

Null Tetrads

食司

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Abstract

感谢 YouTube 博主@TensorCalculusRobertDavie的教学视频Null Vectors and Null Tetrads。鉴于笔者未能找到如此基础的中文资源，特针对其视频整理笔记如下。

1 Null Tetrads

Null Tetrads 是定义在时空上的一组基底切向量 $\{l, n, m, \bar{m}\}$ ，它们会随度规张量 g_{ab} 的不同而不同。

在任意时空中考虑，记其正交归一的四线元余切基矢组为 $\{v, i, j, k\}$ ，满足

$$v \cdot v = \eta_{\alpha\beta} v^\alpha v^\beta = v_\beta v^\beta = -1 \quad (1.1)$$

$$i \cdot i = \eta_{\alpha\beta} i^\alpha i^\beta = i_\beta i^\beta = j \cdot j = k \cdot k = 1 \quad (1.2)$$

且它们两两正交

定义

$$l = \frac{1}{\sqrt{2}}(v + i), \quad n = \frac{1}{\sqrt{2}}(v - i) \quad (1.3)$$

$$m = \frac{1}{\sqrt{2}}(j + ik), \quad \bar{m} = \frac{1}{\sqrt{2}}(j - ik) \quad (1.4)$$

利用基矢的关系很容易检验以下关系成立

- 均为零矢量 $l \cdot l = n \cdot n = m \cdot m = \bar{m} \cdot \bar{m} = 0$
- 正交性 $l \cdot m = l \cdot \bar{m} = n \cdot m = n \cdot \bar{m} = 0$
- 归一性 $l \cdot n = -1, \quad m \cdot \bar{m} = 1$

现在我们可以取切矢量基矢组为 $\{e_0, e_1, e_2, e_3\} = \{l, n, m, \bar{m}\}$

即有

$$e_0^\alpha = l^\alpha = \frac{1}{\sqrt{2}}(v^\alpha + i^\alpha), \quad e_1^\alpha = n^\alpha = \frac{1}{\sqrt{2}}(v^\alpha - i^\alpha) \quad (1.5)$$

$$e_2^\alpha = m^\alpha = \frac{1}{\sqrt{2}}(j^\alpha - ik^\alpha), \quad e_3^\alpha = \bar{m}^\alpha = \frac{1}{\sqrt{2}}(j^\alpha + ik^\alpha) \quad (1.6)$$

先考虑闵氏时空。利用 frame metric γ_{mn} 的性质，可求此基矢下的 frame metric 为

$$\gamma_{mn} = \begin{bmatrix} e_0 \cdot e_0 & e_0 \cdot e_1 & e_0 \cdot e_2 & e_0 \cdot e_3 \\ e_1 \cdot e_0 & e_1 \cdot e_1 & e_1 \cdot e_2 & e_1 \cdot e_3 \\ e_2 \cdot e_0 & e_2 \cdot e_1 & e_2 \cdot e_2 & e_2 \cdot e_3 \\ e_3 \cdot e_0 & e_3 \cdot e_1 & e_3 \cdot e_2 & e_3 \cdot e_3 \end{bmatrix} = \begin{bmatrix} l \cdot l & l \cdot n & l \cdot m & l \cdot \bar{m} \\ n \cdot l & n \cdot n & n \cdot m & n \cdot \bar{m} \\ m \cdot l & m \cdot n & m \cdot m & m \cdot \bar{m} \\ \bar{m} \cdot l & \bar{m} \cdot n & \bar{m} \cdot m & \bar{m} \cdot \bar{m} \end{bmatrix} = \begin{bmatrix} 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \quad (1.7)$$

容易检验 γ_{mn} 是它自己的逆，即有

$$\begin{bmatrix} 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad (1.8)$$

故

$$\gamma_{mn} = \gamma^{mn} \quad (1.9)$$

$$\begin{bmatrix} e^0 \\ e^1 \\ e^2 \\ e^3 \end{bmatrix} = \begin{bmatrix} 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} e_0 \\ e_1 \\ e_2 \\ e_3 \end{bmatrix} = \begin{bmatrix} -e_1 \\ -e_0 \\ e_3 \\ e_2 \end{bmatrix} \quad (1.10)$$

下面给出 frame metric 和 coordinate metric 的关系

$$g_\alpha = e^m_\alpha \gamma_m \quad (1.11)$$

$$g_{\alpha\beta} = e^m_\alpha \gamma_m e^n_\beta \gamma_n = \gamma_{mn} e^m_\alpha e^n_\beta \quad (1.12)$$

$$\gamma_m = e_m^\alpha \partial_\alpha = e_m^\alpha g_\alpha \quad (1.13)$$

特别地我们有

$$ds^2 = g_{\alpha\beta} dx^\alpha dx^\beta \quad (1.14)$$

$$= \gamma_{mn} e^m_\alpha dx^\alpha e^n_\beta dx^\beta \quad (1.15)$$

$$e^m = e^m_\alpha dx^\alpha \quad (1.16)$$

$$e_m = e_m^\alpha \partial_\alpha \quad (1.17)$$

和前面类似的我们有

$$g^{\alpha\beta} = \gamma^{mn} e_m^\alpha e_n^\beta \quad (1.18)$$

现在代入 γ^{mn} 的具体形式

$$g^{\alpha\beta} = \gamma^{mn} e_m^\alpha e_n^\beta \quad (1.19)$$

$$= \gamma^{01} e_0^\alpha e_1^\beta + \gamma^{10} e_1^\alpha e_0^\beta + \gamma^{23} e_2^\alpha e_3^\beta + \gamma^{32} e_3^\alpha e_2^\beta \quad (1.20)$$

$$= -l^\alpha n^\beta - n^\alpha l^\beta + m^\alpha \bar{m}^\beta + \bar{m}^\alpha m^\beta \quad (1.21)$$

以及

$$g_{\alpha\beta} = -l_\alpha n_\beta - n_\alpha l_\beta + m_\alpha \bar{m}_\beta + \bar{m}_\alpha m_\beta \quad (1.22)$$

这就是常见的利用 **Null Tetrads** 表示度规张量的公式，这个公式是具有普遍性的。这个变换在研究对称性等方面具有一些明显的优势。

现在以闵氏时空为例。闵氏时空中有

$$ds^2 = -c^2 dt^2 + dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \quad (1.23)$$

故其余切基矢组为 $\{cdt, dr, rd\theta, r \sin \theta d\phi\}$ 。构造

$$l = \frac{1}{\sqrt{2}}(cdt + dr), \quad n = \frac{1}{\sqrt{2}}(cdt - dr) \quad (1.24)$$

$$m = \frac{1}{\sqrt{2}}(rd\theta + ir \sin \theta d\phi), \quad \bar{m} = \frac{1}{\sqrt{2}}(rd\theta - ir \sin \theta d\phi) \quad (1.25)$$

写成分量的形式

$$l_\alpha = \frac{1}{\sqrt{2}}(1, 1, 0, 0), \quad n_\alpha = \frac{1}{\sqrt{2}}(1, -1, 0, 0) \quad (1.26)$$

$$m_\alpha = \frac{1}{\sqrt{2}}(0, 0, r, ir \sin \theta), \quad \bar{m}_\alpha = \frac{1}{\sqrt{2}}(0, 0, r, -ir \sin \theta) \quad (1.27)$$

由 $g_{\alpha\beta} = -l_\alpha n_\beta - n_\alpha l_\beta + m_\alpha \bar{m}_\beta + \bar{m}_\alpha m_\beta$

$$g_{00} = -l_0 n_0 - n_0 l_0 + m_0 \bar{m}_0 + m_0 \bar{m}_0 \quad (1.28)$$

$$= -\frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} - \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} + 0 + 0 \quad (1.29)$$

$$= -1 \quad (1.30)$$

类似地，可以求出 $g_{11} = 1, g_{22} = r, g_{33} = r \sin \theta$ ，则 $g_{\alpha\beta}$ 确为闵氏度规

$$g_{\alpha\beta} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{bmatrix} \quad (1.31)$$

其逆为

$$g^{\alpha\beta} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \frac{1}{r^2} & 0 \\ 0 & 0 & 0 & \frac{1}{r^2 \sin^2 \theta} \end{bmatrix} \quad (1.32)$$

利用 $g^{\alpha\beta}$ 可对 Null Tetrads 进行指标升降

$$l^\alpha = g^{\alpha\beta} l_\beta = \frac{1}{\sqrt{2}}(-1, 1, 0, 0) \quad (1.33)$$

类似地, 得到

$$n^\alpha = \frac{1}{\sqrt{2}}(-1, -1, 0, 0), \quad m^\alpha = \frac{1}{\sqrt{2}}\left(0, 0, \frac{1}{r}, \frac{i}{r \sin \theta}\right) \quad (1.34)$$

直接取复共轭就有

$$\bar{m}^\alpha = \frac{1}{\sqrt{2}}\left(0, 0, \frac{1}{r}, \frac{-i}{r \sin \theta}\right) \quad (1.35)$$

可以检验切矢形式的 Null Tetrads 仍满足前文所提到的性质。

现在考虑一个更复杂但常用的情况。考虑 Advanced Eddington-Finkelstein 坐标 (u, r, θ, ϕ) 中的史瓦西度规

$$ds^2 = -\left(1 - \frac{2GM}{c^2 r}\right) du^2 + 2du dr + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \quad (1.36)$$

令

$$F = \left(1 - \frac{2GM}{c^2 r}\right) \quad (1.37)$$

度规的矩阵形式可以写为

$$g_{\alpha\beta} = \begin{bmatrix} -F & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{bmatrix} \quad (1.38)$$

注意到

$$\begin{bmatrix} -F & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{bmatrix} \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & F & 0 & 0 \\ 0 & 0 & \frac{1}{r^2} & 0 \\ 0 & 0 & 0 & \frac{1}{r^2 \sin^2 \theta} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad (1.39)$$

不妨先假设

$$l^\alpha = (l^u, l^r, l^\theta, l^\phi) = (0, 1, 0, 0) \quad (1.40)$$

则

$$l_\alpha = g_{\alpha\beta} l^\beta = (g_{ur} l^r, 0, 0, 0) = (1, 0, 0, 0) \quad (1.41)$$

m , \bar{m} 与闵氏时空中的相同。现在目标是解出 n^α 。

利用 $g_{\alpha\beta} = -l_\alpha n_\beta - n_\alpha l_\beta + m_\alpha \bar{m}_\beta + \bar{m}_\alpha m_\beta$

$$g_{uu} = -l_u n_u - n_u l_u + m_u \bar{m}_u + \bar{m}_u m_u \quad (1.42)$$

$$= -2l_u n_u \quad (1.43)$$

$$= -2n_u \quad (1.44)$$

$$= -F \quad (1.45)$$

故

$$n_u = \frac{F}{2} \quad (1.46)$$

$$g_{ur} = -l_u n_r - n_u l_r + m_u \bar{m}_r + \bar{m}_u m_r \quad (1.47)$$

$$= -n_r \quad (1.48)$$

$$= 1 \quad (1.49)$$

故

$$n_r = -1 \quad (1.50)$$

于是得到

$$n_\alpha = \left(\frac{F}{2}, -1, 0, 0\right) \quad (1.51)$$

利用 $g^{\alpha\beta}$ 升指标

$$n^\alpha = \left(-1, -\frac{F}{2}, 0, 0\right) \quad (1.52)$$

由前文

$$m^\alpha = \frac{1}{\sqrt{2}}\left(0, 0, \frac{1}{r}, \frac{i}{r \sin \theta}\right), \quad \bar{m}^\alpha = \frac{1}{\sqrt{2}}\left(0, 0, \frac{1}{r}, \frac{-i}{r \sin \theta}\right) \quad (1.53)$$

综上

$$l^\alpha = (0, 1, 0, 0), \quad l_\alpha = (1, 0, 0, 0) \quad (1.54)$$

$$n^\alpha = \left(-1, -\frac{F}{2}, 0, 0\right), \quad n_\alpha = \left(\frac{F}{2}, -1, 0, 0\right) \quad (1.55)$$

$$m^\alpha = \frac{1}{\sqrt{2}}\left(0, 0, \frac{1}{r}, \frac{i}{r \sin \theta}\right), \quad m_\alpha = \frac{1}{\sqrt{2}}(0, 0, r, ir \sin \theta) \quad (1.56)$$

$$\bar{m}^\alpha = \frac{1}{\sqrt{2}}\left(0, 0, \frac{1}{r}, \frac{-i}{r \sin \theta}\right), \quad \bar{m}_\alpha = \frac{1}{\sqrt{2}}(0, 0, r, -ir \sin \theta) \quad (1.57)$$