



UNIVERSITÄT
BAYREUTH

Bachelor Thesis

Size convergence of the $E \times B$ staircase pattern in flux tube simulations of ion temperature gradient driven turbulence

all nouns with
upper case letters

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(*) 29.05.2005 - (†) 12.05.2023

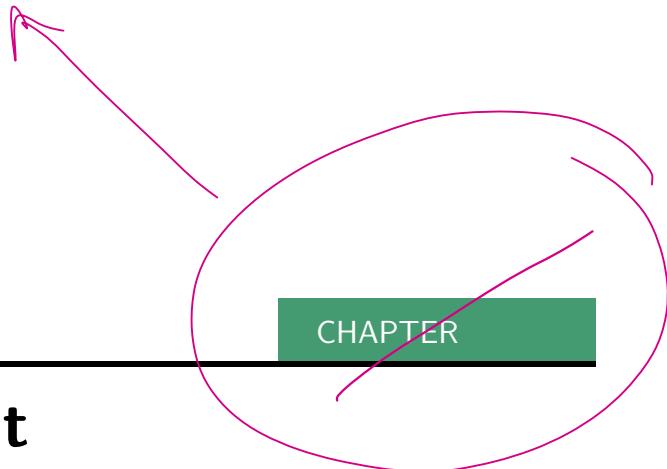
*who almost reached
the age of*

This thesis is dedicated to my cat **Blacky**
~~which get nearly 19 years old~~ ^{*she*} and followed me ~~through life~~
since I was 5 years ~~through my life~~.
You were ^{old.} always a big part of my life and

I will miss you.

Manuel Lippert
May 21, 2023

Use different format
for first parts (only black line)



Acknowledgement

either use:
- express ... for (noun) → e.g. express gratitude for my supervisor's guidance
- extend ... to (person) → e.g. extend gratitude to my supervisor for his guidance

First, I would like to express my greatest gratitude to my supervisor Arthur Peeters and Dr. Florian Rath. I would like to thank both of you for your guidance and support during the submission process of the briefcommunication and the helping hands during the progress writing process of this thesis and the restart script. ↗ two words
development of the

I would like to thank Bernhard Winkler and Markus Hilt for the great technical support regarding my questions to the btrzx1-cluster which helps me to develop the restart script as it is now.

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of

I thank Anna-Maria Pleyer and my sister Cornelia Lippert for proofreading this thesis and briefcommunication, which helps me to get good at English spelling.

Outside academia, I would like to express my gratitude to my parents, my brother, his wife and daughter, my sister and my girlfriend for their encouragement, endurance and support during the time of this thesis. I know sometimes I am a workaholic and I need s to be reminded about the life outside my academic work.

Abstract

Ion temperature gradient driven turbulence (ITG) close to marginal stability exhibits zonal flow pattern formation on mesoscales, so-called $E \times B$ staircase structures. Such pattern formation has been observed in local gradient-driven flux-tube simulations as well as global gradient-driven and global flux-driven studies.

To reduce the computational effort for the simulations lower input parameter of GKW (Gyro Kinetic Workshop) were tested to find the ~~optimal~~ minimum resolution for the ~~optimum of~~

For convenience, a python script (`slurm_monitor.py`) was written which monitors the simulation on the btrzx1-cluster and starts/restarts until the completion criteria is fulfilled. *no findings have been presented to this point, so using "furthermore" carries the wrong implication.
use instead: Simulations show... / Findings show... / or just "It is shown..."*

to
are
plural

Furthermore, it is shown by multiple box size convergence scans that a mesoscale pattern size of $\sim 57 - 76 \rho$ is inherent to ITG driven turbulence with Cyclone Base Case parameters in the local limit. This outcome also implies that a typical scale for avalanche-like transport is inherent to ITG driven turbulence.

Zusammenfassung



*Neues Wort
für Hangman :)*

Ionen-Temperaturgradienten getriebene Turbulenzen (ITG) weisen nahe niedriger Stabilität Zonal Flow Strukturbildungen, sogenannte $E \times B$ Treppenstrukturen, auf Mesoskalen auf. Solche Strukturbildungen wurden sowohl in lokal gradientengetriebene Flusschlauchsimulationen als auch in global gradientengetriebenen und global flussgetriebenen Untersuchungen entdeckt.

Um den Aufwand der Berechnungen der Simulationen zu reduzieren wurden mehrere kleinere Inputparameter für GKW (Gyro Kinetic Workshop) getestet um die optimal kleinste Auflösung für die ausgeführten Simulationen zu finden.

Der halber
Zur Einfachheit wurde ein python-Skript (`slurm_monitor.py`) geschrieben, was Simulationen auf den `btrzx1`-Cluster überwacht und gegebenenfalls startet/neustartet bis das Kriterium zur Vollendung erfüllt ist.

→ Abstract: plural → für einen entscheiden

Weiterhin wurde doch mehrere radiale konvergierende Boxgrößen-Scans gezeigt, dass eine Mesoskalengröße von $\sim 57 - 76 \rho$ inhärent zur ITG getriebenen Turbulenz mit Cyclone Base Parameter in lokalen Limit ist. Dieses Ergebnis impliziert auch, dass eine typische Skala für den lawinenenartig Transport inhärent für die ITG getriebene Turbulenz ist.

Declaration

The author declares that the materials ~~presented~~ ^{are} in this thesis ~~is~~ his own work, unless explicitly stated otherwise. This thesis is based on

LIPPERT, M., RATH, F. & PEETERS, A. G. 2023 Size convergence of the E×B staircase pattern in flux tube simulations of ion temperature gradient driven turbulence. *Phys. of Plasmas* **7** (3), 969–983

and is a further iteration of this publication (brief communication). It provides additional plots and paragraphs that were not ~~necessary to be~~ included in the ~~final~~ publication. The brief communication can be found in Appendix **6.5**.

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CHAPTER

1

Motivation

Ion temperature gradient driven turbulence close to marginal stability exhibits zonal flow pattern formation on mesoscales, so-called $E \times B$ staircase structures^[16]. Such pattern formation has been observed in local gradient-driven flux-tube simulations^{[42][55][45]} as well as global gradient-driven^{[37][51][49]} and global flux-driven^{[16][17][54][28][29]} studies. In global studies, spanning a larger fraction of the minor radius, multiple radial repetitions of staircase structures are usually observed, with a typical pattern size of several ten Larmor radii. By contrast, in the aforementioned local studies the radial size of $E \times B$ staircase structures is always found to converge to the radial box size of the flux tube domain. The above observations lead to the question:

*'always' is a very absolute word for scientific research
↑ "significantly" or something like it would be better
Does the basic pattern size always converges to the box size, or is there a typical mesoscale size inherent to staircase structures also in a local flux-tube description?
too well*

The latter case would imply that it is not necessarily global physics, i.e., profile effects, that set

- (i) the radial size of the $E \times B$ staircase pattern
- (ii) the scale of avalanche-like transport events.

These transport events are usually restricted to $E \times B$ staircase structures and considered as a nonlocal transport mechanism^[16].

In this bachelor thesis the above question is addressed through a box size convergence scan of the same cases close to the nonlinear threshold for turbulence generation as studied in Ref. [42]

*→ case ≠ Fall
use instead:
• experiment iterations
• same conditions*

CHAPTER

Plasma Physics, Zonal Flows and Gyrokinetic Theory

2

2.1 Charged Particle Motion in Magnetic and Electric Field

In magnetic confinement devices like the tokamak reactor, the charged particles experience magnetic and electric fields which results in distinct motion under the associated force. Charged particles can be separated in species, e.g. electrons and ions, which will be later on not displayed in the governing equation. Throughout this thesis the charge q , the mass m or the temperature T indicate the quantities of a specific species.

what do they experience?
A force? A pull?

2.1.1 Particle Motion perpendicular to the magnetic field

Due to the Lorentz force, particles with a velocity component perpendicular to the homogenous magnetic field v_{\perp} undergo a circular motion in the plane perpendicular to the magnetic field. This type^{of} motion has circular frequency, which is often referred^{to} as *cyclotron frequency* and is defined as

$$\omega_c = \frac{|q|B}{m} , \quad (2.1)$$

where m and q are the mass and the charge of the particle and B the strength of the magnetic field. The radius, the so called *Larmor radius*, of this motion is given by

$$r_L = \frac{mv_{\perp}}{|q|B} \quad (2.2)$$

with the center being to which center is often referred as *gyrocenter*. Note that since the Lorentz force depends on the species of the particle, the circulation direction is opposed between electron and ions.

different (?)

Due to the Coulomb collisions the plasma gets thermalized. Together with the Maxwell-Boltzmann distribution the typical thermal velocity is

$$v_{th} = \sqrt{\frac{2T}{m}} , \quad (2.3)$$

where T represents the species temperature. Based on the thermal velocity v_{th} the *thermal Larmor radius* gets introduced as

$$r = \frac{mv_{th}}{|q|B} . \quad (2.4)$$

2.1.2 Particle Motion parallel to the magnetic field

In absence of forces in the parallel direction of the magnetic field the particles can stream flow
(freely)parallel to the homogenous magnetic field. The velocity of this motion is connected
 to the thermal velocity v_{th} and is dominated bei electrons due to their lighter mass
 compared to ions ($v_{\text{th,e}}/v_{\text{th,i}} = 60$). XD

*not so
"freely" if
they are parallel
:P*

When an electric field with a component parallel to the magnetic field E_{\parallel} acts on plasma influences the plasma
 the charged particles gets accelerated by the electric force

↑ are

$$F_{\parallel,E} = qE_{\parallel} . \quad (2.5)$$

The parallel motion follows then from the equation of motion. Here the direction of the motion also depends on the species type.

Since magnetic fields are not always homogenous, an inhomogeneous magnetic field with its gradient ∇B containing a component parallel to the magnetic field which is given by

$$\nabla_{\parallel}B = \frac{\mathbf{B}}{B} \cdot \nabla B \quad (2.6)$$

~~and causes the force~~

$$F_{\parallel,\nabla_{\parallel}B} = -\frac{mv_{\perp}^2}{2B} \nabla_{\parallel}B = \mu \nabla_{\parallel}B ; \quad \mu = \frac{mv_{\perp}^2}{2B} \quad (2.7)$$

with *magnetic moment* μ .] The magnetic moment μ can be an adiabatic invariant (constant of motion) if the variation of the magnetic field over time is smaller than the inverse of the cyclotron frequency ω_c^{-1} and the spatial variation is larger the Larmor radius ρ_L . The *resulting* caused force has its application in the mirror effect where a charged particle gets reflected due to this force. [56]

2.1.3 Drifts in the Gyrocenter

In the presence of a magnetic field (homogenous, inhomogeneous or perturbed) and electric fields the gyrocenter undergoes slow (compared to the thermal velocity v_{th}) drift motions perpendicular to the magnetic field. There are several examples for this drift motion. According to this thesis topic only the main three drift types will be covered in the following.

1. $\mathbf{E} \times \mathbf{B}$ Drift:

If an electric field \mathbf{E} with a perpendicular component together with the magnetic field \mathbf{B} (both fields are homogenous) is present the acting Coulomb force and Lorentz force results to a drift of the gyrocenter with

$$\mathbf{v}_E = \frac{\mathbf{E} \times \mathbf{B}}{B^2} \quad (2.8)$$

Why not bold?

which is called the $\mathbf{E} \times \mathbf{B}$ drift. Since both acting forces direction depends on the species type the direction of the $\mathbf{E} \times \mathbf{B}$ drift is for every species the same.

2. ∇B Drift:

Inhomogeneous magnetic field causes a gradient ∇B of the magnetic field. Because of that gradient the gyrocenter undergoes a ∇B drift defined by

$$\mathbf{v}_{\nabla B} = \frac{mv_{\perp}^2}{2q} \frac{\mathbf{B} \times \nabla B}{B^3} . \quad (2.9)$$

The ∇B drift varies thereby on scales which are larger compared to the Larmor radius. The direction of the ∇B drift depends on the species type.

3. Curvature Drift:

Due to centrifugal force acting on the particle in a curved magnetic field the gyrocenter experiences a curvature drift given by *according to field*

$$\mathbf{v}_{\kappa} = \frac{mv_{\parallel}^2}{q} \frac{\mathbf{B} \times \boldsymbol{\kappa}}{B^2} = \frac{mv_{\parallel}^2}{q} \frac{\mathbf{B} \times \nabla B}{B^3} ; \quad \boldsymbol{\kappa} = -(\mathbf{b} \cdot \nabla) \mathbf{b} = \frac{\nabla B}{B} , \quad (2.10)$$

where \mathbf{b} is the unit vector along the magnetic field. To obtain the result for the curvature $\boldsymbol{\kappa}$ in Eq. (2.10) the plasma pressure has to be small in comparison to the magnetic field strength B . In the form of Eq. (2.10) ∇B and curvature drift can be treated similarly.⁵⁶

2.2 Magnetic Confinement in Tokamak

In tokamak devices strong magnetic fields confine the hot plasma. As mentioned in Chapter 2.1 a magnetic field forces a perpendicular particle motion and a motion which contains the gyro motion and slow perpendicular gyro center drifts. Because of the much smaller size of the Larmor radius compared to the device size R the particle and energy losses are caused by the gyro center drift. To avoid this at the end of the field lines the magnetic field in the tokamak devices is shaped like a torus. This type of geometry has nested surfaces with constant magnetic flux, so-called *flux surfaces*, and ~~with~~ magnetic field lines which appear ~~ing~~ on ~~that~~ surfaces. ~~For stability, reasons~~ the magnetic field has a toroidal and a poloidal component which also holds equivalence of the plasma pressure.^{[50] [56]} The toroidal component is produced by external coils whereas the poloidal component is provided by the toroidal plasma current. Together the components result in a magnetic field which follows helical trajectories [Fig 2.1]. To characterize the quality of confinement the so-called *plasma beta* is used and is given by as

$$\beta = \frac{nT}{\mu_0 B^2/2} .$$

(2.11)

↗ as all forces strive for stability e.g. lowest energy level, this sentence should be restricted: as a way to maintain stability...

with n the plasma density, T as temperature, μ_0 the permeability in vacuum and the magnetic field strength B . Respectively, the plasma beta compares the thermal plasma pressure nT to the ambient magnetic field pressure $\mu_0 B^2/2$ which indicates when the plasma beta is smaller the confinement is better. In a tokamak reactor the plasma beta has a typical order of few percent.^[56]

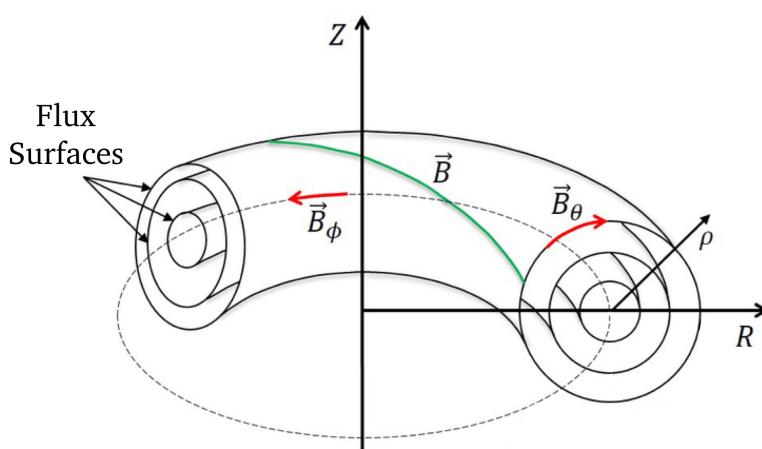


Figure 2.1: Toroidal flux surfaces in tokamak plasma with helical magnetic field (green line) in torus coordinates (ρ (radial), ϕ (toroidal), θ (poloidal)) or cylinder coordinates (Z , R , ϕ).^[3]

↗ might be "cylindrical coordinates" but I'm not sure

2.3 Ion Temperature Gradient (ITG) Driven Instability

The described electrostatic forces acting on the tokamak plasma causes a variety of collective phenomena.^{*The occurrence of*} Microinstabilities is one of these phenomena which are instabilities on the spatial scale of the Larmor radius ρ_L and frequencies much smaller than the cyclotron frequency ω_c .^{*They*} These instabilities are driven by density and temperature gradients that are present in the tokamak devices. Turbulence caused from microinstabilities are considered as the main reason for the loss of particles/energy and momentum from tokamak plasma^{*6[20]26*} which limits the confinement time of tokamak reactors. The most dominant instability in tokamak plasma is the ion temperature gradient (ITG) driven instability.^{*11[12]46*}

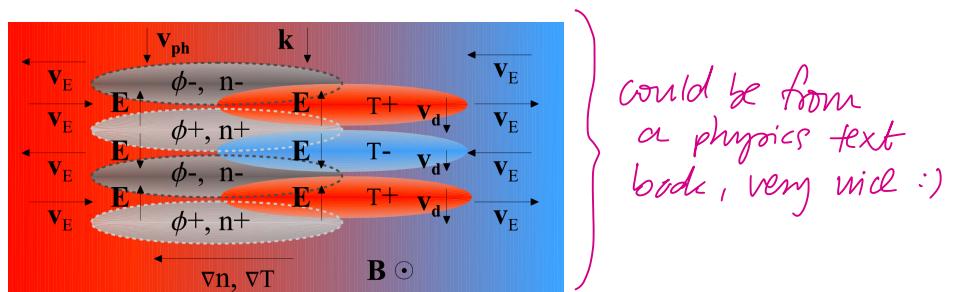


Figure 2.2: Sketch of ion temperature gradient driven instability on the outboard low field side of the tokamak^[10]

The ITG driven instability is present when a sufficiently large ion temperature gradient occurs and is driven by the so-called bad curvature of the outboard low field side of the tokamak. The mechanism will be briefly explained with the help of Fig. 2.2 and Ref. [41013]. An ion temperature perturbation ($T^\pm \gtrless 0$ blue and red cells) on the outboard low field side of the tokamak causes the formation of curvature drift and ∇B drift with velocity \mathbf{v}_d , which compresses the plasma ($n_\pm \gtrless 0$, gray cells). This relates to the perturbation of the potential ($\phi^\pm \gtrless 0$) in the adiabatic electron response. As stated in the equation for the electric field ($\mathbf{E} = -\nabla\phi$) the perturbation of the potential ϕ causes an electric field \mathbf{E} which results in a $\mathbf{E} \times \mathbf{B}$ drift with velocity \mathbf{v}_E . The $\mathbf{E} \times \mathbf{B}$ drift direction transports hot plasma in the outer hot region of perturbation ($n^+ > 0$, red cell) and cooler plasma into inner cold region perturbation ($n^- < 0$, blue cell) which reinforce the process and instability grows.

The described mechanism causes perturbation on the outboard low field side of the tokamak this region has a “bad curvature” and instability happens there. On the inboard high field side the opposing ∇B and ∇T results in a decay of perturbation which is referred to as “good curvature”. Furthermore, such turbulent structures are called *eddies* which have a typical radial scale of several Larmor radius.⁵⁹ Through the explanation above is clear that the initial perturbation is growing when the normalized temperature *if*

gradient overcomes a critical value $R/L_{T,c}$. For nonlinear turbulence the critical value is given by $R/L_{T,c} \sim 6$.^{[18][25]}

2.4 Zonal Flows and Shearing Rate $\omega_{E \times B}$

Zonal flows are linked in the tokamak plasma to the $E \times B$ flows [Eq.(2.8)] that is connected to the on flux-surfaces constant electrostatic potential, also called zonal potential.^[14] Because of the variation only across flux-surfaces, the zonal flow is tangential to those surfaces and do not contribute to the turbulent transport across the flux-surfaces caused by the $E \times B$ motion. Zonal flows are driven nonlinearly through turbulent Reynolds stresses.^[15] For more details about the mechanism behind zonal flow generation the reader is referred to Ref. [6][19][24][31] → *an elegant way to say "I cannot be bothered to explain this to you noobs in detail." XD*

So there
is a point
where they
influence each
other?

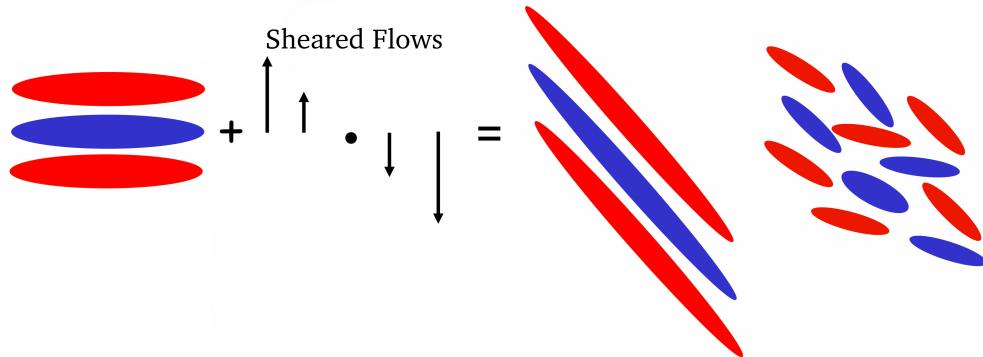


Figure 2.3: Subdued turbulence because *if* stretched eddies, which breaks in smaller structures caused by sheared $E \times B$ drifts^[23]

Furthermore, zonal flows are connected with the radially alternating zonal potential, which introduces a velocity shear, i.e., a radial variation of the zonal flow velocity. This causes a shear deformation of turbulent structures which is considered to play a significant role in the suppression of turbulence^{[5][18]}, especially the ion temperature gradient driven turbulence^{[40][36][35][57][58]}. The nonlinear generation mechanism causes a shift, also known as Dimits shift^[18], of the critical temperature gradient $R/L_{T,c}$ which leads to a later development of ITG driven turbulence. The shearing process tilted and stretches turbulent eddies, caused by the shear flow, until they break into smaller eddies with reduced radial correlation length, [Fig. 2.3 which characterizes] the shearing as a decorrelation process. Turbulence will then be reduced because of this process.^{[5][14][7]}

You often use "which+Vel_ointuitive". like "which characterize".
Alternatively you could use "characterizing" for more flavour :)
→ use at your own convenience

When the zonal flows are strong enough to suppress turbulence completely due to the shearing process, pattern formation occurs in the stabilized tokamak plasma, the so-called $E \times B$ staircase structure. The $E \times B$ staircase pattern has the following properties:

- (1) They occur on a *radial mesoscale* of order of 10ρ which is larger than the Larmor radius but smaller than the machine size R .
- (2) The structures are *quasi stationary* in space and *vary* on time scales much longer than the typical turbulent *time scales*.
- (3) They have an *typical amplitude* in terms of the $E \times B$ shearing rate of the order of $10^{-1}v_{th}/R$.
- (4) Turbulent transport occurs in *avalanches* with propagation strongly linked to the local $E \times B$ shearing rate.^[37]

The $E \times B$ staircase pattern is manifested as radial structure formation in the $E \times B$ shearing rate defined by^{[44][43][42]}

$$\omega_{E \times B} = \frac{1}{2} \frac{\partial^2 \langle \phi \rangle}{\partial x^2}, \quad (2.12)$$

where $\langle \phi \rangle$ is the zonal electrostatic potential normalized by $\rho_* T/e$ ($\rho_* = \rho/R$ is the thermal Larmor radius normalized with the major radius R , T is the temperature, e is the elementary charge). The zonal potential is calculated from the electrostatic potential ϕ on the two-dimensional x - y -plane at the low field side according to^[45]

$$\langle \phi \rangle = \frac{1}{L_y} \int_0^{L_y} dy \phi(x, y, s=0). \quad (2.13)$$

The $E \times B$ shearing rate $\omega_{E \times B}$ is the radial derivative of the advecting zonal flow velocity^{[21][52]} and quantifies the zonal flow induced shearing of turbulent structures^{[5][21][7]}.

As mentioned before the turbulent transport occurs through avalanches which also mediate the turbulence of the ITG driven turbulence in the adabatic electron approximation. Avalanches are *initiated* when the $E \times B$ sharing rate has a zero crossing with a steep flank and propagates outwards with negative shearing and inwards *with* for positive shearing.^{[27][37]} The fully developed avalanche results in almost completely quenched turbulence. Its development results in a discontinuous step in the dependency of turbulent fluxes on the temperature gradient als known as *finite heat flux threshold*.^{[42][55]}

2.5 Gyrokinetic Theory

2.5.1 Vlasov Equation

Because of the large number of particles in the fusion plasma a prediction only using Newton-Maxwell dynamics results in an impossible task for computers, but this problem can be solved with a statistical approach. For that the particle density distribution function $f(\mathbf{x}, \mathbf{v}, t)$ in the six dimensional phase space $\{\mathbf{x}, \mathbf{v}\}$ with the particles position \mathbf{x} and velocity \mathbf{v} is needed. Because collisions are happening at much smaller frequencies than the characteristic frequencies connected to turbulence, the collisionless model is often considered^[20] which gives through evolution of the particle density distribution function in the *Vlasov equation*

$$\frac{\partial f}{\partial t} + \frac{\partial f}{\partial \mathbf{x}} \cdot \frac{d\mathbf{x}}{dt} + \frac{\partial f}{\partial \mathbf{v}} \cdot \frac{d\mathbf{v}}{dt} = 0 . \quad (2.14)$$

To obtain a closed system the Maxwell equations with the particle density n and current density j can be described by the distribution function as follows

$$n = \int d\mathbf{v} f(\mathbf{x}, \mathbf{v}, t) \quad j = q \int d\mathbf{v} v f(\mathbf{x}, \mathbf{v}, t) , \quad (2.15)$$

which are then substituted into the Maxwell equations

$$\begin{aligned} \nabla \cdot \mathbf{B} &= 0 & \nabla \times \mathbf{B} &= \mu_0 \left(\sum_{\text{species}} j + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right) \\ \nabla \cdot \mathbf{E} &= \frac{1}{\epsilon_0} \sum_{\text{species}} qn & \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} . \end{aligned} \quad (2.16)$$

The Vlasov equation (2.14) together with the Maxwell equations (2.15) and (2.16) is the basis of the gyrokinetic model.^[30]

(*) maybe phrase it like: "simulation-based calculation", as there might be a PC out there strong enough)

⊗ a little round-about,
 instead use: "differences between the turbulence ... gyro motion
 led to the need to define spatio-temporal scales in
 categories of gyrokinetic ordering"

2.5.2 Gyrokinetic Ordering

Dynamics in tokamak plasma is described by spatio-temporal scales which is also referred to as *gyrokinetic ordering*. This separation is a comparison between the turbulence frequency ω and the cyclotron frequency ω_c ($\omega/\omega_c \sim 10^{-3}$) due to strong external magnetic fields which causes a fast gyro motion. The length scales of the turbulence are associated with the wavevector \mathbf{k} which can be separated into a perpendicular component $k_\perp = |\mathbf{k} \times \mathbf{b}|$ and a parallel component $k_\parallel = |\mathbf{k} \cdot \mathbf{b}|$ where \mathbf{b} is parallel to the poloidal component of the magnetic field. The perpendicular component k_\perp scales with the thermal Larmor radius $k_\perp^{-1} \sim \rho$ which is significantly smaller than the equilibrium density n_0 , what can be expressed with the gradient length $L_n = |\nabla \ln n_0|^{-1}$. The gradient length compares with the machine size R which leads to the normalized Larmor radius $\rho_* = \rho/R \sim 10^{-3} - 10^{-4}$ with R as major radius of the tokamak. The parallel component k_\parallel on the other hand scales with the machine size R . Experiments also show that in the core plasma the fluctuation amplitude of the density perturbation $\delta n/n_0$ and magnetic field fluctuation $\delta B/B_0$, with B_0 as strength of the magnetic field in equilibrium, is of order $\sim 10^{-4}$. Together this separation result in the following result

$$\frac{\omega}{\omega_c} \sim \frac{k_\parallel}{k_\perp} \sim \frac{\rho}{L_n} \sim \frac{\delta n}{n_0} \sim \frac{\delta B}{B_0} \sim \frac{v_d}{v_{th}} \sim \epsilon_g \quad (2.17)$$

with $\epsilon_g \ll 1$ which applies to the typical dynamics of a fusion core plasma.[6][20] The gyrokinetic ordering allows the reduction of parameters in the Vlasov equation (2.14).

2.5.3 Gyrokinetic Equation

For the description *for* charged particle *in* the tokamak device the *gyrocenter coordinates* *are* will be used [Fig. 2.4]. In this set of coordinates the gyrocenter follows the magnetic field with the parallel velocity v_{\parallel} . The gyro motion *will be* described together with the magnetic moment $\mu = \frac{1}{2}mv_{\perp}/\omega_c$, the gyrocenter \mathbf{X} and the gyro phase ζ which gives a parameter set of six quantities $\{\mathbf{X}, v_{\parallel}, \mu, \zeta\}$.

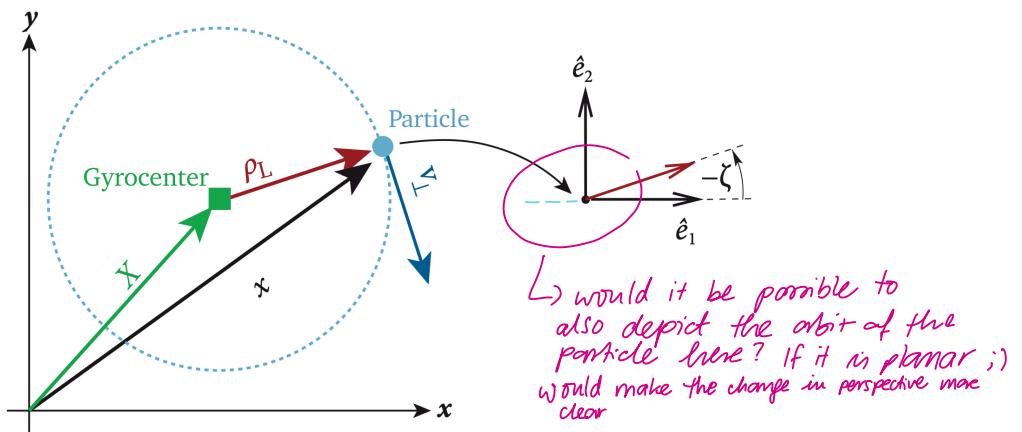


Figure 2.4: Sketch of gyrocenter coordinates where the charged particle performs a circular motion around the gyrocenter.^[31]

In the next step the coordinates of the Vlasov equation (2.14) *will be* transformed into the gyrocenter coordinates. For that the particle position \mathbf{x} *will be* described through the gyrocenter position $\mathbf{X} = \mathbf{x} - \rho_L$ with the position Larmor radius vector ρ_L . The particle velocity \mathbf{v} *gets* decomposed into a parallel component $v_{\parallel} = \mathbf{v} \cdot \mathbf{b}$ and a perpendicular component $v_{\perp} = |\mathbf{v} \times \mathbf{b}|$ which relates to the magnetic moment μ . Through the gyrokinetic ordering a phase space transformation can be performed which is based on the Lie perturbation theory, the Hamiltonian or Lagrangian formalism and an asymptotic expansion in the smallness parameter ϵ_g . In the lowest order is the magnetic moment *um (?)*

as an exact invariant which makes the gyro phase ζ an ignorable variable.^[2089] Taking everything into account the new coordinates of the Vlasov equation are $\{\mathbf{X}, v_{\parallel}, \mu\}$ which transforms the Vlasov equation into the *gyrokinetic equation*

$$\frac{\partial f}{\partial t} + \frac{\partial f}{\partial \mathbf{X}} \cdot \frac{d\mathbf{X}}{dt} + \frac{\partial f}{\partial v_{\parallel}} \cdot \frac{dv_{\parallel}}{dt} = 0 . \quad (2.18)$$

The time derivative of the magnetic moment μ is zero because the magnetic moment is an exact invariant.

2.5.4 δf Approximation and Local Limit

All simulations in this thesis are performed with GKW (Gyro Kinetic Workshop) which uses in application of the gyrokinetic equation the δf approximation and local limit approximation. For the δf approximation the particle density distribution function f gets separated into perturbation δf and equilibrium f_0 (constant in time) as expressed in the gyrokinetic ordering (2.17). This is given by

$$f = f_0 + \delta f \quad \delta f \sim \rho_* f_0 . \quad (2.19)$$

The distribution function in the equilibrium f_0 is set in GKW to the Maxwell distribution function. The perturbed distribution function δf is considered to vary perpendicular to the poloidal magnetic field in order of the Larmor radius while the perturbation changes with the system size which results in the ordering

$$\nabla_{\perp}(\delta f) = \nabla_{\perp}(f_0) . \quad (2.20)$$

The δf approximation applies for global and local description where the latter case is used in this thesis. The *local limit* on the other hand uses the spatial separation of equilibrium and perturbation and implies that all equilibrium and geometry quantities can be taken constant over the radial extent of the simulation volume. Because of the tokamak symmetry all equilibrium constants are also constant in the toroidal direction where the dependency of the equilibrium magnetic field on the poloidal angle is restrained in general.

In addition to this, in GKW the following coordinates get applied with the parameter $\mathbf{X} = \{x, y, s\}$. Here, x is the radial coordinate that labels the flux surfaces normalized by the thermal Larmor radius ρ , y labels the field lines and is an approximate binormal coordinate. Together with the coordinate s , which parameterizes the length along the field lines and is referred to as the parallel coordinate, these quantities form the Hamada coordinates²². This results in the new coordinates $\{x, y, s, v_{\parallel}, \mu\}$. Combining the δf approximation and the Hamada coordinates the gyrokinetic equation can be formulated as follows

$$\frac{\partial g}{\partial t} + \mathbf{v}_x \cdot \nabla g + (v_{\parallel} \mathbf{b} + \mathbf{v}_D) \cdot \nabla(\delta f) - \frac{\mu B}{m} \frac{\mathbf{B} \cdot \nabla B}{B^2} \frac{\partial(\delta f)}{\partial v_{\parallel}} = S , \quad (2.21)$$

where g is a function which contains the perturbation δf and the equilibrium f_0 . Furthermore, \mathbf{v}_D and \mathbf{v}_x are drift velocities which are caused by the inhomogeneous magnetic field and the perturbed magnetic field (\mathbf{v}_D) and through the $E \times B$ drift (\mathbf{v}_x). The source term S combines from the distribution function in the equilibrium f_0 and a correction term for the collisions as well as the energy injection term. Due to periodic boundary conditions in the radial direction the radially gradients cannot change in time which is also referred to as gradient driven approach.²¹

in time for what?

Oh no, the boring part...

CHAPTER

Methods and Material

3

3.1 Simulation Setup

The gyrokinetic simulations are performed with the non-linear flux tube version of Gyrokinetic Workshop (GKW)^[41] with adiabatic electron approximation. In agreement with Ref. [42] Cyclone Base Case (CBC) like parameters are chosen with an inverse background temperature gradient length $R/L_T = 6.0$ and circular concentric flux surfaces. The numerical resolution is compliant to the "Standard resolution with 6th order (S6)" set-up of the aforementioned reference. A summary of the numerical parameters is given in Tab. [4.1] and for more details about the definition of individual quantities the reader is referred to Refs. [41][42].

	N_m	N_x	N_s	$N_{\nu_{ }}$	N_μ	D	ν_d	$D_{\nu_{ }}$	D_x	D_y	Order	$k_y\rho$	$k_x\rho$
S6	21	83	16	64	9	1	$ \nu_{ } $	0.2	0.1	0.1	6	1.4	2.1

Table 3.1: Resolution used in this paper: Number of toroidal modes N_m , number of radial modes N_x , number of grid points along the magnetic field N_s , number of parallel velocity grid points $N_{\nu_{||}}$, number of magnetic moment grid points N_μ , dissipation coefficient used in convection along the magnetic field D , the velocity in the dissipation scheme ν_d , dissipation coefficient used in the trapping term $D_{\nu_{||}}$, damping coefficient of radial modes D_x , damping coefficient of toroidal modes D_y , order of the scheme used for the zonal mode, maximum poloidal wave vector $k_y\rho$, and maximum radial wave vector $k_x\rho$

Consistent with Ref. [42] the turbulence level is quantified by the turbulent heat conduction coefficient χ , which is normalized by $\rho^2 v_{\text{th}}/R$ ($v_{\text{th}} = \sqrt{2T/m}$ is the thermal velocity and m is the mass). Furthermore, quantities ρ , R , T , v_{th} and m are referenced quantities from Ref. [42][41].

3.2 Diagnostics

In order to diagnose the temporal evolution of the staircase pattern and to obtain an estimate of its amplitude the radial Fourier transform of the $E \times B$ shearing rate is considered. It is defined by

applied

$$\omega_{E \times B} = \sum_{k_{ZF}} \hat{\omega}_{E \times B}(k_{ZF}, t) \exp(ik_{ZF}x), \quad (3.1)$$

where $\hat{\omega}_{E \times B}$ is the complex Fourier coefficient and $k_{ZF} = 2\pi n_{ZF}/L_x$ defines the zonal flow wave vector with the zonal flow mode number n_{ZF} ranging in $-(N_x - 1)/2 \leq n_{ZF} \leq (N_x - 1)/2$ and the radial box size L_x . Based on the definitions above, the shear carried by the zonal flow mode with wave vector k_{ZF} is defined by $|\hat{\omega}_{E \times B}|_{n_{ZF}} = 2|\hat{\omega}_{E \times B}(k_{ZF}, t)|$. In general, the zonal flow mode that dominates the $E \times B$ staircase pattern, also referred to as the *basic mode* of the pattern in this work, exhibits the maximum amplitude in the spectrum $|\hat{\omega}_{E \times B}|_{n_{ZF}}$.

3.3 btrzx1 Cluster

Very sexy { The simulation itself will be performed on the btrzx1-cluster of the University Bayreuth. This cluster has a wall time from 24 hours, an total amount of 345 compute nodes with two AMD Epyc Processors (2nd generation) with 16 cores each and 128 GB of main memory each node and as resource manager Slurm.^[48]

The simulations are performed on the /scratch directory because the /home folder has an disk space limit of 80 GB which is not enough to perform long runs of GKW.

3.4 Restart Script for Simulation

As stated in Chapter 3.3 the btrzx1-cluster has a wall time of 24 hours but since simulations typically run longer than 24 hours, one have to restart the simulation after the wall time gets exceeded. Some simulations need even longer so that the simulation itself needs multiple weeks to finish which makes the restart of the simulation a daily task that could be easily forgotten. So to eliminate this minor inconveniences the following restart script were written. *to autsource and simplify the process.* *didn't phrase it like that ;)*

As script language python was chosen because of its easiness to write code and the variety of tools which can interact with the bash shell. bash was also considered but after some issues occurred regarding the mechanism to check whether a file exists or not,

python was selected. Furthermore, the script has no requirements because it runs on the standard python3 library. *and is therefore executable with any python software version.*

The script was named slurm_monitor.py to indicate that this script works with the Slurm resource manager. The source code was pushed to the gkw repository on BitBucket³³ or can be found in Appendix 6.1

3.4.1 How to run Restart Script in Terminal

To run the restart script python3 has to be installed on the system and should be accessible via the command line. The monitoring can be started with the following command:

```
python3 -u slurm_monitor.py --job-name $JOBNAME
```

Additional parser options can be added to the command above which will be summarized in the upcoming Chapter 3.4.2. The script was build to run in the background of the terminal for which there are two approaches:

1. nohup:

Runs command in the background while terminal can still be used or closed.³⁴ The command for this method would be:

```
nohup python3 -u slurm_monitor.py --job-name $JOBNAME &> /dev/null &
```

and to cancel the monitoring:

```
python3 -u slurm_monitor.py --job-name $JOBNAME --kill
```

2. screen:

Opens a virtual terminal session which is detached from the main terminal (like a window) and can be entered and closed.³⁵ Here the command would be:

```
screen -S $SESSION #Create screen session
python3 -u slurm_monitor.py --job-name $JOBNAME --verbose
```

and to cancel the monitoring use the keyboard shortcut **[ctrl]+[C]** or kill screen itself with **[ctrl]+[D]**. To leave the screen use (**[ctrl]+[A]**) + **[D]** and to enter the screen use:

```
screen -r $SESSION
```

3.4.2 Features of Restart Script

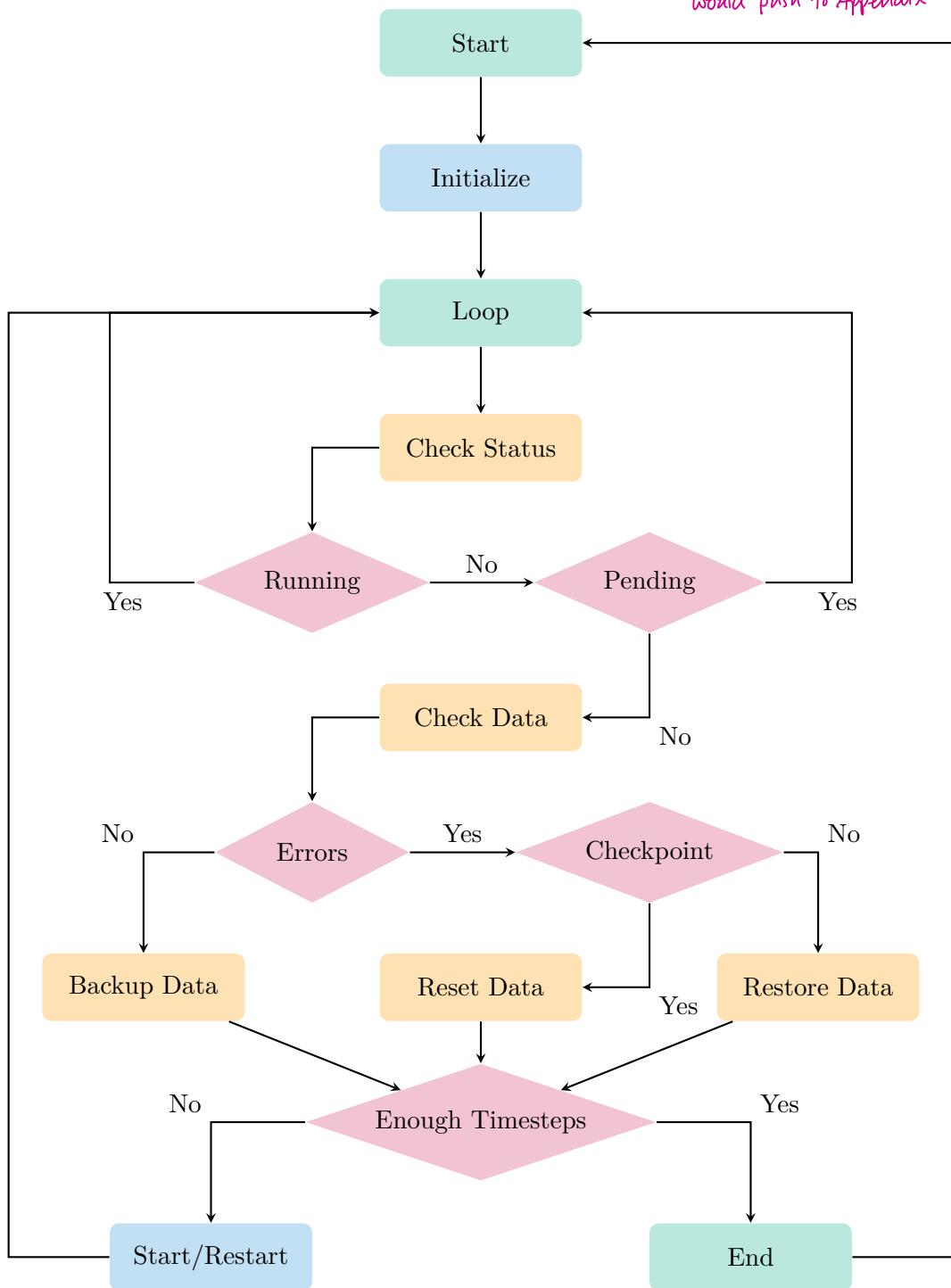
The restart script has several features which can be used via the command line interface as parser options. The following features with the corresponding parser option were implementing:

- **!Required!** Set job name not longer than 8 characters to identify simulation
-j [JOBNAME], --job-name [JOBNAME]
- Creates jobscrip file with defined content
(look into variable `jobscripContent` to add more option)
--jobscrip [JOBSRIPTFILE] (default=jobscrip-create)
--nodes [NODES] (default=3)
--ntask-per-node [TASKS] (default=32)
--walltime [WALLTIME] (d-hh:mm:ss) (default=0-24:00:00)
- Start/Restarts simulation until job criteria is suffused
-n [TIMESTEPS] (default=10000)
--restartfile [RESTARTFILE] (default=FDS.dat)
- Makes backup after each run before Restart and Restore files after failed run
(uses `rsync` command line utility)
Quicklinks: local (simulation folder), home (home folder)
-b [LOCATION], --backup [LOCATION] (default=None)
- Reset option after run fails and dump files were written
(rely on `h5py`, `pandas` and `numpy` which get installed by the script itself)
-r, --reset (default=False)
- Creates status file with current status and progress bars and updates it dynamically
--statusfile [STATUSFILE] (default=status.txt)
- Set form of the output table in status file
Options: fancy (round box), universal (crossplatform), None (No frame)
--format [FORMATTABLE] (default=None)
- Sends mail at the beginning, end and restart with status file as attachment
(`mailx` other equivalent has to be installed, look into `send_mail` function for more info)
--mail [MAILADDRESS] (mail@server.de) (default=None)
--restart-mail (default=False)
- Kills monitoring process by using the PID number
--kill (default=False)
- Set sleep interval after which the status of the current simulation should be checked
--refresh-rate [SLEEPSIME] (default=300)
- Print output of status file table to console
-v, --verbose (default=False)

All boolean parameters are switches, i.e., the praser option only changes the value from False to True.

3.4.3 General Structure

This was surely very hard to create, but I don't see how this schematic adds to your scientific accomplishments ...
would push to Appendix



3.4.4 Status File created from Restart Script

As stated in Chapter 3.4.2 the restart script is writing a status file with the current status and progress bars of the simulation. The output is formatted like a table and can be adjusted in its appearance with the --format parser option [Chapter 3.4.2]. The output types which are written in the status file are listed in Tab. 3.2. To each output prints the restart script the date, time and the duration of the monitoring additionally into the table.

Type	Information
STARTING	Start of monitoring process or Start/Restart of simulation
CONTINUE	If the simulation has performed at least one run
CONTROL	Check of the current time steps of the simulation
SUCCESS	Stop of monitoring because the time steps condition has been fulfilled
IMPORT	Modules h5py, pandas and numpy will be loaded
INSTALL	Script installs the modules h5py, pandas and numpy on the user level
CHECK	Required modules to use the reset option are already installed
WAITING	Simulation is pending in the queue of Slurm
RUNNING	Simulation is running on the server
BACKUP	The script performs the backup of the data to the predefined location
RESTORE	The script reloads the data from the backup folder if errors occur
RESET	The simulation gets reset with the use of dump files
ERROR	During the monitoring does occur an error Types: No executable (gkw.x), walltime, timeout, no config, hdf5, string error

Table 3.2: Output types in status file

To continue, the status file has different style options, e.g., fancy, universal and None which are displayed in Fig. 3.1

fancy	+-----+ universal +-----+	----- None
	 +-----+	

Figure 3.1: Styles of output table in status file

dynamically updated

In addition to that, it writes down the progress of simulation to the end of the status file in form of a progress bar which gets updated dynamically. The progress bar contains the current time step of the simulation in contrast to the required time steps and the progress of the run itself in seconds in relation to the wall time. It also contains the

current output from `squeue`, if the simulation is running or pending, else the important information of the simulation. An example progress bar ist shown in Fig. 3.2

```
+-----+
|   JOBID   NAME      USER ST      TIME NODES  NODELIST (REASON)   |
| 970455 3x1.5 bt712347 CG 3:16:01       3  r03n03, r05n[15-16]  |
+-----+
| PROGRESS  [=====>.....] ( 80%) 20000/25000 |
| RUN 13    [==>.....] ( 13%) 11462/86400 |
+-----+
```

Figure 3.2: Progress bar example

displays all important information (e.g. ..) for the user's convenience.
The progress bar has the advantage that someone can easily access all important information at the glance. To boost productivity the following command could be quite handy: was added:

```
cd $DATA
find . -name $STATUSFILENAME -exec tail -8 {} \;
```

which prints all progress bars of every simulation ^{to} in the data folder. If one wants to see the complete status file change `tail -8` to `cat`. An example status file is displayed in the Appendix 6.2

3.4.5 Support for other Resource Managers

The general idea behind the restart script could be adapted for other resource managers. For example in the early stage of the development the script was changed to support the PBS/Torque resource manager which mainly runs on the btrzx2-cluster of the University of Bayreuth⁴⁷ and supports the research of Dominik Müller for his Bachelor Thesis³⁹ with the RHMD-code³⁸. The script was named `torque_monitor.py` and can be found in the GitHub Repository of this thesis³² or in the Appendix 6.3. But note that due to the early adaptation the code does lack many features of `slurm_monitor.py` and was not maintained actively.

CHAPTER

4

Results and Discussion

The performed simulations for this chapter are documented in Appendix 6.4.

4.1 Variation of Computational Resolution

At the beginning of this work the goal was to estimate the minimal resolution needed to run the simulation without fearing numerical dissipation. Numerical dissipation can therefore result to no formation of zonal flow structures, which cause an permanent turbulent state of the simulation. The goal behind this testing is to reduce calculation time and costs of the simulation.

Because of that, the criteria for the best resolution should be:

- (1) Subdued turbulence after **short** time periods
- (2) Stability for **long** time periods

and the following procedure will be applied for verification:

1. Reduce only one number of grid points and look if criterias (1), (2) are satisfied
2. Reduce to known minimum number of grid points to verify result in general.

4.1.1 Benchmark

Starting from the “Standard resolution with 6th order (S6)” [Tab. 3.1] the first step is to reproduce the result of Ref. 42 in Section IV. Note that, because of selected circular concentric geometry the used inverse background temperature gradient length is $R/L_T = 6.0$ instead of $R/L_T = 7.0$ which was used in Section IV of Ref. 42 for $s\text{-}\alpha$ geometry. In Fig. 4.1 the obtained data is simular to the results from Ref. 42 with subdued turbulence after $\sim 3000 R/v_{\text{th}}$.

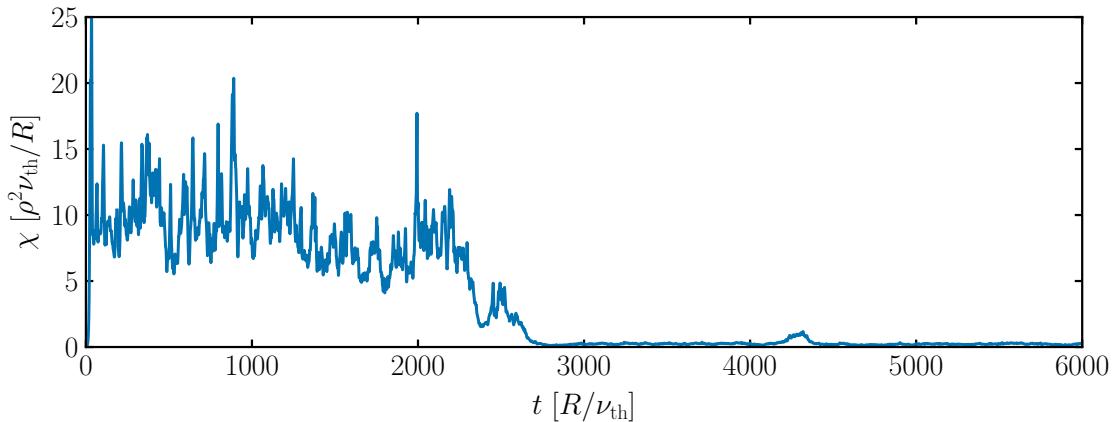


Figure 4.1: Time trace of the heat conduction coefficient χ for $R/L_T = 6.0$ for benchmark

As next step an approach to the finite heat flux threshold was made to verify the selection of the gradient length R/L_T . As in Ref. [42] in Section V concludes, the finite heat flux threshold is approximately located at a gradient length of $R/L_T = 6.3$ for circular geometry [FIG. 4 of Ref. [42]]. Therefore following parameters were used

$$R/L_T \in [6.0, 6.3].$$

As seen in Fig. 4.2 for $R/L_T = 6.3$ no suppression of turbulence occur in the whole time domain, which is in agreement with Ref. [42]

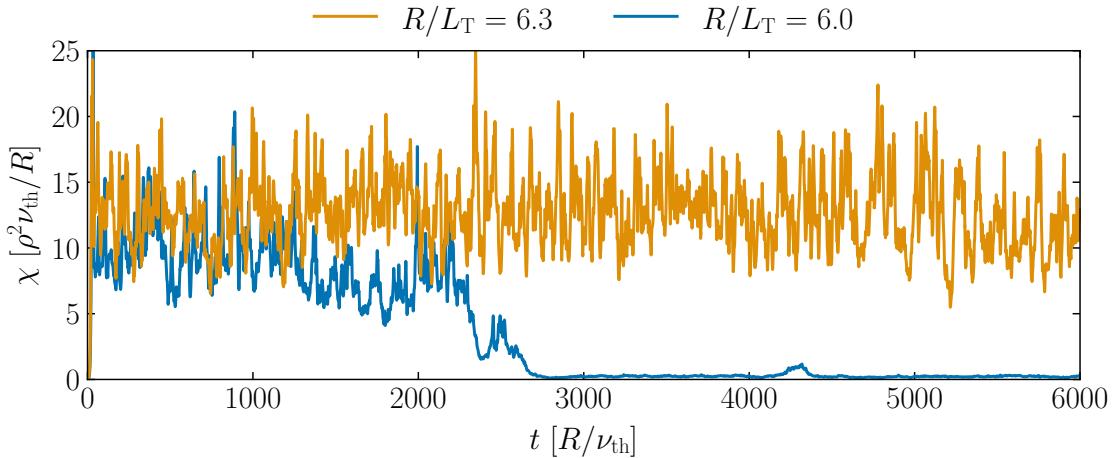


Figure 4.2: Time traces for $R/L_T = 6.0$ and $R/L_T = 6.3$ for benchmark

4.1.2 Reduction of parallel Velocity Grid Points $N_{\nu_{||}}$

In the following the number of grid points for the parallel velocity $N_{\nu_{||}}$ is reduced to

$$N_{\nu_{||}} \in [16, 32, 48, 64].$$

In Fig. 4.3 is clearly visible that the resolution with $N_{\nu_{||}} = 16$ is not suitable for criteria (1) because here the turbulence is not subdued after a long time period. But resolution $N_{\nu_{||}} = 32$ is not acceptable as well according to criteria (2) since the suppressed turbulence state gain instability after $\sim 3000 R/v_{th}$.

So to conclude, only grid points with $N_{\nu_{||}} = 48, 64$ are satisfying the set criteria. Due to criteria (1) the selected resolution will be *is*

$$N_{\nu_{||}} = 48.$$

till too informal

With this number of grid points the time till turbulence subdued is halved compared to the benchmark case, i.e., turbulence suppression occur after $\sim 1500 R/v_{th}$.

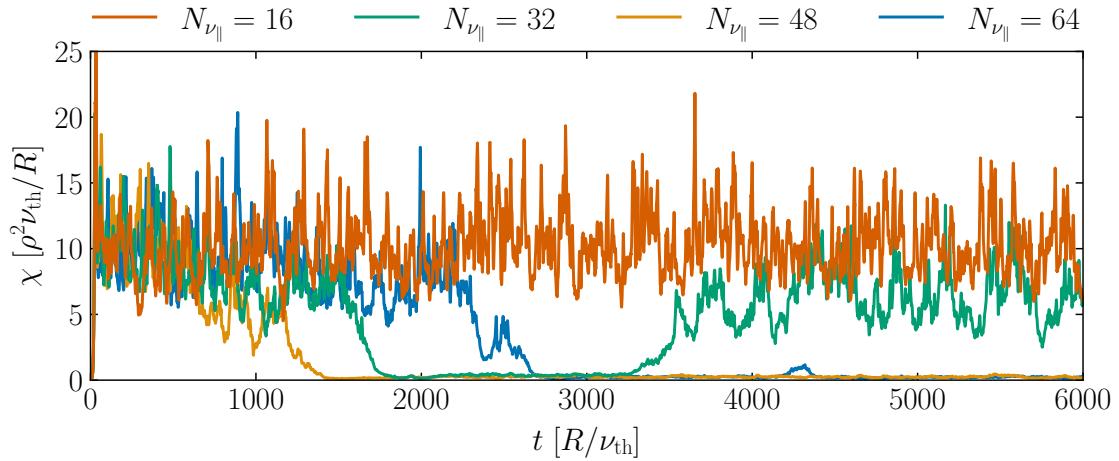


Figure 4.3: Time traces for reduced parallel velocity grid points $N_{\nu_{||}}$

4.1.3 Reduction of Magnetic Moment Grid Points N_{μ}

In the
As next step the number of grid points for the magnetic moment N_{μ} were reduced with

$$N_{\mu} \in [6, 9] .$$

As in Fig. 4.4 shown, the reduction of grid points for the magnetic moment does not significantly impact the suppression of turbulence. The turbulence enters the stationary state in both cases after $\sim 3000 R/v_{th}$.

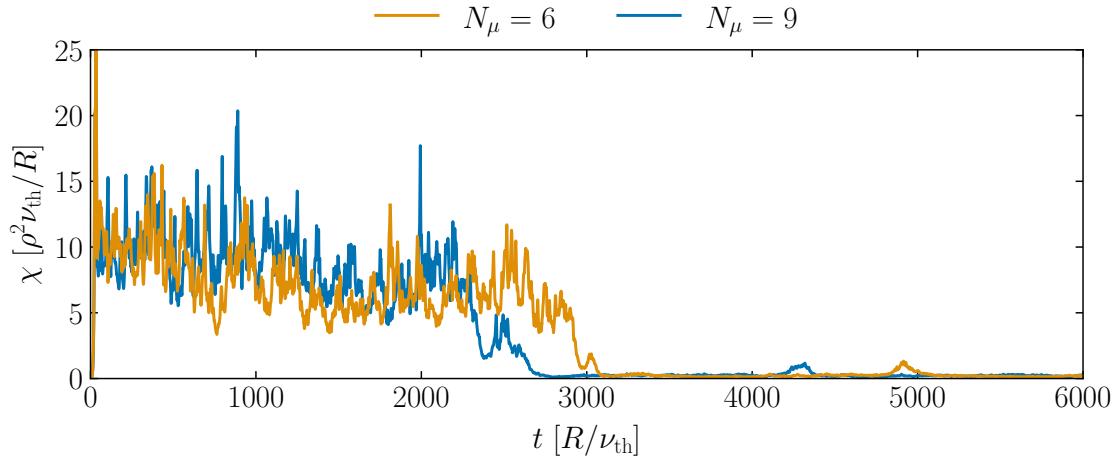


Figure 4.4: Time traces for reduced magnetic moment grid points N_{μ}

To conclude a curial result the number of grids point for the parallel velocity got reduced to $N_{\nu_{\parallel}} = 48$ according to Chapter 4.1.2. In this case the turbulence does not subdue for the resolution $N_{\mu} = 6$ which leads, with the both criteria in mind, to the following resolution

$$N_{\mu} = 9 .$$

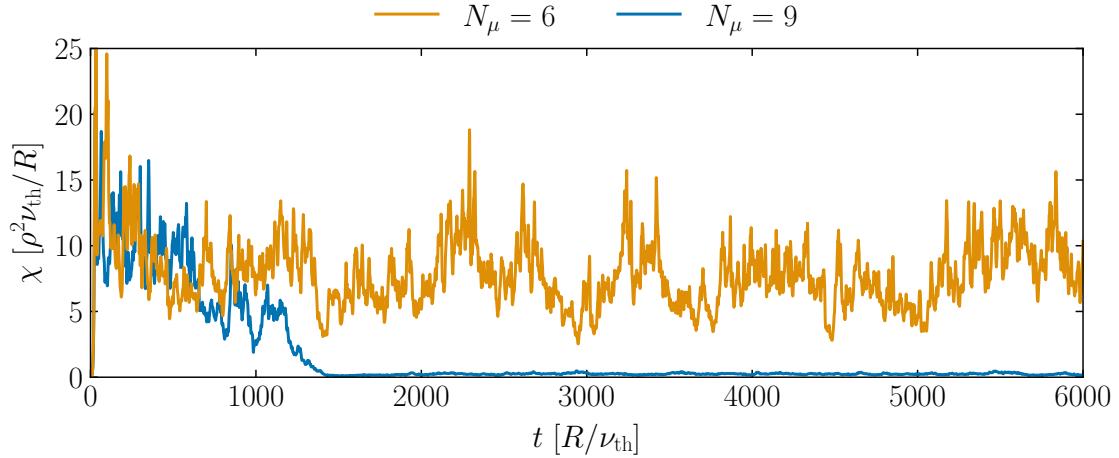


Figure 4.5: Time traces for reduced magnetic moment grid points N_{μ}

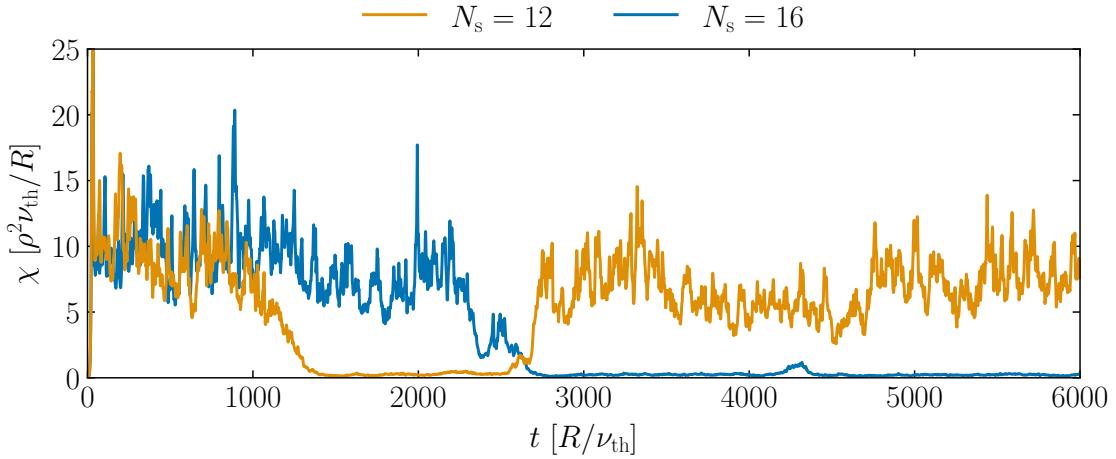
4.1.4 Reduction of Magnetic Field Grid Points N_s

In the final step the number of grid points for the magnetic field N_s gets reduced with the following parameters

$$N_s \in [12, 16] .$$

In Fig. 4.6 can be seen that the reduction to $N_s = 12$ does not satisfy the criteria (2) because the stationary state of the turbulence gets instabil after $\sim 2500 R/v_{th}$. This concludes to the following resolution

$$N_s = 16 .$$


 Figure 4.6: Time traces for reduced magnetic field grid points N_s

4.1.5 Final Resolution for Simulation

Together with the results of Chapter 4.1.2, 4.1.3 and 4.1.4 the final resolution used in the upcoming simulations are displayed in Tab. 4.1, with reduced number of grid points for the parallel velocity $N_{\nu_{||}}$.

	N_m	N_x	N_s	$N_{\nu_{ }}$	N_μ	D	ν_d	$D_{\nu_{ }}$	D_x	D_y	Order	$k_y\rho$	$k_x\rho$
S6	21	83	16	48	9	1	$ \nu_{ } $	0.2	0.1	0.1	6	1.4	2.1

Table 4.1: Resolution used in this paper: Number of toroidal modes N_m , number of radial modes N_x , number of grid points along the magnetic field N_s , number of parallel velocity grid points $N_{\nu_{||}}$, number of magnetic moment grid points N_μ , dissipation coefficient used in convection along the magnetic field D , the velocity in the dissipation scheme ν_d , dissipation coefficient used in the trapping term $D_{\nu_{||}}$, damping coefficient of radial modes D_x , damping coefficient of toroidal modes D_y , order of the scheme used for the zonal mode, maximum poloidal wave vector $k_y\rho$, and maximum radial wave vector $k_x\rho$

4.2 Size Convergence of $E \times B$ Staircase Pattern

In the following the box size ~~is~~^{was} increased relative to the standard box size (L_x, L_y) = (76.27, 89.76) ρ in the radial and binormal direction. The increased box sizes are indicated by the real parameter N_R for radial and N_B for the binormal direction with the nomenclature $N_R \times N_B$ throughout this work. Note that, the number of modes in the respective direction, i.e., N_x and N_m , respectively, ~~is~~^{was} always adapted accordingly to retain a spatial resolution compliant to the standard resolution [Tab. 4.1] and standard box size.

⑩ did you mean "different runs with different $N_R \times N_B$ configurations"?

→ then "iteration" would be better I think

4.2.1 Radial increased Box Size

In the first test the radial box size is increased while the binormal box size ~~is~~^{was} kept fixed to the standard size. The scan covers the realizations ⑩

$$N_R \times N_B \in [1 \times 1, 2 \times 1, 3 \times 1, 4 \times 1].$$

Each realization exhibits an initial quasi-stationary turbulent phase and a second final phase with almost suppressed turbulence [Fig. 4.7(a)]. The latter state is indicative for the presence of a fully developed staircase pattern as depicted in Fig. 4.10. This type of structure is characterized by intervals of almost constant shear with alternating sign satisfying the Waltz criterion $|\omega_{E \times B}| \approx \gamma$ ^{[53][52]} (γ is the growth rate of the most unstable linear ITG driven Eigenmode), connected by steep flanks where $\omega_{E \times B}$ crosses zero. Fig. 4.10(a) shows a striking repetition of the staircase structure, with the number of repetitions equal to N_R . Hence, the basic size of the pattern ~~not~~^{did} only converges with increasing radial box size, the converged radial size turns ~~out to at least~~^{ed} roughly agree ~~with~~^{correspond} with the standard radial box size of Ref. [42].

not clear what this is supposed to mean

You often reference 4.10, therefore you should include this figure earlier, text is hard to follow without the graph visible

⑩ configuration

Due to the lack of a substantial turbulent drive in the final suppressed state, no further zonal flow evolution ~~was~~ observed [Fig. 4.7(b)] and one might critically ask whether the structures shown in Fig. 4.10 represent the real converged pattern in a statistical sense. Note that in the 3×1 case the initial quasi-stationary turbulent state extends up to a few $\sim 10^4 R/v_{th}$. During this period the zonal flow mode with $n_{ZF} = 3$, i.e., the mode that dominates the staircase pattern in final suppressed phase, undergoes a long-term evolution with a typical time scale of several $\sim 10^3 R/v_{th}$.

Hence, several of such cycles are covered by the initial turbulent phase, which is evident from the occurrence of phases with reduced amplitude around $t \approx 8000 R/v_{th}$ and $t \approx 18000 R/v_{th}$. It is the $n_{ZF} = 4$ zonal flow mode, i.e., the next shorter radial scale mode, that dominates the shear spectrum $|\hat{\omega}_{E \times B}|_{n_{ZF}}$ in the latter two phases (not shown). This demonstrates a competition between the $n_{ZF} = 3$ and $n_{ZF} = 4$ modes. Most

importantly, no secular growth of the $n_{\text{ZF}} = 1$ (box scale) zonal flow mode is observed during the entire quasi-stationary turbulent phase [Fig. 4.7(b) dotted line]. The above discussion indicates that although the $n_{\text{ZF}} = 3, 4$ zonal modes compete, the pattern scale does not converge to the radial box scale but rather to a mesoscale of $\sim 57.20 - 76.27 \rho$ (i.e., $n_{\text{ZF}} = 4, 3$ in the 3×1 case).

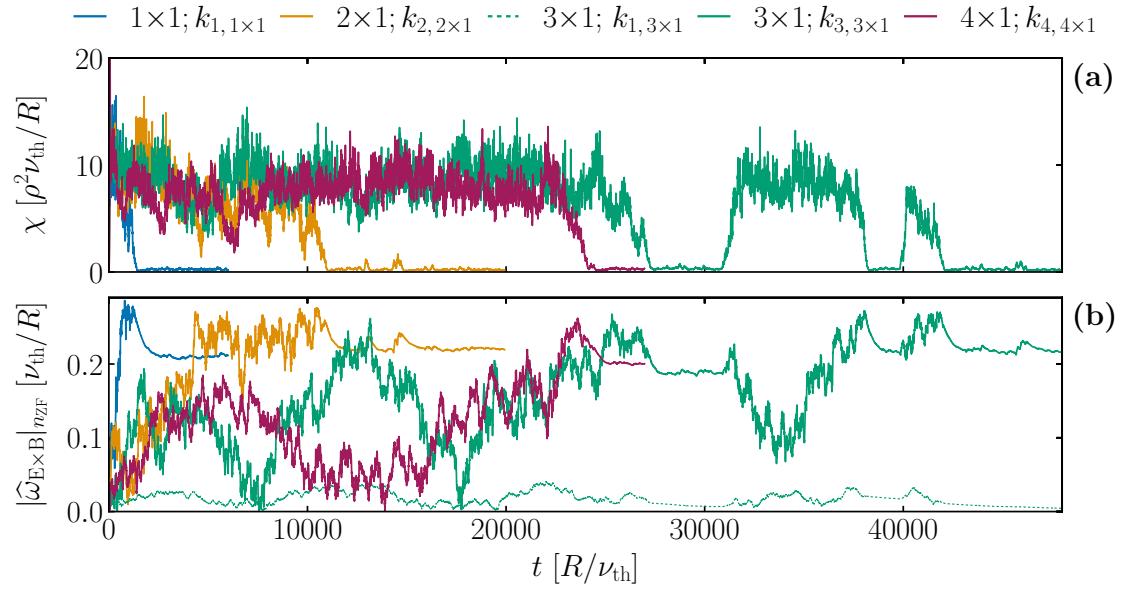


Figure 4.7: (a) Time traces of the heat conduction coefficient χ for $R/L_T = 6.0$ for radial increased box sizes
 (b) Time traces of $|\hat{\omega}_{E \times B}|_{n_{\text{ZF}}}$ for radial increased box sizes

4.2.2 Isotropic increased Box Size

Since the radially elongated simulation domain might inhibit the development of isotropic turbulent structures, in the second test the radial and binormal box size is ~~is~~^{was} increased simultaneously. This scan covers the realizations

$$N_R \times N_B \in [1 \times 1, 2 \times 2, 3 \times 3].$$

Interestingly, suppression of the turbulence by the emergence of a fully developed staircase pattern always occurs ~~ed~~ after $\sim 1000 R/v_{th}$ [Fig. 4.8(a)], i.e., significantly faster compared to the 3×1 and 4×1 realizations. As shown in Fig. 4.10(b) also this test confirms ~~ed~~ the convergence of the staircase pattern size to a typical mesoscale, ~~which~~ is distinct from the radial box size in the $N_R > 1$ realizations.

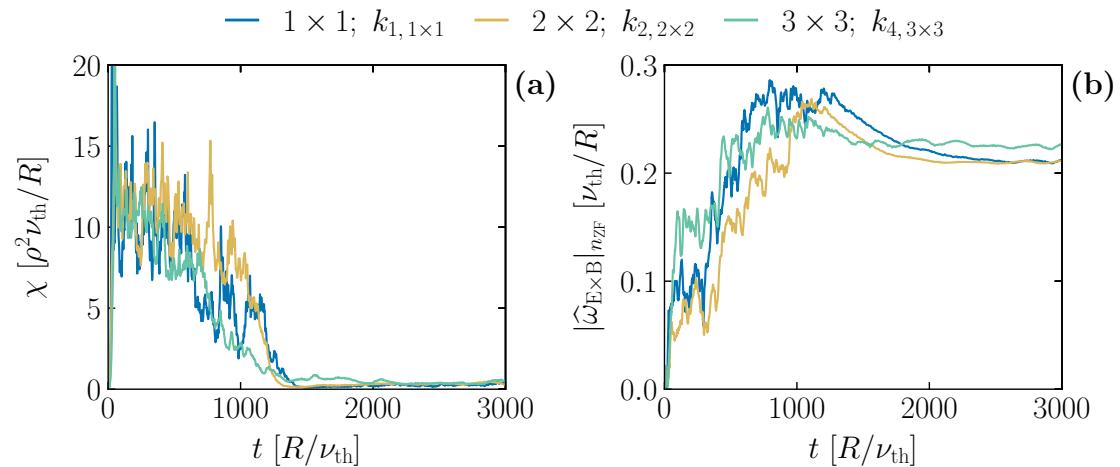


Figure 4.8: (a) Time traces of the heat conduction coefficient χ for $R/L_T = 6.0$ for isotropic increased box sizes
 (b) Time traces of $|\hat{\omega}_{E \times B}|_{n_{ZF}}$ for isotropic increased box sizes

By contrast to the radial box size scan the 3×3 realization shows ~~ed~~ a stationary pattern with four repetitions of the fully developed staircase structure, i.e., a somewhat smaller pattern size [Fig. 4.8(b), Fig. 4.10(b)]. Whether this is related to a possible pattern size dependence on the binormal box size or to the competition between patterns with the two sizes $\lambda \in [57.20, 76.27] \rho$ as observed in the first test is addressed in the next paragraph.

4.2.3 Binormal increased Box Size

In a third test the binormal box size ~~is~~^{was} varied with the radial box size fixed to $N_R = 3$. This test covers the realizations

$$N_R \times N_B \in [3 \times 1.5, 3 \times 2.5, 3 \times 3, 3 \times 5].$$

As in the isotropic scan the turbulence subdued and a fully developed staircase pattern forms ~~it~~ after $\sim 2000 R/v_{th}$ [Fig. 4.9(a)]. The convergence of staircase pattern can be seen in Fig. 4.10(c) and confirms again a size of a typical mesoscale. Fig. 4.10(c) also confirms that indeed a competition between patterns with two sizes $\lambda \in [57.20, 76.27] \rho$ causing the different results for 3×1 and 3×3 . The zonal flow mode number varies between $n_{ZF} = 3, 4$ which can be seen in Fig. 4.10(c) in the 3×2.5 realization. The staircase structure has a pattern between 3 and 4 repetitions which ~~get~~ ^{were produced} represented in the second repetition with no significant plateau at positive shear. Instead, the pattern returns immediately after reaching the maximum shear ($+\gamma$) to the minimum shear ($-\gamma$) of the third repetition in a steep flank. The Fourier analysis of this case yields no definitely basic mode, rather two dominating modes with $n_{ZF} = 3, 4$ with a fraction of the maximum amplitude $|\hat{\omega}_{E \times B}|_{n_{ZF}}$ each [Fig. 4.9(b)].

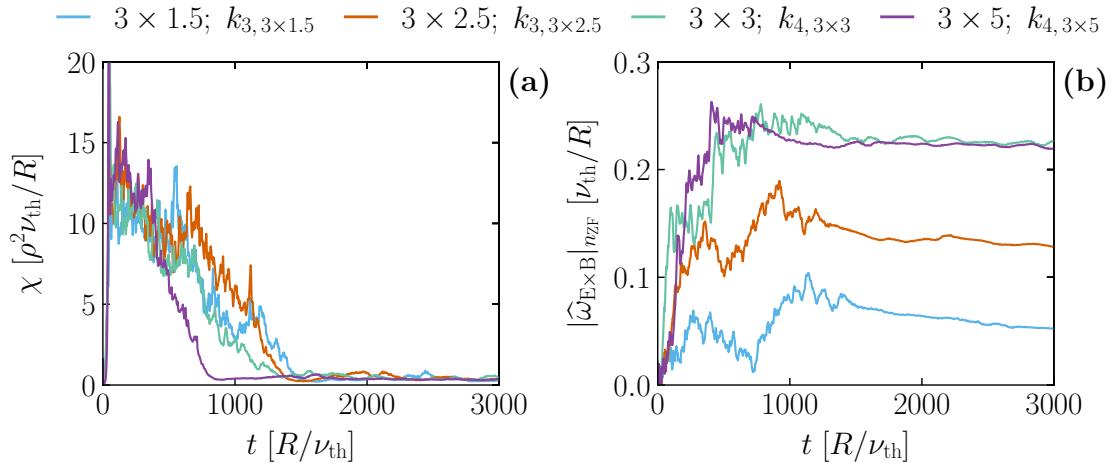


Figure 4.9: (a) Time traces of the heat conduction coefficient χ for $R/L_T = 6.0$ for binormal increased box sizes
 (b) Time traces of $|\hat{\omega}_{E \times B}|_{nZF}$ for binormal increased box sizes

4.2.4 Staircase Structures in Comparison

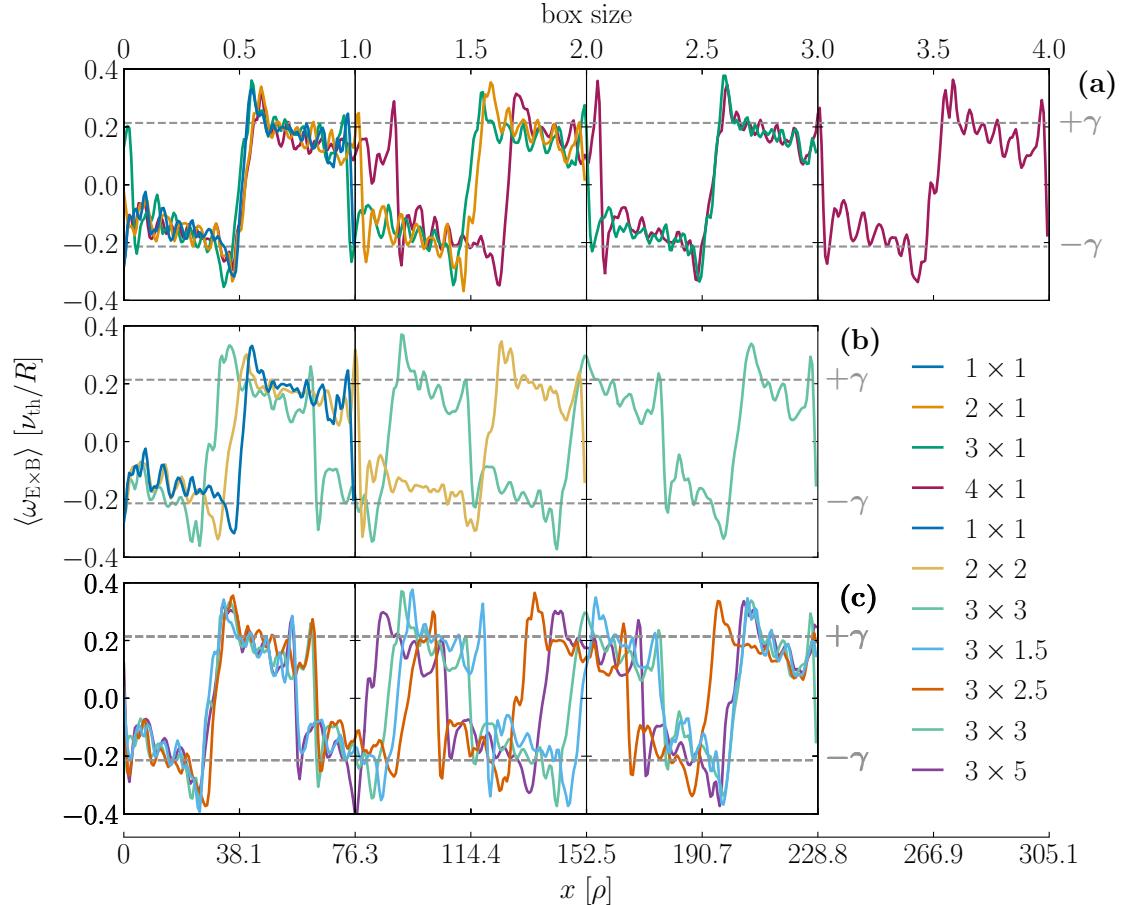


Figure 4.10: Comparison of shearing rate $\langle \omega_{E \times B} \rangle$ for each box size ℓ scan averaged over α given time interval and the growth rate $\pm\gamma$ of the most unstable linear ITG driven Eigenmode. The staircase structures are radially shifted with respect to each other till alignment for better visibility.

(a) **radial:** $t_{1 \times 1} \in [2000, 5000]$, $t_{2 \times 1} \in [15000, 18000]$,
 $t_{3 \times 1} \in [43000, 45000]$, $t_{4 \times 1} \in [26000, 28000]$

(b) **isotropic:** $t_{1 \times 1} \in [2000, 5000]$, $t_{2 \times 2} \in [2000, 3000]$,
 $t_{3 \times 3} \in [2000, 3000]$

(c) **binormal:** $t_{3 \times 1.5} \in [2000, 3000]$, $t_{3 \times 2.5} \in [2000, 3000]$,
 $t_{3 \times 3} \in [2000, 3000]$, $t_{3 \times 5} \in [1000, 3000]$

4.3 The finite Heat Flux Threshold

In the final test the inverse background temperature gradient length R/L_T was varied at fixed 3×3 box size. Since suppression of turbulence usually occurs at later times, when approaching the finite heat flux threshold from below^[42], the analysis aims to lengthen the phase during which the zonal flow varies in time due to turbulent Reynolds stresses. This scan covers realizations with

$$R/L_T \in [6.0, 6.2, 6.4] .$$

In the case of $R/L_T = 6.2$ turbulence suppression was observed for $t > 11000 R/v_{\text{th}}$, while stationary turbulence during the entire simulation time trace of $12000 R/v_{\text{th}}$ was found for $R/L_T = 6.4$ [Fig. 4.11]. The finite heat flux threshold, hence, is

$$R/L_T|_{\text{finite}} = 6.3 \pm 0.1$$

in accordance to Ref. [42]. Although the initial quasi-stationary turbulence in the former case was significantly longer compared to the $R/L_T = 6.2$ realization discussed in the second test, a stationary pattern with basic zonal flow mode $n_{\text{ZF}} = 3$ establishes. Again, the $n_{\text{ZF}} = 1$ (box scale) zonal flow mode does not grow secularly during the entire turbulent phase. Also, this test confirms the statistical soundness of the converged pattern size of $\sim 57.20 - 76.27 \rho$.

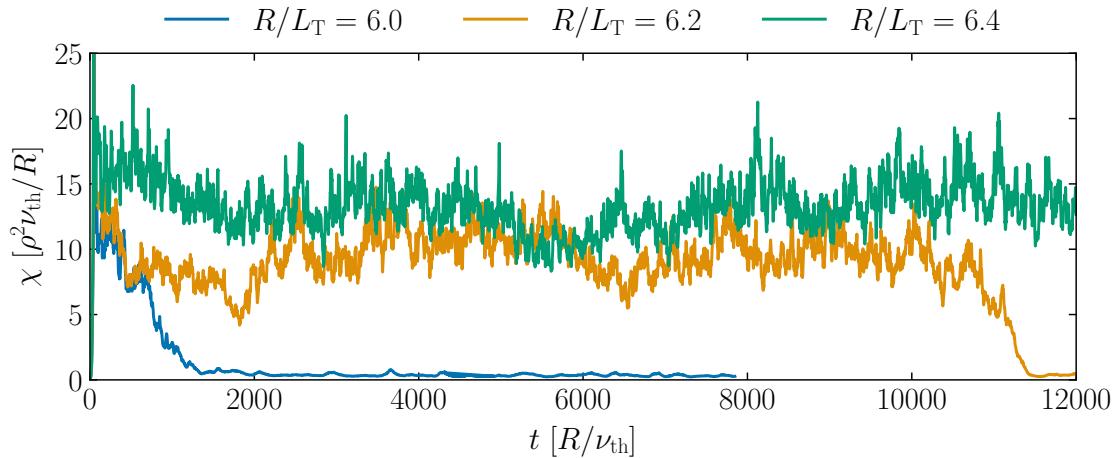


Figure 4.11: Time traces of the heat conduction coefficient χ for different gradient lengths R/L_T

Closure / Conclusion

5

In this thesis the minimal resolution for simulations with GKW in the Cyclone Base parameter were observed in which the number of grid points for the parallel velocity $N_{v_{\parallel}}$ could be reduced from 64 to 48, which halves^{ed} the time until suppression.^{of turbulence.}

Additionally,
Additional~~l~~^{ed} the active development of a restart script in python leads to further convenience during the task of performing ~~a~~ simulations on the btrzx1 cluster.

Through careful tests, this bachelor thesis confirms^{ed} the radial size convergence of the $E \times B$ staircase pattern in local gyrokinetic flux tube simulations of ion temperature gradient (ITG) driven turbulence. A mesoscale pattern size of $\sim 57.20 - 76.27 \rho$ is^{was} found to be intrinsic to ITG driven turbulence for Cyclone Base Case parameters. This length scale is somewhat larger compared to results from global studies with finite ρ_* , which report of a few 10ρ ^[16], which would and has to be considered the proper mesoscale in the local limit $\rho_* \rightarrow 0$. The occurrence of this mesoscale implies that non-locality, in terms of Ref. [16], is inherent to ITG driven turbulence, since avalanches are spatially organized by the $E \times B$ staircase pattern^{[37][16][44][42]}.

CHAPTER

Appendix 6

6.1 slurm_monitor.py

```
#!/usr/bin/env python3

# AUTHOR =====
#   ↪ =====
# Manuel Lippert (GitHub: ManeLippert (https://github.com/ManeLippert))
#   ↪ =====

# MODULES =====
#   ↪ =====

import datetime, time, os, sys, subprocess, math, argparse, pkg_resources

# PARSER =====
#   ↪ =====

description_text = """

===== DESCRIPTION =====
#   ↪ =====

FEATURES:
o NO REQUIREMENTS: Default script runs with standard python3 library
o Creates jobscript file with defined content
  (look into variable "jobscriptContent" to add more option)
o Start/Restarts job until criteria is suffused (default=10000)
o Makes backup after each run before Restart and Restore files after failed
  ↪ run
  (uses rsync command line utility)
o Reset option after run fails and dump files were written
  (rely on "h5py" & "pandas" & "numpy" which get installed by the script)
o Creates status file with current status and progress bars and updates it
  ↪ dynamically
  Progress: Total progress of simulation
  Run X : Progress of current run
o Sends mail at the beginning, end and restart (default=False) with status
  ↪ file
  as attachment
  (mailx has to be installed, look into "send_mail" function for more info)

START SCRIPT IN BACKGROUND:

o WITH NOHUP:
  Command:
    >>> nohup python3 -u slurm_monitor.py --job-name $JOBNAME &> /dev/null &

  Kill Process:
    >>> python3 -u slurm_monitor.py --job-name $JOBNAME --kill
```

6 Appendix

```
o WITH SCREEN (has to be installed):
Create Screen:
>>> screen -S $SESSION

Command:
>>> python3 -u slurm_monitor.py --job-name $JOBNAME --verbose

Leave Screen:
>>> ((Strg + a) + d)

Enter Screen:
>>> screen -r
or
>>> screen -r $SESSION

Kill Process:
>>> ^C (Strg + c)
or
>>> (Strg + d) (kill Screen itself)

OUTPUT STATUS:
Just run the file will create an dynamic output (recommended with using
    ↪ screen) or
>>> cd $DATA && find . -name $STATUSFILENAME -exec tail -8 {} \;

=====
    ↪ =====
"""

parser = argparse.ArgumentParser(description=description_text,
    ↪ formatter_class=argparse.RawTextHelpFormatter)
#parser._action_groups.pop()

required = parser.add_argument_group("required arguments")

required.add_argument("-j", "--job-name", dest="jobname", nargs="?",
    ↪ type=str, required= True,
        help="job name not longer than 8 character")

additional = parser.add_argument_group("additional arguments")

additional.add_argument("-n", dest="timesteps", nargs="?", type=int,
    ↪ default=10000,
        help="required timesteps
    ↪ (default=10000)")

additional.add_argument("-r", "--reset", dest="reset", action="store_true",
    ↪ help="Uses Dumpfiles to reset Simulation
    ↪ (default=False)")

additional.add_argument("-v", "--verbose", dest="verbose",
    ↪ action="store_true",
        help="activate output of script")
```

```

    ↵ (default=False)")

additional.add_argument("--b", "--backup", dest="backup", nargs="?", type=str,
                      help="backup location for files
    ↵ (default=None)\n"+
                          "- local (creates backup in simulation
    ↵ folder)\n"+
                          "- home (creates backup in home folder)")

additional.add_argument("--jobscript", dest="jobscriptFile", nargs="?",
                      type=str, default="jobscript-create",
                      help="jobscript to run SLURM job
    ↵ (default=jobscript-create)")

additional.add_argument("--ntask-per-node", dest="tasks", nargs="?",
                      type=str, default="32",
                      help="MPI task per node
    ↵ (default=32)")

additional.add_argument("--nodes", dest="nodes", nargs="?", type=str,
                      default="3",
                      help="number of nodes
    ↵ (default=3)")

additional.add_argument("--walltime", dest="wallTime", nargs="?", type=str,
                      default="0-24:00:00",
                      help="walltime of server (d-hh:mm:ss)
    ↵ (default=0-24:00:00)")

additional.add_argument("--mail", dest="mail", nargs="?", type=str,
                      help="mail address (mail@server.de)
    ↵ (default=None)")

additional.add_argument("--restart-mail", dest="restartmail",
                      action="store_true",
                      help="mail after every restart
    ↵ (default=False)")

additional.add_argument("--statusfile", dest="statusFile", nargs="?",
                      type=str, default="status.txt",
                      help="file with output from nohup command
    ↵ (default=status.txt)")

additional.add_argument("--restartfile", dest="restartFile", nargs="?",
                      type=str, default="FDS.dat",
                      help="restart file with data
    ↵ (default=FDS.dat)")

additional.add_argument("--format", dest="formattable", nargs="?", type=str,
                      help="format of output table
    ↵ (default=None)\n"+
                          "- fancy (round box)\n"+
                          "- universal (crossplatform)\n"+

```

```

        "-- None (No frame)")

additional.add_argument("--refresh-rate", dest="sleepTime", nargs="?",
                      type=int, default=300,
                      help="time interval to check status in sec
                      ↵ (default=300)")

additional.add_argument("--kill", dest="kill", action="store_true",
                      help="kills monitor process
                      ↵ (default=False)")

args = parser.parse_args()

# VARIABLES =====
# ↵ =====

outputCriteria = ["0", "Run successfully completed"]

slurmFiles = [f for f in os.listdir() if "slurm-" in f]
runCounter = len(slurmFiles)

nTimestepsCurrent = 0

currentTime = "00:00:00"

dataFilename = "gkwdata.h5"

check1Filename, check2Filename = "DM1.dat", "DM2.dat"
check1Bin, check2Bin = check1Filename.replace(".dat", ""),
                      ↵ check2Filename.replace(".dat", "")

# PARSER VARIABLES =====
# ↵ =====

emailAddress = args.mail
RESTARTMAIL = args.restartmail

backupLocation = args.backup

jobName = args.jobname
tasks = args.tasks
nodes = args.nodes
wallTime = args.wallTime

nTimestepsRequired = args.timesteps

jobscriptFilename = args.jobscriptFile
restartFilename = args.restartFile
restartBin = restartFilename.replace(".dat", "")

statusFilename = args.statusFile
#statusFile = open(statusFilename, "r+")

```

```
formatTable = args.formattable
VERBOSE = args.verbose
RESET = args.reset

# Kill process of monitoring
kill = args.kill

# Changing this value can cause problems in writing status file
sleepTime = args.sleepTime

def outputFilename(info):
    return "./slurm-" + info + ".out"

## JOBSRIPT CONTENT =====
→ =====

jobscriptContent = """#!/bin/bash -l

# jobname
#SBATCH --job-name=""" + jobName + """

# MPI tasks
#SBATCH --ntasks-per-node=""" + tasks + """

# number of nodes
#SBATCH --nodes=""" + nodes + """

# walltime
#           d-hh:mm:ss
#SBATCH --time=""" + wallTime + """

# execute the job
time mpirun -np $SLURM_NTASKS ./gkw.x > output.dat

# end
exit 0
"""

jobName = jobName[0:8]

## FLAGS =====
→ =====

startOutputFlag = "Submitted batch job "
restartFlag = "FILE_COUNT"

## SWITCHES =====
→ =====

if emailAddress==None:
    EMAIL = False
else:
    EMAIL = True
```

```

if backupLocation==None:
    BACKUP = False
else:
    BACKUP = True

## PATHS =====
→ =====

user = os.getlogin()
folder = os.getcwd()
path = folder.split(user + "/") [1]

# Set backup location
if BACKUP:
    # If local has been specified as backupLocation, then the backup is
    → copied to the same directory as the simulation folder
    # (simFolder) is located in. The backup is copied to a directory with
    → name simFolder + "-backup".
    if backupLocation == "local":
        simFolder = path.split("/") [-1]
        backupPath = folder + "/.." + simFolder + "-backup"

    # Create backup in home folder
    elif backupLocation == "home":
        simFolder = path.split("/") [-1]
        backupPath = "~/" + simFolder + "-backup"

    # Otherwise, use the backupLocation parsed as argument.
    else:
        if backupLocation[-1] != "/":
            backupLocation += "/"
        backupPath = backupLocation + path

    if not os.path.exists(backupPath):
        os.makedirs(backupPath)

## COMMANDS =====
→ =====

commandJobRunning = "squeue --states=running --name " + jobName
commandJobPending = "squeue --states=pending --name " + jobName

commandMonitorKill = "ps ax | grep " + jobName + " | awk '{print $1}'"

commandJobStarting = "sbatch"

if BACKUP:
    commandBackup = ["rsync", "-a", "", backupPath]
    commandRestore = ["rsync", "-a", "-I", "--exclude=" + statusFilename,
    → backupPath + "/", ""]

# FUNCTIONS =====

```

```

→ =====
## PROGRESSBAR =====
→ =====

def progressbar(required_value, current_value, barsize=42,
                 prefix="",
                 progress_fill="=", progress_fill_top="",
→ progress_fill_bot="",
                 progress_unfill=".",
                 progress_bracket=["[", "]"]):

    x = int(barsize*current_value/required_value)
    percent = int(100*current_value/required_value)

    try:
        percent_format = " (" + (int(math.log10(100)) -
→ int(math.log10(percent)))*" " + "{}%)"
    except ValueError:
        percent_format = " (" + int(math.log10(100))*" " + "{}%)"

    try:
        ratio_format = " " + (int(math.log10(required_value)) -
→ int(math.log10(current_value)))*" " + "{}/{}"
    except ValueError:
        ratio_format = " " + int(math.log10(required_value))*" " + "{}/{}"

    bar_format = "{ } " + progress_bracket[0] + progress_fill_bot + "{ }" +
→ progress_fill_top + "{ }" + progress_bracket[1] + percent_format +
→ ratio_format
    bar = bar_format.format(prefix, progress_fill*x,
→ progress_unfill*(barsize-x), percent, current_value, required_value)

    return bar

## OUTPUT TABLE =====
→ =====

def message(string, add):
    string = string + add + "\n"
    return string

def print_table_row(content,
                    delete_line_index = -10, table_width = 80,
                    output_type = None,
                    TIMEINFO = True, WRITEFILE=True):

    current_value, required_value = nTimestepsCurrent, nTimestepsRequired,
    run_conter = runCounter
    current_time, required_time = currentTime, wallTime

    msg=""

```

```

table_inner_width = table_width - 4
table_content_width = table_inner_width - 47

if formatTable == "fancy":
    table_outline = ["╭", "╮", "╮", "╮", "┫", "┫", "┫", "┫", "┫", "╮"]

if formatTable == "universal":
    table_outline = ["+−", "+−", "+−", "+−", "+−", "+−", "+−", "+−", "+−", "+−"]

if formatTable == None:
    table_outline = ["−−", "−−", "−−", "−−", "−−", "−−", "−−", "−−", "−−", "−−"]

sep_top = table_outline[0] + table_inner_width*table_outline[8] +
    ↪ table_outline[1]
sep_mid = table_outline[4] + table_inner_width*table_outline[8] +
    ↪ table_outline[5]
sep_end = table_outline[2] + table_inner_width*table_outline[8] +
    ↪ table_outline[3]

row_cols    = [2, 10, table_content_width, 13, 11, 13, 2]
row_format = "".join(["{:<" + str(col) + "}" for col in row_cols])

progress_cols    = [2, table_inner_width, 2]
progress_format = "".join(["{:<" + str(col) + "}" for col in
    ↪ progress_cols])

progressbar_content = [progressbar(required_value, current_value,
    ↪ progress_fill_top=">",
                                prefix="PROGRESS")]
progressbar_content.insert(0, table_outline[6])
progressbar_content.insert(len(progressbar_content), table_outline[7])

required_time = get_time_in_seconds(required_time)
current_time  = get_time_in_seconds(current_time)

jobStatus = subprocess.getoutput("squeue --name " +
    ↪ jobName).strip().split("\n")

try:
    jobStatusHeader = [" " + jobStatus[0]]
    jobStatusInfo   = [jobStatus[1][11:11+table_inner_width]]
except IndexError:
    jobStatus_cols    = [12, 12, 10, table_inner_width- 71, 13, 11, 13]
    jobStatus_format = "".join(["{:<" + str(col) + "}" for col in
        ↪ jobStatus_cols])

    jobStatusHeader = ["OUTPUT", "NAME", "USER", "", "DATE", "TIME",
    ↪ "W:DD:HH:MM:SS"]
    jobStatusInfo   = [content[0], jobName, user, "", time_date(),
    ↪ time_time(), time_duration(startTime, pastTime)]

    jobStatusHeader    = [jobStatus_format.format(*jobStatusHeader)]
    jobStatusInfo      = [jobStatus_format.format(*jobStatusInfo)]

```

```

jobStatusHeader.insert(0, table_outline[6])
jobStatusHeader.insert(len(jobStatusHeader), table_outline[7])
jobStatusInfo.insert(0, table_outline[6])
jobStatusInfo.insert(len(jobStatusInfo), table_outline[7])

try:
    progressbartime_content = [progressbar(required_time, current_time,
    ↪ progress_fill_top=">",
                                prefix="RUN " +
    ↪ str(run_conter) + (2 - int(math.log10(run_conter)))*" " + "")]
except ValueError:
    progressbartime_content = [progressbar(required_time, current_time,
    ↪ progress_fill_top=">",
                                prefix="RUN " +
    ↪ str(run_conter) + 2*" " + "")]
progressbartime_content.insert(0, table_outline[6])
progressbartime_content.insert(len(progressbartime_content),
    ↪ table_outline[7])

content[1] = content[1][-(table_content_width-1):]
content.insert(0, table_outline[6])
if TIMEINFO:
    if output_type == "header":
        content.append("DATE")
        content.append("TIME")
        content.append("W:DD:HH:MM:SS")
    else:
        content.append(time_date())
        content.append(time_time())
        content.append(time_duration(startTime, pastTime))

content.insert(len(content), table_outline[7])

if output_type == "header":
    delete_line_index = -1

msg += sep_top + "\n"
msg += row_format.format(*content) + "\n"
msg += sep_end + "\n"

elif output_type == "middle":
    if VERBOSE:
        sys.stdout.write("\x1b[1A"*(-delete_line_index + 1))

    msg += sep_mid + "\n"
    msg += row_format.format(*content) + "\n"
    msg += sep_end + "\n"

elif output_type == "update":
    if VERBOSE:
        sys.stdout.write("\x1b[1A"*(-delete_line_index + 1))

```

```

msg += sep_end + "\n"

else:
    if VERBOSE:
        sys.stdout.write("\x1b[1A"*(-delete_line_index + 1))

    msg += row_format.format(*content) + "\n"
    msg += sep_end + "\n"

msg += " "*table_width + "\n" + " "*table_width + "\n"

msg += sep_top + "\n"
msg += progress_format.format(*jobStatusHeader) + "\n"
msg += progress_format.format(*jobStatusInfo) + "\n"
msg += sep_mid + "\n"
msg += progress_format.format(*progressbar_content) + "\n"
msg += progress_format.format(*progressbar_time_content) + "\n"
msg += sep_end + "\n"

if VERBOSE:
    print(msg, flush=True)
if WRITEFILE:
    delete_write_line_to_file(statusFilename, msg, end=delete_line_index)

## PIP INSTALL =====
→ =====

def pip_install(modules):
    required = modules
    installed = {pkg.key for pkg in pkg_resources.working_set}
    missing = required - installed

    if missing:
        subprocess.check_call([sys.executable, "-m", "pip", "install",
        → "--user", *missing],
        → stdout=subprocess.DEVNULL,
        → stderr=subprocess.DEVNULL)

        print_table_row(["INSTALL", "Required Modules installed"])
    else:
        print_table_row(["CHECK", "Modules already installed"])

## FILE =====
→ =====

def get_value_of_variable_from_file(filename, file_index, relative_index,
    → string):
    try:
        content = [i.strip().split() for i in open(filename).readlines()]
        index = [idx for idx, s in enumerate(content) if string in
        → s][file_index]
        value = content[index][relative_index]
        return value
    
```

```

except IndexError:
    print_table_row(["ERROR", "String not found in file"])
    quit()

def find_string_in_file(filename, string):
    with open(filename) as f:
        if string in f.read():
            return True
        else:
            return False

def delete_write_line_to_file(filename, add = "", start=None, end=None):
    with open(filename, "r") as file:
        try:
            lines = file.readlines()[start:end]
            content = "".join(lines)

        except IndexError:
            pass

        content += add

        write_file(filename, content)

def write_file(filename, content):
    with open(filename, "w") as file:
        file.write(content)
        file.flush()

# Reset gkwdata.h5 with the use of dump files DM1, DM2
# AUTHOR: Florian Rath
# IMPORT: gkw_reset_checkpoint.py
# ↳ (https://bitbucket.org/gkw/gkw/src/develop/python/gkw_reset_checkpoint.py)
def reset_simulation(SIM_DIR, NTIME=None, use_ntime=False):

    # Function that deletes all data in interval [nt_reset:nt_broke].
    # ncol considers, if data series is ordered by a multiple interger of
    # ntime.
    def reset_time_trace(indata, dim, ncol, nt_reset):

        # Shift dimension dim to 0.
        data_shifted = np.moveaxis(indata, dim, 0)

        # Reset data.
        data_shifted_trimmed = data_shifted[0:int(nt_reset*ncol),]

        # Shift dimension back.
        out = np.moveaxis(data_shifted_trimmed, 0, dim)

```

```

# free memory
del data_shifted
del data_shifted_trimmed

return out

# Checks if file has binary format.
def is_binary(filename):
    try:
        with open(filename, 'tr') as check_file:
            check_file.read()
        return False
    except:
        return True


# Check for specific files that are no ordinary data files.
def is_file_exception(filename):

    # substrings that have to be checked
    check_list = ['geom.dat', 'DM1.dat', 'DM2.dat', 'FDS.dat', '.o',
    ↪ 'FDS',
    ↪ 'input.dat', 'perform_first.dat', 'perform.dat',
    ↪ 'output.dat',
    ↪ 'gkwdata.meta', 'gkw_hdf5_errors.txt',
    ↪ 'kx_connect.dat',
    ↪ 'jobscript', 'Poincare1.mat', 'perfloop_first.dat',
    ↪ 'par.dat',
    ↪ 'input_init.dat', 'sgrid', 'gkw', '.out', 'status.txt']
    for key in check_list:
        if key in filename:
            return True

    return False


# Check if given file is a PBS or SLURM jobscript.
def is_jobscript(filename):
    with open(filename, 'r') as file:
        for line in file:
            if '#PBS -l' in line:
                return True
            if '#SBATCH' in line:
                return True

    return False


def get_timestep(FILE):
    if(os.path.isfile(SIM_DIR+'/'+FILE+'.dat')):
        EXISTS = True
        with open(SIM_DIR+'/'+FILE+'.dat', 'r') as file:
            for line in file:

```

```

        if 'NT_REMAIN' in line:
            expr = line.replace(' ', '')
            expr = expr.replace(',', '')
            expr = expr.replace('\n', '')
            REMAIN = int(expr.split('=')[1])
        if 'NT_COMPLETE' in line:
            expr = line.replace(' ', '')
            expr = expr.replace(',', '')
            expr = expr.replace('\n', '')
            COMPLETE = int(expr.split('=')[1])

    else:
        REMAIN, COMPLETE, EXISTS = None, None, False

    return REMAIN, COMPLETE, EXISTS

HDF5_FILENAME = "gkwdata.h5"
RESTARTFILE, DUMPF1, DUMPF2, = "FDS", "DM1", "DM2"

# Change to simulation directory.
if(not os.path.isdir(SIM_DIR)):
    return
else:
    os.chdir(SIM_DIR)

# Check if hdf5-file exists.
if(not os.path.isfile(HDF5_FILENAME)):
    return

# First, read hdf5 file and determine the number of big time steps NTIME,
# requested in the input.dat file.
f = h5py.File(HDF5_FILENAME, "r+")

# Get requested big time steps from the /control group in the hdf5-file.
if(NTIME==None):
    NTIME = int(f['input/control/ntime'][::])

# Get number of big time steps after which simulation broke.
# If time.dat exists read this file to obtain number of time steps after
# which simulation broke.
if(os.path.isfile('time.dat')):
    tim = pd.read_csv(SIM_DIR+'time.dat', header=None, sep='\s+').values
    NT_BROKE = tim.shape[0]
# Else, get time from hdf5-file.
else:
    NT_BROKE = f['diagnostic/diagnos_growth_freq/time'].shape[1]

# Set NT_BROKE for output files holding temporal derivates and therefore
# one timestep less.
NT_BROKE_DERIV = NT_BROKE-1

```

6 Appendix

```
# Close the hdf5-file again.  
f.close()  
  
# Get the number of remaining big time steps NT_REMAIN from checkpoint  
# files FDS.dat. This is used lateron to determine the most recent  
# checkpoint file.  
NT_REMAIN, NT_COMPLETE, FDS_EXISTS = get_timestep(RESTARTFILE)  
  
# Get the number of remaining big time steps NT_REMAIN[1/2] from  
# checkpoint  
# files DM[1/2]. This is used lateron to determine the most recent  
# checkpoint file.  
NT_REMAIN1, NT_COMPLETE1, DM1_EXISTS = get_timestep(DUMPF1LE1)  
NT_REMAIN2, NT_COMPLETE2, DM2_EXISTS = get_timestep(DUMPF1LE1)  
  
# Check if FDS is the most recent checkpoint file. In this case  
# resetting the simulation makes no sense.  
if(FDS_EXISTS):  
    DM1_OLD, DM2_OLD = False, False  
    if(DM1_EXISTS):  
        if(NT_COMPLETE1 < NT_COMPLETE):  
            DM1_OLD = True  
    if(DM2_EXISTS):  
        if(NT_COMPLETE2 < NT_COMPLETE):  
            DM2_OLD = True  
    if(DM1_OLD and DM2_OLD):  
        return  
  
# Now determine which checkpoint file is the recent one.  
if(DM1_EXISTS and DM2_EXISTS):  
    if(NT_REMAIN1 > NT_REMAIN2):  
        NT_REMAIN, NT_COMPLETE, DUMPF1LE = NT_REMAIN2, NT_COMPLETE2,  
        ↪ DUMPF1LE2  
    else:  
        NT_REMAIN, NT_COMPLETE, DUMPF1LE = NT_REMAIN1, NT_COMPLETE1,  
        ↪ DUMPF1LE1  
  
    elif(DM1_EXISTS and not DM2_EXISTS):  
        NT_REMAIN, NT_COMPLETE, DUMPF1LE = NT_REMAIN1, NT_COMPLETE1,  
        ↪ DUMPF1LE1  
  
    elif(DM2_EXISTS and not DM1_EXISTS):  
        NT_REMAIN, NT_COMPLETE, DUMPF1LE = NT_REMAIN2, NT_COMPLETE2,  
        ↪ DUMPF1LE2  
  
else:  
    return  
  
if(not use_ntime):  
    # Find the total number ob big time steps the simulation time trace
```

```

    ↪ should have,
        # when considering the big time steps already completed as well as
    ↪ the big time
        # steps that remain. Can be different from NTIME, since the
    ↪ simulation could
        # have been restarted several times such that NTIME > NT_COMPLETE.
    NTOT = NT_COMPLETE + NT_REMAIN
    N_REQUEST = NTOT

        # Determine the time steps to which the time trace files have to be
    ↪ reset.
        # Use NTOT here, since NTIME could have been changed at some point,
    ↪ or NT_COMPLETE
        # could be larger than NTIME.
    NT_RESET = NTOT-NT_REMAIN
else:
    # Determine the time steps to which the time trace files have to be
    ↪ reset.
    NT_RESET = NTIME-NT_REMAIN
    N_REQUEST = NTIME

# Same for files holding time derivatives.
NT_RESET_DERIV = NT_RESET-1

# -----
# Cycle over all nodes of hdf5-file and reset time trace datasets.

# Check if hdf5-file exists.
if(os.path.isfile(HDF5_FILENAME)):

    # Find all possible keys items, i.e. both groups and datasets
    f = h5py.File(HDF5_FILENAME, "a")
    h5_keys = []
    f.visit(h5_keys.append)

    # Cycle over all keys items and check, if any dimension has size
    ↪ NT_BROKE,
        # i.e., it is a time trace file.
        for item in enumerate(h5_keys):

            data = f.get(item)

            # Consider datasets only.
            if(isinstance(data, h5py.Dataset)):

                #Check if any dimension is an integer multiple of NT_BROKE,
                ↪ by checking
                    # the residual of the division.
                res = [None]*len(data.shape)
                for i in range(len(data.shape)):
                    res[i] =
    ↪ data.shape[i]/NT_BROKE-np.floor(data.shape[i]/NT_BROKE)

```

```

if 0.0 in res:

    # Check which dimension is integer multiple of NT_BROKE
    # and save dimension as well as integer.
    new_shape = data.shape
    for i in range(len(data.shape)):
        ncol = data.shape[i]/NT_BROKE
        res = ncol - np.floor(ncol)
        if(res == 0):
            dim = i
            ncol = int(ncol)
            # adjust new shape to ncol*NT_RESET
            y = list(new_shape)
            y[dim] = int(ncol*NT_RESET)
            new_shape = tuple(y)
            break

    # Reset dataset (.resize discards data with indices
    ↪ larger than
    # ncol*NT_RESET along dimension dim).
    dset = f[item]
    dset.resize(int(ncol*NT_RESET),dim)

# After having repaired all datasets, close the hdf5-file again.
f.close()

# -----
# Cycle over all csv-files and reset time trace.
for filename in os.listdir(SIM_DIR):

    # First perform some checks on files; cycle if file is binary, an
    ↪ exception
    # or a jobscript.
    if(is_binary(filename)):
        continue
    if(is_file_exception(filename)):
        continue
    if(is_jobscript(filename)):
        continue

    # no file
    if(not os.path.isfile(filename)):
        continue

    # Load csv file.
    data = pd.read_csv(SIM_DIR+'/'+filename, header=None,
    ↪ sep='\s+').values

    #Check if any dimension is an integer multiple of NT_BROKE, by
    ↪ checking
    #the residual of the division.

```

```

res = [None]*len(data.shape)

# residual for output that holds time derivatives and therefore nt-1
→ datapoints
res_deriv = [None]*len(data.shape)
for i in range(len(data.shape)):
    res[i] = data.shape[i]/NT_BROKE-np.floor(data.shape[i]/NT_BROKE)
    res[i] =
→ data.shape[i]/(NT_BROKE_DERIV)-np.floor(data.shape[i]/(NT_BROKE_DERIV))

# ordinary files
if 0.0 in res:

    #Check which dimension is integer multiple of NT_BROKE
    #and save dimension as well as integer.
    for i in range(len(data.shape)):
        ncol = data.shape[i]/NT_BROKE
        res = ncol - np.floor(ncol)
        if(res == 0):
            dim = i
            ncol = int(ncol)
            break

    # Load original dataset.
    original_data = data
    # print('\t Original shape: \t'+str(original_data.shape))

    # Reset time trace.
    reset_data = reset_time_trace(original_data, dim, ncol, NT_RESET)
    # print('\t Reset shape: \t' +str(reset_data.shape))

    # Save resetted data.
    pd.DataFrame(reset_data).to_csv(filename, sep='\t', header=None,
→ index=None)

    # files holding time derivatives
    if 0.0 in res_deriv:

        # Check which dimension is integer multiple of NT_BROKE
        # and save dimension as well as integer.
        for i in range(len(data.shape)):
            ncol = data.shape[i]/NT_BROKE_DERIV
            res = ncol - np.floor(ncol)
            if(res == 0):
                dim = i
                ncol = int(ncol)
                break

        # Load original dataset.
        original_data = data

        # Reset time trace.
        reset_data =

```

```

→ reset_time_trace(original_data, dim, ncol, NT_RESET_DERIV)

    # Save resetted data.
    pd.DataFrame(reset_data).to_csv(filename, sep='\t', header=None,
→ index=None)

# -----
# Finally, copy most recent dump file to FDS[/.dat].

# Copy the most recent dump file to FDS[/.dat].
copyfile(SIM_DIR+'/' +DUMPFFILE, SIM_DIR+'/' +'FDS')
copyfile(SIM_DIR+'/' +DUMPFFILE+'.dat', SIM_DIR+'/' +'FDS.dat')

# First line of the so produced FDS.dat has to be modified.
old_text = '!Dump filename: '+DUMPFFILE
new_text = '!Dump filename: +'+'FDS'

# Replace first line in FDS.dat to set the correct file name.
with fileinput.input(SIM_DIR+'/' +'FDS.dat', inplace=True) as f:
    for line in f:
        line.replace(old_text, new_text)

def check_and_delete_file(filenames):

    for filename in filenames:
        if(os.path.isfile(filename)):
            os.remove(filename)

def check_checkpoint_files():
    DM1, DM2 = False, False

    if(os.path.isfile(check1Filename) and os.path.isfile(check1Bin)):
        DM1 = True
    if(os.path.isfile(check2Filename) and os.path.isfile(check2Bin)):
        DM2 = True

    return DM1, DM2

def get_timestep_from_restartfile(filename, flag):
    return int(get_value_of_variable_from_file("./" + filename, 0, 2,
→ flag).replace(", ", ""))

```

def get_ntimestepCurrent(filenames):

```

ntimestep = 0

for f in filenames:
    if(os.path.isfile(f)):
        timestepFile = get_timestep_from_restartfile(f, restartFlag)

        if timestepFile > ntimestep:

```

```

        ntimestep = ntimestepFile

    return ntimestep

def get_time_from_statusfile(filename, line_index):
    with open(filename, "r") as file:
        line = file.readlines()[line_index]

        content = line.split(" ")
        time = content[-2]
        time_sec = get_time_in_seconds(time)

    return time_sec

## TIME =====
→ =====

def format_num(time):
    if time < 10:
        return "0" + str(time)
    else:
        return str(time)

def get_time_as_string(sec):

    mins, sec = sec // 60, sec % 60
    hours, mins = mins // 60, mins % 60
    days, hours = hours // 24, hours % 24
    weeks, days = days // 7, days % 7

    timeConvertedString = (str(int(weeks)) + ":" +
                           format_num(int(days)) + ":" +
                           format_num(int(hours)) +
                           → ":" +
                           format_num(int(mins)) + ":" +
                           format_num(int(sec)))

    return timeConvertedString

def get_time_in_seconds(time):

    # Format D-HH:MM:SS or HH:MM:SS or MM:SS

    time = time.replace("-", ":")
    time_split = time.split(":")

    seconds = [7*24*60*60, 24*60*60, 60*60, 60, 1]
    seconds = seconds[-len(time_split):]

    time_sec = sum([a*b for a,b in zip(seconds, map(int,time_split))])

    return time_sec

def time_date():
    e = datetime.datetime.now()

```

```

    return "%s-%s-%s" % (format_num(e.year), format_num(e.month),
↪ format_num(e.day))

def time_time():
    e = datetime.datetime.now()
    return "%s:%s:%s" % (format_num(e.hour), format_num(e.minute),
↪ format_num(e.second))

def time_duration(startTime, pastTime):
    stop = time.time()
    return get_time_as_string(stop - startTime + pastTime)

## STATUS =====
↪ =====

def get_job_status():

    jobStatusRunning =
    ↪ subprocess.getoutput(commandJobRunning).strip().split()
    jobStatusPending =
    ↪ subprocess.getoutput(commandJobPending).strip().split()

    return jobStatusRunning, jobStatusPending

def set_output_type():
    jobStatusRunning, jobStatusPending = get_job_status()

    if jobName in jobStatusRunning:
        outputType = "running"
    elif jobName in jobStatusPending:
        outputType = "pending"
    else:
        outputType ="no Output"

    return outputType

def get_error_type(filename):

    slurm_errors = {"executable": ["error on file ./gkw.x (No such file or
↪ directory)", "No executable found"],
                    "walltime": ["process killed (SIGTERM)", "Exceeded wall
↪ time"],
                    "timeout": ["DUE TO TIME LIMIT", "Exceeded time limit"],
                    "config": ["couldn't open config directory", "Config not
↪ loading"],
                    "hdf5": ["HDF5-DIAG", "Writing h5 file failed"]}

    for key in slurm_errors:
        if find_string_in_file(filename, slurm_errors[key][0]):
            return slurm_errors[key][1]

    return "Unknown error occurred"

```

```

## MAIL =====
→ =====

def send_mail(recipient, subject, body = None):

    recipient = recipient.encode("utf_8")

    subject = subject.replace(" ", "_")
    subject = subject.encode("utf_8")

    if body == None:
        body = "For futher information open attachment"

    body = body.encode("utf_8")

    attachmentPath = folder + "/" + statusFilename
    attachment = attachmentPath.encode("utf_8")

    process = subprocess.Popen(["ssh", "master", "/usr/bin/mailx", "-s",
                               → subject, "-a", attachment, recipient],
                               stdin=subprocess.PIPE)
    process.communicate(body)

## KILL JOB =====
→ =====

PID = subprocess.getoutput(commandMonitorKill).split("\n") [0]
if kill:
    subprocess.run(["kill", PID])
    quit()

# START/RESTART JOB =====
→ =====

startTime = time.time()
outputType = set_output_type()
jobStatusRunning, jobStatusPending = get_job_status()

# Set current Time for progress bar
if outputType == "running":
    jobStatusRunningNameIndex = [idx for idx, s in
                                → enumerate(jobStatusRunning) if jobName in s][0]
    currentTime = jobStatusRunning[jobStatusRunningNameIndex + 3]
elif outputType == "pending":
    jobStatusPendingNameIndex = [idx for idx, s in
                                → enumerate(jobStatusPending) if jobName in s][0]
    currentTime = jobStatusPending[jobStatusPendingNameIndex + 3]
else:
    currentTime = "00:00:00"

# Set pastTime and create status file if necessary. When status file exist
→ append next lines
WRITEHEADER = True

```

```

if not os.path.isfile(statusFilename):
    pastTime = 0
    write_file(statusFilename, "")
else:
    try:
        pastTime = get_time_from_statusfile(statusFilename, -11)
        WRITEHEADER = False
    except IndexError:
        pastTime = 0

print_table_row(["OUTPUT", "INFO"], output_type="header",
    ↪ WRITEFILE=WRITEHEADER)

## BEGIN =====
    ↪ =====

# Check if timesteps criterion is satisfied, send mail and end monitoring
# else continue monitoring and set output type accordingly
nTimestepsCurrent = get_ntimestepCurrent([restartFilename, check1Filename,
    ↪ check2Filename])

if nTimestepsCurrent >= nTimestepsRequired:
    print_table_row(["CONTROL", "Current Timesteps " +
    ↪ str(nTimestepsCurrent)], output_type="middle")
    print_table_row(["SUCCESS", "Stop monitoring " + jobName])

    if EMAIL:
        send_mail(emailAddress, "Ended Job " + jobName)

    quit()

# Continue monitoring and send mail
else:

    if outputType != "no Output":
        print_table_row(["CONTINUE", "Continue monitoring " + jobName],
    ↪ output_type="middle")
        print_table_row(["CONTROL", "Current Timesteps " +
    ↪ str(nTimestepsCurrent)])

        if EMAIL:
            send_mail(emailAddress, "Continued Job " + jobName)

    elif os.path.isfile(restartFilename):
        print_table_row(["CONTINUE", "Continue monitoring " + jobName],
    ↪ output_type="middle")

        if EMAIL:
            send_mail(emailAddress, "Continued Job " + jobName)

    else:
        print_table_row(["STARTING", "Start monitoring " + jobName],
    ↪ output_type="middle")

```

```

if BACKUP:
    print_table_row(["BACKUP", backupPath])
    subprocess.run(commandBackup)

if EMAIL:
    send_mail(emailAddress, "Started Job " + jobName)

if RESET:
    pip_install({"h5py", "pandas", "numpy"})

    import h5py, fileinput
    import pandas as pd
    import numpy as np
    from shutil import copyfile

    print_table_row(["IMPORT", "Load numpy, pandas, h5py"])

else:
#    pip_install({"h5py"})
#    import h5py
#    print_table_row(["IMPORT", "Load module h5py"])

## MONITOR ROUTINE =====
# =====

while True:

    # Check job status to monitor current state
    jobStatusRunning, jobStatusPending = get_job_status()

    # Job running
    if jobName in jobStatusRunning:

        jobStatusRunningNameIndex = [idx for idx, s in
        ↪ enumerate(jobStatusRunning) if jobName in s][0]
        jobID = jobStatusRunning[jobStatusRunningNameIndex - 2]

        currentTime = jobStatusRunning[jobStatusRunningNameIndex + 3]
        nTimestepsCurrent = get_ntimestepCurrent([restartFilename,
        ↪ check1filename, check2filename])

        if outputType == "running":
            print_table_row(["RUNNING", "Job is executed"])
            outputType = "no Output"
        else:
            print_table_row(["RUNNING", "Job is executed"],
            ↪ output_type="update")

        time.sleep(sleepTime)

    # Job pending
    elif jobName in jobStatusPending:

```

```

        jobStatusPendingNameIndex = [idx for idx, s in
    ↵ enumerate(jobStatusPending) if jobName in s][0]
        jobID = jobStatusPending[jobStatusPendingNameIndex - 2]

        currentTime = jobStatusPending[jobStatusPendingNameIndex + 3]
        nTimestepsCurrent = get_ntimestepCurrent([restartFilename,
    ↵ check1Filename, check2Filename])

        if outputType == "pending":
            print_table_row(["WAITING", "Job is pending"])
            outputType = "running"
        else:
            print_table_row(["WAITING", "Job is pending"],
    ↵ output_type="update"))

        time.sleep(sleepTime)

# Job start/restart
else:

    # Check errors and making Backup
    while True:
        try:
            outputContent =
    ↵ open(outputFilename(jobID)).readlines()[-5].replace("\n", "")

            # Create Backup if run is successful
            # If scan of output.dat is needed: Scans for string "Run
    ↵ Successful in output.dat and returns bool value"
            #runSuccess = find_string_in_file("output.dat",
    ↵ outputCriteria[1])
            #if outputCriteria[0] in outputContent and runSuccess:

            if outputCriteria[0] in outputContent:

                # Check if h5 file is closed before start/restart
    ↵ simulation
                    # than check if FDS/FDS.dat is updated
                    try:
                        f = open(dataFilename)
                        #f = h5py.File(dataFilename)
                        f.close()

                        # Check if FDS/FDS.dat is updated after run and has
    ↵ equally time stamp as gkwdata.h5
                        timestamp_data =
    ↵ int(os.path.getmtime(dataFilename))
                        timestamp_restart =
    ↵ int(os.path.getmtime(restartFilename))

                        wallTime_sec = get_time_in_seconds(wallTime)
                        timestamp_remain = timestamp_data - timestamp_restart

```

```

# FDS/FDS.dat does not get written at the same time
↪ as gkwdata.h5

# For that a time interval have to be considered
# To be certain the half wall time is set aus time
↪ interval
    if timestamp_remain > wallTime_sec/2:

        print_table_row(["ERROR", "FDS/FDS.dat not
↪ updated"])

        # Reset simulation and save as backup
        if RESET:

            DM1, DM2 = check_checkpoint_files()

            if (DM1 or DM2):
                print_table_row(["RESET", "Reset to last
↪ checkpoint."])
                reset_simulation(folder)

                # Update backup
                if BACKUP:
                    print_table_row(["BACKUP",
↪ backupPath])
                    subprocess.run(commandBackup)

                # Restore backup to rerun simulation
                elif BACKUP:
                    print_table_row(["RESTORE", backupPath])
                    subprocess.run(commandRestore)

                elif BACKUP:
                    print_table_row(["BACKUP", backupPath])
                    subprocess.run(commandBackup)

                break
            except OSError:
                time.sleep(sleepTime)

        else:
            print_table_row(["ERROR",
↪ get_error_type(outputFilename(jobID))])

            # Reset simulation and save as backup
            if RESET:

                DM1, DM2 = check_checkpoint_files()

                if (DM1 or DM2):
                    print_table_row(["RESET", "Reset to last
↪ checkpoint."])
                    reset_simulation(folder)

```

```

        # Update backup
        if BACKUP:
            print_table_row(["BACKUP", backupPath])
            subprocess.run(commandBackup)

        # Restore backup to rerun simulation
        elif BACKUP:
            print_table_row(["RESTORE", backupPath])
            subprocess.run(commandRestore)

        break

    # Wait sleepTime and check output file again
    except (IndexError, FileNotFoundError):
        time.sleep(sleepTime)

    # If jobID undefined break cycle
    except NameError:
        break

    # Check if timesteps criterion is satisfied, send mail and end
    ↪ monitoring
    nTimestepsCurrent = get_ntimestepCurrent([restartFilename,
    ↪ check1Filename, check2Filename])
    print_table_row(["CONTROL", "Current Timesteps " +
    ↪ str(nTimestepsCurrent)])

    if nTimestepsCurrent >= nTimestepsRequired:
        print_table_row(["SUCCESS", "Stop monitoring " + jobName])

        if EMAIL:
            send_mail(emailAddress, "Ended Job " + jobName)
        break

    # Delete checkpoint files
    check_and_delete_file([check1Bin, check1Filename, check2Bin,
    ↪ check2Filename])

    # Create jobscript
    if jobscriptFilename == "jobscript-create":
        jobscriptFilename = "jobscript"
        write_file(jobscriptFilename, jobscriptContent)

    # Start Job and send restart mail (if activated)
    startOutput = subprocess.check_output([commandJobStarting,
    ↪ jobscriptFilename]).decode("utf-8").replace("\n", "")
    jobID = startOutput.split(startOutputFlag)[1]

    runCounter += 1

    print_table_row(["STARTING", startOutput], output_type="middle")

```

```
try:
    if RESTARTMAIL and EMAIL:
        send_mail(emailAddress, "Restart Job " + jobName)
except NameError:
    continue

time.sleep(30)

# Set output type to running or pending
outputType = set_output_type()

## RESTART =====
→ =====
```

6.2 Status File

OUTPUT	INFO	DATE	TIME	W:DD:HH:MM:SS
STARTING	Start monitoring 3x1.5	2023-01-23	09:58:18	0:00:00:00:00
BACKUP	/scratch/bt712347/backup/	2023-01-23	09:58:18	0:00:00:00:00
STARTING	Submitted batch job 903536	2023-01-23	09:58:20	0:00:00:00:01
RUNNING	Job is executed	2023-01-23	09:58:50	0:00:00:00:31
BACKUP	/scratch/bt712347/backup/	2023-01-24	09:38:08	0:00:23:39:49
CONTROL	Current Timesteps 2086	2023-01-24	09:38:49	0:00:23:40:30
STARTING	Submitted batch job 905108	2023-01-24	09:38:49	0:00:23:40:30
RUNNING	Job is executed	2023-01-24	09:39:19	0:00:23:41:01
BACKUP	/scratch/bt712347/backup/	2023-01-25	09:18:41	0:01:23:20:22
CONTROL	Current Timesteps 4179	2023-01-25	09:19:35	0:01:23:21:16
STARTING	Submitted batch job 906542	2023-01-25	09:19:35	0:01:23:21:16
RUNNING	Job is executed	2023-01-25	09:20:05	0:01:23:21:46
BACKUP	/scratch/bt712347/backup/	2023-01-26	08:59:21	0:02:23:01:03
CONTROL	Current Timesteps 6289	2023-01-26	09:00:38	0:02:23:02:19
STARTING	Submitted batch job 907237	2023-01-26	09:00:38	0:02:23:02:19
RUNNING	Job is executed	2023-01-26	09:01:08	0:02:23:02:49
BACKUP	/scratch/bt712347/backup/	2023-01-27	08:40:56	0:03:22:42:37
CONTROL	Current Timesteps 8411	2023-01-27	08:42:23	0:03:22:44:04
STARTING	Submitted batch job 910236	2023-01-27	08:42:23	0:03:22:44:04
RUNNING	Job is executed	2023-01-27	08:42:53	0:03:22:44:34
BACKUP	/scratch/bt712347/backup/	2023-01-28	02:41:23	0:04:16:43:05
CONTROL	Current Timesteps 10000	2023-01-28	02:42:50	0:04:16:44:31
SUCCESS	Stop monitoring 3x1.5	2023-01-28	02:42:50	0:04:16:44:31
CONTINUE	Continue monitoring 3x1.5	2023-02-09	18:22:46	0:04:16:44:31
BACKUP	/scratch/bt712347/backup/	2023-02-09	18:22:46	0:04:16:44:31
CONTROL	Current Timesteps 10000	2023-02-09	18:22:47	0:04:16:44:31
STARTING	Submitted batch job 936758	2023-02-09	18:22:47	0:04:16:44:31
WAITING	Job is pending	2023-02-09	18:23:17	0:04:16:45:01
CONTINUE	Continue monitoring 3x1.5	2023-02-11	00:12:32	0:04:16:45:01
BACKUP	/scratch/bt712347/backup/	2023-02-11	00:12:32	0:04:16:45:01
CONTROL	Current Timesteps 10000	2023-02-11	00:12:34	0:04:16:45:02
STARTING	Submitted batch job 937665	2023-02-11	00:12:34	0:04:16:45:02
WAITING	Job is pending	2023-02-11	00:13:04	0:04:16:45:32
RUNNING	Job is executed	2023-02-11	01:58:08	0:04:18:30:36
BACKUP	/scratch/bt712347/backup/	2023-02-12	01:38:54	0:05:18:11:22
CONTROL	Current Timesteps 14497	2023-02-12	01:41:22	0:05:18:13:50

6 Appendix

```

| STARTING Submitted batch job 937946 2023-02-12 01:41:22 0:05:18:13:50 |
| RUNNING Job is executed 2023-02-12 01:41:52 0:05:18:14:20 |
| BACKUP /scratch/bt712347/backup/ 2023-02-13 01:22:37 0:06:17:55:06 |
| CONTROL Current Timesteps 18985 2023-02-13 01:25:40 0:06:17:58:08 |
+-----+
| STARTING Submitted batch job 938797 2023-02-13 01:25:40 0:06:17:58:08 |
| RUNNING Job is executed 2023-02-13 01:26:10 0:06:17:58:38 |
| BACKUP /scratch/bt712347/backup/ 2023-02-13 06:51:20 0:06:23:23:48 |
| CONTROL Current Timesteps 20000 2023-02-13 06:54:11 0:06:23:26:40 |
| SUCCESS Stop monitoring 3x1.5 2023-02-13 06:54:11 0:06:23:26:40 |
+-----+
| CONTINUE Continue monitoring 3x1.5 2023-02-13 15:34:37 0:06:23:26:40 |
| BACKUP /scratch/bt712347/backup/ 2023-02-13 15:34:37 0:06:23:26:40 |
| CONTROL Current Timesteps 20000 2023-02-13 15:34:38 0:06:23:26:40 |
+-----+
| STARTING Submitted batch job 939062 2023-02-13 15:34:38 0:06:23:26:41 |
| RUNNING Job is executed 2023-02-13 15:35:08 0:06:23:27:11 |
| BACKUP /scratch/bt712347/backup/ 2023-02-14 11:31:13 1:00:19:23:15 |
| CONTROL Current Timesteps 20000 2023-02-14 11:38:17 1:00:19:30:19 |
+-----+
| STARTING Submitted batch job 952358 2023-02-14 11:38:17 1:00:19:30:19 |
| RUNNING Job is executed 2023-02-14 11:38:47 1:00:19:30:49 |
| ERROR SLURM Job failed to execute 2023-02-14 15:38:58 1:00:23:31:01 |
| RESTORE /scratch/bt712347/backup/ 2023-02-14 15:38:58 1:00:23:31:01 |
| CONTROL Current Timesteps 20000 2023-02-14 15:44:44 1:00:23:36:46 |
+-----+
| STARTING Submitted batch job 956235 2023-02-14 15:44:44 1:00:23:36:46 |
| RUNNING Job is executed 2023-02-14 15:45:14 1:00:23:37:17 |
| ERROR SLURM Job failed to execute 2023-02-15 11:01:35 1:01:18:53:38 |
| RESTORE /scratch/bt712347/backup/ 2023-02-15 11:01:44 1:01:18:53:46 |
| CONTROL Current Timesteps 20000 2023-02-15 11:06:52 1:01:18:58:54 |
+-----+
| STARTING Submitted batch job 969292 2023-02-15 11:06:52 1:01:18:58:54 |
| RUNNING Job is executed 2023-02-15 11:07:22 1:01:18:59:24 |
| ERROR SLURM Job failed to execute 2023-02-16 10:48:07 1:02:18:40:09 |
| RESTORE /scratch/bt712347/backup/ 2023-02-16 10:48:07 1:02:18:40:09 |
| CONTROL Current Timesteps 20000 2023-02-16 10:51:17 1:02:18:43:20 |
+-----+
| STARTING Submitted batch job 970455 2023-02-16 10:51:17 1:02:18:43:20 |
| WAITING Job is pending 2023-02-16 10:51:47 1:02:18:43:50 |
| RUNNING Job is executed 2023-02-16 12:56:50 1:02:20:48:53 |
| BACKUP /scratch/bt712347/backup/ 2023-02-16 16:11:56 1:03:00:03:59 |
+-----+
+-----+
| JOBID PARTITION NAME USER ST TIME NODES NODELIST(REAISON) |
| 970455 normal 3x1.5 bt712347 CG 3:16:01 3 r03n03,r05n[15-16] |
+-----+
| PROGRESS [=====>.....] ( 80%) 20000/25000 |
| RUN 13 [====>.....] ( 13%) 11462/86400 |
+-----+

```

6.3 torque_monitor.py

```
#!/usr/bin/env python3

# AUTHOR: Manuel Lippert (GitHub: ManeLippert
#         ↪ (https://github.com/ManeLippert))
#
# DESCRIPTION:
# Script that starts a given job (shell script) with Sun Grid Engine until
# ↪ the previous defined timestep are completed.
# It also sends an mail alert when the job gets started and ended.

# IMPORTANT:
# Take a look in the Variable Section but overall everthing should work
# ↪ automatically

# START SCRIPT:
# To start the script in the background following command is needed:
#
# >>> nohup python3 -u monitor_job.py &> status.txt &
#
# Output:
#
# >>> [1] 10537
#
# This will write every output in the file jobstatus.txt that will be send
# ↪ as attachment to the defined mail address.

# LIST PRGRESS:
# To see which progress is in background running following command is needed:
#
# >>> ps ax | grep monitor_job.py
#
# Output:
#
# >>> 10537 pts/1      S          0:00 python3 -u monitor_job.py
# >>> 23426 pts/1      S+         0:00 grep --color=auto monitor_job.py
#
# This will give you the ID to kill monitoring script with the command:
#
# >>> kill 10537
#
# This will kill the monitor script.

# OUTPUT STATUS:
# To see all status.txt files with one command follwing command is needed:
#
# >>> cd $DATA && find . -name status.txt -exec tail --lines=+10 {} \;

'''
```

```
#####
#                               #
#          Modules           #
#                               #
#####
'''
```

```
import datetime
import time
from time import sleep
import os
import subprocess
```

```
'''
```

```
#####
#                               #
#          VARIABLES          #
#                               #
#####
'''
```

```
#####
# ADDITIONAL #####
#emailAddress = 'Dominik.Mueller@uni-bayreuth.de'
#backupLocation = '/scratch/bt712347/backup'
```

```
#####
# FILENAMES #####

```

```
# Finds File that ends with .sh
for file in os.listdir():
    if file.endswith('.sh'):
        jobscriptFilename = file
```

```
# Rename jobscript
dirName = os.getcwd().split('/')[-1]
os.rename(jobscriptFilename, dirName + '.sh')
jobscriptFilename = dirName + '.sh'
```

```
monitorFilename = 'status.txt'
```

```
# Finds File that ends with .json
for file in os.listdir():
    if file.endswith('.json'):
        inputFilename = file
```

```
restartFilename = inputFilename
```

```
# Declared as function for dynamic changes
def outputFilename(info):
    return jobscriptFilename + '.o' + info
```

```
#####
# FLAGS #####

```

```

walltimeFlag = 'time='
startOutputFlag = '.mgmt'
outputCriteria = 'WARNING'
inputFlag = '"Number'

restartFlag = 'ETA:'
restartString = '           "Restart from step": '

##### SLEEP TIME #####
sleepTime = 5*60

##### COMMANDS #####
commandJobStatus = 'qstat -u'
commandJobStarting = 'qsub'

'''
#####
#          FUNCTIONS #
#
#####
'''


#####
#INFORMATIONS#
#####

def read_file_to_string(file):
    content = ''.join(open(file).readlines())

    return content

def get_value_of_variable_from_input_file(file, string, idx):
    try:
        content = [i.strip().split() for i in open(file).readlines()]
        index = [idx for idx, s in enumerate(content) if string in s][idx]
        value = int(content[index][4].split(',') [0])
        return value
    except IndexError:
        print('! String not in input file !')
        quit()

def get_value_of_variable_from_output_file(file, string, idx):
    try:
        content = [i.strip().split() for i in open(file).readlines()]
        index = [idx for idx, s in enumerate(content) if string in s][idx]
        value = int(content[index - 1][0])
        return value
    except IndexError:
        print('! String not in input file !')

```

```

quit()

def get_job_information_from_jobscrip_flag(content, flag):
    index = [idx for idx, s in enumerate(content) if flag in s][0]
    info = content[index].split(flag, 1)[1]

    return info

#####
# FILE #####
#####

def write_add_string_into_file(file, substring, add, comment = None):
    # open file
    with open(file, 'r') as f:
        # read a list of lines into data
        data = f.readlines()

    try:
        index = [idx for idx, s in enumerate(data) if substring in s][0]
        data[index] = substring + add + '\n'

    except IndexError:
        index = [idx for idx, s in enumerate(data) if '\n' in s][1]
        data.insert(index, '\n')
        data.insert(index + 1, comment)
        data.insert(index + 2, substring + add + '\n')

    # replace line
    with open(file, 'w') as f:
        f.writelines(data)

#####
# TIME #####
#####

def get_time_in_seconds(time):
    # Check time format of time
    if len(time) < 5:
        ## d:hh:mm:ss
        if len(time) == 4:
            timeSeconds = int(time[0])*24*60*60 + int(time[1])*60*60 +
        ↵ int(time[2])*60 + int(time[3])
        ## hh:mm:ss
        elif len(time) == 3:
            timeSeconds = int(time[0])*60*60 + int(time[1])*60 + int(time[2])
        ## mm:ss
        elif len(time) == 2:
            timeSeconds = int(time[0])*60 + int(time[1])
        ## ss
        elif len(time) == 1:
            timeSeconds = int(time[0])
    else:
        print('! Time format is not supported !')
        quit()

    return timeSeconds

```

```

def format_num(time):
    if time < 10:
        return '0' + str(time)
    else:
        return str(time)

def get_time_converted(sec):
    mins, sec = sec // 60, sec % 60
    hours, mins = mins // 60, mins % 60
    days, hours = hours // 24, hours % 24
    weeks, days = days // 7, days % 7

    timeConverted = (str(int(weeks)) + ':' +
                      format_num(int(days)) + ':' + format_num(int(hours)) + ':' +
                      format_num(int(mins)) + ':' + format_num(int(sec)))

    return timeConverted

def time_date():
    e = datetime.datetime.now()
    return '%s-%s-%s' % (format_num(e.year), format_num(e.month),
                         format_num(e.day))

def time_time():
    e = datetime.datetime.now()
    return '%s:%s:%s' % (format_num(e.hour), format_num(e.minute),
                         format_num(e.second))

def time_duration(startTime):
    stop = time.time()
    return get_time_converted(stop - startTime)

#####
# TABLE #####
#####

def table_row_format(content):
    if len(content) == 7:
        cols = [2, 10, 29, 13, 11, 13, 2]
    elif len(content) == 6:
        cols = [2, 10, 29, 13, 24, 2]
    elif len(content) == 5:
        cols = [2, 10, 25, 41, 2]
    elif len(content) == 4:
        cols = [2, 10, 66, 2]
    else:
        cols = [2, 76, 2]

    i, sep = 0, []
    while i < len(cols):
        if i == 0:
            sep.append('o' + (cols[i]-1)*'-')
        elif i == (len(cols)-1):
            sep.append('o' + (cols[i]-1)*'-')
        else:
            sep.append('o' + (cols[i]-1)*'-')
        i += 1
    return sep

```

```

        sep.append((cols[i]-1)*'-' + 'o')
    else:
        sep.append(cols[i]*'-')

    i += 1

    row_format = "".join(["{:<" + str(col) + "}" for col in cols])

    return row_format, sep

def print_table_row(content, output_type = None, time_info = True):

    if isinstance(content[0], list):

        i = 0
        while i < len(content):
            content[i].insert(0, '| ')
            if time_info:
                if output_type == 'header':
                    content[i].append('DATE')
                    content[i].append('TIME')
                    content[i].append('W:DD:HH:MM:SS')
                else:
                    content[i].append(time_date())
                    content[i].append(time_time())
                    content[i].append(time_duration(startTime))
            content[i].insert(len(content[i]), '| ')
            i += 1

        row_format, sep = table_row_format(content[0])

        if output_type == 'header':
            print('\n')
            print(row_format.format(*sep))
            for row in content:
                print(row_format.format(*row))
            print(row_format.format(*sep))

        elif output_type == 'end':
            for row in content:
                print(row_format.format(*row))
            print(row_format.format(*sep))

        else:
            for row in content:
                print(row_format.format(*row))
    else:
        content.insert(0, '| ')
        if time_info:
            if output_type == 'header':
                content.append('DATE')
                content.append('TIME')

```

```
        content.append('W:DD:HH:MM:SS')
    else:
        content.append(time_date())
        content.append(time_time())
        content.append(time_duration(startTime))
    content.insert(len(content), ' | ')

    row_format, sep = table_row_format(content)

    if output_type == 'header':
        print('\n')
        print(row_format.format(*sep))
        print(row_format.format(*content))
        print(row_format.format(*sep))

    elif output_type == 'end':
        print(row_format.format(*content))
        print(row_format.format(*sep))

    else:
        print(row_format.format(*content))

#####
# STATUS #####
# MAIL #####
#  
  
def get_job_status(user):
    jobStatus = subprocess.getoutput(commandJobStatus + user).strip().split()

    return jobStatus

def set_output_type(user):
    jobStatus = get_job_status(user)

    if jobName in jobStatus:
        outputType = 'running'
    else:
        outputType ='no Output'

    return outputType

def send_mail(recipient, subject, body = None):

    recipient = recipient.encode('utf_8')

    subject = "''" + subject + "''"
    subject = subject.encode('utf_8')

    if body == None:
        body = 'For futher information open attachment'

    body = body.encode('utf_8')
```

6 Appendix

```
attachmentPath = folder + '/' + monitorFilename
attachment = attachmentPath.encode('utf_8')

process = subprocess.Popen(['ssh', 'master', '/usr/bin/mailx', '-s',
                           ↳ subject, '-a', attachment, recipient],
                           stdin=subprocess.PIPE)
process.communicate(body)

'''

#####
#          #
#          JOB INFORMATIONS          #
#          #
#####


#####
#          OUTPUT STRING          #
#          #

jobInformations = []

#####
#          START WATCH          #
#          #



startTime = time.time()

#####
#          OUTPUT INIT          #
#          #



print_table_row(['OUTPUT', 'JOB INITIALIZE'], output_type='header')

#####
#          USER          #
#          #



user = os.getlogin()

#####
#          JOB NAME          #
#          #



jobName = jobscripfilename

if len(jobName) > 16:
    jobName = jobName[0:16]

jobInformations.append(['INFO', 'Name', jobName])

#####
#          FILE PATH          #
#          #



folder = os.path.dirname(os.path.abspath(__file__))
path = os.path.dirname(os.path.abspath(__file__)).split(user + '/') [1]

#####
#          MAIL ADDRESS          #
#          #



try:
    emailNotification = True
    jobInformations.append(['INFO', 'E-Mail', emailAddress])
except NameError:
    emailNotification = False
```

```

#####
# JOB SCRIPT #####
#####

try:
    jobsript = [filename for filename in os.listdir('.') if
    ↪ filename.startswith(jobsriptFilename)][0]
    print_table_row(['SUCCESS', 'Found ' + 'jobsript file'])

except IndexError:
    print_table_row(['ERROR', 'No jobsript found'])
    # Send error mail
    if emailNotification:
        send_mail(emailAddress, 'Failed Job ' + jobName)
    quit()

jobsriptContent = open(jobsript, 'r').read().splitlines()

#####
# Timesteps #####
#####

# Number of required timesteps
try:
    nTimestepsRequired = 0

    for i in range(10):
        nTimestepsRequired += get_value_of_variable_from_input_file('./' +
    ↪ inputFilename, inputFlag, i)

    print_table_row(['SUCCESS', 'Found ' + 'input file'], output_type='end',
    ↪ time_info=True)
except FileNotFoundError:
    print_table_row(['ERROR', 'No input file found'], output_type='end')
    # Send error mail
    if emailNotification:
        send_mail(emailAddress, 'Failed Job ' + jobName)
    quit()

jobInformations.append(['INFO', 'Required Timesteps', nTimestepsRequired])

#####
# WALLTIME #####
#####

walltime = get_job_information_from_jobsript_flag(jobsriptContent,
    ↪ walltimeFlag).replace('-', ':').split(':')
walltimeSeconds = get_time_in_seconds(walltime)

jobInformations.append(['INFO', 'Walltime/s', walltimeSeconds])

#####
# BACKUP PATH #####
#####

try:
    # Backup location
    if backupLocation[-1] != '/':
        backupLocation += '/'

```

```

backupPath = backupLocation + path

# Create backup directory if do not exist
if not os.path.exists(backupPath):
    os.makedirs(backupPath)

# BackUp switch
backup = True

jobInformations.append(['INFO', 'Backup Path', backupLocation])

except NameError:
    # BackUp switch
    backup = False

#####
# OUTPUT INFO #####
print_table_row(['OUTPUT', 'JOB INFORMATIONS', 'VALUE'],
               ↪ output_type='header', time_info=False)
print_table_row(jobInformations, output_type='end', time_info=False)

'''

#####
# START/RESTART JOB #####
#           #
#           #
#####

'''

#####
# OUTPUT MONITOR #####
print_table_row(['OUTPUT', 'JOB MONITORING'], output_type='header')

#####
# BEGIN #####
outputType = set_output_type(user)

# read restart file
## If gkw has run requiered timesteps stop already here
while True:
    try:

        identity = []
        # find last output file
        for file in os.listdir():
            if '.sh.o' in file:
                identity.append(file.split('.sh.o')[1])

        nTimestepsCurrent = get_value_of_variable_from_output_file('..' +
        ↪ outputfilename(max(identity)), restartFlag, -1)
        nTimestepsCurrent = int((nTimestepsCurrent // 1e4) * 1e4)

```

```

# Output current timesteps
if outputType == 'running':
    print_table_row(['CONTROL', 'Current Timesteps ' +
    ↪ str(nTimestepsCurrent)]) 

# Check if gkw has run requiered timesteps
if nTimestepsCurrent >= nTimestepsRequired:
    print_table_row(['SUCCESS', 'Stop monitoring'],
    ↪ output_type='end')

# Send end email
if emailNotification:
    send_mail(emailAddress, 'Ended Job ' + jobName)

quit()

# Continue
else:
    print_table_row(['CONTINUE', 'Continue monitoring'])

# Send continue mail
if emailNotification:
    send_mail(emailAddress, 'Continued Job ' + jobName)
break

# Start
except (FileNotFoundException, NameError):
    print_table_row(['STARTING', 'Start monitoring'])

# Making backup
if backup:
    print_table_row(['BACKUP', backupLocation], output_type='end')
    subprocess.run(['rsync', '-a', '', backupPath])

# Send start mail
if emailNotification:
    send_mail(emailAddress, 'Started Job ' + jobName)
break

##### MONITOR ROUTINE #####
while True:

    jobStatus = get_job_status(user)

    # Job running
    if jobName in jobStatus:
        # Job ID Running
        jobStatusNameIndex = [idx for idx, s in enumerate(jobStatus) if
    ↪ jobName in s][0]
        jobID = jobStatus[jobStatusNameIndex -3].split(startOutputFlag)[0]

    # Set output type

```

```

if outputType == 'running':
    print_table_row(['RUNNING', 'Job is executed'])
    outputType = 'no Output'

sleep(sleepTime)

# Job start/restart
else:

    # Check error and making Backup
    while True:
        try:
            outputContent = read_file_to_string('..' +
→ outputfilename(jobID))

            if outputCriteria in outputContent:

                print_table_row(['ERROR', 'NaN Value in surf_dens'])

                quit()

            else:

                # Making backup
                if backup:
                    print_table_row(['BACKUP', backupLocation],
→ output_type='end')
                    subprocess.run(['rsync', '-a', '', backupPath])

                outputType = 'restart'

                break

        # If jobID is not defined or file is not generated
        except (IndexError, FileNotFoundError):
            sleep(30)
        except NameError:
            break

    # Check Timesteps
    try:
        nTimestepsCurrent = get_value_of_variable_from_output_file('..'
→ + outputfilename(jobID), restartFlag, -1)
        nTimestepsCurrent = int((nTimestepsCurrent // 1e4) * 1e4)

        print_table_row(['CONTROL', 'Current Timesteps ' +
→ str(nTimestepsCurrent)])

        # Check if gkw has run requiered timesteps
        if nTimestepsCurrent >= nTimestepsRequired:
            print_table_row(['SUCCESS', 'Stop monitoring'],
→ output_type='end')

```

```
# Send end email
if emailNotification:
    send_mail(emailAddress, 'Ended Job ' + jobName)
break

# write restart timestep in input file
else:
    write_add_string_into_file(restartFilename, restartString,
→ str(nTimestepsCurrent))

except (FileNotFoundException, NameError):
    pass

# Start Job
startOutput = subprocess.check_output([commandJobStarting,
→ jobscriptFilename]).decode('utf-8').replace('\n', '')
jobID = startOutput.split(startOutputFlag)[0]

print_table_row(['STARTING', startOutput])

# Set output type
if outputType == 'restart':
    # Send end email
    if emailNotification:
        send_mail(emailAddress, 'Restart Job ' + jobName)

    sleep(30)

outputType = set_output_type(user)

##### RESTART #####
```

6.4 Simulation parameter

	R/L_T	box size	N_s	$N_{\nu_{ }}$	N_μ	d_{tim}	$k\rho^{\max}$	N_{mod}	N_x
1	6.0	1x1	12	16	6	0.02	1.4	21	83
2	6.0	1x1	12	32	6	0.02	1.4	21	83
3	6.0	1x1	12	48	9	0.02	1.4	21	83
4	6.0	1x1	12	64	9	0.02	1.4	21	83
5	6.0	1x1	16	16	9	0.02	1.4	21	83
6	6.0	1x1	16	32	9	0.02	1.4	21	83
7	6.0	1x1	16	48	6	0.02	1.4	21	83
8	6.0	1x1	16	48	9	0.02	1.4	21	83
9	6.0	1x1	16	48	9	0.025	1.4	21	83
10	6.0	1x1	16	48	9	0.02	0.7	11	83
11	6.0	1x1	16	48	9	0.02	1.4	21	43
12	6.0	1x1	16	48	9	0.02	1.4	21	63
13	6.0	1x1	16	64	6	0.02	1.4	21	83
14	6.0	1x1	16	64	9	0.02	1.4	21	83
15	6.0	2x1	16	48	9	0.02	1.4	21	83
16	6.0	2x2	16	48	9	0.02	1.4	21	83
17	6.0	3x1	16	48	9	0.02	1.4	21	83
18	6.0	3x1.5	16	48	9	0.02	1.4	21	83
19	6.0	3x1.5	16	48	9	0.02	1.4	21	83
20	6.0	3x2.5	16	48	9	0.02	1.4	21	83
21	6.0	3x3	16	48	9	0.02	1.4	21	83
22	6.0	3x5	16	48	9	0.02	1.4	21	83
23	6.0	4x1	16	48	9	0.02	1.4	21	83
24	6.2	2x2	16	64	9	0.02	1.4	21	83
25	6.2	3x3	16	48	9	0.02	1.4	21	83
26	6.3	1x1	16	64	9	0.02	1.4	21	83
27	6.4	3x3	16	48	9	0.02	1.4	21	83

Table 6.1: Parameters of simulations performed for this thesis

	R/L_T	box size	time	timestep	error_index	stable	n_{ZF}	backup
1	6.0	1x1	6000.0	10000		False	0	True
2	6.0	1x1	12000.0	20000		False	0	True
3	6.0	1x1	6000.0	10000		False	0	True
4	6.0	1x1	6000.0	10000		False	0	True
5	6.0	1x1	12000.0	20000		False	0	True
6	6.0	1x1	12000.0	20000		False	0	True
7	6.0	1x1	6000.0	10000		False	0	True
8	6.0	1x1	6000.0	10000		True	1	True
9	6.0	1x1	7125.0	10000		True	1	True
10	6.0	1x1	6000.0	10000		True	1	True
11	6.0	1x1	6000.0	10000		False	0	True
12	6.0	1x1	6000.0	10000		False	0	True
13	6.0	1x1	12000.0	20000		True	1	True
14	6.0	1x1	12000.0	20000		True	1	True
15	6.0	2x1	18000.0	30000		True	2	True
16	6.0	2x2	18000.0	30000		True	2	True
17	6.0	3x1	48000.0	80000		True	3	True
18	6.0	3x1.5	16840.3	28068		True	4	True
19	6.0	3x1.5	14170.0	23618	23618	True	3	True
20	6.0	3x2.5	6000.0	10000		True	3, 4	True
21	6.0	3x3	7847.0	13083		True	4	True
22	6.0	3x5	4769.0	7958		True	4	True
23	6.0	4x1	30000.0	50000	47084	True	4	True
24	6.2	2x2	6000.0	10000		True	2	True
25	6.2	3x3	14682.7	24475	14473-16132	True	3	True
26	6.3	1x1	6000.0	10000		False	0	True
27	6.4	3x3	12000.0	20000		False	0	True

Table 6.2: Time steps with index which should be exclude in data analysis. Additional the stable column which indicates if turbulence subdued in simulation with the corresponding zonal flow mode number n_{ZF} (0 stands for turbulence not stable) and if the data is saved on the NAS of TPV (backup column)

	path
1	./data/S6_rlt6.0/boxsize1x1/Ns12/Nvpar16/Nmu6
2	./data/S6_rlt6.0/boxsize1x1/Ns12/Nvpar32/Nmu6
3	./data/S6_rlt6.0/boxsize1x1/Ns12/Nvpar48/Nmu9
4	./data/S6_rlt6.0/boxsize1x1/Ns12/Nvpar64/Nmu9
5	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar16/Nmu9
6	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar32/Nmu9
7	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar48/Nmu6
8	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar48/Nmu9
9	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar48/Nmu9/dtim0.025
10	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar48/Nmu9/krhomax0.70/Nmod11
11	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar48/Nmu9/Nx43
12	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar48/Nmu9/Nx63
13	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar64/Nmu6
14	./data/S6_rlt6.0/boxsize1x1/Ns16/Nvpar64/Nmu9
15	./data/S6_rlt6.0/boxsize2x1/Ns16/Nvpar48/Nmu9
16	./data/S6_rlt6.0/boxsize2x2/Ns16/Nvpar48/Nmu9
17	./data/S6_rlt6.0/boxsize3x1/Ns16/Nvpar48/Nmu9
18	./data/S6_rlt6.0/boxsize3x1.5/Ns16/Nvpar48/Nmu9
19	./data/S6_rlt6.0/boxsize3x1.5/Ns16/Nvpar48/Nmu9/Broken
20	./data/S6_rlt6.0/boxsize3x2.5/Ns16/Nvpar48/Nmu9
21	./data/S6_rlt6.0/boxsize3x3/Ns16/Nvpar48/Nmu9
22	./data/S6_rlt6.0/boxsize3x5/Ns16/Nvpar48/Nmu9
23	./data/S6_rlt6.0/boxsize4x1/Ns16/Nvpar48/Nmu9
24	./data/S6_rlt6.2/boxsize2x2/Ns16/Nvpar64/Nmu9
25	./data/S6_rlt6.2/boxsize3x3/Ns16/Nvpar48/Nmu9
26	./data/S6_rlt6.3/boxsize1x1/Ns16/Nvpar64/Nmu9
27	./data/S6_rlt6.4/boxsize3x3/Ns16/Nvpar48/Nmu9

Table 6.3: Data location for each simulation

Size convergence of the $E \times B$ staircase pattern in flux tube simulations of ion temperature gradient driven turbulence

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(Dated: 21 May 2023)

The radial size convergence of the $E \times B$ staircase pattern is addressed in local gradient-driven flux tube simulations of ion temperature gradient (ITG) driven turbulence. It is shown that a mesoscale pattern size of $\sim 57 - 76\rho$ is inherent to ITG driven turbulence with Cyclone Base Case parameters in the local limit.

Ion temperature gradient driven turbulence close to marginal stability exhibits zonal flow pattern formation on mesoscales, so-called $E \times B$ staircase structures¹. Such pattern formation has been observed in local gradient-driven flux-tube simulations^{2–4} as well as global gradient-driven^{5–7} and global flux-driven^{1,8–11} studies. In global studies, spanning a larger fraction of the minor radius, multiple radial repetitions of staircase structures are usually observed, with a typical pattern size of several ten Larmor radii. By contrast, in the aforementioned local studies the radial size of $E \times B$ staircase structures is always found to converge to the radial box size of the flux tube domain. The above observations lead to the question: *Does the basic pattern size always converges to the box size, or is there a typical mesoscale size inherent to staircase structures also in a local flux-tube description?* The latter case would imply that it is not necessarily global physics, i.e., profile effects, that set (i) the radial size of the $E \times B$ staircase pattern and (ii) the scale of avalanche-like transport events. These transport events are usually restricted to $E \times B$ staircase structures and considered as a nonlocal transport mechanism¹. In this brief communication the above question is addressed through a box size convergence scan of the same cases close to the nonlinear threshold for turbulence generation as studied in Ref. 2.

The gyrokinetic simulations are performed with the nonlinear flux tube version of Gyrokinetic Workshop (GKW)¹² with adiabatic electron approximation. In agreement with Ref. 2, Cyclone Base Case (CBC) like parameters are chosen with an inverse background temperature gradient length $R/L_T = 6.0$ and circular concentric flux surfaces. The numerical resolution is compliant to the "Standard resolution with 6th order (S6)" set-up of the aforementioned reference, with a somewhat lowered number of parallel velocity grid points. It has been carefully verified that this modification preserves the same physical outcome as the original study. A summary of the numerical parameters is given in Tab. I and for more details about the definition of individual quantities the reader is referred to Ref. 2 and 12.

	N_m	N_x	N_s	$N_{V\parallel}$	N_μ	D	v_d	$D_{V\parallel}$	D_x	D_y	Order	$k_y\rho$	$k_x\rho$
S6	21	83	16	48	9	1	$ v_\parallel $	0.2	0.1	0.1	6	1.4	2.1

TABLE I: Resolution used in this paper for further information the author links to Ref. 2.

In the following the box size is increased relative to the standard box size $(L_x, L_y) = (76, 89)\rho$ in the radial and binormal direction. Here, x is the radial coordinate that labels the flux surfaces normalized by the thermal Larmor radius ρ , y labels the field lines and is an approximate binormal coordinate. Together with the coordinate s which parameterizes the length along the field lines and is referred to as the parallel coordinate these quantities form the Hamada coordinates¹³. The increased box sizes are indicated by the real parameter N_R for radial and N_B for the binormal direction with the nomenclature $N_R \times N_B$ throughout this work. Note that, the number of modes in the respective direction, i.e., N_x and N_m , respectively, is always adapted accordingly to retain a spatial resolution compliant to the standard resolution [Tab. I] and standard box size.

The $E \times B$ staircase pattern is manifest as radial structure formation in the $E \times B$ shearing rate defined by^{2,14,15}

$$\omega_{E \times B} = \frac{1}{2} \frac{\partial^2 \langle \phi \rangle}{\partial x^2}, \quad (1)$$

where $\langle \phi \rangle$ is the zonal electrostatic potential normalized by $\rho_* T/e$ ($\rho_* = \rho/R$ is the thermal Larmor radius normalized with the major radius R , T is the temperature, e is the elementary charge). The zonal potential is calculated from the electrostatic potential ϕ on the two-dimensional x - y -plane at the low field side according to⁴

$$\langle \phi \rangle = \frac{1}{L_y} \int_0^{L_y} dy \phi(x, y, s=0). \quad (2)$$

The $E \times B$ shearing rate $\omega_{E \times B}$ is the radial derivative of the advecting zonal flow velocity^{16,17} and quantifies the zonal flow induced shearing of turbulent structures^{16,18,19}.

Consistent with Ref. 2 the turbulence level is quantified by the turbulent heat conduction coefficient χ , which is normalized by $\rho^2 v_{th}/R$ ($v_{th} = \sqrt{2T/m}$ is the thermal velocity and m is the mass). Furthermore, quantities ρ , R , T , v_{th} and m are referenced quantities from Ref. 2 and 12.

^{a)}Repository of this work:

<https://github.com/ManeLippert/Bachelorthesis-Shearingrate-Convergence>

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In order to diagnose the temporal evolution of the staircase pattern and to obtain an estimate of its amplitude the radial Fourier transform of the $E \times B$ shearing rate is considered. It is defined by

$$\omega_{E \times B} = \sum_{k_{ZF}} \hat{\omega}_{E \times B}(k_{ZF}, t) \exp(i k_{ZF} x), \quad (3)$$

where $\hat{\omega}_{E \times B}$ is the complex Fourier coefficient and $k_{ZF} = 2\pi n_{ZF}/L_x$ defines the zonal flow wave vector with the zonal flow mode number n_{ZF} ranging in $-(N_x - 1)/2 \leq n_{ZF} \leq (N_x - 1)/2$. Based on the definitions above, the shear carried by the zonal flow mode with wave vector k_{ZF} is defined by $|\hat{\omega}_{E \times B}|_{n_{ZF}} = 2|\hat{\omega}_{E \times B}(k_{ZF}, t)|$. In general, the zonal flow mode that dominates the $E \times B$ staircase pattern, also referred to as the *basic mode* of the pattern in this work, exhibits the maximum amplitude in the spectrum $|\hat{\omega}_{E \times B}|_{n_{ZF}}$.

In the first test the radial box size is increased while the binormal box size is kept fixed to the standard size. The scan covers the realizations $N_R \times N_B \in [1 \times 1, 2 \times 1, 3 \times 1, 4 \times 1]$. Each realization exhibits an initial quasi-stationary turbulent phase and a second final² phase with almost suppressed turbulence [Fig. 1 (a)]. The latter state is indicative for the presence of a fully developed staircase pattern as depicted in Fig. 2. This type of structure is characterized by intervals of almost constant shear with alternating sign satisfying the Waltz criterion $|\omega_{E \times B}| \approx \gamma^{17,20}$ (γ is the growth rate of the most unstable linear ITG driven Eigenmode), connected by steep flanks where $\omega_{E \times B}$ crosses zero. Fig. 2(a) shows a striking repetition of the staircase structure, with the number of repetitions equal to N_R . Hence, the basic size of the pattern not only converges with increasing radial box size, the converged radial size turns out to at least roughly agree with

the standard radial box size of Ref. 2.

Due to the lack of a substantial turbulent drive in the final suppressed state no further zonal flow evolution is observed [Fig. 1 (b)] and one might critically ask whether the structures shown in Fig. 2 represent the real converged pattern in a statistical sense. Note that in the 3×1 case the initial quasi-stationary turbulent state extends up to a few $\sim 10^4 R/v_{th}$. During this period the zonal flow mode with $n_{ZF} = 3$, i.e., the mode that dominates the staircase pattern in final suppressed phase, undergoes a long-term evolution with a typical time scale of several $\sim 10^3 R/v_{th}$. Hence, several of such cycles are covered by the initial turbulent phase, which is evident from the occurrence of phases with reduced amplitude around $t \approx 8000 R/v_{th}$ and $t \approx 18000 R/v_{th}$. It is the $n_{ZF} = 4$ zonal flow mode, i.e., the next shorter radial scale mode, that dominates the shear spectrum $|\hat{\omega}_{E \times B}|_{n_{ZF}}$ in the latter two phases (not shown). This demonstrates a competition between the $n_{ZF} = 3$ and $n_{ZF} = 4$ modes. Most importantly, no secular growth of the $n_{ZF} = 1$ (box scale) zonal flow mode is observed during the entire quasi-stationary turbulent phase [Fig. 1 (b) dotted line]. The above discussion indicates that although the $n_{ZF} = 3, 4$ zonal modes compete, the pattern scale does not converge to the radial box scale but rather to a mesoscale of $\sim 57 - 76 \rho$ (i.e., $n_{ZF} = 4, 3$ in the 3×1 case).

Since the radially elongated simulation domain might inhibit the development of isotropic turbulent structures, in the second test the radial and binormal box size is increased simultaneously. This scan covers the realizations $N_R \times N_B \in [1 \times 1, 2 \times 2, 3 \times 3]$. Interestingly, suppression of the turbulence by the emergence of a fully developed staircase pattern always occurs after $\sim 1000 R/v_{th}$ [Fig. 3], i.e., significantly faster compared to the 3×1 and 4×1 realizations. As shown in Fig. 2 (b) also this test confirms the convergence of the staircase pattern size to a typical mesoscale that is distinct from the radial box size in the $N_R > 1$ realizations.

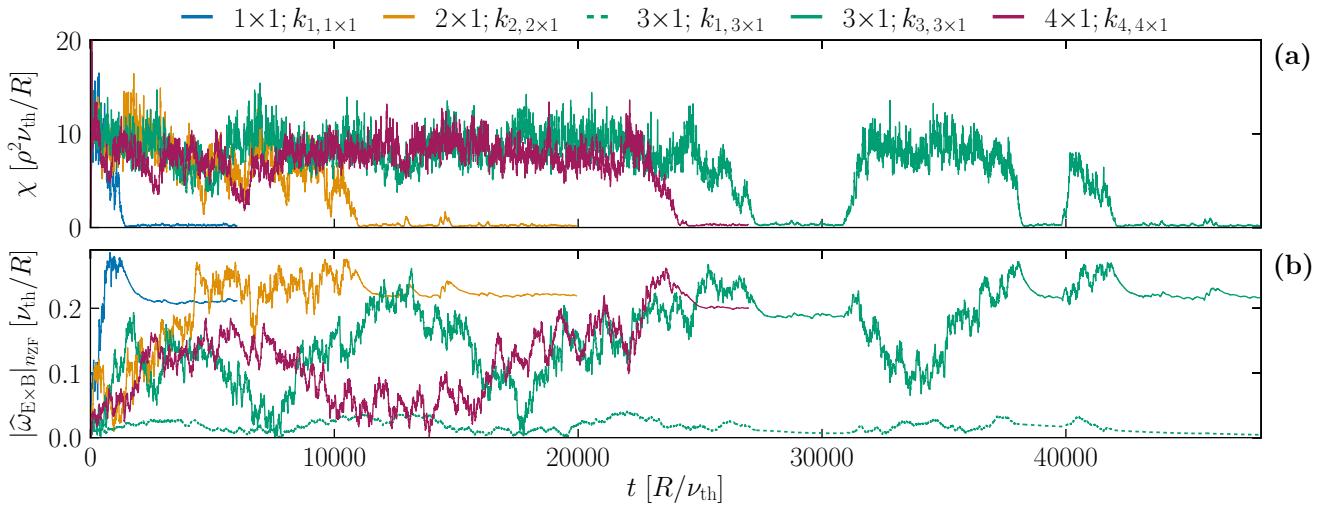


FIG. 1: (a) Time traces of the heat conduction coefficient χ for $R/L_T = 6.0$ for radial increased box sizes
(b) Time traces of $|\hat{\omega}_{E \times B}|_{n_{ZF}}$ for radial increased box sizes

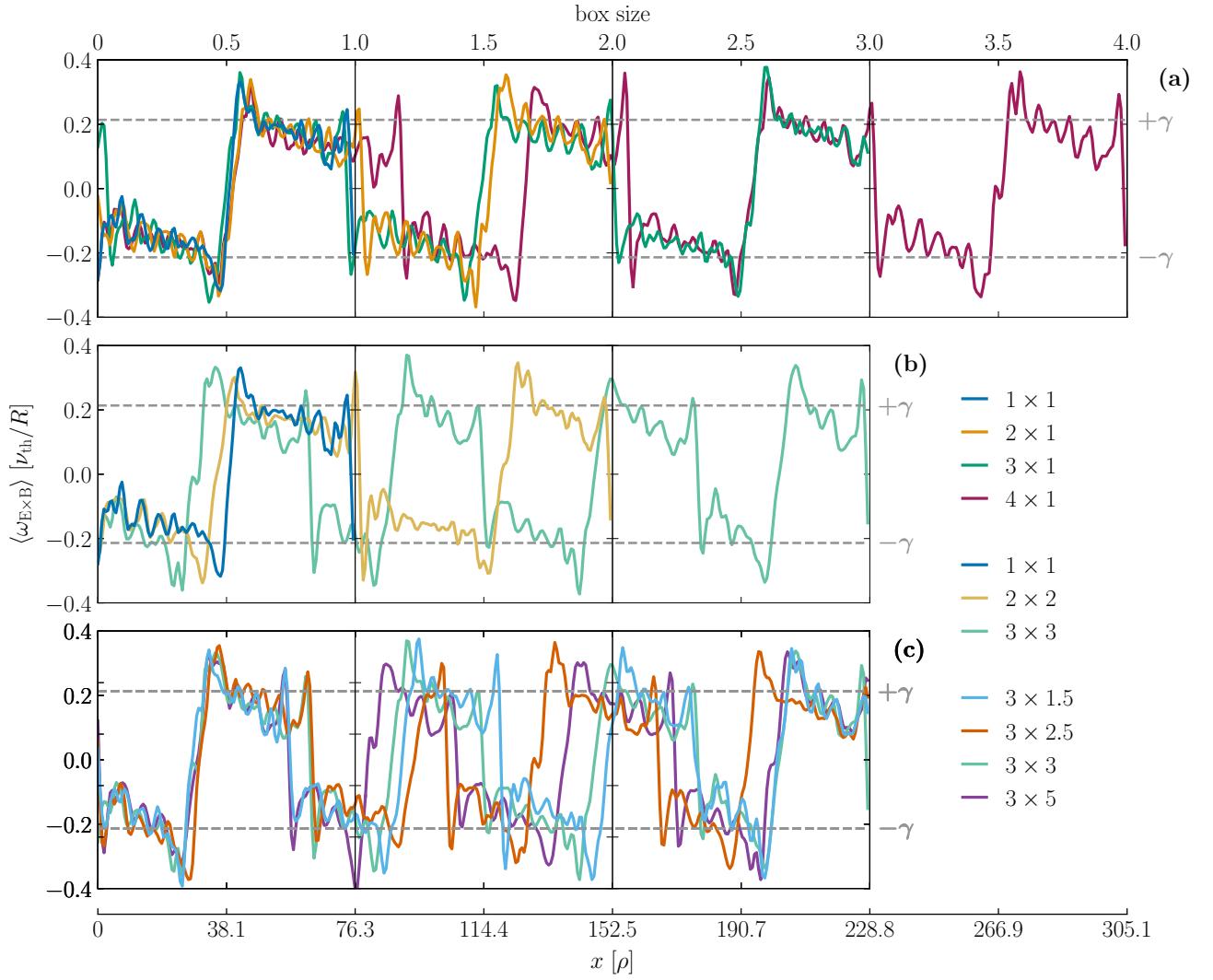


FIG. 2: Comparison of shearing rate $\omega_{E \times B}$ for each box sizes scan averaged over given time interval and the growth rate $\pm\gamma$ of the most unstable linear ITG driven Eigenmode. The staircase structures are radially shifted with respect to each other till alignment for better visibility.

(a) **radial:** $t_{1 \times 1} \in [2000, 5000]$, $t_{2 \times 1} \in [15000, 18000]$, $t_{3 \times 1} \in [43000, 45000]$, $t_{4 \times 1} \in [26000, 28000]$

(b) **isotropic:** $t_{1 \times 1} \in [2000, 5000]$, $t_{2 \times 2} \in [2000, 3000]$, $t_{3 \times 3} \in [2000, 3000]$

(c) **binormal:** $t_{3 \times 1.5} \in [2000, 3000]$, $t_{3 \times 2.5} \in [2000, 3000]$, $t_{3 \times 3} \in [2000, 3000]$, $t_{3 \times 5} \in [1000, 3000]$

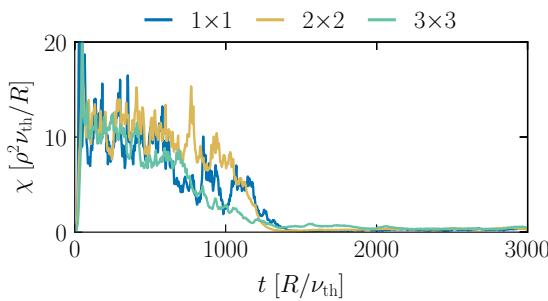


FIG. 3: Time traces of the heat conduction coefficient χ for $R/L_T = 6.0$ for isotropic increased box sizes

By contrast to the radial box size scan the 3×3 realization shows a stationary pattern with four repetitions of the fully developed staircase structure, i.e., a somewhat smaller pattern size. Whether this is related to a possible pattern size dependence on the binormal box size or to the competition between patterns with the two sizes $\lambda \in [57, 76]\rho$ as observed in the first test is addressed in the next paragraph.

In a third test the binormal box size is varied with the radial box size fixed to $N_R = 3$. This test covers the realizations $N_R \times N_B \in [3 \times 1.5, 3 \times 2.5, 3 \times 3, 3 \times 5]$. As in the isotropic scan the turbulence subdued and a fully developed staircase pattern forms after $\sim 2000 R/v_{th}$ [Fig. 4]. The convergence of staircase pattern can be seen in Fig. 2(c) and confirms again a size of a typical mesoscale. Fig. 2(c) also confirms that indeed a competition between patterns with two sizes $\lambda \in [57, 76]\rho$ causing the different results for 3×1 and 3×3 . The zonal flow mode number varies between $n_{ZF} = 3, 4$ which can be seen in Fig. 2(c) in the 3×2.5 realization. The staircase structure has a pattern between 3 and 4 repetitions which get represented in the second repetition with no significant plateau at positive shear. Instead the pattern returns immediately after reaching the maximum shear ($+\gamma$) to the minimum shear ($-\gamma$) of the third repetition in a steep flank. The Fourier analysis of this case yields no definitely basic mode rather two dominating modes with $n_{ZF} = 3, 4$ with a fraction of the maximum amplitude $|\hat{\omega}_{E \times B}|_{n_{ZF}}$ each (not shown).

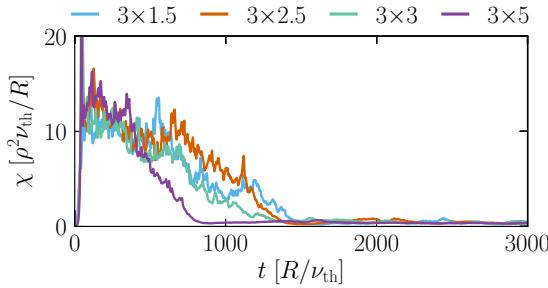


FIG. 4: Time traces of the heat conduction coefficient χ for $R/L_T = 6.0$ for binormal increased box sizes

In the final test the inverse background temperature gradient length R/L_T is varied at fixed 3×3 box size. Since suppression of turbulence usually occurs at later times when approaching the finite heat flux threshold from below², the analysis aims to lengthen the phase during which the zonal flow varies in time due to turbulent Reynolds stresses. This scan covers realizations with $R/L_T \in [6.0, 6.2, 6.4]$. In the case of $R/L_T = 6.2$ turbulence suppression is observed for $t > 11000 R/v_{th}$, while stationary turbulence during the entire simulation time trace of $12000 R/v_{th}$ is found for $R/L_T = 6.4$. The finite heat flux threshold, hence, is $R/L_T|_{\text{finite}} = 6.3 \pm 0.1$ in accordance to Ref. 2. Although the initial quasi-stationary turbulence in the former case is significantly longer compared to the $R/L_T = 6.2$ realization discussed in the second test, a stationary pattern with basic zonal flow mode $n_{ZF} = 3$ establishes. Again, the $n_{ZF} = 1$ (box scale) zonal flow mode does not grow secularly during the entire turbulent phase. Also, this test confirms the statistical soundness of the converged pattern size of $\sim 57 - 76\rho$.

Through careful tests this brief communication confirms the radial size convergence of the $E \times B$ staircase pattern in local gyrokinetic flux tube simulations of ion temperature gradient (ITG) driven turbulence. A mesoscale pattern size of $\sim 57 - 76\rho$ is found to be intrinsic to ITG driven turbulence for Cyclone Base Case parameters. This length scale is somewhat larger compared to results from global studies with finite ρ_* , which report of a few $10\rho^1$, and has to be considered the proper mesoscale in the local limit $\rho_* \rightarrow 0$. The occurrence of this mesoscale implies that non-locality, in terms of Ref. 1, is inherent to ITG driven turbulence, since avalanches are spatially organized by the $E \times B$ staircase pattern^{1,2,5,14}.

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

- ¹G. Dif-Pradalier, P. H. Diamond, V. Grandgirard, Y. Sarazin, J. Abiteboul, X. Garbet, P. Ghendrih, A. Strugarek, S. Ku, and C. S. Chang, Phys. Rev. E **82**, 025401 (2010).
- ²A. G. Peeters, F. Rath, R. Buchholz, Y. Camenen, J. Candy, F. J. Casson, S. R. Grosshauser, W. A. Hornsby, D. Sintzhi, and A. Weikl, Phys. Plasmas **23**, 082517 (2016).
- ³A. Weikl, A. G. Peeters, F. Rath, S. R. Grosshauser, R. Buchholz, W. A. Hornsby, F. Seiferling, and D. Sintzhi, Phys. Plasmas **24**, 102317 (2017).
- ⁴F. Rath, A. G. Peeters, and A. Weikl, Phys. Plasmas **28**, 072305 (2021).
- ⁵B. F. McMillan, S. Jolliet, T. M. Tran, L. Villard, A. Bottino, and P. Angelino, Physics of Plasmas **16**, 022310 (2009).
- ⁶L. Villard, P. Angelino, A. Bottino, S. Brunner, S. Jolliet, B. F. McMillan, T. M. Tran, and T. Vernay, Plasma Physics and Controlled Fusion **55**, 074017 (2013).
- ⁷J. Seo, H. Jhang, and J.-M. Kwon, Physics of Plasmas **29**, 052502 (2022).
- ⁸G. Dif-Pradalier, G. Hornung, P. Ghendrih, Y. Sarazin, F. Clairet, L. Vermaire, P. H. Diamond, J. Abiteboul, T. Cartier-Michaud, C. Ehrlacher, D. Estève, X. Garbet, V. Grandgirard, O. D. Gürcan, P. Hennequin, Y. Kosuga, G. Latu, P. Maget, P. Morel, C. Norscini, R. Sabot, and A. Storelli, Phys. Rev. Lett. **114**, 085004 (2015).
- ⁹W. Wang, Y. Kishimoto, K. Imadera, H. Liu, J. Li, M. Yagi, and Z. Wang, Nuclear Fusion **60**, 066010 (2020).
- ¹⁰Y. J. Kim, K. Imadera, Y. Kishimoto, and T. S. Hahm, Journal of the Korean Physical Society **81**, 636 (2022).
- ¹¹Y. Kishimoto, K. Imadera, A. Ishizawa, W. Wang, and J. Q. Li, Philosophical Transactions of the Royal Society A: Mathematical, Physical and Engineering Sciences **381**, 20210231 (2023).
- ¹²A. G. Peeters, Y. Camenen, F. J. Casson, W. A. Hornsby, A. P. Snodin, D. Sintzhi, and G. Szepesi, Comput. Phys. Commun. **180**, 2650 (2009).
- ¹³S. Hamada, Kakuyugo Kenkyu **1**, 542 (1958).
- ¹⁴F. Rath, A. G. Peeters, R. Buchholz, S. R. Grosshauser, P. Migliano, A. Weikl, and D. Sintzhi, “Comparison of gradient and flux driven gyrokinetic turbulent transport,” Phys. Plasmas **23**, 052309 (2016).
- ¹⁵M. J. Pueschel, M. Kammerer, and F. Jenko, Physics of Plasmas **15**, 102310 (2008).
- ¹⁶T. S. Hahm and K. H. Burrell, Phys. Plasmas **2**, 1648–1651 (1995).
- ¹⁷R. E. Waltz, R. L. Dewar, and X. Garbet, Phys. Plasmas **5**, 1784–1792 (1998).
- ¹⁸H. Biglari, P. H. Diamond, and P. W. Terry, Phys. Fluids B: Plasma Physics **2**, 1–4 (1990).
- ¹⁹K. H. Burrell, Phys. Plasmas **4**, 1499–1518 (1997).
- ²⁰R. E. Waltz, G. D. Kerbel, and J. Milovich, Phys. Plasmas **1**, 2229 (1994).

7

Bibliography

- [1] 2018 nohup. URL <https://wiki.ubuntuusers.de/nohup/> – Zugriffsdatum: 2023-04-15.
- [2] 2021 Screen. URL <https://wiki.ubuntuusers.de/Screen/> – Zugriffsdatum: 2023-04-15.
- [3] BARTON, JUSTIN E., WEHNER, WILLIAM P., SCHUSTER, EUGENIO, FELICI, FEDERICO & SAUTER, OLIVIER 2015 Simultaneous closed-loop control of the current profile and the electron temperature profile in the tcv tokamak. URL <https://ieeexplore.ieee.org/document/7171844>.
- [4] BEER, M.A. 1994 *Gyrofluid models of turbulent transport in tokamaks*. PhD thesis, Princeton University.
- [5] BIGLARI, H., DIAMOND, P. H. & TERRY, P. W. 1990 Influence of sheared poloidal rotation on edge turbulence. *Phys. Fluids B: Plasma Physics* **2** (1), 1–4.
- [6] BRIZARD, A. J. & HAHM, T. S. 2007 Foundations of nonlinear gyrokinetic theory. *Rev. Mod. Phys.* **79**, 421–468.
- [7] BURRELL, K. H. 1997 Effects of $e \times b$ velocity shear and magnetic shear on turbulence and transport in magnetic confinement devices. *Phys. Plasmas* **4** (5), 1499–1518.
- [8] CARY, JOHN R. 1981 Lie transform perturbation theory for Hamiltonian systems. *Physics Reports* **79** (2), 129–159.
- [9] CARY, JOHN R & LITTLEJOHN, ROBERT G 1983 Noncanonical Hamiltonian mechanics and its application to magnetic field line flow. *Annals of Physics* **151** (1), 1–34.

- [10] CASSON, F.J. 2011 *Turbulent transport in rotating tokamak plasmas*. PhD thesis, University of Warwick.
- [11] COPPI, B., ROSENBLUTH, M. N. & SAGDEEV, R. Z. 1967 Instabilities due to temperature gradients in complex magnetic field configurations. *The Physics of Fluids* **10** (3), 582–587.
- [12] COWLEY, S. C., KULSRUD, R. M. & SUDAN, R. 1991 Considerations of ion-temperature-gradient-driven turbulence. *Physics of Fluids B: Plasma Physics* **3** (10), 2767–2782.
- [13] DANNERT, T. 2005 *Gyrokinetische Simulation von Plasmaturbulenz mit gefangenen Teilchen und Elektromagnetischen Effekten*. PhD thesis, Technische Universität München.
- [14] DIAMOND, P. H., ITOH, S.-I., ITOH, K. & HAHM, T. S. 2005 Zonal flows in plasma—a review. *Plasma Phys. Controlled Fusion* **47**, R35.
- [15] DIAMOND, P. H. & KIM, Y.-B. 1991 Theory of mean poloidal flow generation by turbulence. *Physics of Fluids B: Plasma Physics* **3** (7), 1626–1633.
- [16] DIF-PRADALIER, G., DIAMOND, P. H., GRANDGIRARD, V., SARAZIN, Y., ABITEBOUL, J., GARBET, X., GHENDRIH, Ph., STRUGAREK, A., KU, S. & CHANG, C. S. 2010 On the validity of the local diffusive paradigm in turbulent plasma transport. *Phys. Rev. E* **82**, 025401.
- [17] DIF-PRADALIER, G., HORNUNG, G., GHENDRIH, Ph., SARAZIN, Y., CLAIRET, F., VERMARE, L., DIAMOND, P. H., ABITEBOUL, J., CARTIER-MICHAUD, T., EHRLACHER, C., ESTÈVE, D., GARBET, X., GRANDGIRARD, V., GÜRCAN, Ö. D., HENNEQUIN, P., KOSUGA, Y., LATU, G., MAGET, P., MOREL, P., NORSCINI, C., SABOT, R. & STORELLI, A. 2015 Finding the elusive $\mathbf{E} \times \mathbf{B}$ staircase in magnetized plasmas. *Phys. Rev. Lett.* **114**, 085004.
- [18] DIMITS, A. M., BATEMAN, G., BEER, M. A., COHEN, B. I., DORLAND, W., HAMMETT, G. W., KIM, C., KINSEY, J. E., KOTSCHENREUTHER, M., KRITZ, A. H., LAO, L. L., MANDREKAS, J., NEVINS, W. M., PARKER, S. E., REDD, A. J., SHUMAKER, D. E., SYDORA, R. & WEILAND, J. 2000 Comparisons and physics basis of tokamak transport models and turbulence simulations. *Phys. of Plasmas* **7** (3), 969–983.
- [19] DUBIN, DANIEL H. E., KROMMES, JOHN A., OBERMAN, C. & LEE, W. W. 1983 Nonlinear gyrokinetic equations. *The Physics of Fluids* **26** (12), 3524–3535.
- [20] GARBET, X., IDOMURA, Y., VILLARD, L. & WATANABE, T. H. 2010 Gyrokinetic simulations of turbulent transport. *Nuclear Fusion* **50** (4).
- [21] HAHM, T. S. & BURRELL, K. H. 1995 Flow shear induced fluctuation suppression in finite aspect ratio shaped tokamak plasma. *Phys. Plasmas* **2** (5), 1648–1651.

- [22] HAMADA, S. 1958 *Kakuyugo Kenkyu* **1**, 542.
- [23] HAMMETT, GREG ???? The Ion Temperature Gradient (ITG) Instability. CM-PD/CMSO Winter School, UCLA, 1/09/2009.
- [24] HASEGAWA, AKIRA & MIMA, KUNIOKI 1978 Pseudo-three-dimensional turbulence in magnetized nonuniform plasma. *The Physics of Fluids* **21** (1), 87–92.
- [25] H.ISLIKER, PISOKAS, TH., STRINTZI, D. & VLAHOS, L. 2010 A self-organized criticality model for ion temperature gradient mode driven turbulence in confined plasma. *Physics of Plasmas* **17**.
- [26] HORTON, W. 1999 Drift waves and transport. *Rev. Mod. Phys.* **71**, 735–778.
- [27] IDOMURA, Y., URANO, H., AIBA, N. & TOKUDA, S. 2009 Study of ion turbulent transport and profile formations using global gyrokinetic full- f vlasov simulation. *Nuclear Fusion* **49** (6), 065029.
- [28] KIM, Y. J., IMADERA, K., KISHIMOTO, Y. & HAHM, T. S. 2022 Transport events and $\mathbf{e} \times \mathbf{b}$ staircase in flux-driven gyrokinetic simulation of ion temperature gradient turbulence. *Journal of the Korean Physical Society* **81**, 636.
- [29] KISHIMOTO, Y., IMADERA, K., ISHIZAWA, A., WANG, W. & LI, J. Q. 2023 Characteristics of constrained turbulent transport in flux-driven toroidal plasmas. *Philosophical Transactions of the Royal Society A: Mathematical, Physical and Engineering Sciences* **381** (2242), 20210231.
- [30] KROMMES, JOHN A. 2012 The Gyrokinetic Description of Microturbulence in Magnetized Plasmas. *Annual Review of Fluid Mechanics* .
- [31] KROMMES, JOHN A. & KIM, CHANG-BAE 2000 Interactions of disparate scales in drift-wave turbulence. *Phys. Rev. E* **62**, 8508–8539.
- [32] LIPPERT, M. 2022 torque_monitor.py. URL https://github.com/ManeLippert/Bachelorthesis-Shearingrate-Convergence/blob/main/python/torque_monitor.py – Zugriffsdatum: 2023-04-14.
- [33] LIPPERT, M. & RATH, F. 2023 slurm_monitor.py. URL https://bitbucket.org/gkw/gkw/src/develop/python/slurm_monitor.py – Zugriffsdatum: 2023-04-12.
- [34] LIPPERT, M., RATH, F. & PEETERS, A. G. 2023 Size convergence of the $\mathbf{E} \times \mathbf{B}$ staircase pattern in flux tube simulations of ion temperature gradient driven turbulence. *Phys. of Plasmas* **7** (3), 969–983.
- [35] MAEYAMA, S., ISHIZAWA, A., WATANABE, T.-H., NAKATA, M., MIYATO, N., YAGI, M. & IDOMURA, Y. 2014 Comparison between kinetic-ballooning-mode-driven turbulence and ion-temperature-gradient-driven turbulence. *Physics of Plasmas* **21** (5), 052301.

- [36] MAKWANA, K. D., TERRY, P. W., PUESCHEL, M. J. & HATCH, D. R. 2014 Subdominant modes in zonal-flow-regulated turbulence. *Phys. Rev. Lett.* **112**, 095002.
- [37] McMILLAN, B. F., JOLLIET, S., TRAN, T. M., VILLARD, L., BOTTINO, A. & ANGELINO, P. 2009 Avalanche-like bursts in global gyrokinetic simulations. *Physics of Plasmas* **16** (2), 022310.
- [38] MITTENDORF, J., SCHOBERT, B. & MÜLLER, D. 2023 Rmhd-code. URL https://bitbucket.org/astro_bayreuth/rmhdcod – Zugriffsdatum: 2023-04-14.
- [39] MÜLLER, D. 2023 Numerical simulations of exor events in protoplanetary disks: Numerical stability and growth of ring structures in the surface density. Bachelorthesis, University of Bayreuth.
- [40] NAKATA, M., WATANABE, T.-H. & SUGAMA, H. 2012 Nonlinear entropy transfer via zonal flows in gyrokinetic plasma turbulence. *Phys. Plasmas* **19**, 022303.
- [41] PEETERS, A. G., CAMENEN, Y., CASSON, F. J., HORNSBY, W. A., SNODIN, A. P., STRINTZI, D. & SZEPESI, G. 2009 The nonlinear gyro-kinetic flux tube code gkw. *Comput. Phys. Commun.* **180**, 2650.
- [42] PEETERS, A. G., RATH, F., BUCHHOLZ, R., CAMENEN, Y., CANDY, J., CASSON, F. J., GROSSHAUSER, S. R., HORNSBY, W. A., STRINTZI, D. & WEIKL, A. 2016 Gradient-driven flux-tube simulations of ion temperature gradient turbulence close to the non-linear threshold. *Phys. Plasmas* **23** (8), 082517.
- [43] PUESCHEL, M. J., KAMMERER, M. & JENKO, F. 2008 Gyrokinetic turbulence simulations at high plasma beta. *Physics of Plasmas* **15** (10), 102310.
- [44] RATH, F., PEETERS, A. G., BUCHHOLZ, R., GROSSHAUSER, S. R., MIGLIANO, P., WEIKL, A. & STRINTZI, D. 2016 Comparison of gradient and flux driven gyro-kinetic turbulent transport. *Phys. Plasmas* **23** (5), 052309.
- [45] RATH, F., PEETERS, A. G. & WEIKL, A. 2021 Analysis of zonal flow pattern formation and the modification of staircase states by electron dynamics in gyrokinetic near marginal turbulence. *Phys. Plasmas* **28** (7), 072305.
- [46] RUDAKOV, L.I. & SAGDEEV, R.Z. 1961 On the instability of a nonuniform rarefied plasma in a strong magnetic field. *Dokl. Akad. Nauk. SSSR* **138** (3), 581–583.
- [47] SCHELTER, DR.RER.NAT. INGO 2016 btrzx2 (2016). URL https://www.bzhpc.uni-bayreuth.de/de/keylab/Cluster/btrzx2_page/index.html – Zugriffsdatum: 2023-04-14.
- [48] SCHELTER, DR.RER.NAT. INGO 2020 btrzx1 (2020). URL https://www.bzhpc.uni-bayreuth.de/de/keylab/Cluster/btrzx1_page/index.html – Zugriffsdatum: 2023-04-12.

- [49] SEO, JANGHOON, JHANG, HOGUN & KWON, JAE-MIN 2022 Effects of light impurities on zonal flow activities and turbulent thermal transport. *Physics of Plasmas* **29** (5), 052502.
- [50] STROTH, ULRICH 2011 *Plasmaphysik*. Wiesbaden: Viewg+Teubner.
- [51] VILLARD, L, ANGELINO, P, BOTTINO, A, BRUNNER, S, JOLLIET, S, McMILLAN, B F, TRAN, T M & VERNAY, T 2013 Global gyrokinetic ion temperature gradient turbulence simulations of iter. *Plasma Physics and Controlled Fusion* **55** (7), 074017.
- [52] WALTZ, R. E., DEWAR, R. L. & GARBET, X. 1998 Theory and simulation of rotational shear stabilization of turbulence. *Phys. Plasmas* **5** (5), 1784–1792.
- [53] WALTZ, R. E., KERBEL, G. D. & MILOVICH, J. 1994 Toroidal gyro-landau fluid modelturbulence simulations in a nonlinearballooning mode representation with radial modes. *Phys. Plasmas* **1**, 2229.
- [54] WANG, W., KISHIMOTO, Y., IMADERA, K., LIU, H.R., LI, J.Q., YAGI, M. & WANG, Z.X. 2020 Statistical study for itg turbulent transport in flux-driven tokamak plasmas based on global gyro-kinetic simulation. *Nuclear Fusion* **60** (6), 066010.
- [55] WEIKL, A., PEETERS, A. G., RATH, F., GROSSHAUSER, S. R., BUCHHOLZ, R., HORNSBY, W. A., SEIFERLING, F. & STRINTZI, D. 2017 Ion temperature gradient turbulence close to the finite heat flux threshold. *Phys. Plasmas* **24** (10), 102317.
- [56] WESSON, JOHN 2011 *Tokamaks*. Oxford: Oxford University Press.
- [57] WHELAN, G. G., PUESCHEL, M. J. & TERRY, P. W. 2018 Nonlinear electromagnetic stabilization of plasma microturbulence. *Phys. Rev. Lett.* **120**, 175002.
- [58] WHELAN, G. G., PUESCHEL, M. J., TERRY, P. W., CITRIN, J., MCKINNEY, I. J., GUTTENFELDER, W. & DOERK, H. 2019 Saturation and nonlinear electromagnetic stabilization of itg turbulence. *Physics of Plasmas* **26** (8), 082302.
- [59] W.M.NEWINS, J.CANDY, S.COWLEY, T.DANNERT, A.DIMITS, W.DORLAND, C.ESTRADA-MILA, G.W.HAMMET, F.JENKO, M.J.PUESCHEL & D.E.SHUMAKER 2006 Characterizing electron temperature gradient turbulence via numerical simulations. *Physics of Plasmas* **13**.

CHAPTER

Eidesstattliche Erklärung

Hiermit erkläre ich, Manuel Lippert, dass ich die vorliegende Arbeit selbstständig und ohne Benutzung anderer als der angegebenen Hilfsmittel angefertigt habe. Alle Stellen, die wörtlich oder sinngemäß aus veröffentlichten oder nicht veröffentlichten Schriften entnommen wurden, sind als solche kenntlich gemacht. Die Arbeit hat in gleicher oder ähnlicher Form noch keiner anderen Prüfungsbehörde vorgelegen.

Bayreuth, den May 21, 2023

Manuel Lippert