

1 System: Surface ship

The nonlinear surge speed equation of a surface ship can be written¹

$$M_u \dot{u} + h(u) = \tau_u \quad (1)$$

where τ_u is the surge force control input, $M_u = m - X_{\dot{u}} > 0$, and $h(u) = -X_u u - X_{|u|u}|u|$ u is monotonically increasing, $h(0) = 0$, and $h(u)u > 0; \forall u \neq 0$. Let u_{ref} be a constant reference speed and choose the feedforward control law

$$\tau_u = h(u_{\text{ref}}). \quad (2)$$

1.1 Task: Function properties

Define

$$g(u) := h(u) - h(u_{\text{ref}}). \quad (3)$$

1. Show that $g(u_{\text{ref}}) = 0$.
2. Show that $g(u)(u - u_{\text{ref}}) > 0, \forall u \neq u_{\text{ref}}$.

1.2 Task: Lyapunov analysis

Show by using the Lyapunov function

$$V(u) = \frac{M_u}{2} (u - u_{\text{ref}})^2 + M_u \int_{u_{\text{ref}}}^u g(y) dy \quad (4)$$

that its time derivative along the solutions of the closed-loop system is given by

$$\dot{V}(u) = -(u - u_{\text{ref}})g(u) - g(u)^2,$$

and that the equilibrium $u - u_{\text{ref}} = 0$ is GAS.

2 System: Pendulum

The system at hand consists of a stand with a bearing. A rod is attached to the outer bearing race. A bob is attached to the other end of the rod. See Figure 1a.

Table 1 summarizes the parameters and variables of the installation.

The control plant model is based on that the angular acceleration $\dot{\omega}$ of the rotating part is proportional to the sum of torques:

$$\dot{\omega} = \frac{1}{J} \sum \tau, \quad (5)$$

where J is the system inertia.

Considering the bob as a point mass, the inertia is

$$J = ml^2. \quad (6)$$

¹Fossen, Handbook of Marine Craft Hydrodynamics and Motion Control, Eq. 7.32

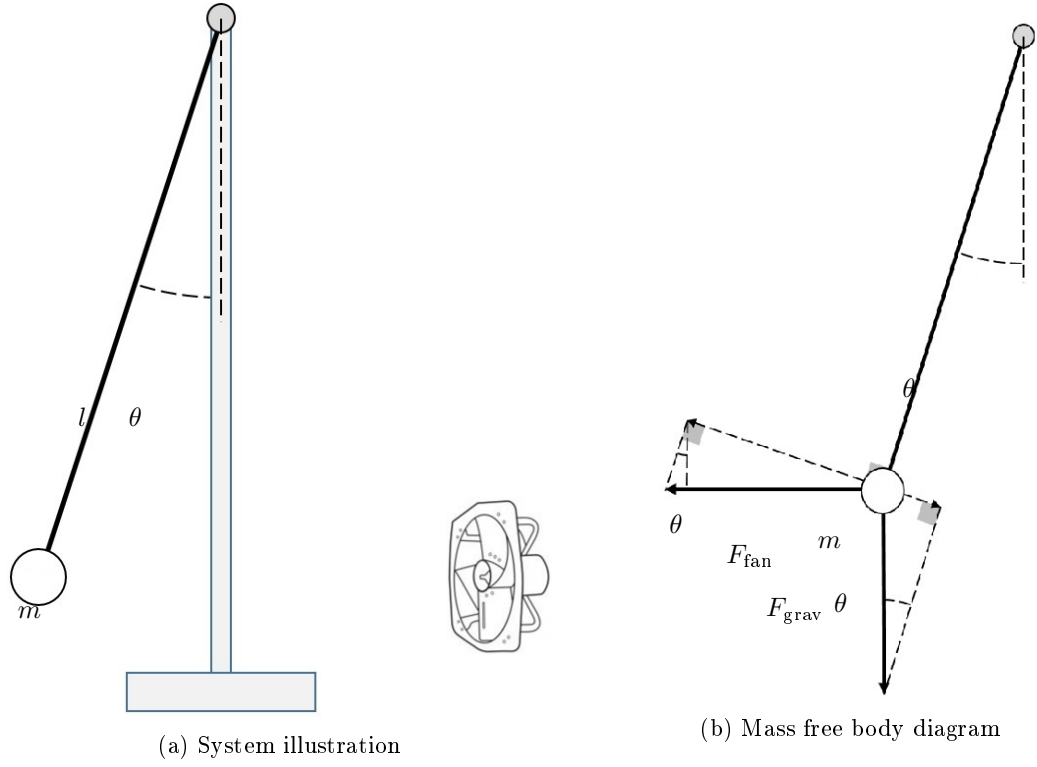


Figure 1: Pendulum

m	0.1 kg	mass of the bob
l	0.5 m	length of the rod attached to the bob
g	9.81 m/s^2	
k	0.01	friction coefficient
θ	rad	angle between the vertical stand and the bob rod
ω	rad/s	angular velocity

Table 1: Parameters and variables

For simplicity, this is chosen as the inertia of the whole system, thus disregarding the mass of the rod.

Figure 1b holds a diagram of the forces acting on the system. Since the forces parallel to the rod are counteracted by the latter, only the tangent components are of interest here:

- The torque τ_{grav} due to gravity is

$$\begin{aligned}\tau_{\text{grav}} &= -lF_{\text{grav}} \sin(\theta) \\ &= -lmg \sin(\theta).\end{aligned}\tag{7}$$

- The torque τ_{fan} due to the fan pressure on the rod is

$$\begin{aligned}\tau_{\text{fan}} &= lF_{\text{fan}} \sin\left(\frac{\pi}{4} - \theta\right), \\ &= lF_{\text{fan}} \cos(\theta)\end{aligned}\tag{8}$$

where F_{fan} is the force from the fan.

Additionally, the friction in the bearing is modeled by a torque τ_{fric} proportional to the velocity:

$$\tau_{\text{fric}} = -k\omega.\tag{9}$$

Inserting (7)-(9) and substituting (6) in (5) yields

$$\begin{aligned}\dot{\omega} &= \frac{1}{ml^2} (-lmg \sin(\theta) - k\omega + lF_{\text{fan}} \cos(\theta)) \\ &= -\frac{g}{l} \sin(\theta) - \frac{k}{ml^2} \omega + \frac{F_{\text{fan}}}{ml} \cos(\theta)\end{aligned}\tag{10}$$

2.1 Task: State equations

1. Write the state equation $\dot{x} = f(x, u)$ using the state vector $x = \begin{bmatrix} \theta \\ \omega \end{bmatrix}$ and $u = F_{\text{fan}}$.
2. Program a corresponding Simulink model.

2.2 Task: System properties

Assuming $|F_{\text{fan}}|$ is bounded, explain why the system is or isn't:

1. Forward complete.
2. Backward complete.
3. Complete
4. Locally Lipschitz.
5. Globally Lipschitz.

2.3 Task: Simple pendulum equilibrium point

Assume the fan is off, i.e. $F_{\text{fan}} = 0$ N.

1. Show that $x^* = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$ is an equilibrium point of the unforced system.
2. Explain why x^* is or isn't
 - (a) a unique equilibrium point,
 - (b) an isolated equilibrium point.

3. Describe the physical situation(s) the equilibrium point(s) correspond(s) to.
4. Simulate the system with initial condition $x(0) = x^*$ to confirm the behavior at the equilibrium point.

2.4 Task: Linearized simple pendulum model

The linearized state equations are

$$\begin{bmatrix} \dot{\theta} \\ \dot{\omega} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{g}{l} & -\frac{k}{ml^2} \end{bmatrix} \begin{bmatrix} \theta \\ \omega \end{bmatrix}.$$

1. Show that x^* is also an equilibrium point for the linearized system.
2. Explain why x^* is or isn't:
 - (a) Locally stable.
 - (b) Globally stable.

2.5 Task: Equilibrium point with fan

Assume that the fan is again running, with $F_{\text{fan}} = 0.56638$ N.

1. Calculate the angle at which the bob now stabilizes.
2. Confirm through simulation.

2.6 Task: Angle control

In order to set the fan to stabilize the bob at $\theta = 60^\circ$,

1. choose a change of variables such that the equilibrium is shifted to this angle,
2. write the state equations using the new states, and
3. determine F_{fan} necessary to the new equilibrium.
4. Confirm through simulation.

3 Lyapunov function

Consider the differential equations

$$\dot{x}_1 = u_1 \tag{11a}$$

$$\dot{x}_2 = u_2, \tag{11b}$$

and a function

$$V(x_1, x_2) = c_1 x_1^2 + c_2 x_1 x_2 + c_3 x_2^2, \tag{12}$$

$$V = \begin{bmatrix} x_1 & x_2 \end{bmatrix} \begin{bmatrix} c_1 & \frac{c_2}{2} \\ \frac{c_2}{2} & c_3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \tag{13}$$

where c_1, c_2, c_3 are positive scalars.

3.1 Task: Differentiation

Do the following:

1. Differentiate $V(x_1, x_2)$ with respect to time (we say, “along the solutions of (11)”) and set up the resulting expression in terms of the states (x_1, x_2) and the inputs (u_1, u_2) .
2. Let $x := \text{col}(x_1, x_2)$ and $u := \text{col}(u_1, u_2)$. Show that V can be written as $V(x) = x^\top P x$ where $P = P^\top$ (symmetric).
3. Give conditions on (c_1, c_2, c_3) for P to be a positive definite matrix ($P = P^\top > 0$).
4. Show that taking the vector differentiation of $V(x)$ gives $\dot{V} = 2x^\top P \dot{x}$ and that this equals the answer in the Subtask 1.
5. Let $u_1 = -x_1 + x_2$ and $u_2 = -x_2$ such that the closed-loop system becomes

$$\dot{x} = Ax \tag{14}$$

$$A = \begin{bmatrix} -1 & 1 \\ 0 & -1 \end{bmatrix}. \tag{15}$$

Choose values for (c_1, c_2, c_3) such that $P = P^\top > 0$ and $PA + A^\top P = -Q$ where $Q > 0$ is a diagonal positive matrix.

6. Differentiate again $V(x)$ along the solutions of (14) and show that $\dot{V} = -x^\top Q x$ for your chosen values of (c_1, c_2, c_3) .

	$h(u) - h(u_{\text{ref}})$	$u - u_{\text{ref}}$
$u > u_{\text{ref}}$	+	+
$u < u_{\text{ref}}$	-	-

Table 2: Sign of terms given ratio of u and u_{ref}

1 Solution: Surface ship

1.1 Task: Function properties

Inserting $u = u_{\text{ref}}$ in (3) yields

$$\begin{aligned} g(u_{\text{ref}}) &= h(u_{\text{ref}}) - h(u_{\text{ref}}) \\ &= 0. \end{aligned} \tag{16}$$

Also,

$$g(u)(u - u_{\text{ref}}) = (h(u) - h(u_{\text{ref}}))(u - u_{\text{ref}}). \tag{17}$$

Since $h(\cdot)$ is monotonically increasing,

$$\begin{aligned} u > u_{\text{ref}} &\Leftrightarrow h(u) > h(u_{\text{ref}}), \text{ and} \\ u < u_{\text{ref}} &\Leftrightarrow h(u) < h(u_{\text{ref}}) \end{aligned}$$

by definition. By inspection, summarized in Table 2, the two terms of (17) have the same sign when $u \neq u_{\text{ref}}$ and the product is thus positive.

1.2 Task: Lyapunov analysis

Differentiation of (4) yields

$$\dot{V}(u) = M_{\dot{u}}(u - u_{\text{ref}})(\dot{u} - \dot{u}_{\text{ref}}) + M_{\dot{u}}g(u)\dot{u}. \tag{18}$$

Rearranging (1), then substituting (2), then (3) gives

$$\begin{aligned} \dot{u} &= \frac{\tau_u - h(u)}{M_u}. \\ &= \frac{h(u_{\text{ref}}) - h(u)}{M_u}. \\ &= -\frac{g(u)}{M_u}. \end{aligned} \tag{19}$$

Also, since u_{ref} is constant,

$$\dot{u}_{\text{ref}} = 0. \tag{20}$$

Applying (19), (20) and (16) to (18) yields

$$\begin{aligned} \dot{V}(u) &= M_{\dot{u}}(u - u_{\text{ref}})\left(-\frac{g(u)}{M_u} - 0\right) + M_{\dot{u}}g(u)\left(-\frac{g(u)}{M_u}\right). \\ &= -(u - u_{\text{ref}})g(u) - g(u)^2 =: -\alpha(|u|) < 0, \quad \forall |u| \neq 0 \end{aligned} \tag{21}$$

The speed equilibrium is at $\dot{u} = 0$. From (19) it follows that this requires $g(u) = 0$. (16) reveals $g(u_{\text{ref}})$ as one equilibrium. The Lyapunov function is

- $\dot{V}(u_{\text{ref}}) = 0$, by (21) with (16), and
- $\dot{V}(u) < 0 \forall u \neq u_{\text{ref}}$, by the inspection of (17) and since $g(u)^2 \geq 0$,

thus the equilibrium is globally asymptotically stable.

2 Solution: Pendulum

2.1 Task: State equations

$$\begin{aligned} \dot{x} &= f(x, u) \\ \begin{bmatrix} \dot{\theta} \\ \dot{\omega} \end{bmatrix} &= \begin{bmatrix} \omega \\ -\frac{g}{l} \sin(\theta) - \frac{k}{ml^2} \omega + \frac{1}{ml} \cos(\theta) F_{\text{fan}} \end{bmatrix} \end{aligned}$$

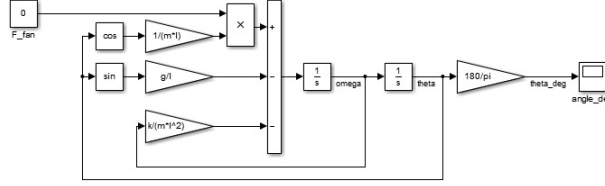


Figure 2: Simulink diagram

2.2 Task: System properties

Assuming $|F_{\text{fan}}|$ is bounded, explain why the system is or isn't:

1. Forward complete: This means that the solutions exists for all forward time. We get for the system that

$$\begin{aligned} |\dot{\theta}| &\leq |\omega| \leq |x| \\ |\dot{\omega}| &\leq \left| -\frac{g}{l} \sin(\theta) - \frac{k}{ml^2} \omega + \frac{1}{ml} \cos(\theta) F_{\text{fan}} \right| \\ &\leq \frac{g}{l} + \frac{k}{ml^2} |\omega| + \frac{1}{ml} |F_{\text{fan}}| \leq a + b |x| \end{aligned}$$

We see that the derivatives have linear sector bounds in the states, and it follows that the solutions can only grow to infinity by an exponential rate - which takes infinite time². Hence, the solutions must be forward complete.

2. Backward complete: Assume for simplicity that $F_{\text{fan}}(t) = 0$ for $t < 0$. Substitute $s = -t$ as a new free variable running backwards in time, and $z = x(s(t))$ is the corresponding solution of x backwards in time. Then

$$\dot{z} = \frac{dx}{ds} \frac{ds}{dt} = -f(x(s), 0) = -f(z, 0).$$

By the same argument as above, this system is forward complete. Since the solutions of $z(t)$ forward in time is the same as the solutions of $x(t)$ backwards in time, it follows that the pendulum system is backwards complete.

3. Complete: Since the system is both forward complete and backwards complete, it is complete.
4. Locally Lipschitz: Since all p -norms are equivalent, it is easiest to use the 1-norm (note: $|x|_2 \leq$

²This is formally proved in Proposition A.1 in Skjetne et al. (2005).

$|x|_1 \leq \sqrt{2}|x|_2$). Assume $u_{\max} := \|F_{\text{fan}}\|_{\infty}$. We check

$$\begin{aligned}
& |f(x, u) - f(y, u)|_2 \leq |f(x, u) - f(y, u)|_1 \\
&= \left\| \begin{bmatrix} -\frac{g}{l} (\sin(\theta_x) - \sin(\theta_y)) - \frac{k}{ml^2} (\omega_x - \omega_y) + \frac{1}{ml} (\cos(\theta_x) - \cos(\theta_y)) F_{\text{fan}} \end{bmatrix} \right\|_1 \\
&= \left\| \begin{bmatrix} -2\frac{g}{l} \cos\left(\frac{\theta_x + \theta_y}{2}\right) \sin\left(\frac{\theta_x - \theta_y}{2}\right) - \frac{k}{ml^2} (\omega_x - \omega_y) - 2\frac{1}{ml} \sin\left(\frac{\theta_x + \theta_y}{2}\right) \sin\left(\frac{\theta_x - \theta_y}{2}\right) F_{\text{fan}} \end{bmatrix} \right\|_1 \\
&= |\omega_x - \omega_y| \\
&\quad + \left| -2\frac{g}{l} \cos\left(\frac{\theta_x + \theta_y}{2}\right) \sin\left(\frac{\theta_x - \theta_y}{2}\right) - \frac{k}{ml^2} (\omega_x - \omega_y) - 2\frac{1}{ml} \sin\left(\frac{\theta_x + \theta_y}{2}\right) \sin\left(\frac{\theta_x - \theta_y}{2}\right) F_{\text{fan}} \right| \\
&\leq |\omega_x - \omega_y| + \frac{g}{l} |\theta_x - \theta_y| + \frac{k}{ml^2} |\omega_x - \omega_y| + \frac{u_{\max}}{ml} |\theta_x - \theta_y| \\
&\leq \max \left\{ 1, \frac{g}{l}, \frac{k}{ml^2}, \frac{u_{\max}}{ml} \right\} |x - y|_1 \leq L |x - y|_2 \\
&\quad L = \sqrt{2} \max \left\{ 1, \frac{g}{l}, \frac{k}{ml^2}, \frac{u_{\max}}{ml} \right\}
\end{aligned}$$

5. Globally Lipschitz: Since the above derived L is valid for all $x \in \mathbb{R}^2$, the above derivation in fact shows that

2.3 Task: Simple pendulum equilibrium point

The equilibrium requires, i.e.

$$0 = \omega \tag{22}$$

$$0 = -\frac{g}{l} \sin(\theta) - \frac{k}{ml^2} \omega + \frac{0}{ml} \cos(\theta). \tag{23}$$

From (22) it follows that $0 = \omega$. Substitution in (23) gives $0 = -\frac{g}{l} \sin(\theta)$, thus $0 = \sin(\theta)$ since g and l are constants. Equilibrium points are thus at

$$(\theta, \omega) = (\pm n\pi, 0) \quad \forall n \in \mathbb{Z}.$$

For even n (including $n = 0$), the bob stands still in the lower position with the rod counteracting gravity and no friction effect. For odd n , the bob stands still upright. The infinite amount of equilibrium points corresponds to the bob assuming the equilibrium angles after any number of rotations.

2.4 Task: Linearized simple pendulum model

Setting $\dot{\theta} = \dot{\omega} = 0$ for the linearized equations shows that $x^* = \text{col}(0, 0)$ is an equilibrium point also for the linearized system.

1. Explain why x^* is or isn't:

(a) Locally stable: We find that the poles of the linearized A -matrix becomes

$$\begin{aligned}
\det(sI - A) &= \det \begin{bmatrix} s & -1 \\ \frac{g}{l} & s + \frac{k}{ml^2} \end{bmatrix} = s^2 + \frac{k}{ml^2} s + \frac{g}{l} \\
s_0 &= -\frac{k}{2ml^2} \pm \frac{1}{2} \sqrt{\left(\frac{k}{ml^2}\right)^2 - 4\frac{g}{l}}
\end{aligned}$$

We see that the real part of the poles will always be negative for $k, m, l > 0$, and hence the linearized system is exponentially stable. For the nonlinear system we can then conclude that $x^* = \text{col}(0, 0)$ is locally stable.

- (b) Globally stable: The linearized system, by itself, is globally stable. However, since the non-linear system has multiple equilibria, for $\theta^* = k\pi$, $k \in \mathbb{Z}$, it cannot be globally stable. In fact we will find that every equilibria $(\theta^*, \omega^*) = (2k\pi, 0)$, $k \in \mathbb{Z}$ (pendulum hanging downright), is locally stable, and $(\theta^*, \omega^*) = ((2k+1)\pi, 0)$, $k \in \mathbb{Z}$ (pendulum staying upright), is locally unstable.

2.5 Task: Equilibrium point with fan

(22)-(23) yields

$$\frac{g}{l} \sin(\theta) = \frac{F_{\text{fan}}}{ml} \cos(\theta)$$

which in turn gives

$$\begin{aligned} \frac{\sin(\theta)}{\cos(\theta)} &= \frac{F_{\text{fan}}}{mg} \\ \Downarrow \\ \theta &= \tan^{-1}\left(\frac{F_{\text{fan}}}{mg}\right) \end{aligned} \tag{24}$$

which corresponds to $\theta = 0.5236 \text{ rad} = 30^\circ$.

2.6 Task: Angle control

Introducing a new variable

$$\tilde{\theta} = \theta - \frac{\pi}{3}$$

gives

$$\begin{aligned} \theta &= \tilde{\theta} + \frac{\pi}{3}, \text{ and} \\ \dot{\tilde{\theta}} &= \dot{\theta}. \end{aligned}$$

The state equations are then

$$\begin{aligned} \dot{\tilde{\theta}} &= \omega \\ \dot{\omega} &= -\frac{g}{l} \sin\left(\tilde{\theta} + \frac{\pi}{3}\right) - \frac{k}{ml^2} \omega + \frac{F_{\text{fan}}}{ml} \cos\left(\tilde{\theta} + \frac{\pi}{3}\right) \end{aligned}$$

The equilibrium requires $\dot{\tilde{\theta}} = \dot{\omega} = 0$, also $\tilde{\theta} = 0$ here, i.e.

$$\begin{aligned} 0 &= \omega \\ 0 &= -\frac{g}{l} \sin\left(0 + \frac{\pi}{3}\right) - \frac{k}{ml^2} \omega + \frac{F_{\text{fan}}}{ml} \cos\left(0 + \frac{\pi}{3}\right) \end{aligned}$$

$$\begin{aligned} \frac{g}{l} \sin\left(\frac{\pi}{3}\right) + \frac{k}{ml^2} 0 &= \frac{F_{\text{fan}}}{ml} \cos\left(\frac{\pi}{3}\right) \\ \Downarrow \\ F_{\text{fan}} &= mg \frac{\sin\left(\frac{\pi}{3}\right)}{\cos\left(\frac{\pi}{3}\right)} \\ &= mg \tan\left(\frac{\pi}{3}\right) \end{aligned}$$

Note: The task could also have been solved by solving (24) for F_{fan} , however revising changes of variables was part of the intention here.

3 Solution: Lyapunov function

3.1 Task: Differentiation

1. Differentiating $V(x_1, x_2)$ along the solutions of (11) results in

$$\dot{V} = \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 = (2c_1x_1 + c_2x_2)u_1 + (c_2x_1 + 2c_3x_2)u_2.$$

2. We get

$$\begin{aligned} V(x_1, x_2) &= c_1x_1^2 + c_2x_1x_2 + c_3x_2^2 \\ &= \begin{bmatrix} x_1 & x_2 \end{bmatrix} \begin{bmatrix} c_1 & \frac{c_2}{2} \\ \frac{c_2}{2} & c_3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \\ P &= \begin{bmatrix} c_1 & \frac{c_2}{2} \\ \frac{c_2}{2} & c_3 \end{bmatrix} \end{aligned}$$

3. For P to be a positive definite its leading principal minors must all be positive. This means that $c_1 > 0$ and $c_1c_3 - \frac{1}{4}c_2^2 > 0$. It follows that $c_3 > 0$ and $c_2 < 2\sqrt{c_1c_3}$.
4. Differentiating $V(x)$ gives

$$\begin{aligned} \dot{V} &= \dot{x}^\top Px + x^\top P\dot{x} = 2x^\top P\dot{x} = 2x^\top Pu \\ &= 2 \begin{bmatrix} x_1 & x_2 \end{bmatrix} \begin{bmatrix} c_1 & \frac{c_2}{2} \\ \frac{c_2}{2} & c_3 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix} \\ &= (2c_1x_1 + c_2x_2)u_1 + (c_2x_1 + 2c_3x_2)u_2 \quad \text{Q.E.D.} \end{aligned}$$

5. We have from above that $c_1 > 0$, $c_3 > 0$, and $c_2 < 2\sqrt{c_1c_3}$. Calculating $PA + A^\top P$ gives

$$\begin{aligned} &PA + A^\top P \\ &= \begin{bmatrix} c_1 & \frac{c_2}{2} \\ \frac{c_2}{2} & c_3 \end{bmatrix} \begin{bmatrix} -1 & 1 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} c_1 & \frac{c_2}{2} \\ \frac{c_2}{2} & c_3 \end{bmatrix} \\ &= - \begin{bmatrix} 2c_1 & -c_1 + c_2 \\ -c_1 + c_2 & 2c_3 - c_2 \end{bmatrix} =: -Q. \end{aligned}$$

Hence, for Q to be diagonal we need $c_1 = c_2 > 0$. Moreover, for it to be positive we need $c_1 > 0$ and $\det Q > 0$. This gives

$$\det Q = -c_1^2 + 4c_3c_1 - c_2^2 = 4c_3c_1 - 2c_1^2 = 2c_1(2c_3 - c_1) > 0.$$

Choosing for instance $c_3 = c_1 = c_2 > 0$ satisfies this, e.g. $c_1 = c_2 = c_3 = 1$.

6. Differentiating $V(x)$ gives

$$\dot{V} = \dot{x}^\top Px + x^\top P\dot{x} = (x^\top A^\top) Px + x^\top P(Ax) = x^\top (PA + A^\top P)x = -x^\top Qx$$

References

Skjetne, R., Fossen, T. I., and Kokotović, P. V. (2005). Adaptive maneuvering, with experiments, for a model ship in a marine control laboratory. *Automatica*, 41(2):289–298.