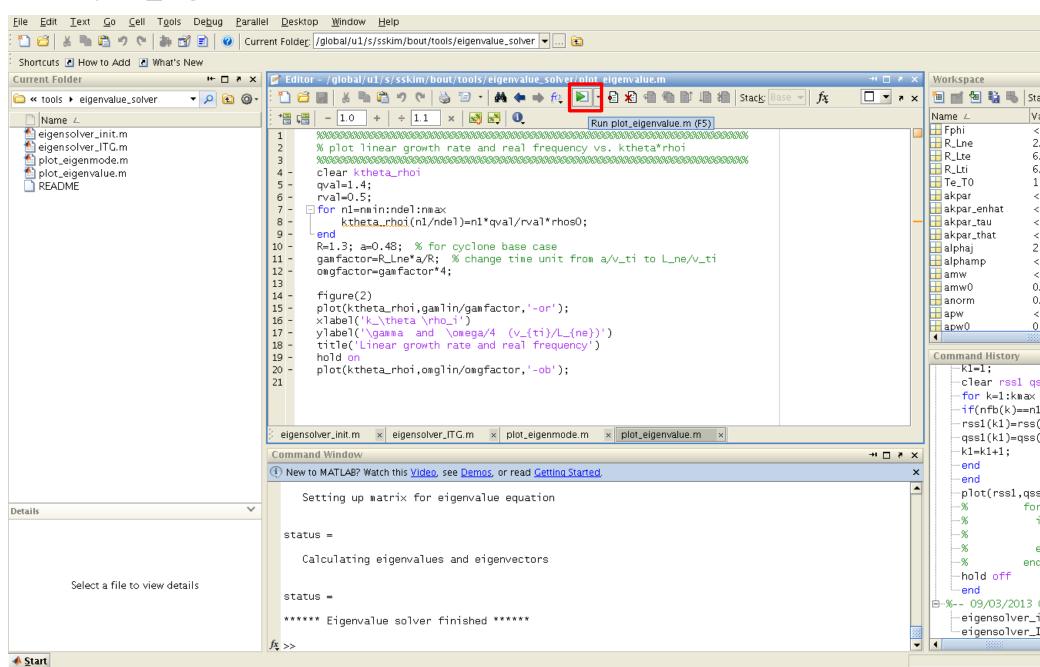
Run eigensolver_ITG.m;

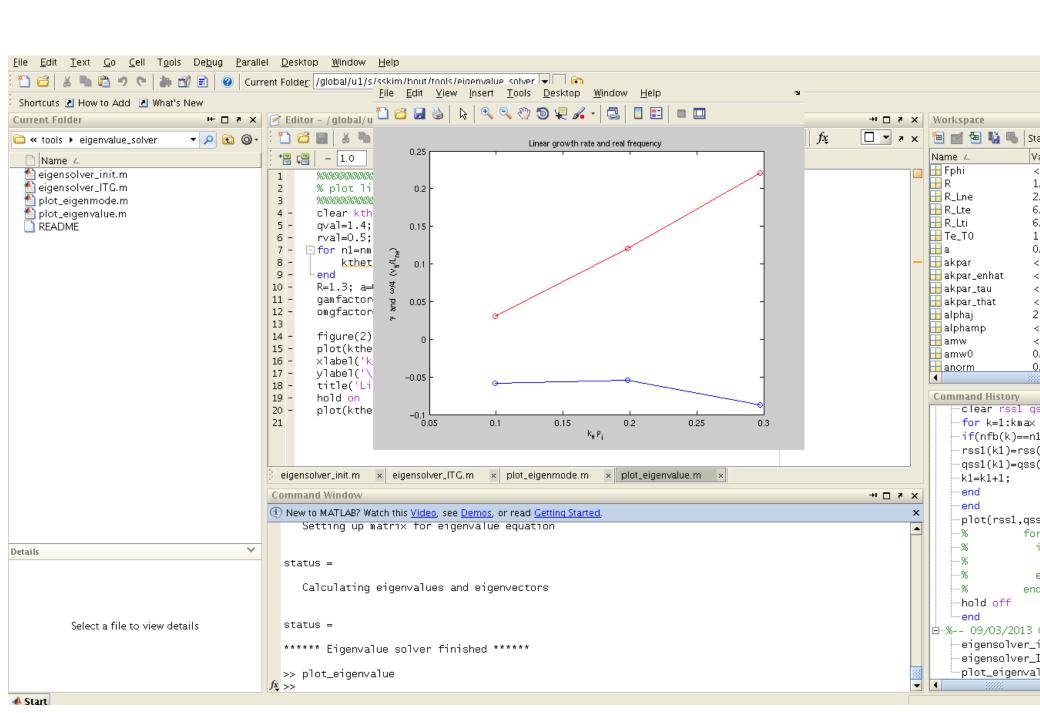


Running eigensolver_ITG.m ... : it takes about 10 min.

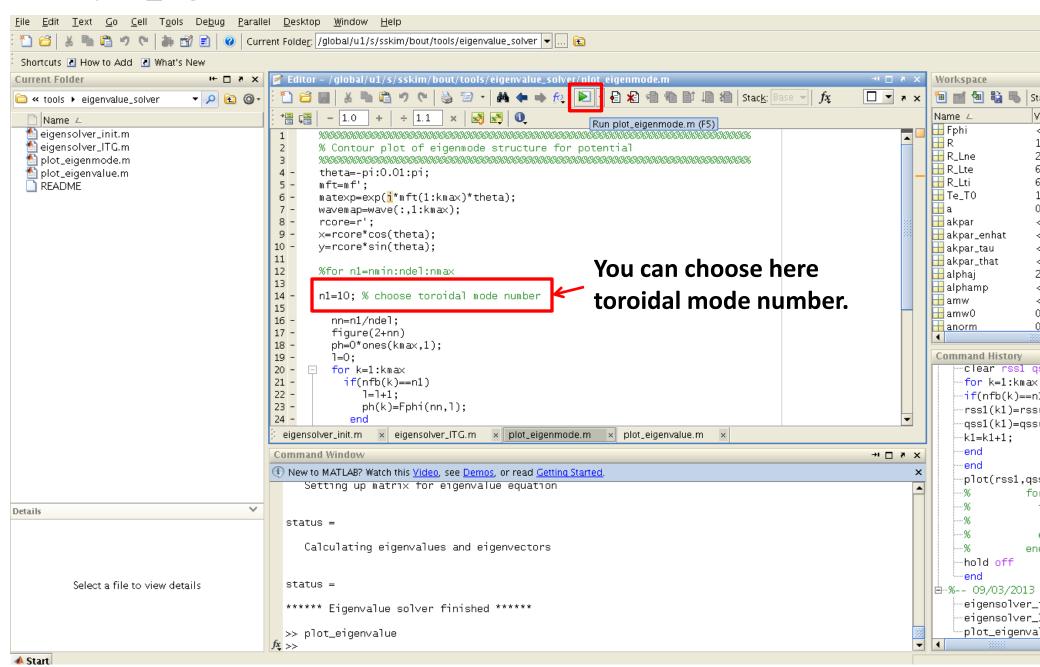


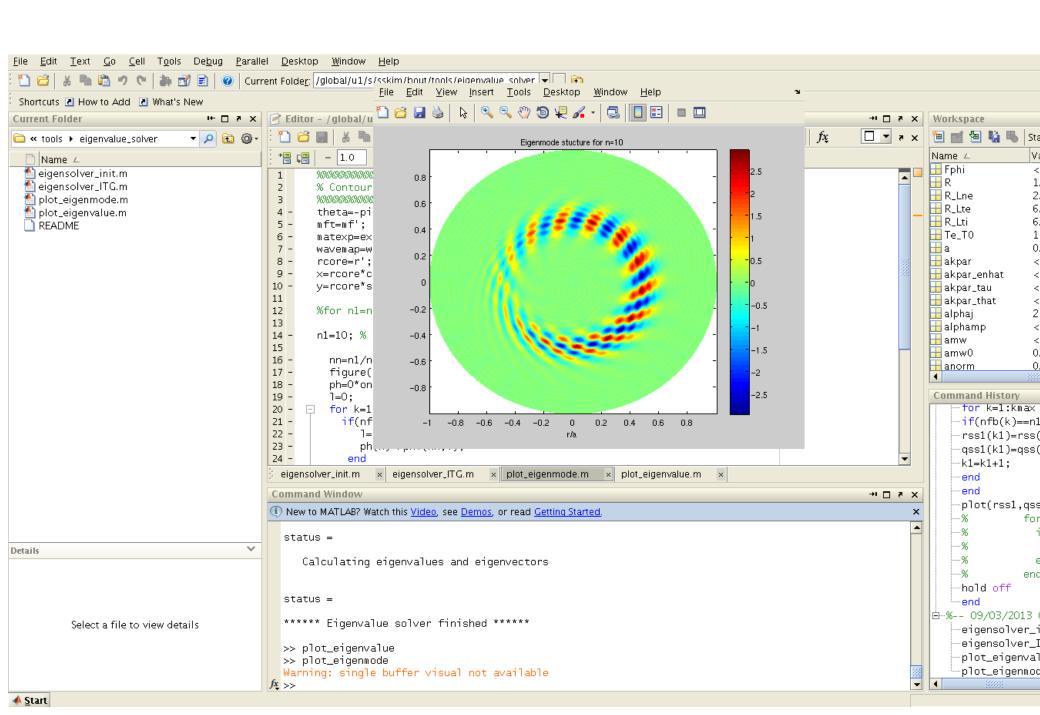
Run plot_eigenvalue.m;





Run plot_eigenmode.m:





Exit matlab:

