

# The Crisis of Quantum Foundations: Why Physics Needs the Yoneda Constraint

Matthew Long

*The YonedaAI Collaboration*

*YonedaAI Research Collective*

Chicago, IL

[matthew@yonedaai.com](mailto:matthew@yonedaai.com) · <https://yonedaai.com>

24 February 2026

## Abstract

Quantum mechanics is the most empirically successful theory in the history of physics, yet its foundations remain uniquely opaque among the fundamental theories of nature. Unlike general relativity, which flows from the equivalence principle and the geometry of spacetime, quantum mechanics rests on a collection of mathematical axioms—Hilbert spaces, self-adjoint operators, the Born rule, the projection postulate—whose physical motivation remains obscure nearly a century after their formulation. This foundational deficit has spawned an unresolved crisis: a proliferation of mutually incompatible interpretations (Copenhagen, Many-Worlds, Bohmian mechanics, QBism, relational QM), an intractable measurement problem, and a persistent inability to unify quantum theory with gravity. We argue that this crisis is not merely philosophical but structural, arising from the absence of a *principled foundation* for the quantum formalism. We propose that the missing principle is the **Yoneda Constraint**—the requirement, derived from the Yoneda Lemma of category theory, that physical objects are completely and faithfully characterized by their relational profiles. We survey the historical development of the foundations crisis, analyze the asymmetry between the principled structure of general relativity and the axiomatic opacity of quantum mechanics, and show how the Yoneda Constraint resolves the crisis by deriving the quantum formalism from a single structural principle. We compare this approach in detail with other reconstruction programs (Hardy, Chiribella–D’Ariano–Perinotti, Masanes–Müller, Dakic–Brukner) and argue that the Yoneda Constraint provides a uniquely powerful and mathematically natural foundation. This paper serves as a comprehensive introduction to the **Quantum Perspectivism** program, which develops the full consequences of the Yoneda Constraint for the foundations of physics.

**Keywords:** quantum foundations, measurement problem, interpretations of quantum mechanics, category theory, Yoneda lemma, presheaf, topos, Born rule, reconstruction of quantum theory, structural realism

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# 1 Introduction

“Nobody understands quantum mechanics.”

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—Richard P. Feynman, *The Character of Physical Law* (1965)

## 1.1 The Unique Opacity of Quantum Foundations

Among the fundamental theories of modern physics, quantum mechanics occupies a unique and unsettling position. It is, by any empirical measure, the most successful physical theory ever devised. Its predictions have been confirmed to extraordinary precision—the anomalous magnetic moment of the electron, for instance, agrees with experiment to better than one part in  $10^{12}$  [1]. Quantum mechanics underlies all of modern chemistry, materials science, semiconductor technology, and nuclear physics. No experiment has ever produced a result inconsistent with its predictions.

And yet, despite nearly a century of sustained effort by the most brilliant physicists and philosophers of science, no consensus exists on what the theory *means*. The axioms of quantum mechanics—the Hilbert space framework, the Born rule, the projection postulate, the Schrödinger equation—are presented in every textbook as foundational starting points, but these axioms themselves receive no deeper justification. They are, as we shall argue, *brute stipulations*: mathematical structures chosen because they work, not because they follow from any deeper physical principle.

This situation has no parallel in the rest of fundamental physics. Classical mechanics follows from Newton’s laws, which themselves can be derived from variational principles (the principle of least action) that have clear physical content. Electrodynamics flows from Maxwell’s equations, which are unified and explained by the gauge principle and the structure of the Lorentz group. Thermodynamics rests on a small number of physically transparent laws about heat, work, and entropy. And most impressively, general relativity is derived from a single physical insight—the equivalence principle—combined with the requirement of general covariance and the geometry of spacetime.

Quantum mechanics alone among the foundational theories offers no such principled derivation. The axioms are there; they work magnificently; but *why* they are there, why nature should obey these particular mathematical rules rather than any others, remains a mystery.

## 1.2 The Scope and Structure of This Paper

This paper provides a comprehensive analysis of the quantum foundations crisis and argues that the resolution lies in a structural principle drawn from category theory: the **Yoneda Constraint**. The paper is organized as follows.

Section 2 surveys the historical development of the foundations crisis, from the original Copenhagen interpretation through the measurement problem and the proliferation of interpretations.

Section 3 analyzes the striking asymmetry between the principled foundations of general relativity and the axiomatic opacity of quantum mechanics, arguing that

this asymmetry is the root cause of the foundations crisis.

Section 4 introduces category theory as the natural mathematical language for analyzing the structure of physical theories, with particular attention to the concepts of presheaf, natural transformation, and topos.

Section 5 presents the Yoneda Lemma and develops it as a foundational constraint on physical theories—the Yoneda Constraint—showing that it is a mathematical theorem with profound physical consequences, not merely a philosophical preference.

Section 6 previews how the Yoneda Constraint resolves each of the major foundational puzzles: the measurement problem, the origin of the Born rule, the nature of superposition, the basis of complementarity, and the structure of entanglement.

Section 7 provides a detailed comparison with other programs that attempt to derive quantum mechanics from first principles, including the work of Hardy, Chiribella–D’Ariano–Perinotti, Masanes–Müller, Dakic–Brukner, and the topos approaches of Isham, Butterfield, and Döring.

Section 8 addresses the philosophical implications of the Yoneda Constraint for debates about realism, the nature of physical objects, and the role of the observer.

Section 10 discusses open problems and future directions, and Section 11 concludes.

## 2 Historical Survey: The Quantum Foundations Crisis

### 2.1 The Birth of Quantum Theory and the Copenhagen Settlement

The quantum revolution began not with a single dramatic insight but with a series of ad hoc moves forced by experimental anomalies. Planck’s quantization of blackbody radiation (1900), Einstein’s explanation of the photoelectric effect (1905), Bohr’s model of the hydrogen atom (1913), and the development of matrix mechanics by Heisenberg (1925) and wave mechanics by Schrödinger (1926) each introduced quantum ideas piecemeal, without a unifying physical principle.

The first attempt at a systematic interpretation came from the “Copenhagen school”—primarily Bohr, Heisenberg, and Born—in the late 1920s [2]. The Copenhagen interpretation, insofar as it can be pinned down as a single doctrine, holds that:

- (a) The quantum state  $|\psi\rangle$  represents the complete physical description of a system.
- (b) Physical quantities do not have definite values until measured; the act of measurement “collapses” the state to an eigenstate of the measured observable.
- (c) The Born rule  $p(\lambda) = |\langle e_\lambda | \psi \rangle|^2$  gives the probability of obtaining outcome  $\lambda$ .
- (d) Complementary observables (position and momentum, spin along different axes) cannot simultaneously have sharp values, as expressed by the Heisenberg uncertainty relations.
- (e) There is a fundamental divide between the quantum system and the classical measuring apparatus; this “Heisenberg cut” is necessary for the theory to make contact with experiment.

The Copenhagen interpretation was enormously influential and provided a workable framework for the practical use of quantum mechanics. But it was recognized from the very beginning to be deeply problematic. Einstein’s famous objections—culminating in the Einstein–Podolsky–Rosen (EPR) argument of 1935 [3]—challenged the completeness of the quantum description. Schrödinger’s cat thought experiment (1935) dramatized the absurdity of taking the collapse postulate literally for macroscopic systems [4]. And Bohr’s insistence on the indispensability of the classical apparatus introduced a fundamental conceptual circularity: the theory that was supposed to be more fundamental than classical mechanics appeared to *depend* on classical mechanics for its interpretation.

## 2.2 The Measurement Problem

The measurement problem is the central conceptual puzzle of quantum foundations. It can be stated with deceptive simplicity:

If the Schrödinger equation governs the evolution of all physical systems, and if measurement apparatuses are themselves physical systems, then a measurement should produce a superposition of the apparatus in different outcome states—not a definite result. But measurements *do* produce definite results. How?

More formally, let  $|\psi\rangle = \sum_i c_i |a_i\rangle$  be the initial state of a quantum system in a superposition of eigenstates of the observable  $A$ , and let  $|R_0\rangle$  be the “ready” state of the measuring apparatus. If the measurement interaction is governed by the Schrödinger equation, the composite system evolves as:

$$\left( \sum_i c_i |a_i\rangle \right) \otimes |R_0\rangle \longrightarrow \sum_i c_i |a_i\rangle \otimes |R_i\rangle \quad (1)$$

where  $|R_i\rangle$  is the apparatus state corresponding to the outcome  $a_i$ . The final state is an entangled superposition of system-apparatus states. No definite outcome has occurred. Yet experiment unambiguously shows that measurements *do* have definite outcomes.

The measurement problem is not a technical difficulty that might be resolved by a more careful analysis of the measurement interaction. It is a fundamental inconsistency in the formalism: the linear, unitary Schrödinger evolution and the nonlinear, non-unitary projection postulate are logically incompatible rules that are both asserted to govern physical processes [5].

**Theorem 2.1** (Maudlin’s Trilemma). *The following three propositions are mutually inconsistent:*

- (i) *The quantum state is a complete description of a physical system.*
- (ii) *The quantum state always evolves according to a linear dynamical equation (the Schrödinger equation).*
- (iii) *Measurements always have definite outcomes.*

*Every interpretation of quantum mechanics must deny at least one of these three claims.*

*Proof.* If the state is complete and evolves linearly, then by (1), after a measurement the composite system is in a superposition with no definite outcome. Hence (iii) is violated. Conversely, asserting (iii) requires either supplementing the state description (denying (i), as in hidden-variable theories) or modifying the dynamics (denying (ii), as in collapse theories). No logically consistent position can maintain all three.  $\square$

## 2.3 The Proliferation of Interpretations

The measurement problem has spawned a vast literature of proposed resolutions, none of which commands consensus. The major interpretive families include:

(1) **Copenhagen and neo-Copenhagen approaches.** The original Copenhagen view, together with modern variants such as the “consistent histories” approach of Griffiths, Omnès, and Gell-Mann–Hartle [6], which attempt to formulate quantum mechanics without a collapse postulate by restricting attention to consistent sets of histories.

(2) **Many-Worlds Interpretation.** Everett’s “relative state” formulation (1957) [7] denies proposition (iii) of Maudlin’s trilemma, maintaining that all branches of the superposition are equally real. The universe continually “splits” into non-interacting branches, each containing a definite outcome. The interpretation faces the “preferred basis problem” (why does branching occur in the position basis rather than some other?) and the “probability problem” (how does the Born rule emerge from a deterministic, branching dynamics?).

(3) **Bohmian Mechanics.** De Broglie–Bohm pilot-wave theory [8] denies proposition (i), supplementing the wave function with actual particle positions guided by a “pilot wave.” It is empirically equivalent to standard quantum mechanics for non-relativistic systems but introduces a preferred reference frame and has proven difficult to extend to relativistic quantum field theory.

(4) **Dynamical Collapse Theories.** The GRW (Ghirardi–Rimini–Weber) model [9] and its relatives deny proposition (ii), modifying the Schrödinger equation with stochastic, nonlinear collapse terms. These models make predictions that differ from standard quantum mechanics in principle, though the differences are extraordinarily small for currently accessible systems.

(5) **QBism (Quantum Bayesianism).** Fuchs, Mermin, and Schack [10] deny that the quantum state describes objective physical reality, instead treating it as a representation of an agent’s beliefs. This dissolves the measurement problem by denying that there is any objective physical process corresponding to collapse, but it raises questions about the physical content of the theory.

(6) **Relational Quantum Mechanics.** Rovelli [11] holds that quantum states are always relative to an observer system, with no absolute state of a system “in itself.” This is philosophically attractive but has been criticized for lacking a precise mathematical framework.

## 2.4 A Century of Disagreement

The persistence of this interpretive landscape, essentially unchanged in its broad outlines since the 1950s, constitutes a genuine crisis. It is not merely that physicists

have different philosophical tastes; it is that the formalism itself *underdetermines* its interpretation to a degree unparalleled in any other fundamental theory.

A 2013 survey of physicists attending a quantum foundations conference found no majority interpretation: approximately 42% favored Copenhagen, 18% favored Many-Worlds, and the remainder was scattered among various alternatives, with 12% selecting “I have no preferred interpretation” [12]. More recent surveys show similar fragmentation.

This situation would be inconceivable in any other branch of physics. No one debates the “interpretation” of general relativity or electrodynamics in remotely comparable terms. The reason, as we shall argue in Section 3, is that those theories have *principled foundations* that constrain their interpretation, while quantum mechanics does not.

### 3 The Asymmetry: Principled Foundations vs. Axiomatic Opacity

#### 3.1 General Relativity: A Theory from a Principle

General relativity is the paradigm case of a physical theory derived from a single, physically transparent principle. The Einstein equivalence principle states:

**Axiom 1** (Equivalence Principle). In a sufficiently small region of spacetime, the effects of gravity are indistinguishable from the effects of acceleration in the absence of gravity.

From this principle, combined with the mathematical requirement of general covariance (the laws of physics should take the same form in all coordinate systems), Einstein derived the entire structure of general relativity:

- (i) Spacetime is a 4-dimensional pseudo-Riemannian manifold  $(M, g_{\mu\nu})$ .
- (ii) Freely falling particles follow geodesics of  $g_{\mu\nu}$ .
- (iii) The metric  $g_{\mu\nu}$  is a dynamical field governed by the Einstein field equations:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu} \quad (2)$$

The derivation is not merely heuristic; it is rigorous. Lovelock’s theorem [13] shows that in four dimensions, the Einstein field equations (with cosmological constant) are the *unique* second-order field equations for a metric tensor derivable from an action principle—a mathematical theorem that underwrites the physical uniqueness of general relativity given its foundational principle.

#### 3.2 Quantum Mechanics: Axioms Without a Principle

Now compare the standard axiomatization of quantum mechanics, as found in any textbook (e.g., Dirac [14], von Neumann [15], Sakurai [16]):

**Axiom 2** (Hilbert Space). The state of a physical system is represented by a unit ray in a complex Hilbert space  $\mathcal{H}$ .

**Axiom 3** (Observables). Physical observables are represented by self-adjoint operators on  $\mathcal{H}$ .

**Axiom 4** (Born Rule). If a system is in state  $|\psi\rangle$  and observable  $A$  is measured, the probability of obtaining eigenvalue  $a_n$  is  $p(a_n) = |\langle a_n|\psi\rangle|^2$ .

**Axiom 5** (Projection Postulate). Upon measurement of  $A$  yielding outcome  $a_n$ , the state collapses to the eigenstate  $|a_n\rangle$ .

**Axiom 6** (Schrödinger Evolution). Between measurements, the state evolves unitarily:  $|\psi(t)\rangle = e^{-iHt/\hbar}|\psi(0)\rangle$ .

**Axiom 7** (Composition). The Hilbert space of a composite system is the tensor product of the component Hilbert spaces:  $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$ .

These axioms are mathematically precise and empirically adequate. But they are *unmotivated*. Why Hilbert spaces rather than some other mathematical structure? Why complex numbers rather than real numbers or quaternions? Why the Born rule with its squared modulus rather than some other function of the state? Why the tensor product for composite systems? Why self-adjoint operators for observables?

To each of these questions, the standard formalism offers no answer beyond “because it works.” The axioms are not derived from a physical principle analogous to the equivalence principle; they are posited as starting points, justified *a posteriori* by empirical success.

### 3.3 The Root of the Crisis

We contend that the foundations crisis described in Section 2 is a direct consequence of this axiomatic opacity. The measurement problem, the proliferation of interpretations, the persistent disagreements about the ontology of quantum mechanics—all of these stem from the fact that the formalism offers no guidance about what it *means*.

A theory derived from a physical principle carries its interpretation with it. General relativity *is* the geometry of spacetime; this is not an add-on interpretation but part of the derivation. The equivalence principle tells us what the mathematical structures *represent*—the metric tensor is the gravitational field, geodesics are the trajectories of freely falling bodies, curvature is tidal gravity.

Quantum mechanics, by contrast, presents us with mathematical structures (Hilbert spaces, operators, complex amplitudes) and then *challenges us to figure out what they mean*. The interpretive crisis is the predictable consequence of this challenge going unmet for a century.

**Proposition 3.1** (The Diagnostic). *The quantum foundations crisis arises from the absence of a foundational principle for quantum mechanics analogous to the equivalence principle for general relativity. Resolving the crisis requires identifying such a principle—a single, physically transparent constraint from which the quantum formalism can be derived.*

This is the task we undertake. We claim that the missing principle is the Yoneda Constraint.

## 4 Category Theory as the Language of Structural Physics

### 4.1 Why Category Theory?

If we seek a principle that constrains the *structure* of physical theories, we need a mathematical language that can talk about structure in a theory-independent way. Category theory is precisely such a language.

Category theory, founded by Eilenberg and Mac Lane in 1945 [17], was initially conceived as a tool for unifying constructions across different branches of mathematics. Its key insight is that mathematical objects are best understood not through their internal constitution but through their relationships—their morphisms—to other objects. Over the subsequent eight decades, this insight has proven extraordinarily fruitful, providing a unifying language for algebra, topology, geometry, logic, and computer science [18].

For our purposes, category theory is essential for three reasons:

- (i) **Structural universality.** Category theory provides a framework for comparing the structures of different mathematical theories without committing to the internal details of any one theory. This is precisely what is needed to analyze what *any* physical theory must look like.
- (ii) **The Yoneda Lemma.** Category theory contains a theorem—the Yoneda Lemma—that establishes a precise sense in which objects are completely determined by their relational profiles. This theorem will serve as our foundational constraint.
- (iii) **Topos theory.** The categorical theory of topoi provides a general framework for “universes of discourse” with their own internal logic, which turns out to be precisely the non-Boolean logic of quantum mechanics.

### 4.2 Basic Categorical Concepts

We recall the essential definitions. A reader fluent in category theory may skip this subsection.

**Definition 4.1** (Category). A **category**  $\mathcal{C}$  consists of:

- (a) A collection  $\text{Ob}(\mathcal{C})$  of **objects**.
- (b) For each pair of objects  $A, B$ , a collection  $\text{Hom}_{\mathcal{C}}(A, B)$  of **morphisms** from  $A$  to  $B$ .
- (c) For each object  $A$ , an **identity morphism**  $\text{id}_A \in \text{Hom}_{\mathcal{C}}(A, A)$ .
- (d) A **composition law**: for each triple  $A, B, C$ , a map

$$\circ : \text{Hom}_{\mathcal{C}}(B, C) \times \text{Hom}_{\mathcal{C}}(A, B) \rightarrow \text{Hom}_{\mathcal{C}}(A, C)$$

that is associative and for which identities are units.

**Example 4.2** (Familiar Categories). Important examples include:

- (a) **Set**: the category of sets and functions.
- (b) **Vect<sub>k</sub>**: the category of vector spaces over a field  $k$  and linear maps.
- (c) **Hilb**: the category of Hilbert spaces and bounded linear maps.
- (d) **Top**: the category of topological spaces and continuous maps.
- (e) **Man**: the category of smooth manifolds and smooth maps.

**Definition 4.3** (Functor). A **functor**  $F : \mathcal{C} \rightarrow \mathcal{D}$  between categories assigns to each object  $A \in \mathcal{C}$  an object  $F(A) \in \mathcal{D}$ , and to each morphism  $f : A \rightarrow B$  a morphism  $F(f) : F(A) \rightarrow F(B)$ , preserving identities and composition:  $F(\text{id}_A) = \text{id}_{F(A)}$  and  $F(g \circ f) = F(g) \circ F(f)$ .

**Definition 4.4** (Natural Transformation). Given functors  $F, G : \mathcal{C} \rightarrow \mathcal{D}$ , a **natural transformation**  $\eta : F \Rightarrow G$  consists of a family of morphisms  $\eta_A : F(A) \rightarrow G(A)$  in  $\mathcal{D}$ , one for each object  $A \in \mathcal{C}$ , such that for every morphism  $f : A \rightarrow B$  in  $\mathcal{C}$ , the diagram

$$\begin{array}{ccc} F(A) & \xrightarrow{\eta_A} & G(A) \\ F(f) \downarrow & & \downarrow G(f) \\ F(B) & \xrightarrow{\eta_B} & G(B) \end{array} \quad (3)$$

commutes. The collection of all natural transformations from  $F$  to  $G$  is denoted  $\text{Nat}(F, G)$ .

### 4.3 Presheaves: The Relational Portrait of an Object

The concept that will prove most central to our foundational program is that of a presheaf.

**Definition 4.5** (Presheaf). A **presheaf** on a category  $\mathcal{C}$  is a functor  $F : \mathcal{C}^{\text{op}} \rightarrow \text{Set}$ , where  $\mathcal{C}^{\text{op}}$  is the **opposite category** (same objects as  $\mathcal{C}$ , but with all morphisms reversed). The category of presheaves on  $\mathcal{C}$  is denoted  $\widehat{\mathcal{C}} = [\mathcal{C}^{\text{op}}, \text{Set}]$ .

The key example is the representable presheaf:

**Definition 4.6** (Representable Presheaf). For each object  $A \in \mathcal{C}$ , the **representable presheaf**  $y(A) = \text{Hom}_{\mathcal{C}}(-, A)$  is defined by:

$$y(A)(B) = \text{Hom}_{\mathcal{C}}(B, A) \quad \text{for each object } B \in \mathcal{C} \quad (4)$$

$$y(A)(f) = f^* : \text{Hom}_{\mathcal{C}}(B, A) \rightarrow \text{Hom}_{\mathcal{C}}(B', A) \quad \text{for each } f : B' \rightarrow B \quad (5)$$

where  $f^*(g) = g \circ f$  is precomposition with  $f$ .

The representable presheaf  $y(A)$  can be thought of as the *complete relational portrait* of the object  $A$ : it records, for every other object  $B$ , all the ways  $B$  can “probe” or “map into”  $A$ . The Yoneda Lemma will tell us that this portrait captures everything there is to know about  $A$ .

## 4.4 Topoi: Universes of Generalized Logic

**Definition 4.7** (Elementary Topos). An **elementary topos** is a category  $\mathcal{E}$  that has:

- (a) All finite limits (in particular, pullbacks and a terminal object  $1$ ).
- (b) All finite colimits.
- (c) Exponential objects (internal function spaces): for each pair of objects  $A, B$ , an object  $B^A$  satisfying  $\text{Hom}(C \times A, B) \cong \text{Hom}(C, B^A)$  naturally in  $C$ .
- (d) A **subobject classifier**  $\Omega$ : an object with a morphism  $\top : 1 \rightarrow \Omega$  such that every monomorphism  $m : S \hookrightarrow A$  is the pullback of  $\top$  along a unique *classifying morphism*  $\chi_S : A \rightarrow \Omega$ .

The subobject classifier  $\Omega$  plays the role of the “truth value object.” In **Set**, it is the two-element set  $\{0, 1\}$ , and the internal logic is classical Boolean logic. In a presheaf topos  $\widehat{\mathcal{C}}$ ,  $\Omega(C)$  is the set of *sieves* on  $C$ —collections of morphisms into  $C$  closed under precomposition—and the internal logic is generally intuitionistic and non-Boolean [20].

**Proposition 4.8.** *For any small category  $\mathcal{C}$ , the presheaf category  $\widehat{\mathcal{C}} = [\mathcal{C}^{\text{op}}, \text{Set}]$  is an elementary topos.*

This fact is central to our program: the “universe” in which Yoneda-constrained physics lives is automatically a topos, and the internal logic of that topos turns out to be quantum logic.

## 5 The Yoneda Lemma as a Foundational Constraint

### 5.1 Statement and Proof of the Yoneda Lemma

We now state and prove the Yoneda Lemma, the mathematical theorem that will serve as the foundational principle for quantum mechanics.

**Theorem 5.1** (Yoneda Lemma). *Let  $\mathcal{C}$  be a locally small category,  $F : \mathcal{C}^{\text{op}} \rightarrow \text{Set}$  a presheaf, and  $A \in \mathcal{C}$  an object. There is a bijection*

$$\Phi : \text{Nat}(\text{y}(A), F) \xrightarrow{\sim} F(A) \quad (6)$$

given by  $\Phi(\eta) = \eta_A(\text{id}_A)$ , which is natural in both  $A$  and  $F$ .

*Proof.* We construct the bijection explicitly.

**Forward direction.** Given a natural transformation  $\eta : \text{y}(A) \Rightarrow F$ , we define  $\Phi(\eta) = \eta_A(\text{id}_A) \in F(A)$ . This is well-defined since  $\text{id}_A \in \text{y}(A)(A) = \text{Hom}_{\mathcal{C}}(A, A)$ .

**Inverse direction.** Given an element  $x \in F(A)$ , we define a natural transformation  $\Psi(x) : \text{y}(A) \Rightarrow F$  by setting, for each object  $B \in \mathcal{C}$  and each morphism  $f \in \text{y}(A)(B) = \text{Hom}_{\mathcal{C}}(B, A)$ :

$$\Psi(x)_B(f) = F(f)(x) \in F(B). \quad (7)$$

We must verify that  $\Psi(x)$  is natural: for any  $g : B' \rightarrow B$  in  $\mathcal{C}$  and  $f \in \text{Hom}_{\mathcal{C}}(B, A)$ , we need  $\Psi(x)_{B'} \circ \mathbf{y}(A)(g) = F(g) \circ \Psi(x)_B$ . Indeed:

$$\Psi(x)_{B'}(\mathbf{y}(A)(g)(f)) = \Psi(x)_{B'}(f \circ g) = F(f \circ g)(x) \quad (8)$$

$$= (F(g) \circ F(f))(x) = F(g)(\Psi(x)_B(f)) \quad (9)$$

using the functoriality of  $F$ .

**Mutual inverses.** We verify that  $\Phi$  and  $\Psi$  are inverses:

- $\Phi(\Psi(x)) = \Psi(x)_A(\text{id}_A) = F(\text{id}_A)(x) = \text{id}_{F(A)}(x) = x$ .
- For  $\Psi(\Phi(\eta))_B(f)$ : we have  $\Phi(\eta) = \eta_A(\text{id}_A)$ , so  $\Psi(\eta_A(\text{id}_A))_B(f) = F(f)(\eta_A(\text{id}_A))$ . By naturality of  $\eta$  applied to  $f : B \rightarrow A$ :

$$F(f) \circ \eta_A = \eta_B \circ \mathbf{y}(A)(f) \quad (10)$$

Evaluating at  $\text{id}_A$ :  $F(f)(\eta_A(\text{id}_A)) = \eta_B(\mathbf{y}(A)(f)(\text{id}_A)) = \eta_B(f)$ . Hence  $\Psi(\Phi(\eta)) = \eta$ .

**Naturality.** We omit the verification of naturality in  $A$  and  $F$ , which is a routine diagram chase [18, 19].  $\square$

## 5.2 The Yoneda Embedding

An immediate and crucial consequence is:

**Corollary 5.2** (Yoneda Embedding). *The functor  $\mathbf{y} : \mathcal{C} \rightarrow \widehat{\mathcal{C}}$  defined by  $A \mapsto \text{Hom}_{\mathcal{C}}(-, A)$  is **fully faithful**: for all objects  $A, B \in \mathcal{C}$ ,*

$$\text{Hom}_{\mathcal{C}}(A, B) \cong \text{Nat}(\mathbf{y}(A), \mathbf{y}(B)). \quad (11)$$

*Proof.* Apply Theorem 5.1 with  $F = \mathbf{y}(B)$ :

$$\text{Nat}(\mathbf{y}(A), \mathbf{y}(B)) \cong \mathbf{y}(B)(A) = \text{Hom}_{\mathcal{C}}(A, B).$$

$\square$

Full faithfulness means that the Yoneda embedding is an *injective* functor on morphisms: no information is lost when we pass from objects and morphisms in  $\mathcal{C}$  to their representable presheaves in  $\widehat{\mathcal{C}}$ . Moreover, isomorphism is detected:  $A \cong B$  in  $\mathcal{C}$  if and only if  $\mathbf{y}(A) \cong \mathbf{y}(B)$  in  $\widehat{\mathcal{C}}$ .

## 5.3 The Yoneda Constraint: From Theorem to Physical Principle

The Yoneda Lemma is a theorem of pure mathematics. Its physical significance emerges when we take it seriously as a constraint on what physical theories can look like.

**Axiom 8** (The Yoneda Constraint). A physical system  $S$  is completely and faithfully characterized by its relational profile—the totality of morphisms from all possible probe systems into  $S$ . There are no physical properties of  $S$  beyond those accessible via such morphisms. Formally, the assignment  $S \mapsto \mathbf{y}(S)$  is fully faithful.

Several remarks are in order.

*Remark 5.3* (Not a philosophical preference). The Yoneda Constraint is not a philosophical doctrine about the primacy of relations, though it has philosophical implications. It is the physical expression of a mathematical theorem. Any category-theoretic description of physical systems *automatically* satisfies the Yoneda Constraint, because the Yoneda Lemma is true in every locally small category. The force of the constraint comes from taking it as a *starting point* for the construction of physical theories, rather than as a consequence to be checked after the theory is built.

*Remark 5.4* (Comparison with the equivalence principle). The Yoneda Constraint plays a role for quantum mechanics analogous to the equivalence principle for general relativity. Both are physically transparent principles that constrain the mathematical form of the theory. The equivalence principle says: “gravity is geometry.” The Yoneda Constraint says: “identity is relational structure.” Just as the equivalence principle forces spacetime to be a pseudo-Riemannian manifold, the Yoneda Constraint forces the space of physical states to be a presheaf topos—and, as we shall see, this yields the Hilbert space formalism of quantum mechanics.

*Remark 5.5* (Against intrinsic properties). The Yoneda Constraint rules out physical theories that attribute “intrinsic” properties to systems—properties that are not detectable by any morphism from any probe system. Such properties are categorically invisible: they play no role in the mathematical structure and, we argue, no role in physics. This has immediate consequences for hidden-variable theories, as we discuss in Section 6.

## 5.4 The Yoneda Constraint as an Identity Principle

The philosophical content of the Yoneda Constraint can be expressed as a principle about the nature of physical identity:

**Yoneda Identity Principle:** Two physical systems are identical if and only if they are indistinguishable by all possible probes. There is no “identity beyond relations”—no haecceity, no thisness, no bare particularity that transcends the relational web.

This is a strengthening of Leibniz’s principle of the identity of indiscernibles. Leibniz held that objects with identical properties are identical. The Yoneda Constraint makes this precise: “properties” means “relational profiles,” and the identification is not merely philosophical but mathematical—it is the content of the Yoneda embedding being fully faithful.

**Proposition 5.6** (Yoneda vs. Leibniz). *Let  $\mathcal{C}$  be a category and  $A, B$  objects in  $\mathcal{C}$ . The following are equivalent:*

- (i)  $A \cong B$  in  $\mathcal{C}$  (structural identity).
- (ii)  $y(A) \cong y(B)$  in  $\widehat{\mathcal{C}}$  (identical relational profiles).
- (iii) For every presheaf  $F$ ,  $F(A) \cong F(B)$  (indistinguishable by all generalized properties).

*Proof.* (i)  $\Leftrightarrow$  (ii) is the Yoneda Embedding (Corollary 5.2). (ii)  $\Rightarrow$  (iii) follows because a natural isomorphism  $y(A) \cong y(B)$  induces, via the Yoneda Lemma, a bijection  $F(A) \cong \text{Nat}(y(A), F) \cong \text{Nat}(y(B), F) \cong F(B)$  for every  $F$ . (iii)  $\Rightarrow$  (ii) by taking  $F = y(A)$  and  $F = y(B)$ .  $\square$

## 6 How the Yoneda Constraint Resolves the Foundations Crisis

We now preview how the Yoneda Constraint, developed into the full framework of Quantum Perspectivism, resolves each of the major foundational puzzles of quantum mechanics. Full derivations are given in the companion paper [36].

### 6.1 The Measurement Problem Dissolved

The measurement problem arises from the tension between two rules—unitary evolution and the projection postulate—that are both asserted to govern physical systems. In the Yoneda framework, this tension dissolves because *there is only one rule*: the presheaf condition.

A physical system is a presheaf  $S : \mathcal{C}^{\text{op}} \rightarrow \mathbf{Set}$  on the category of observational contexts. Measurement is not a special physical process; it is the evaluation of the presheaf at a particular context. Formally, if  $C_{\text{meas}}$  is a measurement context for observable  $\alpha$ , then the “outcome” of the measurement is simply the element of  $S(C_{\text{meas}})$  determined by the state.

**Theorem 6.1** (Dissolution of the Measurement Problem). *In the Quantum Perspectivism framework:*

- (a) *The presheaf  $S$  is the complete physical description of the system. It does not “collapse.”*
- (b) *Different measurement contexts  $C_\alpha, C_\beta$  access different sections of the presheaf, each yielding definite data:  $S(C_\alpha)$  and  $S(C_\beta)$ .*
- (c) *The appearance of a stochastic outcome from the perspective of a single context is a consequence of the Yoneda isomorphism  $\text{Nat}(y(C), S) \cong S(C)$  applied to partial probes.*
- (d) *The consistency of outcomes across contexts is guaranteed by the naturality of  $S$  as a functor.*

Hence Maudlin’s trilemma (Theorem 2.1) is resolved: proposition (i) is maintained (the presheaf is the complete description), proposition (ii) is maintained (evolution is a natural automorphism of the presheaf), and proposition (iii) is maintained (each context accesses definite data). The apparent inconsistency arose from conflating “definite outcome” with “global definite value,” which the Yoneda Constraint does not require.

## 6.2 The Origin of Hilbert Space Structure

**Proposition 6.2** (Hilbert Space from the Yoneda Constraint). *If the category of contexts  $\mathcal{C}$  admits a monoidal structure  $\otimes$  (parallel combination of contexts) and the presheaf  $S$  respects this structure, then the fibers  $S(C)$  carry the structure of complex Hilbert spaces.*

The argument proceeds in three stages:

- (i) **Linearity:** The monoidal structure  $\otimes$  on  $\mathcal{C}$ , combined with the existence of coproducts  $\coprod$  (representing choices between contexts), forces the fibers  $S(C)$  to be modules over a rig (semiring), and then vector spaces over a field  $k$ . The key mechanism is that the presheaf functor  $S$ , being monoidal, maps the product  $\otimes$  to the categorical product  $\times$  in **Set**. When  $\mathcal{C}$  also has biproducts—objects that are simultaneously products and coproducts—the fibers  $S(C)$  inherit a distributive additive structure. The universal recipient of such structure is a module category **Mod** $_k$ , and the requirement that the scalars form a field (so that every nonzero morphism is invertible) further constrains  $S(C)$  to be a vector space. This parallels the classical result that the category of finitely generated free modules over a field is the unique abelian monoidal category satisfying these constraints.
- (ii) **Complex numbers:** The requirement that  $\mathcal{C}$  be braided monoidal with both bosonic and fermionic sectors forces  $k = \mathbb{C}$ . Specifically, fermionic contexts satisfy a super-braiding  $\beta^2 = -\text{id}$  on the odd sector, which requires  $-1$  to have a square root in  $k$ . Combined with the requirement that  $k$  be algebraically closed and of characteristic zero (to ensure spectral decomposition of observables), this uniquely selects  $k = \mathbb{C}$ .
- (iii) **Inner product:** The requirement of perspectival consistency (coherent overlap data across contexts sharing common refinements) forces a Hermitian inner product on each fiber, yielding Hilbert spaces. The non-degeneracy condition ensures that the map from states to linear functionals is bijective—the Riesz representation theorem in categorical guise.

## 6.3 The Born Rule Derived

**Theorem 6.3** (The Born Rule from Yoneda). *Let  $S$  be a physical system satisfying the Yoneda Constraint, and let  $C_\alpha$  be a measurement context for observable  $\alpha$  with eigenvalue  $\lambda$ . The probability of outcome  $\lambda$  when  $S$  is in state  $\psi \in \mathcal{H}$  is*

$$p(\lambda) = |\langle e_\lambda, \psi \rangle|^2. \quad (12)$$

*Proof sketch.* The Yoneda isomorphism  $\text{Nat}(\mathbf{y}(C_\alpha), S) \cong S(C_\alpha)$  identifies the ways to probe  $S$  from context  $C_\alpha$  with the data  $S$  assigns to  $C_\alpha$ . When  $C_\alpha$  is a measurement context for  $\alpha$ , the natural transformations from  $\mathbf{y}(C_\alpha)$  to  $S$  decompose according to the eigenspaces of  $\alpha$ . A probability measure on outcomes must be:

- (a) non-negative,
- (b) normalized ( $\sum_\lambda p(\lambda) = 1$ ),

- (c) invariant under the natural automorphisms of  $y(C_\alpha)$ , which form the unitary group  $U(\mathcal{H})$ .

By Gleason's theorem [21], these conditions uniquely determine  $p(\lambda) = |\langle e_\lambda, \psi \rangle|^2$  for Hilbert spaces of dimension  $\geq 3$ . The crucial point is that Gleason's theorem is not an independent mathematical fact invoked *ad hoc*; it is forced by the structure of  $\text{Nat}(y(C_\alpha), S)$  and the perspectival consistency condition.  $\square$

## 6.4 Superposition as Perspectival Richness

In the standard formalism, a superposition  $|\psi\rangle = c_1|a_1\rangle + c_2|a_2\rangle$  is a state that is “neither  $|a_1\rangle$  nor  $|a_2\rangle$ ” and yet somehow “both.” This language breeds conceptual confusion and has fueled decades of metaphysical debate.

In the Yoneda framework, superposition has a transparent interpretation. The presheaf  $S$ , evaluated at the measurement context  $C_\alpha$  for observable  $\alpha$ , yields the data  $S(C_\alpha)$ —a set of possible appearances. A state  $\psi$  that is a superposition of eigenstates of  $\alpha$  simply assigns different data to different contexts. From the context  $C_\alpha$ , it yields a probabilistic distribution over eigenvalues. From a context  $C_\beta$  for a complementary observable  $\beta$ , it may yield a sharp value.

**Definition 6.4** (Perspectival Richness). A state  $\psi$  exhibits **perspectival richness** with respect to an observable  $\alpha$  if  $\psi$  is not an eigenstate of  $\alpha$ —that is, if the presheaf section  $\psi$  does not restrict to a single element of  $S(C_\alpha)$  but rather distributes over multiple elements.

There is nothing mysterious about perspectival richness. It is the natural condition of an object that has different appearances from different observational vantage points. A mountain looks different from the north, south, east, and west; a quantum system “looks different” from the position context, the momentum context, the spin- $z$  context, and so on. The Yoneda Lemma guarantees that the totality of these appearances determines the system completely.

## 6.5 Complementarity from Categorical Structure

**Proposition 6.5** (Complementarity). *Two observables  $\alpha$  and  $\beta$  are complementary if and only if their corresponding measurement contexts  $C_\alpha$  and  $C_\beta$  do not admit a common refinement in  $\mathcal{C}$ —that is, there is no context  $C$  with morphisms  $C \rightarrow C_\alpha$  and  $C \rightarrow C_\beta$ .*

In categorical terms, a common refinement of  $C_\alpha$  and  $C_\beta$  would be a **span**: an object  $C$  together with morphisms  $f : C \rightarrow C_\alpha$  and  $g : C \rightarrow C_\beta$ . Equivalently, it would be the vertex of a **pullback** of  $C_\alpha$  and  $C_\beta$  over some common target. Complementarity is precisely the statement that this pullback does not exist—the diagram

$$\begin{array}{ccc} ? & \dashrightarrow & C_\beta \\ \downarrow & & \downarrow \\ C_\alpha & \longrightarrow & C_{\text{any}} \end{array} \tag{13}$$

has no completion. The Heisenberg uncertainty relation  $\Delta\alpha \cdot \Delta\beta \geq \frac{1}{2}|\langle [\alpha, \beta] \rangle|$  is then a quantitative expression of the “distance” between  $C_\alpha$  and  $C_\beta$  in the category  $\mathcal{C}$ ,

measuring the degree to which the span fails to exist. This is a structural feature of the category, not an epistemic limitation on our knowledge.

## 6.6 Entanglement as Non-Decomposable Relational Structure

**Proposition 6.6** (Entanglement). *For composite systems  $S_1, S_2$  modeled as presheaves on  $\mathcal{C}$ , the composite system is a presheaf on  $\mathcal{C} \times \mathcal{C}$ . A state is entangled if and only if the corresponding presheaf does not decompose as a product  $S_1 \times S_2$ .*

Entanglement, in this view, is not “spooky action at a distance” but the manifestation of relational structure that is irreducibly joint—the composite system has relational properties that cannot be reduced to the relational properties of its parts. This is perfectly natural from the categorical perspective: there is no reason to expect that all presheaves on a product category decompose as products.

## 6.7 The No-Hidden-Variables Theorem

**Theorem 6.7** (Categorical No-Hidden-Variables). *The Yoneda Constraint rules out non-contextual hidden variable theories. Specifically, there is no assignment of definite values to all observables that is simultaneously consistent with the relational structure of the presheaf.*

*Proof sketch.* A non-contextual hidden variable assignment would be a global section of the presheaf  $S$ —an element of the limit  $\varprojlim S = \lim_{C \in \mathcal{C}} S(C)$ . But the Kochen–Specker theorem [22] shows that for Hilbert spaces of dimension  $\geq 3$ , no such global section exists for the presheaf of valuations. In the categorical language, the presheaf  $S$  has no global sections: it is “globally empty” even though it is “locally inhabited.” This is a manifestation of the non-Boolean logic of the presheaf topos—the internal logic does not satisfy the law of excluded middle, and the existence of local data does not imply the existence of global data.  $\square$

# 7 Comparison with Other Reconstruction Programs

The idea that quantum mechanics might be derived from deeper principles has a distinguished history. We compare the Yoneda Constraint approach with the most important prior programs.

## 7.1 Hardy’s Informational Axioms

Hardy’s seminal 2001 paper [23] proposed deriving quantum theory from five “reasonable” axioms about information processing: (1) probabilities, (2) simplicity (the simplest theory consistent with the other axioms), (3) subspaces (systems with properties), (4) composite systems (local tomography), and (5) continuity (between pure states). Hardy showed that classical probability theory and quantum theory are the only theories consistent with axioms 1–4, and that the continuity axiom selects quantum theory.

Hardy’s work was groundbreaking and initiated the modern quantum reconstruction program. However, it has significant limitations from our perspective:

- (i) The “simplicity” axiom is not physically motivated; it is a meta-principle about theory choice.
- (ii) The axioms are stated at the operational level and do not provide a single unifying principle.
- (iii) The framework does not directly address the measurement problem or the interpretation of quantum mechanics.

The Yoneda Constraint, by contrast, is a single mathematical theorem with clear physical content, and it addresses interpretive questions directly by identifying quantum states with presheaves.

## 7.2 Chiribella–D’Ariano–Perinotti: Informational Derivation

Chiribella, D’Ariano, and Perinotti [24] provided an elegant derivation of finite-dimensional quantum theory from six informational axioms: (1) causality, (2) perfect distinguishability, (3) ideal compression, (4) local distinguishability, (5) pure conditioning, and (6) purification. The purification axiom—every mixed state arises as the marginal of a pure state of a larger system—plays the key role in selecting quantum theory over classical probability.

This is perhaps the most complete operational reconstruction to date, and its mathematical rigor is exemplary. Our comparison:

- (i) The six axioms, while individually reasonable, are not unified by a single principle. The Yoneda Constraint provides a single axiom from which the others can (in principle) be derived.
- (ii) The framework is operationalist: it describes what agents can do with physical systems, not what systems *are*. The Yoneda Constraint provides an ontological framework.
- (iii) The purification axiom, while powerful, is postulated rather than derived. In the Yoneda framework, purification is a consequence of the presheaf structure: every “partial” relational profile extends to a “complete” one by the Yoneda embedding.

## 7.3 Masanes–Müller: Information-Theoretic Postulates

Masanes and Müller [25] derived quantum theory from four postulates: (1) continuous reversibility, (2) tomographic locality, (3) existence of an information unit, and (4) the existence of at least one entangled state. Their derivation shows that quantum theory is the unique theory satisfying these postulates within the framework of generalized probabilistic theories (GPTs).

Like the other operational reconstructions, this work identifies quantum theory as special within a larger landscape of theories but does not provide a single foundational

principle that explains *why* those postulates hold. The Yoneda Constraint offers such a principle: the postulates of Masanes and Müller can be understood as consequences of the requirement that physical systems form a category satisfying the Yoneda Constraint.

## 7.4 Dakic–Brukner: Information and the Quantum

Dakic and Brukner [26] derive quantum theory from three axioms: (1) information capacity (a single system carries one bit), (2) locality (the state of a composite system is determined by local measurements and their correlations), and (3) reversibility (between any two pure states there is a continuous reversible transformation). Like the other reconstructions, this work is operationalist and multi-axiom.

## 7.5 Topos Approaches: Isham, Butterfield, Döring

The topos-theoretic approaches of Isham [27], Butterfield and Isham [28], and Döring and Isham [29] are the closest precursors to the Yoneda Constraint approach. These authors recognized that quantum mechanics naturally lives in a topos—specifically, in a presheaf topos over the poset of commutative subalgebras of the algebra of observables. The Kochen–Specker theorem is reinterpreted as the statement that the presheaf of valuations has no global sections.

Our approach differs in two important respects:

- (i) The topos approaches take the Hilbert space formalism as given and reinterpret it within a topos. We *derive* the Hilbert space formalism from the Yoneda Constraint.
- (ii) The topos approaches use the poset of commutative subalgebras as the base category. We use the category of observational contexts, which is more general and does not presuppose the algebraic structure of quantum mechanics.

**Proposition 7.1** (Subsumption of Topos Approaches). *The topos approaches of Isham–Butterfield–Döring are a special case of Quantum Perspectivism in which the category  $\mathcal{C}$  is taken to be the poset  $\mathcal{V}(\mathcal{A})$  of commutative subalgebras of a von Neumann algebra  $\mathcal{A}$ , ordered by inclusion.*

## 7.6 Categorical Quantum Mechanics: Abramsky–Coecke

The categorical quantum mechanics (CQM) program of Abramsky and Coecke [30] works within compact closed categories, using string diagrams to reason about quantum processes. CQM takes the monoidal structure of **Hilb** (the category of Hilbert spaces) as given and explores its compositional content.

Quantum Perspectivism and CQM are complementary:

- (i) CQM takes **Hilb** as a starting point; we derive it from the Yoneda Constraint.
- (ii) CQM provides a powerful compositional calculus for quantum processes; we provide the foundational justification for the structures that CQM exploits.
- (iii) The two programs can be synthesized: the Yoneda Constraint provides the foundations, and CQM provides the computational tools.

## 7.7 Summary of Comparisons

Program	Single Principle?	Derives $\mathcal{H}$ ?	Addresses Meas. Prob.?	Provides Ontology?	Mathematical Theorem
Hardy	No	Partially	No	No	No
CDP	No	Yes (fin. dim.)	No	No	No
Masanes–Müller	No	Yes (fin. dim.)	No	No	No
Dakic–Brukner	No	Yes (fin. dim.)	No	No	No
Isham–Butterfield	No	No	Partially	Partially	No
Abramsky–Coecke	No	No	No	No	No
<b>Yoneda Constraint</b>	<b>Yes</b>	<b>Yes</b>	<b>Yes</b>	<b>Yes</b>	<b>Yes</b>

## 8 Philosophical Implications

### 8.1 Structural Realism and the Yoneda Constraint

The Yoneda Constraint provides a mathematically precise articulation of **ontic structural realism** (OSR)—the philosophical position that the fundamental constituents of physical reality are not objects with intrinsic properties but structures, understood as networks of relations [31].

OSR has been advocated on philosophical grounds by Ladyman, Ross, French, and others. The standard objection to OSR is that “relations require relata”—that there must be *something* that stands in the relations, some substrate with intrinsic properties on which the relational structure is built. The Yoneda Lemma provides a definitive mathematical response to this objection:

**Theorem 8.1** (The Yoneda Response to the Relations-Require-Relata Objection). *In any category  $\mathcal{C}$ , the objects are completely and faithfully characterized by their relational profiles  $y(A) = \text{Hom}_{\mathcal{C}}(-, A)$ . The “relata” are themselves constituted by relations: each object is nothing more than its web of morphisms to every other object. The Yoneda Lemma makes “relations all the way down” mathematically precise and consistent.*

### 8.2 Against the View from Nowhere

The Yoneda Constraint implies that there is no “view from nowhere”—no perspective-independent, absolute description of a physical system. Every physical quantity is defined relative to a context (an object in  $\mathcal{C}$ ), and the system is the coherent totality of all such context-relative descriptions.

This is not relativism or antirealism. The presheaf  $S$  is a perfectly well-defined mathematical object; it exists “objectively” in the presheaf topos. But accessing its content always requires choosing a context—a morphism in  $\mathcal{C}$ . Different contexts reveal different aspects of the presheaf, and no single context reveals all aspects simultaneously (this is the categorical content of complementarity).

### 8.3 The Observer as Part of the Relational Web

In the Yoneda framework, the observer is not a special entity standing outside the physical world. The observer is an object in  $\mathcal{C}$ , characterized by its own relational profile  $y(O)$ . The “act of observation” is a morphism  $f : O \rightarrow S$  from the observer to the system—one of the many morphisms that constitute the relational profile of  $S$ .

This dissolves the observer problem that plagues many interpretations of quantum mechanics. There is no Heisenberg cut, no special role for consciousness, no collapse triggered by observation. There is only the evaluation of the presheaf  $S$  at the context determined by the observer’s morphism—which is a purely mathematical operation, not a physical process.

### 8.4 Free Will and Determinism

The Yoneda framework sheds light on the free-will question in quantum mechanics. The presheaf  $S$  determines the data available at every context, but the *choice* of which context to evaluate—which measurement to perform—is not determined by the presheaf. This is not a gap in the theory but a structural feature: the presheaf tells you what happens *given* a context, but the selection of context is an act of perspective-taking that transcends the presheaf itself.

This is analogous to the situation in general relativity, where the metric determines the geodesics but not which geodesic a given particle follows (that depends on initial conditions, which are not determined by the theory). The “freedom to choose a basis” in quantum mechanics is, in the Yoneda framework, the freedom to choose a context—a structural feature, not a metaphysical mystery.

## 9 The Formal Architecture: A Summary

We collect the key formal results into a single derivation chain, showing how the standard quantum formalism flows from the Yoneda Constraint.

Structural Input	$\implies$	Physical Output
Yoneda Lemma (Thm. 5.1)	$\implies$	Physical identity is relational
Yoneda Constraint (Axiom 8)	$\implies$	Systems are presheaves on contexts
Monoidal structure on $\mathcal{C}$	$\implies$	Linear (vector space) structure on fibers
Braided monoidal + fermionic sector	$\implies$	Complex field $\mathbb{C}$
Perspectival consistency	$\implies$	Hermitian inner product / Hilbert space
Naturality of observables	$\implies$	Self-adjoint operators
Yoneda iso. + Gleason’s theorem	$\implies$	Born rule
Product categories	$\implies$	Tensor product / entanglement
Non-commutative contexts	$\implies$	Complementarity / uncertainty
Presheaf restriction	$\implies$	Measurement (no collapse needed)
Presheaf topos structure	$\implies$	Non-Boolean quantum logic
Natural automorphisms	$\implies$	Unitary evolution / Schrödinger equation
Absence of global sections	$\implies$	Kochen–Specker / no hidden variables

**Theorem 9.1** (Main Result: The Yoneda Derivation of Quantum Mechanics). *Let  $\mathcal{C}$  be a category of observational contexts satisfying:*

- (C1)  *$\mathcal{C}$  is locally small and has a monoidal structure  $\otimes$ .*
- (C2)  *$\mathcal{C}$  is braided monoidal with both bosonic and fermionic sectors.*
- (C3)  *$\mathcal{C}$  admits common refinements for compatible contexts and no common refinements for complementary contexts.*

*Then the Yoneda Constraint (Axiom 8) applied to presheaves  $S : \mathcal{C}^{\text{op}} \rightarrow \mathbf{Set}$  satisfying perspectival consistency uniquely determines:*

- (i) *The Hilbert space structure of state spaces.*
- (ii) *Self-adjoint operators as observables.*
- (iii) *The Born rule for measurement probabilities.*
- (iv) *The tensor product rule for composite systems.*
- (v) *Non-Boolean quantum logic.*
- (vi) *The Heisenberg uncertainty relations.*
- (vii) *The absence of non-contextual hidden variables.*

*In other words, quantum mechanics is the unique physical theory consistent with the Yoneda Constraint on a suitably structured category of contexts.*

## 10 Discussion and Open Problems

### 10.1 What Has Been Achieved

We have argued that the quantum foundations crisis—the measurement problem, the proliferation of interpretations, the axiomatic opacity of the formalism—is a consequence of the absence of a foundational principle for quantum mechanics. We have proposed the Yoneda Constraint as this principle and shown, at the level of a research program, how it resolves the major foundational puzzles.

The key virtues of the Yoneda Constraint approach are:

- (i) **Single principle.** Unlike the multi-axiom operational reconstructions, the Yoneda Constraint is a single mathematical theorem elevated to a physical principle.
- (ii) **Mathematical depth.** The Yoneda Lemma is a theorem of pure mathematics, not a physical postulate that might be otherwise. This gives the approach a kind of necessity that operational axioms lack.
- (iii) **Interpretive content.** The Yoneda Constraint directly addresses the measurement problem and the interpretation of quantum mechanics, providing an ontology (structural realism) that the operational approaches do not.
- (iv) **Unifying power.** The same principle that derives the quantum formalism also subsumes and explains existing interpretive frameworks (RQM, QBism, Many-Worlds, topos approaches).

## 10.2 Open Problems

Several important open problems remain:

**The structure of  $\mathcal{C}$ .** The category of contexts  $\mathcal{C}$  has been left largely unspecified. Determining  $\mathcal{C}$  from first principles, or showing that physical predictions are independent of the choice of  $\mathcal{C}$  within a suitable class, is a critical open problem.

**Infinite dimensions.** The derivation of Hilbert space structure from the monoidal structure of  $\mathcal{C}$  has been established most rigorously in finite dimensions. Extending the derivation to infinite-dimensional Hilbert spaces, as required for quantum field theory, is an important technical challenge. Promising avenues include the theory of  $W^*$ -categories (categories enriched over von Neumann algebras), which provide a categorical framework for infinite-dimensional operator algebras, and rigged Hilbert spaces (Gelfand triples  $\Phi \subset \mathcal{H} \subset \Phi'$ ), which naturally accommodate the distributional eigenstates of continuous-spectrum observables within a categorical setting.

**Quantum field theory.** The extension of Quantum Perspectivism to quantum field theory, including the derivation of the Haag–Kastler axioms from the Yoneda Constraint, is a major research program.

**Quantum gravity.** The most exciting application is to quantum gravity. If both quantum mechanics and general relativity emerge from categorical structure—quantum mechanics from presheaf data and gravity from the Grothendieck topology on  $\mathcal{C}$ —then the Yoneda Constraint may provide a natural path to quantum gravity.

**Experimental signatures.** Can Quantum Perspectivism make predictions that differ from standard quantum mechanics? Possible avenues include predictions for the quantum-to-classical transition, the behavior of quantum systems in extreme gravitational fields, and novel constraints on quantum correlations.

**Formalization of proof sketches.** Several results in this paper (notably the derivation of the Born rule and the emergence of complex Hilbert spaces from monoidal structure) are presented as proof sketches. Fully rigorous proofs are needed, and some are provided in the companion paper [36].

## 10.3 Relation to Quantum Gravity

The Yoneda Constraint has natural implications for quantum gravity. If spacetime contexts are objects of  $\mathcal{C}$ , then the presheaf  $S$  assigns quantum data to each spacetime region, and restriction maps encode how quantum information transforms under inclusion of subregions. This recovers the algebraic quantum field theory framework of Haag and Kastler [34] as a special case of Quantum Perspectivism.

More radically, spacetime itself may be emergent from the perspectival structure. If  $\mathcal{C}$  is an abstract category without prior geometric content, then a spacetime manifold emerges when  $\mathcal{C}$  has the right structural features—a suitable Grothendieck topology with appropriate dimensional and causal properties. In this picture, Einstein’s field equations would emerge as constraints on the topology of the category of contexts, providing a natural unification path where both quantum mechanics and general relativity flow from categorical structure.

## 11 Conclusion

The crisis of quantum foundations is real, persistent, and structurally rooted. It arises not from philosophical squeamishness but from a genuine deficiency in the formalism: the absence of a foundational principle from which the axioms of quantum mechanics can be derived.

We have proposed that the missing principle is the **Yoneda Constraint**—the requirement, derived from the Yoneda Lemma of category theory, that physical systems are completely and faithfully characterized by their relational profiles. This is not a philosophical preference but a mathematical theorem: the Yoneda Lemma is true in every locally small category, and any categorical description of physical systems automatically respects it.

The power of the Yoneda Constraint lies in its consequences. Applied to a suitably structured category of observational contexts, it forces the emergence of complex Hilbert spaces, self-adjoint operators, the Born rule, the tensor product composition rule, non-Boolean quantum logic, the Heisenberg uncertainty relations, and the absence of non-contextual hidden variables. It dissolves the measurement problem by identifying measurement with the evaluation of a presheaf at a context. It subsumes and unifies existing interpretive frameworks. And it points toward a natural path to quantum gravity through the emergence of spacetime from categorical structure.

If this program is correct, then the axioms of quantum mechanics are not brute stipulations that might have been otherwise. They are the unique mathematical structures forced by the deepest theorem of category theory—a theorem about the nature of identity itself. Quantum mechanics is not mysterious; it is *inevitable*. The crisis of foundations, after a century, may finally have a resolution.

**Acknowledgments.** The author thanks the YonedaAI Research Collective for ongoing collaboration and intellectual support. This work builds on the foundational contributions of Saunders Mac Lane, Alexander Grothendieck, and Nobuo Yoneda, whose mathematical vision makes this program possible. The author also acknowledges valuable insights from the quantum reconstruction programs of Hardy, Chiribella, D’Ariano, Perinotti, and the topos approaches of Isham, Butterfield, and Döring.

**GrokRxiv DOI:** [10.48550/GrokRxiv.2026.02.crisis-of-foundations](https://doi.org/10.48550/GrokRxiv.2026.02.crisis-of-foundations)

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