## Computational Optimal Transport - Project report:

# Regularized Optimal Transport methods for solving variational Mean-Field Games

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November 11, 2019

## 1 General setting: variational mean-field games

A mean-field game [8, 9] is a strategic decision-making problem with a very large, continuously-distributed number of interacting agents inside a state space: the overall theory developed by Lasry and Lions can be used as a means to model large, computationally intractable games. In the continuous-time setting explored in [9], each agent evolves according to some dynamics – leading to a so-called differential game – with the response to his choices depends on other agents' states and actions through a mean-field effect.

The general setup of a MFG has every agent penalize a running cost on the control and mean-field interaction, as well as a terminal cost on the its final position and the overall final distribution of agents (see [9]). The framework of [3, 4] focuses on games with agent dynamics  $dX_t = \alpha_t dt + dW_t$  where the running cost of the control  $\alpha$  is a quadratic function.

The (Nash) equilibrium agent-control dynamics can be summarized by the system of coupled nonlinear partial differential equations:

$$-\partial_t u - \frac{1}{2}\Delta u + \frac{1}{2}|\nabla u|^2 = f[\rho_t] \quad (t, x) \in (0, T) \times \Omega$$
 (1.1a)

$$\partial_t \rho_t - \frac{1}{2} \Delta \rho_t - \operatorname{div}(\rho_t \nabla u) = 0$$
 (1.1b)

$$\rho_0$$
 given (1.1c)

$$u(T, \cdot) = g[\rho_T] \tag{1.1d}$$

where and  $t \mapsto \rho_t$  is a trajectory in the space of measures, and  $\Omega$  a subset of the Euclidean space  $\mathbb{R}^d$ . The applications  $f[\mu]$  and  $g[\mu]$  are supposed to be derivatives of some real-valued functionals F and G on the space of measures. For instance, if  $G(\mu) = \int_{\Omega} \Psi \, d\mu(x)$  then its derivative is  $g[\mu](x) = \Psi(x)$ .

Equations (1.1a) and (1.1b) form a coupled system of control (Hamilton-Jacobi-Bellman) and diffusion (Fokker-Planck) partial differential equations. They can be solved in some cases using finite-difference methods (see Achdou, Camilli, and Capuzzo-Dolcetta [1]).

## 2 Variational formulations for the quadratic MFG

The first idea of [3] is to cast the MFG partial differential equations to a variational problem over an appropriate function space. Denote  $\mathbb{W}_2(\Omega) = (\mathcal{P}_2(\Omega), \mathcal{W}_2)$  the set of probability

measures with finite second moment, equipped with the Wasserstein metric

$$W_2(\mu,\nu)^2 = \inf_{\gamma \in \Pi(\mu,\nu)} \int |x-y|^2 d\gamma \tag{2.1}$$

where  $\Pi(\mu,\nu)=\{\gamma\in\mathcal{P}_2(\Omega\times\Omega):P_\#^1\gamma=\mu,\ P_\#^2\gamma=\nu\}$  is the set of transport plants from  $\mu$  to  $\nu$ . Then,  $\mathcal{C}([0,T],\mathbb{W}_2(\Omega))$  is the Wiener space of continuous  $\mathbb{W}_2$ -valued trajectories. Benamou, Carlier, and Santambrogio [3] show that the MFG be reformulated to the following variational problem:

$$\inf_{\rho,v} J(\rho,v) = \frac{1}{2} \int_0^T \int_{\Omega} |v_t|^2 d\rho_t(x) dt + \int_0^T F(\rho_t) dt + G(\rho_T)$$
s.t.  $\partial_t \rho_t - \frac{1}{2} \Delta \rho_t + \operatorname{div}(\rho_t v) = 0$ 

$$\rho_0 \in \mathbb{W}_2(\Omega)$$
(2.2)

where  $\rho = (\rho_t)_{t \in [0,T]} \in \mathcal{C}([0,T], \mathbb{W}_2(\Omega))$  is a trajectory in  $\mathbb{W}_2$  and v is a sufficiently regular function on  $[0,T] \times \Omega$ , most likely lying in a Sobolev space – see [3] for further discussion on regularity.

This point of view [3] is called *Eulerian*: we minimize over both the velocity v and the time-trajectory of the agents' density  $\rho$ . This can be solved by introducing Lagrange multipliers, exploiting duality, and using a finite element method, as shown in [3].

Benamou, Carlier, and Santambrogio [3] and Benamou et al. [4] also introduce a *Lagrangian* point of view, which allows to use tools from optimal transport theory: the variational problem is changed to optimize over the space of probability distributions on the space of agent trajectories.

## 2.1 Lagrangian formulation

#### 2.1.1 Wiener space and measure

This new point of view involves a change in function spaces. We denote  $\mathcal{X} = \mathcal{C}([0,T],\Omega)$  the Wiener space of (agents') trajectories  $[0,T] \to \Omega$ . Following [3, 2], we equip it with the Wiener measure (the law of a Wiener process with any starting point x)

$$R = \int_{\Omega} \delta_{x+W} \, dx$$

where W is a standard Wiener process in  $\mathbb{R}^d$ . It is an analogue in the space  $\mathcal{X}$  to the usual finite-dimensional Lebesgue measure<sup>1</sup>.

Measures  $Q \in \mathcal{P}(\mathcal{X})$  can also be seen as trajectories  $(Q_t)_{t \in [0,T]} \in \mathcal{C}([0,T],\mathcal{P}(\Omega))$ , with

$$Q_t = e_{t\#}Q \in \mathcal{P}(\Omega)$$

the push-forward of Q by the evaluation map  $e_t : \xi \in \mathcal{X} \longmapsto \xi(t)$ . This naturally defines an injection  $\underline{i} : \mathcal{P}(\mathcal{X}) \to \mathcal{C}([0,T],\mathcal{P}(\Omega))$ . We also introduce the more general marginals  $Q_{t_1,\ldots,t_n} = (e_{t_1},\ldots,e_{t_n})_{\#}Q$  for  $0 \leq t_1 < \cdots < t_N \leq T$ .

https://en.wikipedia.org/wiki/Infinite-dimensional\_Lebesgue\_measure

Marginals of the Wiener measure R. In particular,  $R_t$  is the Lebesgue measure  $\mathcal{L}^d$  on  $\mathbb{R}^d$ , and

$$R_{s,t}(dx, dy) = G_{t-s}(x - y) dx dy. (2.3)$$

Where  $G_t$  is the standard heat kernel

$$G_t(u) = \frac{1}{(2\pi t)^{d/2}} \exp\left(-\frac{|u|^2}{2t}\right)$$

Benamou, Carlier, and Santambrogio [3] and Benamou and Carlier [2] then re-cast the Eulerian variational game (2.2) into a Lagrangian optimization problem over the set of Borel probability measures. This new problem is also solved in [3] using a finite element method, which is computationally expensive.

#### 2.1.2 Entropic Lagrangian

Instead, Benamou et al. [4] propose using an entropy minimization approach to allow for a more computationally efficient method adapted from the Sinkhorn algorithm introduced by Cuturi [6].

This method, just like the Sinkhorn for OT between histograms (discrete measures), introduces some sort of entropic regularization [4], but this time on the measure over the trajectory space  $\mathcal{X}$ . The resulting numerical algorithm becomes a regularization of the Lagrangian from [3, 2].

The entropic Lagrangian variational problem is

$$\inf_{Q \in \mathcal{P}(\mathcal{X})} H(Q|R) + \int_0^T F(Q_t) \, dt + G(Q_T) \quad \text{s.t. } Q_0 = \rho_0$$
 (2.4)

Intuitively, this is the same as fixing the marginals  $\rho_t$ , finding the optimal bridge Q between them that has minimal entropy relative to the Wiener measure, and then optimizing over the  $\rho_t$ .

#### 2.2 Viscosity and the deterministic limit

We change the MFG problem to one following the agent-level dynamics  $dX_t = \alpha_t dt + \sigma dW_t$  with a diffusion coefficient  $\sigma$ . The MFG equilibrium equations become

$$-\partial_t u - \frac{\sigma^2}{2} \Delta u + \frac{1}{2} |\nabla u|^2 = f[\rho_t]$$

$$\partial_t \rho - \frac{\sigma^2}{2} \Delta \rho - \operatorname{div}(\rho \nabla u) = 0$$
(2.5)

This can be used to approximate first-order MFGs by setting a low viscosity parameter  $\sigma$ . Denoting  $\varepsilon = \sigma^2$ , the entropic variational problem (2.4) becomes

$$\inf_{Q \in \mathcal{P}(\mathcal{X})} \varepsilon H(Q|\mathbf{R}_{\varepsilon}) + \int_0^T F(Q_t) dt + G(Q_T) \quad \text{s.t. } Q_0 = \rho_0$$

where  $R_{\varepsilon}$  is the Wiener measure associated with Wiener processes scaled by  $\varepsilon$ .

## 3 Numerical algorithm

#### 3.1 Time discretization

Let N be the number of discrete steps for the time discretization of the problem, and h = T/N the time step.

Benamou et al. [4] propose a discretization of (2.4) obtained by connecting the marginals through a multimarginal OT problem:

$$S(\mu_0, \dots, \mu_N) = \inf_{\gamma \in \Pi(\mu_0, \dots, \mu_N)} H(\gamma | R^N)$$
(3.1)

where  $t_k = kh$ ,  $R^N = R_{t_0,...,t_N}$  and the marginals  $\mu_k \in \mathcal{P}_2(\Omega)$ . Then, define

$$\mathcal{U}(\mu_0, \dots, \mu_N) = h \sum_{k=1}^{N-1} F(\mu_k) + G(\mu_N).$$
 (3.2)

Thus, the discretized entropy minimization problem can be written as

$$\inf \{ S(\mu_0, \dots, \mu_N) + \mathcal{U}(\mu_0, \dots, \mu_N) : \mu_k \in \mathcal{P}_2(\Omega), \ \mu_0 = \rho_0 \}.$$

Expanding the inf-within-inf leads to the following convex optimization problem:

$$\inf_{\gamma \in \mathcal{P}(\Omega^{N+1})} H(\gamma | R^N) + \iota_{\rho_0}(\mu_0) + \sum_{k=1}^{N-1} F(\mu_k) + G(\mu_N)$$
s.t.  $\mu_k = P_{\#}^k \gamma$  (3.3)

where  $i_{\rho_0}(\mu) = +\infty$  if  $\mu \neq \rho_0$  and 0 otherwise is the convex indicatrix of the measure  $\rho_0$ . This is a generalized multimarginal optimal transport problem.

Benamou et al. [4] provide the corresponding dual problem involving the convex conjugates and potential functions, by using a multimarginal generalization of a result from Chizat et al. [5]:

$$\sup_{u} -i_{\rho_0}^*(-u_0) - \sum_{k=1}^{N-1} F^*(-u_k) - G^*(-u_N) - \int_{\Omega^{N+1}} \left( \exp\left(\bigoplus_{k=0}^N u_k\right) - 1 \right) dR^N$$
 (3.4)

where the supremum is taken over  $u = (u_0, \dots, u_N) \in L^{\infty}(\Omega)^{N+1}$ .

Benamou et al. [4] introduce a Sinkhorn-like iterative algorithm to solve the above dual problem. We rewrite it more explicitly with slightly different notations inspired by [5]

**Algorithm 1.** Denote for k = 0, ..., N and  $(a_i)_{i \neq k}$ 

$$\mathcal{I}_k((a_j)_{j \neq k})(z_k) = \int_{\Omega^N} \prod_{j \neq k} a_j(x_j) dR^N(x_{0:k-1}, z_k, x_{k+1:N})$$

the integral of functions  $a_j, j \neq k$  with respect to  $R^N$  and variables  $x_j, j \neq k$ . For convenience we use the shorthand

$$\mathcal{I}_k^{(n)} = \mathcal{I}_k \left( \left( e^{u_j^{(n+1)}} \right)_{j < k}, \left( e^{u_j^{(n)}} \right)_{j > k} \right)$$

for the nth iterate.

Then we compute the dual potentials iteratively:

$$\begin{cases} u_0^{(n+1)} = \underset{v \in L^{\infty}}{\operatorname{argmax}} \int_{\Omega} (1 - e^{-v(x_0)}) \mathcal{I}_0^{(n)} dx_0 - \imath_{\rho_0}^*(v) \\ u_k^{(n+1)} = \underset{v \in L^{\infty}}{\operatorname{argmax}} \int_{\Omega} (1 - e^{-v(x_k)}) \mathcal{I}_k^{(n)} dx_k - h F^*(v), \quad 1 \le k < N \\ u_N^{(n+1)} = \underset{v \in L^{\infty}}{\operatorname{argmax}} \int_{\Omega} (1 - e^{-v(x_N)}) \mathcal{I}_N^{(n)} dx_N - G^*(v) \end{cases}$$
(3.5)

until convergence.

Using duality, we find that the iterates  $u_k^{(n)}$  satisfy

$$a_k^{(n)} = \exp(-u_k^{(n)}) = \frac{\operatorname{prox}_{F_k}^{\mathrm{KL}}(\mathcal{I}_k^{(n)})}{\mathcal{I}_{L}^{(n)}}$$
 (3.6)

where

$$\operatorname{prox}_{F}^{\operatorname{KL}}(z) = \operatorname{argmin} F(s) + \operatorname{KL}(s|z).$$

**Remark 1** (Some convex conjugates). In practice, the convex conjugates of the cost functions are difficult to compute. For some of the examples in the paper, we have closed-form conjugates.

- The conjugate of the convex indicatrix  $\iota_{\nu}$  of any measure  $\nu$  is given by  $\iota_{\nu}^{*}(u) = \langle u, \nu \rangle$ .
- The hard congestion constraint

$$C(\rho) = \begin{cases} 0 & \text{if } \rho \le \bar{m} \\ +\infty & \text{otherwise} \end{cases}$$

has convex conjugate (on the domain  $\rho \geq 0$ )

$$C^*(u) = \sup_{\rho \le \bar{m}} \langle u, \rho \rangle = \langle u^+, \bar{m} \mathbb{1} \rangle$$

• Obstacle constraints, given by

$$F(\rho) = \int_{\Omega} V(x) \, d\rho(x) = \begin{cases} 0 & \text{if } \rho = 0 \text{ on } \mathscr{O} \\ +\infty & \text{otherwise} \end{cases} = \imath_0(\mathbb{1}_{\mathscr{O}}\rho)$$

where V is the convex indicatrix of the complement  $\Omega \backslash \mathcal{O}$  of the obstacles. Its conjugate is given by

$$F^*(u) = \begin{cases} 0 & \text{if } u \le 0 \text{ on } \Omega \backslash \mathscr{O} \\ +\infty & \text{otherwise} \end{cases}$$

#### 3.2 Full discretization

For full numerical implementation, all measures are replaced by multi-dimensional arrays representing discrete histograms over a fixed grid of points in  $\mathbb{R}^d$  of dimensionality  $M = N_1 \times \cdots \times N_d$ . Integration is exchanged with summation.

In the general case, the KL-projections in the Sinkhorn iterations can be solved using the Python library CVXPY<sup>2,3</sup>. Some can be computed explicitly.

**Proposition 1.** The KL-projection under the hard congestion constraint of a measure  $\beta \in \mathbb{R}^M$  is given by

$$\operatorname{prox}_{C}^{\operatorname{KL}}(\beta) = \min(\beta, \overline{m}) \tag{3.7}$$

where the minimum is taken element-wise.

If we also add the obstacle constraint on a set  $\mathcal O$  of points in the grid, then the proximal operator reads

$$\operatorname{prox}_{F}^{\operatorname{KL}}(\beta) = \min(\beta, \bar{m} \mathbb{1}_{\Omega \setminus \mathscr{O}}). \tag{3.8}$$

## 4 Examples

#### 4.1 Transport with a soft target

#### 4.1.1 Two-marginal case

We start with a very simplified approximation of the crowd displacement problem on  $\Omega = [0,1]^2$ , with only the first step (with initial agent distribution) and final step decided by the terminal penalty function G.

We set G to be the obstacle constraint related to a subset  $\mathscr O$  of  $\Omega$  as well as a potential  $\Psi(x)=d(x,\mathscr A)^\beta$  for some  $\beta>0$ , related to the distance to a target subset  $\mathscr A$  (see fig. 1):

$$G(\mu) = \int_{\Omega} \Psi \, d\mu + \imath_0(\mu \mathbb{1}_{\mathscr{O}})$$

Thus, the agents engage in a one-round mean-field game where they are only concerned with moving to regions with lower potential  $\Psi$  – as close as possible to  $\mathscr{A}$  – whilst obeying physical constraints related to the obstacles.

The discretized MFG problem with viscosity parameter  $\varepsilon = \sigma^2$  can be written as the following transport problem:

$$\inf_{\gamma} \langle \Psi, \gamma^T \mathbb{1} \rangle + \varepsilon H(\gamma | R_{\varepsilon})$$
s.t.  $\gamma \mathbb{1} = \rho_0, \quad \gamma^T \mathbb{1} \odot \mathbb{1}_{\mathscr{O}} = 0$  (4.1)

The interesting aspect of this problem is observing what kind of optimal distribution  $\rho_1^* = (\gamma^*)^T \mathbb{1}$  the agents reach.

**Proposition 2.** Problem (4.1) can be solved in closed form: the Lagrange multiplier  $u_0^*$  for the marginal law constraint satisfies

$$e^{-u_0^*} = \frac{\rho_0}{R_\varepsilon a_1^*}$$

where  $a_1^* = e^{-\Psi/\varepsilon} \odot \mathbb{1}_{\Omega \setminus \mathscr{O}}$ , and the optimal coupling is

$$\gamma^* = R_{\varepsilon} \odot (e^{-u^*} \otimes \hat{\varphi})$$

It satisfies, as expected, that  $\gamma_{i,j}^* = 0$  for all  $j \in \mathscr{O}$ .

<sup>&</sup>lt;sup>2</sup>https://github.com/cvxgrp/cvxpy

<sup>&</sup>lt;sup>3</sup>Steven Diamond and Stephen Boyd. "CVXPY: A Python-Embedded Modeling Language for Convex Optimization". In: *Journal of Machine Learning Research* 17.83 (2016), pp. 1–5.

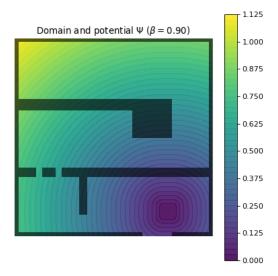


Figure 1: Computational domain of the game  $\Omega$  with set of obstacles  $\mathscr{O}$  (transparent grey), and contour of the potential function  $\Psi(x) = d(x, \mathscr{A})^{\beta}$ .

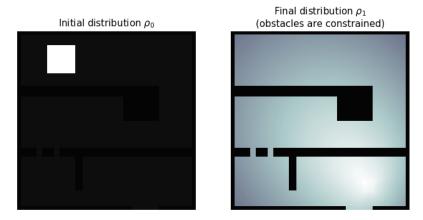
Numerical experiment We ran a numerical experiment by implementing the solutions given by proposition 2 to the discrete MFG (4.1). fig. 2a provides a representation of both . We also checked the results when removing the constraints on the obstacles (essentially setting  $\mathcal{O}$ ), and when lowering the viscosity parameter  $\sigma = \sqrt{\varepsilon}$  (see figs. 2b and 2c).

Of course, replacing G by a hard marginal constraint turns the problem into a classical regularized OT problem, and the above proposition leads to usual Sinkhorn iterations on the grid. The matrix-vector product in the Lagrange multiplier potentially becomes a computational bottleneck, but the separability of the heat kernel  $R_{\varepsilon}$  allows for fast computation using simple 1D convolutions.

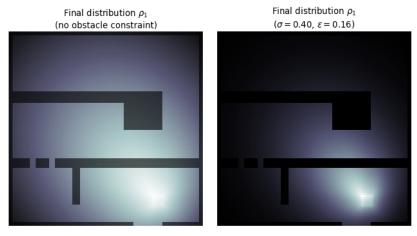
Remark 2 (A "smarter" (more realistic) potential for crowd dynamics). The results shown fig. 2 are satisfactory for the given potential  $\Psi$  – as expected the agents try to stay near the low-potential regions. However, for modeling of crowd dynamics they would be deeply nonphysical because the potential is inadequate. In a room evacuation scenario, for instance, agents would seek to minimize the time-to-exit: the literature shows this leads to the Eikonal equation, a kind of Hamilton-Jacobi PDE. We computed the adequate potential shown fig. 3a using the Fast Sweeping method [10], as well as the discrete MFG fig. 3b.

## 4.1.2 One intermediate time step

We now go up to three marginals  $(\rho_0, \rho_1, \rho_2)$ . We assign to the single intermediate marginal  $\rho_1$  the congestion constraint  $\rho_1 \leq \bar{m}$  and the obstacle constraint. The primal problem then

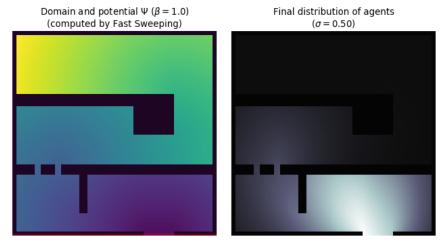


(a) Result of the "fuzzy" transport problem with enforcement of the obstacle constraints and viscosity parameter  $\sigma=1.$ 



(b) Final distribution  $\rho_1^*$  without en- (c) Obstacles are constrained as in forcing the obstacles. The mass of the fig. 2a, but with a lower viscosity padistribution "bleeds" through the ob- rameter  $\sigma=0.4$ . stacles.

Figure 2: Marginal distributions of the solution of two-step MFG or "fuzzy transport" problem (4.1), with a few variations.



(a) Domain and potential associated (b) Optimal terminal distribution  $\rho_1^*$  of with the fastest path distance. the discrete MFG with the potential from fig. 3a.

Figure 3: Setup and solution for the discrete MFG using the time-to-exit potential discussed in remark 2.

reads

$$\inf_{\gamma,\rho_{1},\rho_{2}} \langle \Psi, \rho_{2} \rangle + \varepsilon H(\gamma | R_{\varepsilon})$$
s.t.  $P_{\#}^{k} \gamma = \rho_{k}, \ k = 1, 2$ 

$$\rho_{1} \leq \bar{m}$$

$$\rho_{1} \odot \mathbb{1}_{\mathscr{O}} = 0$$

$$\rho_{2} \odot \mathbb{1}_{\mathscr{O}} = 0$$
(4.2)

**Proposition 3.** The Lagrange multipliers  $u_i^*$  at the optimum satisfy the fixed-point conditions:

$$\begin{split} a_0^* &= \frac{\rho_0}{R_\varepsilon[\cdot, a_1^*, a_2^*]} \\ a_1^* &= \min\left(\frac{\bar{m}}{R_\varepsilon[a_0^*, \cdot, a_1^*]}\right) \end{split}$$

where  $a_i^* = \exp(-u_i^*)$ ,  $a_1^*$  is supported on  $\Omega \setminus \mathscr{O}$  and  $a_2^* = e^{-\Psi/\varepsilon} \odot \mathbb{1}_{\Omega \setminus \mathscr{O}}$ , and we denote  $R_{\varepsilon}[\cdot, \cdot, \cdot]$  the tensor product by  $R_{\varepsilon}$ .

The fixed point can then computed using an iterative algorithm à la generalized Sinkhorn, just as in the Algorithm 1 suggested by [4].

The issue of computational efficiency is more pronounced here than before due to the tensor product and need for multiple iterations until convergence.

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