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to the pressure, p, of the confined gas

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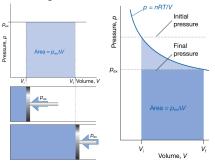
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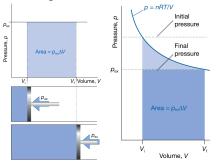
If, on the other hand, the external pressure differs measurably from the internal pressure, then changing $p_{\rm ex}$ infinitesimally will not decrease it below the pressure of the gas, so will not change the direction of the process. Such a system is **not** in mechanical equilibrium

Calculating work:



For a perfect gas, reversible work of expansion =

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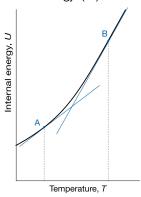
For a perfect gas, reversible work of expansion =

$$-\int_{V_{i}}^{V_{f}} p_{\text{ext}} dV = -\int_{V_{i}}^{V_{f}} p dV$$

$$= -nRT \ln \frac{V_{f}}{V_{i}}$$

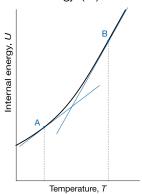
for spontaneous expansion : $work = -p_{ext} (V_f - V_i)$

Internal energy (U) as a function of temperature : U(V, T)



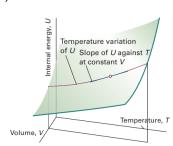
internal energy of a system increases as temperature is raised

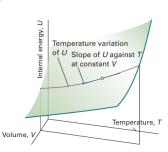
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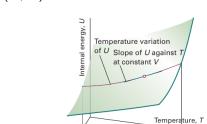
internal energy of a system increases as temperature is raised slope of tangent at any temperature =heat capacity, *C* in the Fig.,

 ${\it C}$ at A < ${\it C}$ at B



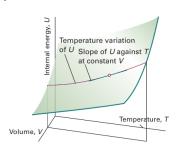


$$C_V = \left(\frac{dq}{dT}\right)_V = \left(\frac{\partial U}{\partial T}\right)_V$$



Volume, V

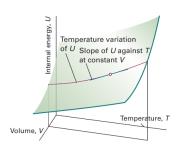
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If heat capacity is independent of temperature over range of temperatures of interest



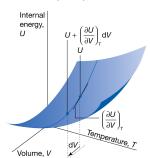
$$C_V = \left(\frac{dq}{dT}\right)_V = \left(\frac{\partial U}{\partial T}\right)_V$$

$$\implies dU = C_V dT$$
If heat capacity is independent

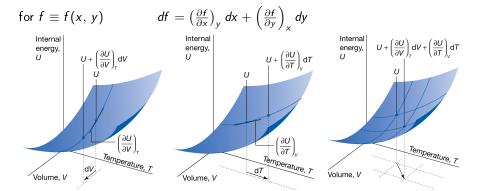
If heat capacity is independent of temperature over range of temperatures of interest then $\Delta U = C_V \Delta T$

for $f \equiv f(x, y)$ Internal

 $df = \left(\frac{\partial f}{\partial x}\right)_y dx + \left(\frac{\partial f}{\partial y}\right)_x dy$



for $f \equiv f(x, y)$ $df = \left(\frac{\partial f}{\partial x}\right)_{y} dx + \left(\frac{\partial f}{\partial y}\right)_{x} dy$ Internal energy, $U + \left(\frac{\partial U}{\partial V}\right)_{\tau} dV$ $U + \left(\frac{\partial U}{\partial V}\right)_{\tau} dV$ $U + \left(\frac{\partial U}{\partial T}\right)_{v} dT$ $U + \left(\frac{\partial U}{\partial T}\right)_{v} dT$ Volume, VVolume, V



•	large heat capacity ⇒ small increase in temper	y of energy transferi	red as heat, a

- ▶ large heat capacity ⇒ for given quantity of energy transferred as heat, a small increase in temperature
 - phase transition : heat capacity of a sample is infinite

large heat capacity \Rightarrow for given quantity of energy transferred as heat, a small increase in temperature

phase transition : heat capacity of a sample is infinite mean energy of a molecule due to its translational motion $=\frac{3}{2}k_BT$

 \therefore contribution to molar energy = $\frac{3}{2}RT$

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$$U_m(T) = U_m(0) + \frac{3}{2}RT$$

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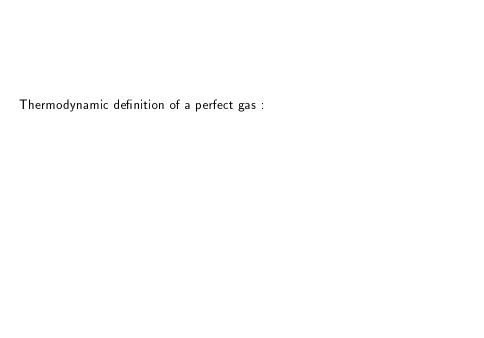
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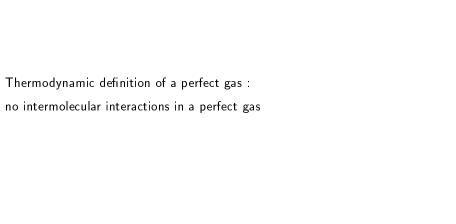
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e.g., for CO_2 , $U_m(T) = U_m(0) + \frac{3}{2}RT + RT + \cdots$





Thermodynamic definition of a perfect gas :
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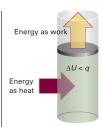
⇒ internal energy of a perfect gas is independent of the volume it occupies

Enthalpy

 $\Delta \text{energy} \neq \text{energy transferred as heat}$ when system is free to change its volume

Under these circumstances some of the

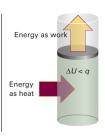
energy supplied as heat to system is returned to surroundings as expansion work, so dU < dq



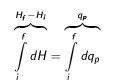
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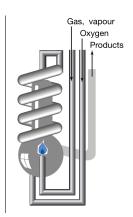


def.
$$H = U + pV$$
, a state function $dH = d(U + pV) = dU + pdV + Vdp = dq + Vdp \implies dq_p = dH$ def. $c_p = \left(\frac{dq}{dT}\right)_p = \left(\frac{\partial H}{\partial T}\right)_p \implies dH = C_p dT$



$$\int_{i}^{H_{f}-H_{i}} dH = \int_{i}^{q_{p}} dq_{p}$$

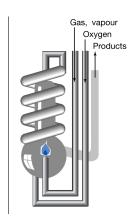
Part of a constant-pressure adiabatic flame calorimeter it is immersed in a stirred waterbath



$$\int_{i}^{H_{f}-H_{i}} dH = \int_{i}^{q_{p}} dq_{p}$$

Part of a constant-pressure adiabatic flame calorimeter

it is immersed in a stirred waterbath



Combustion occurs as a known amount of reactant is passed through to fuel the flame, and the rise of temperature is monitored.

when 1.0 mol CaCO₃ as calcite converts to aragonite, $\Delta U = 0.21$ kJ

Calculate $\Delta H - \Delta U$ when p=1 bar given densities of solids : 2.71 g cm⁻³ and

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► :
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usually justifiable to ignore $\Delta H - \Delta U$ for condensed phases, except at very high pressures, when ρV is no longer negligible

Ex. Calculate $\Delta H - \Delta U$ when 1.0 mol Sn(s, grey) of density 5.75 g cm⁻³ changes to Sn(s, white) of density 7.31 g cm⁻³ at 10.0 bar. At 298 K, $\Delta H = 2.1$ kJ

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 V_m for Sn (118.71 g)= ${20.65(g) \over 16.24(w)}$ cm³

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 \rightarrow $\Delta H - \Delta U = -0.44 \text{ J} \ (\approx 0.2\% \text{ of } \Delta U)$

 V_m for Sn (118.71 g)= ${20.65(g) \atop 16.24(w)}$ cm³

enthalpy of a perfect gas:

 $H = U + pV = U + nRT \implies \Delta H = \Delta U + \Delta nRT$

Water is heated to boiling under p = 1.0 atm. When an electric current of 0.50 A

from a 12 V supply is passed for 300 s through a resistance in thermal contact

with it,

0.798 g of water is vaporized. Calculate ΔU_m and ΔH_m at the boiling point

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$$\Delta H_m = q_p = rac{0.50 imes 12 imes 300}{rac{0.798}{100}} = 41 \; {
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$$\Delta H_m = q_p = rac{0.50 imes 12 imes 300}{rac{0.798}{18}} = 41 \ {
m kJ \ mol^{-1}}$$

$$p|^{-1} \left[\because p \left(V_{\sigma} - V_{I} \right) \approx p V_{\sigma} = RT \right]$$

$$\Delta U_m = \Delta H_m - RT = 38 \text{ kJ mol}^{-1} \left[:: p\left(V_g - V_l\right) \approx pV_g = RT \right]$$

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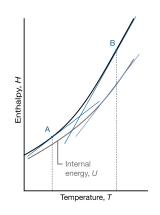
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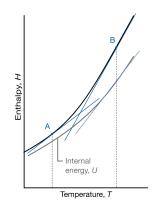
Notice $\Delta U_m < \Delta H_m$ because energy has been used to drive back surrounding atmosphere to make room for vapour

Enthalpy (H) as a function of temperature :



internal energy and enthalpy of system increase as temperature is raised

Enthalpy (H) as a function of temperature :



internal energy and enthalpy of system increase as temperature is raised at constant pressure, slope of tangent at any temperature =heat capacity, C_p

For gases, at a given temperature the slope is steeper than that of internal energy versus temperature

$$C_{p, m} > C_{V, m}$$

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of a reference material which does not undergo a physical or chemical change

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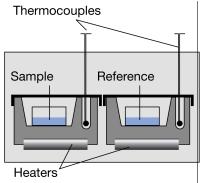
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'scanning': temperatures are increased, or scanned, during the analysis

during the analysis

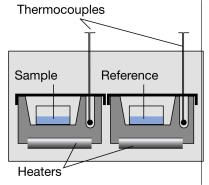
Differential scanning calorimeter (DSC) measures energy transferred as heat to or from a sample at constant pressure during a physical or chemical change 'differential' refers to the fact that the behaviour of the sample is compared to that of a reference material which does not undergo a physical or chemical change during the analysis

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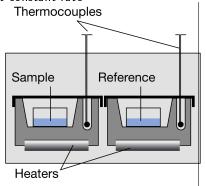
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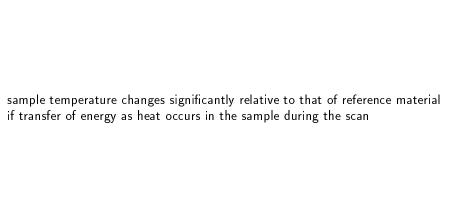
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'scanning': temperatures are increased, or scanned, during the analysis DSC consists of two small, separate but identical compartments heated electrically at constant rate



output : difference in power needed to maintain heat sinks at equal temperatures as temperature rises. $T(t) = T_0 + \alpha t$, where $T_0 =$ initial temperature $\alpha =$ temperature scan rate (in Ks^{-1})

A computer controls electrical power to maintain same temperature in sample and reference compartments



sample temperature changes significantly relative to that of reference material if transfer of energy as heat occurs in the sample during the scan
To maintain same temperature in both compartments, excess energy is transferred as heat to or from sample during the process

an endothermic process lowers sample temperature relative to that of reference and, as a result, the sample must be heated more strongly than reference in order to maintain equal temperatures if transfer of energy as heat occurs in the sample during the scan

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heat capacity $= \mathcal{C}_p + \mathcal{C}_{p,\, ext{ex}}$ and $q_p + q_{p,\, ext{ex}} = (\mathcal{C}_p + \mathcal{C}_{p,\, ext{ex}})\,\Delta T$

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heat capacity
$$=$$
 C_{p} $+$ $C_{p,\,ex}$ and q_{p} $+$ $q_{p,\,ex}$ $=$ $(C_{p}$ $+$ $C_{p,\,ex})\Delta T$

 $\therefore C_{p, ex} = \frac{q_{p, ex}}{\Delta T} = \frac{q_{p, ex}}{\alpha t} = \frac{P_{ex}}{\alpha}; \quad P_{ex} = \text{excess electrical power necessary to}$ equalize the temperature of sample and reference compartments

DSC trace - a thermogram : plot of P_{ex} or C against T

Broad peaks in thermogram indicate processes requiring transfer of energy as heat

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$$\Delta H = \int_{T_1}^{T_2} C_{p, ex} dT$$

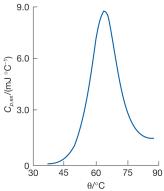
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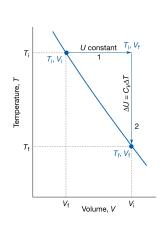
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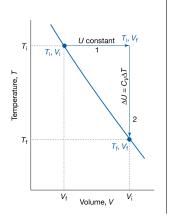
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thermogram for protein ubiquitin at pH = 2.45

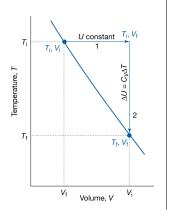
native structure up to $\approx 45^{\circ}\text{C}$ and then endothermic conformational change



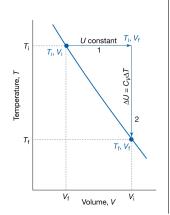




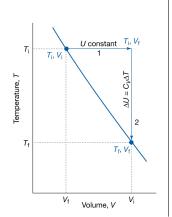
when a perfect gas expands adiabatically



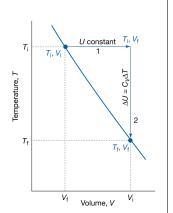
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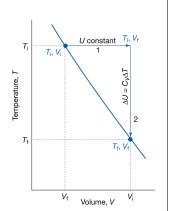
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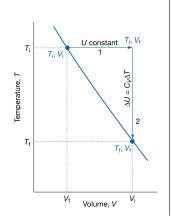


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 ΔU for temperature change from T_i to T_f and the volume change from V_i to V_f can be seen as the sum of two steps

1st step : only volume changes and temperature is constant

U for perfect gas is independent of volume of molecules, the overall ΔU arises solely from second step, the change in temperature at

$$\Delta U = C_v (T_f - T_i) = w_{ad}$$

constant volume

$$w_{ad} = C_v \Delta T$$

 $w_{ad} = C_v \Delta T$ reversible adiabatic expansion : $C_v dT = -pdV$

$$w_{ad}=C_v\Delta T$$

reversible adiabatic expansion : $C_vdT=-pdV$
For perfect gas, $C_v\frac{dT}{T}=-\frac{nRdV}{V}$

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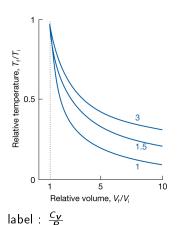
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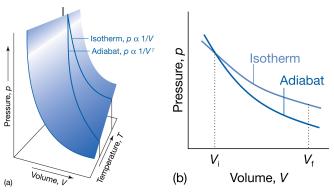
or, $\ln\left(\frac{T_2}{T_1}\right)^{\frac{C_V}{nR}} = \ln\left(\frac{V_1}{V_2}\right)$ or, $\left(\frac{T_2}{T_1}\right)^{\frac{C_V}{nR}} = \frac{V_1}{V_2}$

show: $pV^{\gamma} = \text{const}$

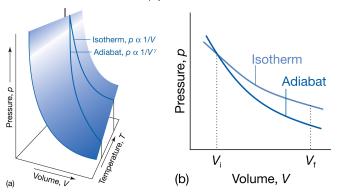
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p-V-T plots: isotherms and adiabats pressure declines more steeply for an adiabat than for an isotherm



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in isothermal expansion, energy flows into the system as heat and maintains the temperature; as a result, pressure does not fall as much as in adiabatic expansion

adiabatic, reversible expansion of 0.020 mol Ar

ex :

initially at 25°C, from 0.50 dm³ to 1.00 dm³

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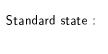
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exothermic process at constant pressure : $\Delta H < 0$ endothermic process at constant pressure : $\Delta H > 0$



Standard state : standard state of a substance at a specified temperature is its pure form at $1\ \mathrm{base}$	ar

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enthalpy of products $_{\rm standard\,states}$ — enthalpy of reactants $_{\rm standard\,states}$ standard enthalpy of vaporization : $\Delta_{\rm vap} H^{\odot} = {\rm enthalpy}$ change per mole when a pure liquid at 1 bar vaporizes to a

 $\Delta_{\text{vap}}H^{\ominus}=$ enthalpy change per mole when a pure liquid at 1 bar vaporizes to a gas at 1 bar $H_2O(I)\longrightarrow H_2O(g)$ $\Delta_{\text{vap}}H^{\ominus}(373\text{K})=40.66 \text{ kJmol}^{-1}$

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 $\Delta_{\sf vap} H^{\odot} = {\sf enthalpy}$ change per mole when a pure liquid at 1 bar vaporizes to a gas at 1 bar

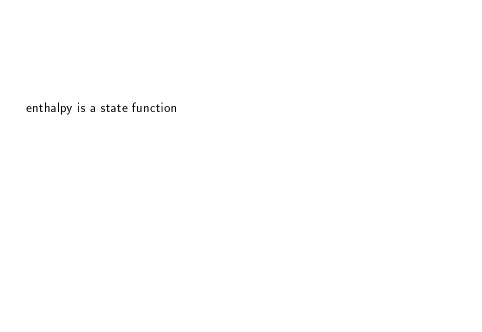
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standard enthalpies may be reported for any temperature



enthalpy is a state function

states

 ΔH is independent of path between the two

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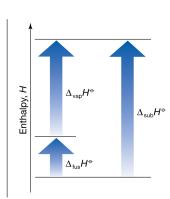
 ΔH is independent of path between the two states

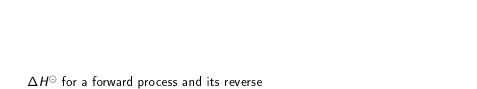
same ΔH^{\odot} for change brought about between same initial and final states

e.g., conversion of a solid to a vapour either as

occurring by sublimation

$$\begin{array}{ll} \mathsf{H}_2\mathsf{O}(\mathsf{s}) &\longrightarrow \mathsf{H}_2\mathsf{O}(\mathsf{I}) & \Delta_\mathsf{fus} H^\circleddash \\ \mathsf{H}_2\mathsf{O}(\mathsf{I}) &\longrightarrow \mathsf{H}_2\mathsf{O}(\mathsf{g}) & \Delta_\mathsf{vap} H^\circleddash \\ \mathsf{H}_2\mathsf{O}(\mathsf{s}) &\longrightarrow \mathsf{H}_2\mathsf{O}(\mathsf{g}) & \Delta_\mathsf{sub} H^\circleddash = (\Delta_\mathsf{fus} H^\circleddash + \Delta_\mathsf{vap} H^\circleddash) > \Delta_\mathsf{vap} H^\circleddash \end{array}$$



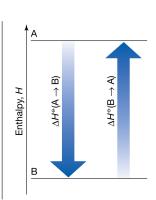


differ in sign

 ΔH^{\odot} for a forward process and its reverse differ in sign

e.g., conversion of a solid to a vapour either as occurring by sublimation

$$\Delta H^{\odot}(A \rightarrow B) = -\Delta H^{\odot}(B \rightarrow A)$$



Transition	Process	Symbol
Transition	Phase $\alpha \rightarrow$ phase β	$\Delta_{ m trs} H$
Fusion	$s \rightarrow 1$	$\Delta_{ m fus} H$
Vaporization	$l \rightarrow g$	$\Delta_{ m vap} H$
Sublimation	$s \rightarrow g$	$\Delta_{ m sub} H$
Mixing	Pure → mixture	$\Delta_{ ext{mix}}H$
Solution	$Solute \rightarrow solution$	$\Delta_{ m sol} H$
Hydration	$X^{\pm}(g) \to X^{\pm}(aq)$	$\Delta_{ m hyd} H$
Atomization	$Species(s, l, g) \rightarrow atoms(g)$	$\Delta_{\mathrm{at}}H$
Ionization	$X(g) \rightarrow X^+(g) + e^-(g)$	$\Delta_{\mathrm{ion}}H$
Electron gain	$X(g) + e^{-}(g) \longrightarrow X^{-}(g)$	$\Delta_{ m eg} H$
Reaction	Reactants \rightarrow products	$\Delta_{\mathbf{r}}H$
Combustion	$Compounds(s, l, g) + O_2(g) \rightarrow CO_2(g), H_2O(l, g)$	$\Delta_c H$
Formation	Elements \rightarrow compound	$\Delta_{ m f} H$
Activation	Reactants \rightarrow activated complex	$\Delta^{\ddagger}H$

Standard reaction	enthalov :	

change in enthalpy for

Pure, separate reactants in standard states \longrightarrow

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Ex. $CH_4(g)+2O_2(g) \longrightarrow CO_2(g)+2H_2O(I)$ $\Delta_r H^{\odot}=-890 \text{ kJ}$

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$$\Delta_r H^{\odot} = \left\{ 3H_m^{\odot} \left(\mathsf{C} \right) + H_m^{\odot} \left(\mathsf{D} \right) \right\} - \left\{ 2H_m^{\odot} \left(\mathsf{A} \right) + H_m^{\odot} \left(\mathsf{B} \right) \right\}$$

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$$\Delta_r H^{\odot} = \{3H_m^{\odot}(C) + H_m^{\odot}(D)\} - \{2H_m^{\odot}(A) + H_m^{\odot}(B)\}$$

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generally,
$$\Delta_r H^{\odot} = \sum_{\text{products}} \nu H_m^{\odot} - \sum_{\text{reactants}} \nu H_m^{\odot}$$

Standard reaction enthalpy : change in enthalpy for

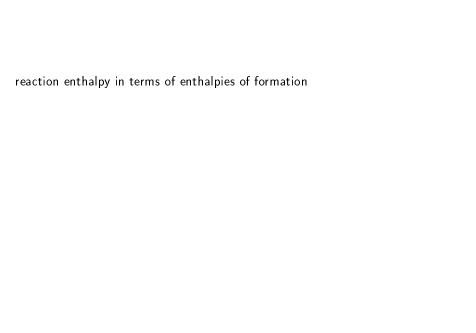
Pure, separate reactants in standard states \longrightarrow

pure, separate products in standard states Ex. $CH_4(g)+2O_2(g) \longrightarrow CO_2(g)+2H_2O(I)$ $\Delta_r H^{\odot} = -890 \text{ kJ}$

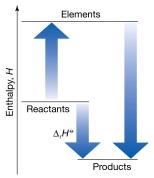
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standard enthalpy of combustion, $\Delta_C H^{\odot} = \text{standard reaction enthalpy for the complete oxidation of an organic compound to CO₂ gas and liquid water if the compound contains C, H, and O, and to N₂ gas if N is also present$



reaction enthalpy in terms of enthalpies of formation



$$\Delta_r H^{\odot} = \sum_{\rm products} \nu \Delta_f H^{\odot} - \sum_{\rm reactants} \nu \Delta_f H^{\odot}$$



Can we construct $\Delta_f H^{\odot}$ from a knowledge of chemical constitution of the species?

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 $dH = C_p dT \implies H(T_2) = H(T_1) + \int_{T_1}^{T_2} C_p dT$ assuming no phase transition in

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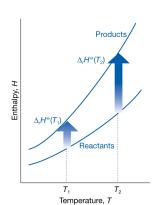
temperature range of interest

 $dH = C_p dT \implies H(T_2) = H(T_1) + \int_{T_2}^{T_2} C_p dT$ assuming no phase transition in

temperature range of interest

Kirchhoff's law:

where
$$\Delta_r C_p^{\odot} = \sum_{\text{products}} \nu C_{\text{p, m}}^{\odot} - \sum_{\text{reactants}} \nu C_{\text{p, m}}^{\odot}$$



Estimate $\Delta_f H^{\odot}$ at 100° C

 $C_{p,m}(O_2,g) = 29.37 \text{J K}^{-1} \text{ mol}^{-1}$

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given $C_{p, m}(H_2O,g) = 33.58 \text{ J K}^{-1}\text{mol}^{-1}$; $C_{p, m}(H_2,g) = 28.84 \text{ JK}^{-1} \text{ mol}^{-1}$;

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 $\Delta_r C_p^{\odot} = C_{p, m} (H_2O, g) - \{C_{p, m}(H_2, g) + \frac{1}{2}C_{p, m}(O_2, g) = -9.94 \text{J K}^{-1} \text{ mol}^{-1}\}$

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 $= -241.82 - 75 \times 9.94 \times 10^{-3} = -242.6 \text{k l mol}^{-1}$

 $\therefore \Delta_r H^{\odot}(T_2) = \Delta_r H^{\odot}(T_1) + (T_2 - T_1) \Delta_r C_n^{\odot}$

 $C_{\rm p,m}(O_2,g) = 29.37 \,\rm J \, K^{-1} \, mol^{-1}$