

MHD Waves in Asymmetric Slab Waveguides

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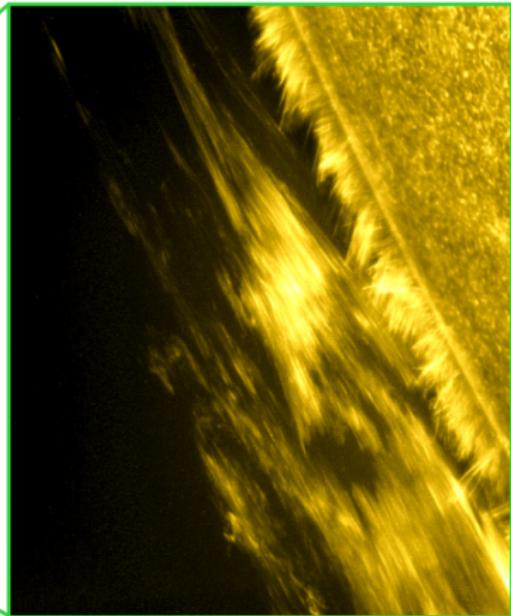
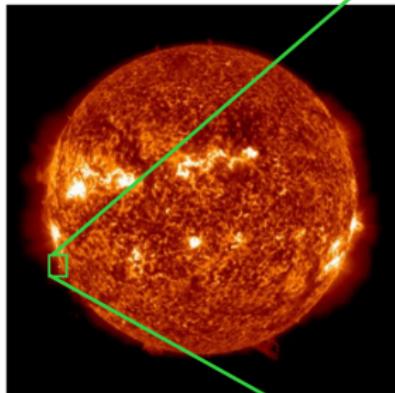
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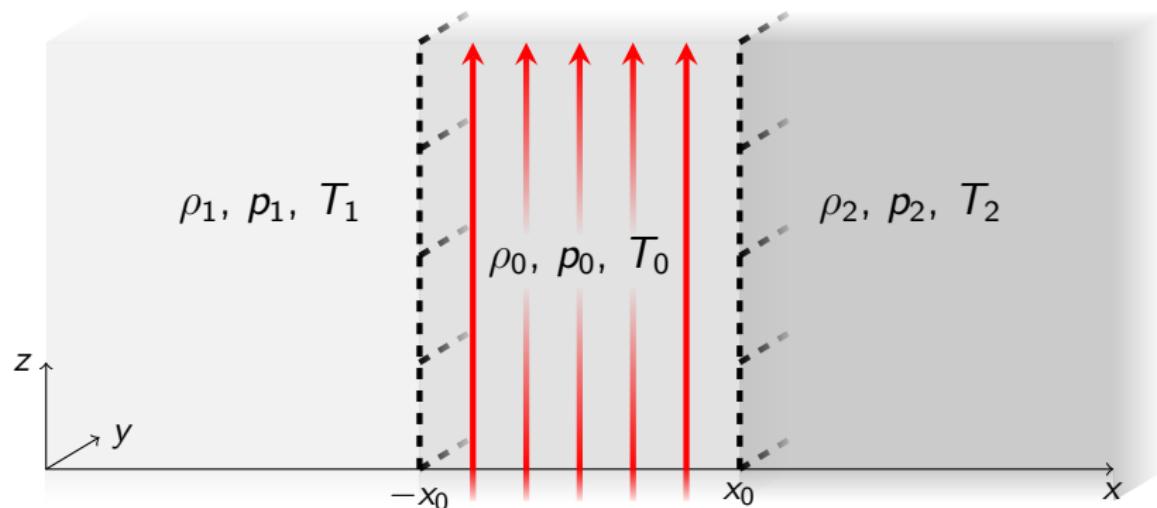


Motivation



- **Solar magneto-seismology,**
- **Numerical simulations,**
- **Coronal heating.**

Equilibrium conditions



- Uniform vertical magnetic field in the slab.
- Non-magnetised plasma outside.
- **Different** density and pressure on each side.

Governing equations

Ideal MHD equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0,$$

$$\rho \frac{D \mathbf{v}}{D t} = -\nabla p - \frac{1}{\mu} \mathbf{B} \times (\nabla \times \mathbf{B}),$$

$$\frac{D}{D t} \left(\frac{p}{\rho^\gamma} \right) = 0,$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}),$$

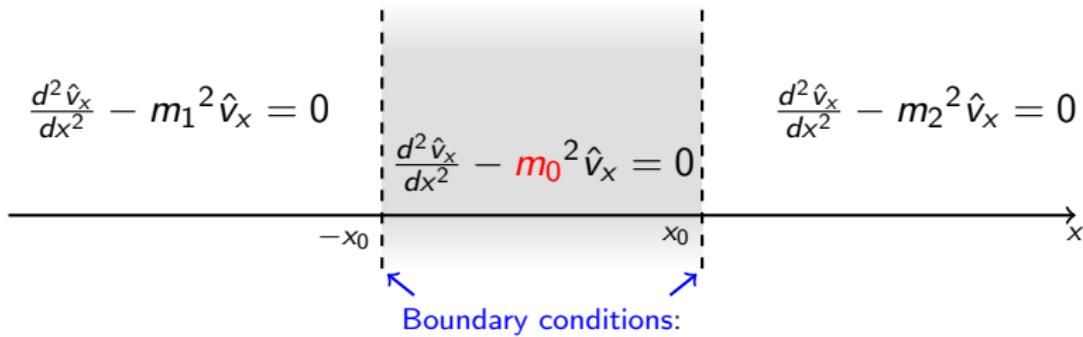
\mathbf{v} = plasma velocity, \mathbf{B} = magnetic field strength, ρ = density,
 p = pressure, μ = magnetic permeability, γ = adiabatic index.

Fourier decomposition

Look for *plane wave* solutions of the form:

$$v_x(\mathbf{x}, t) = \hat{v}_x(x) e^{i(kz - \omega t)}, \quad v_y(\mathbf{x}, t) = 0, \quad v_z(\mathbf{x}, t) = \hat{v}_z(x) e^{i(kz - \omega t)},$$

to arrive at the following ODEs:



pressure and velocity continuous across boundaries.

$$m_0^2 = \frac{(k^2 v_A^2 - \omega^2)(k^2 c_0^2 - \omega^2)}{(c_0^2 + v_A^2)(k^2 c_T^2 - \omega^2)}, \quad m_{1,2}^2 = k^2 - \frac{\omega^2}{c_{1,2}^2},$$

$$c_T^2 = \frac{c_0^2 v_A^2}{c_0^2 + v_A^2}, \quad v_A = \frac{B_0}{\sqrt{\mu \rho_0}},$$

Dispersion relation

The conditions for the existence of non-trivial solutions yield the dispersion relation (DR):

Dispersion relation

$$\frac{\omega^4 m_0^2}{k^2 v_A^2 - \omega^2} + \frac{\rho_0}{\rho_1} m_1 \frac{\rho_0}{\rho_2} m_2 (k^2 v_A^2 - \omega^2) - \frac{1}{2} m_0 \omega^2 \left(\frac{\rho_0}{\rho_1} m_1 + \frac{\rho_0}{\rho_2} m_2 \right) (\tanh m_0 x_0 + \coth m_0 x_0) = 0,$$

Approximate dispersion relation

When each side of the slab have **approximately equal** parameters, the asymmetric slab DR decouples to:

Approximate dispersion relation

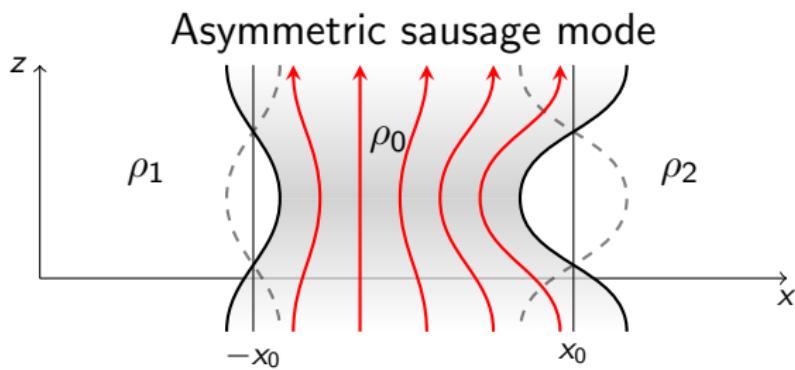
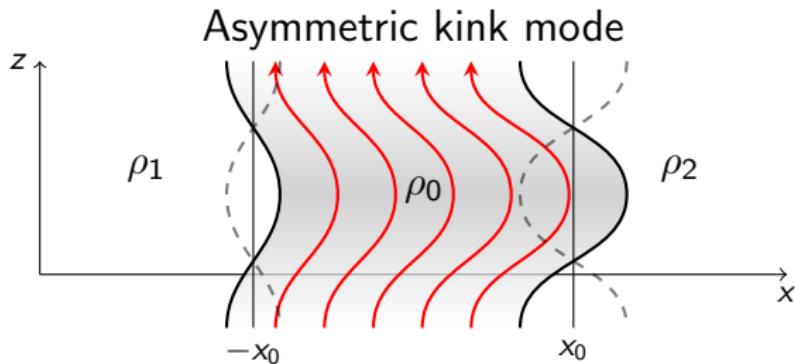
$$(k^2 v_A^2 - \omega^2) \left(\frac{\rho_0}{\rho_1} m_1 + \frac{\rho_0}{\rho_2} m_2 \right) = 2\omega^2 m_0 \left(\frac{\tanh}{\coth} \right) (m_0 x_0).$$

Compare with the symmetric slab DR, *Roberts 1981*:

Symmetric slab dispersion relation

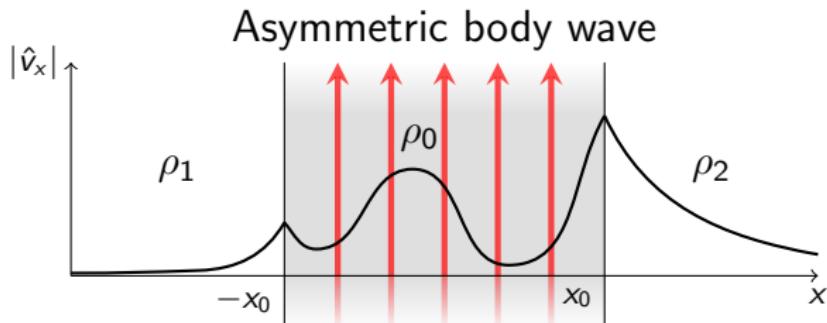
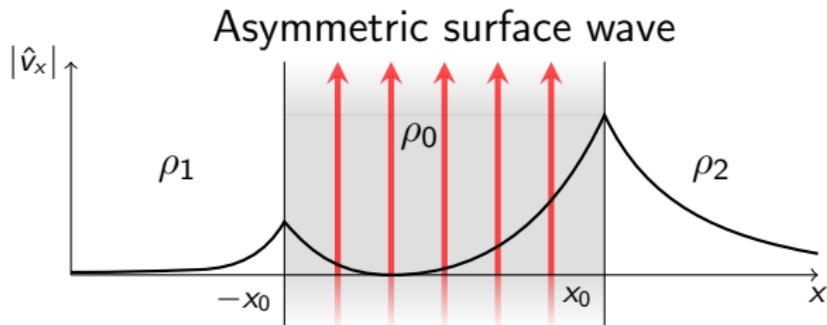
$$(k^2 v_A^2 - \omega^2) \frac{\rho_0}{\rho_e} m_e = \omega^2 m_0 \left(\frac{\tanh}{\coth} \right) (m_0 x_0).$$

Asymmetric sausage and kink modes

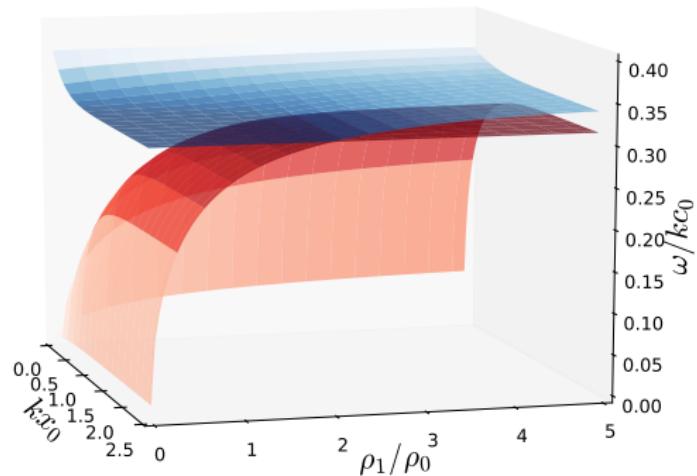


Asymmetric surface and body modes

- When $m_0^2 > 0$, the solution is **exponential**: *surface wave*.
- When $m_0^2 < 0$, the solution is **oscillatory**: *body wave*.

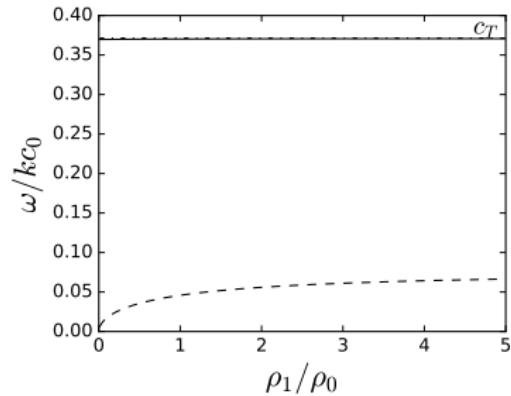


Density variation

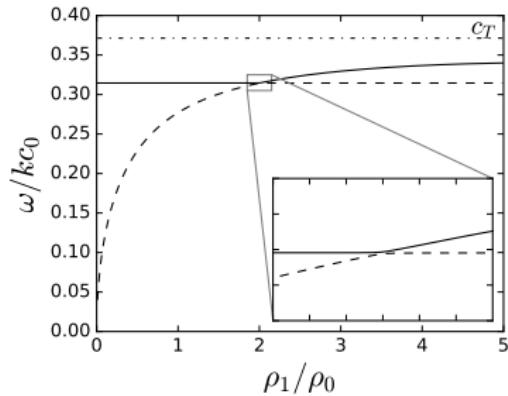


Slow surface modes of the **asymmetric** slab varying slab width (kx_0) and density ratio (ρ_1/ρ_0).

Density variation



(a) Small kx_0

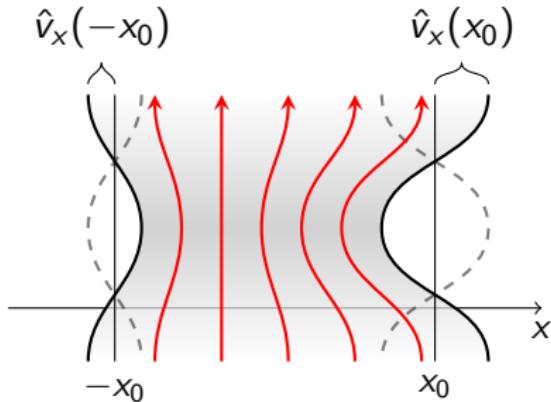


(b) Large kx_0

Avoided crossing

Can energy be transferred between sausage and kink surface modes of an asymmetric slab?

Cross-slab amplitude ratio



Cross-slab amplitude ratio

$$R := \frac{\hat{v}_x(x_0)}{\hat{v}_x(-x_0)} = \frac{\rho_1 m_2}{\rho_2 m_1} \frac{(k^2 v_A^2 - \omega^2) m_1 \frac{\rho_0}{\rho_1} - \omega^2 m_0 \tanh(m_0 x_0)}{(k^2 v_A^2 - \omega^2) m_2 \frac{\rho_0}{\rho_2} - \omega^2 m_0 \tanh(m_0 x_0)}$$

Solar magneto-seismology:

- Measure: R , ρ_i , c_i , ω , k , x_0 .
- Solve to find: v_A and hence B_0 .

Summary

- **Slab structuring** - kink and sausage modes.
- **Asymmetric** slab DR - does not decouple into sausage and kink eigenmodes.
- **Asymmetric** external plasma - difference in amplitude on each boundary.
- **Density ratio variation** - mode conversion between surface modes.
- **Cross-slab amplitude ratio** - potential use as a diagnostic tool to determine the magnetic field strength.

Future work

- Generalise the asymmetric slab further:
 - Shear flow
 - External magnetic field
 - Gravity
- Apply the **cross-slab amplitude ratio** method to **solar observations**.

