

# ① Recombination

Saha equation for ionisation fraction  $X_e$ :  $\left(\frac{1-X_e}{X_e^2}\right) = \frac{2\zeta(3)}{\pi^2} \eta \left(\frac{2\pi T}{m_e}\right)^{3/2} e^{E_I/T}$

a.) Derive this.

$$X_e = \frac{n_e}{n_p + n_H} \leftarrow \text{fraction of electrons left behind}$$

$$n_H = n_p - n_n$$

$$1 - X_e = 1 - \frac{n_e}{n_p + n_H}$$

$$\frac{1-X_e}{X_e^2} = \frac{1 - \frac{n_e}{n_p + n_H}}{\frac{n_e^2}{(n_p + n_H)^2}}$$

$$= \frac{n_p + n_H - n_e}{n_p + n_H} \cdot \frac{(n_p + n_H)^2}{n_e^2}$$

$$= \frac{n_p + n_H - n_e}{n_e^2} (n_p + n_H), \text{ where Universe is uncharged so } n_p = n_e$$

$$\frac{1-X_e}{X_e^2} = \frac{n_H}{n_e^2} (n_p + n_H), \text{ where } \eta = \frac{n_B}{n_\gamma}$$

$$= \frac{n_H}{n_e^2} \eta n_\gamma, \text{ where } n_b = \eta n_\gamma = \eta \times \frac{2\zeta(3)}{\pi^2} T^3$$

$$= \eta \frac{2\zeta(3)}{\pi^2} T^3 \frac{n_H}{n_e^2}$$

$$-E_I = m_H - m_p - m_e$$

$$\left(\frac{n_H}{n_e n_p}\right)_{eq} = \frac{g_H}{g_e g_p} \left(\frac{m_H}{m_e m_p} \frac{2\pi}{T}\right)^{3/2} e^{(m_p + m_e - m_H)/T}$$

$\downarrow n_e = n_p$  (no charge)     $\downarrow g_e = g_p = 2, g_H = 4 \rightarrow \frac{4}{2 \cdot 2} = 1$      $\downarrow m_H \approx m_p$      $\downarrow = E_I$

$$\left(\frac{n_H}{n_e^2}\right)_{eq} = \left(\frac{2\pi}{m_e T}\right)^{3/2} e^{E_I/T}$$

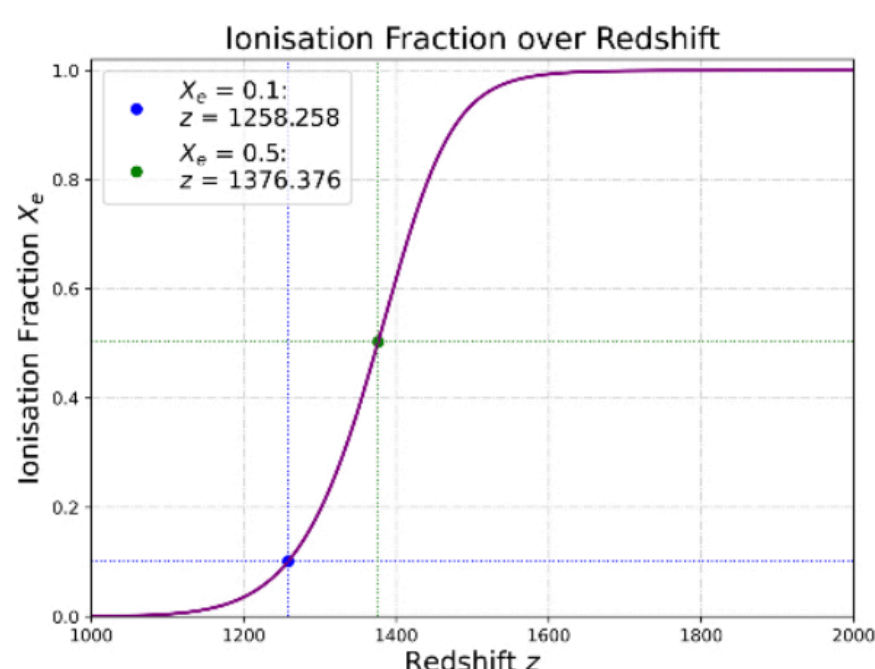
$$= \frac{2\zeta(3)}{\pi^2} \eta T^3 \left(\frac{2\pi}{m_e T}\right)^{3/2} e^{E_I/T}, \zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s}$$

$$\frac{1-X_e}{X_e^2} = \frac{2\zeta(3)}{\pi^2} \eta \left(\frac{2\pi T}{m_e}\right)^{3/2} e^{E_I/T} \leftarrow \text{Saha equation}$$

b.) Solve for  $X_e$  as function of  $z$ . Plot.

(Did not solve by hand, but rearranged for quadratic formula.)

$$X_e = \frac{-1 + \sqrt{1+4f}}{2f}, \text{ where } f(T, \eta) = \frac{2\zeta(3)}{\pi^2} \eta \left(\frac{2\pi T}{m_e}\right)^{3/2} e^{E_I/T}$$



c.) From figure, find  $z$  at  $X_e = 0.1$ . How different from  $z(X_e = 0.5)$ ?

See plot markers and labels:

• for  $X_e = 0.1 \rightarrow z = 1258.258$

• for  $X_e = 0.5 \rightarrow z = 1376.376$

Difference of over  $z = 100$

d.) Calculate age of Universe at this epoch. Assume matter-dominated.

$$t = \frac{2}{3H} \quad \downarrow \quad \text{Friedmann with } \Omega_m = 1 \rightarrow H = H_0 \Omega_m a^{-3/2}$$

$$t = \frac{2}{3} \frac{a^{+3/2}}{H_0 \Omega_m}$$

$$= \frac{2}{3} \frac{1}{H_0 \Omega_m (1+z)^{3/2}}$$

$$= \frac{2}{3} \frac{\text{Gyr}}{0.07 | 1 | (1+1258.258)^{3/2}}$$

$$= 2.131 \times 10^{-4} \text{ Gyr}$$

$$t = 0.2131 \text{ Myr} \rightarrow \text{Universe was very young at this ionisation fraction!}$$

e.) Using this ionisation fraction, estimate  $z_{dec}$  where Thomson scattering rate = Hubble rate.

• Estimate Universe age at this  $z_{dec}$ . Find ionisation frac at this  $z_{dec}$

Interaction rate  $\Gamma_\gamma \approx n_e \sigma_T$ , where  $\sigma_T \approx 2 \times 10^{-3} \text{ MeV}^{-2}$

At decoupling,  $\Gamma_\gamma(T_{dec}) \approx H(T_{dec})$

Substitute for  $n_b$  to find  $\Gamma_\gamma(T_{dec})$ :

$$\Gamma_\gamma(T_{dec}) = n_b X_e(T_{dec}) \sigma_T, \text{ where from 1a, } n_b = \eta n_\gamma = \eta \frac{2\zeta(3)}{\pi^2} T^3$$

$$= \frac{2\zeta(3)}{\pi^2} \eta \sigma_T X_e(T_{dec}) T_{dec}^3$$

Solve for  $H(T_{dec})$ :

$$H(T_{dec}) = H_0 \sqrt{\Omega_m} \left(\frac{T_{dec}}{T_0}\right)^{3/2}$$

Set equal  $\Gamma_\gamma$  and  $H$ :

$$\frac{2\zeta(3)}{\pi^2} \eta \sigma_T X_e(T_{dec}) T_{dec}^3 = H_0 \sqrt{\Omega_m} \left(\frac{T_{dec}}{T_0}\right)^{3/2}$$

$$X_e(T_{dec}) T_{dec}^{3/2} = \frac{\pi^2}{2\zeta(3)} \frac{H_0 \sqrt{\Omega_m}}{\eta \sigma_T T_0^{3/2}}$$

$$T_{dec} = \left[ \frac{10}{2} \frac{\pi^2}{1.2021} \frac{0.07 \text{ Gyr} \text{ MeV}}{7.60 \times 10^{36}} \frac{1}{6 \times 10^{-10}} \frac{1}{2 \times 10^{-3}} \frac{1}{(2.3525 \times 10^{-10})^{3/2}} \right]^{2/3}$$

$$T_{dec} = 2.579 \times 10^{-7} \text{ MeV}$$

• Estimate age:

$$T_{dec} = T_0 (1+z_{dec})$$

$$z_{dec} = \frac{T_{dec}}{T_0} - 1$$

$$z_{dec} = \frac{2.480 \times 10^{-7} \text{ MeV}}{2.3525 \times 10^{-10} \text{ MeV}} - 1$$

$$z_{dec} = 1095 \approx 1100$$

From my plot,  $X_e(z=1095) = 0.004$

$$t_{dec} = \frac{2}{3} \frac{a^{+3/2}}{H_0 \Omega_m} = \frac{2}{3} \frac{1}{H_0 \Omega_m (1+z_{dec})^{3/2}} = \frac{2}{3} \frac{\text{Gyr}}{0.07 | 1 | (1+1095)^{3/2}}$$

$$t_{dec} = 262 \text{ Myr}$$

Universe would be 262 Myr with redshift 1095 and ionisation fraction 0.004.



## ② "What if?" BBN

Neutron fraction defined as  $X_n \equiv \frac{n_n}{n_n + n_p}$

a.) Derive equilibrium abundance  $X_n = \frac{e^{-Q/T}}{1 + e^{-Q/T}}$

where  $Q = m_n - m_p = 1.3 \text{ MeV}$

From Baumann (3.126):

$$\left(\frac{n_n}{n_p}\right)_{eq} = \left(\frac{m_n}{m_p}\right)^{3/2} e^{-(m_n - m_p)/T} \quad \text{Using equilibrium number densities as detailed in Baumann 3.73}$$

Because  $m_n$  is very close to  $m_p$ , we disregard  $\left(\frac{m_n}{m_p}\right)^{3/2}$ :

$$\left(\frac{n_n}{n_p}\right)_{eq} = e^{-(m_n - m_p)/T}, \quad \text{where } Q = m_n - m_p$$

$$\left(\frac{n_n}{n_p}\right)_{eq} = e^{-Q/T}$$

Putting this into the given  $X_n$  expression:

$$X_n \equiv \frac{n_n}{n_n + n_p} \rightarrow \text{dividing by } n_p \text{ in all terms}$$

$$X_n = \frac{\frac{n_n}{n_p}}{\frac{n_n}{n_p} + \frac{n_p}{n_p}}$$

$$X_n = \frac{e^{-Q/T}}{1 + e^{-Q/T}} \quad \leftarrow \text{intended expression for neutron fraction } X_n$$

b.) Estimate freeze-out abundance of neutrons  $X_n$ .

Assume "freeze-out" temperature is  $0.8 \text{ MeV}$

$$X_n = \frac{e^{-Q/T}}{1 + e^{-Q/T}}, \quad \text{where } Q = 1.3 \text{ MeV}, T = 0.8 \text{ MeV}$$

$$= \frac{e^{-1.3 \text{ MeV}/0.8 \text{ MeV}}}{1 + e^{-1.3 \text{ MeV}/0.8 \text{ MeV}}}$$

$$= \frac{e^{-1.625}}{1 + e^{-1.625}}$$

$$= \frac{0.19691}{1.19691}$$

$$X_n = 0.1645 \quad \leftarrow \text{freeze-out neutron abundance}$$

c.) Calculate mass fraction of helium  $Y_p = \frac{4n_{\text{He}}}{n_H}$ .

Assume all neutrons from 4b are converted to helium-4.

No neutron decay, so all  $n \rightarrow \text{He}$

\* Important: The wording was unclear to me (and others), so my interpretation of 'assume all neutrons are converted' is that we are being told to not account for decay. I also think we're being told to use this  $Y_p$  expression and not derive it.

So that's what I solve here:

$$Y_p = \frac{4n_{\text{He}}}{n_H}, \quad \text{where } n_{\text{He}} = \frac{1}{2}n_n \text{ and } n_H \approx n_p - n_n$$

$$= \frac{4 \cdot \frac{1}{2}n_n}{n_p - n_n}$$

$$= \frac{2 \frac{n_n}{n_p}}{\frac{n_p}{n_p} - \frac{n_n}{n_p}}$$

$$= \frac{2X_n}{1 - X_n}$$

$$= \frac{2(0.1645)}{1 - 0.1645}$$

$$Y \approx 0.39 \quad \leftarrow$$

Mass fraction of Helium for freeze-out neutron abundance of  $X_n = 0.1645$ .

d) If we maintained freeze-out temperature but doubled  $Q$ , how is He abundance impacted?

$$T = 0.8 \text{ MeV}$$

$$Q = m_n - m_p = \cancel{1.3 \text{ MeV}} \quad 2.6 \text{ MeV}$$

$$X_n = \frac{e^{-Q/T}}{1 + e^{-Q/T}},$$

$$= \frac{e^{-2.6 \text{ MeV}/0.8 \text{ MeV}}}{1 + e^{-2.6 \text{ MeV}/0.8 \text{ MeV}}}$$

$$= \frac{e^{-3.25}}{1 + e^{-3.25}}$$

$$= \frac{0.03877}{1.03877}$$

$$X_n = 0.0373 \quad \leftarrow \text{freeze-out neutron abundance}$$

Now use this to calculate helium mass fraction:

$$Y_p = \frac{4n_{\text{He}}}{n_H}$$

$$= \frac{4 \cdot \frac{1}{2}n_n}{n_p - n_n}$$

$$= \frac{2 \frac{n_n}{n_p}}{\frac{n_p}{n_p} - \frac{n_n}{n_p}}$$

$$= \frac{2X_n}{1 - X_n}$$

$$= \frac{2(0.0373)}{1 - 0.0373}$$

$$Y \approx 0.078$$



### 3. Freeze-In DM

Huterer 6.4

DM particle  $\chi$  produced by  $\sigma \rightarrow \chi\chi$  and  $\delta n_\chi = -2\delta n_\sigma$  with  $\Gamma \equiv \Gamma_{\sigma \rightarrow \chi\chi}$  (no  $n$  dependence; only on  $\sigma$ ).

For small enough  $\Gamma$ ,  $n_\chi$  increases then flattens when  $n_\sigma$  becomes Boltzmann suppressed:  $n_\sigma \propto e^{-m_\sigma/T} \ll 1$

#### a) FREEZE-OUT

Boltzmann for freezeout:

$$\frac{1}{a^3} \frac{d(na^3)}{dt} = -\langle \sigma v \rangle \left[ n_1 n_2 - \left( \frac{n_1 n_2}{n_3 n_4} \right)_{eq} n_3 n_4 \right] = -\langle \sigma v \rangle [n^2 - n_{eq}^2]$$

For freezeout Boltzmann eqs:  $\frac{dY}{dx} = -\frac{1}{x^2} [Y^2 - Y_{eq}^2]$

#### FREEZE-IN

Boltzmann for freeze-in:

$$\frac{1}{a^3} \frac{d(na^3)}{dt} = 2\Gamma h(t) n_{\sigma,eq}(t), \text{ where } h(x) \approx \frac{x}{x+2}, x \equiv \frac{m_\sigma}{T}$$

Show for freeze-in Boltzmann:  $\frac{dY}{dx} = \lambda \times h(x) Y_{eq}(x)$

Solution:

$$\Gamma = n_2 \langle \sigma v \rangle, s \equiv \frac{S}{V} = \frac{p+p}{T} = \frac{2\pi^2}{45} g_{*s} T^3 \quad (H 4.32)$$

$$Y \equiv \frac{n}{T^3} \propto \frac{n}{s}, x \equiv \frac{M}{T}, \lambda \equiv \frac{2\pi^2}{45} g_{*s} \frac{m^3 \langle \sigma v \rangle}{H(T=m)}$$

$$\begin{aligned} \frac{1}{a^3} \frac{d(na^3)}{dt} &= 2\Gamma h(t) n_{\sigma,eq}(t) & \frac{dx}{dt} &= Hx & \frac{\lambda}{s} &= \frac{m^3 \langle \sigma v \rangle}{T^3 H(T=m)} \\ &= 2n_2 \langle \sigma v \rangle \frac{t}{t+2} n_{\sigma,eq}(t) & T &\propto a^{-1} & \langle \sigma v \rangle &= \frac{\lambda T^3 H(T=m)}{s m^3} \\ &= 2n_2 \langle \sigma v \rangle \frac{\frac{m_\sigma}{T}}{\frac{m_\sigma}{T}+2} n_{\sigma,eq}(t) \end{aligned}$$

$$T^3 \frac{dY}{dt} = 2n_2 \frac{\lambda T^3 H(T=m)}{s m^3} \frac{\frac{m_\sigma}{T}}{\frac{m_\sigma}{T}+2} n_{\sigma,eq}(t)$$

$$\boxed{\frac{dY}{dx} = \lambda \times h(x) Y_{eq}(x)}$$

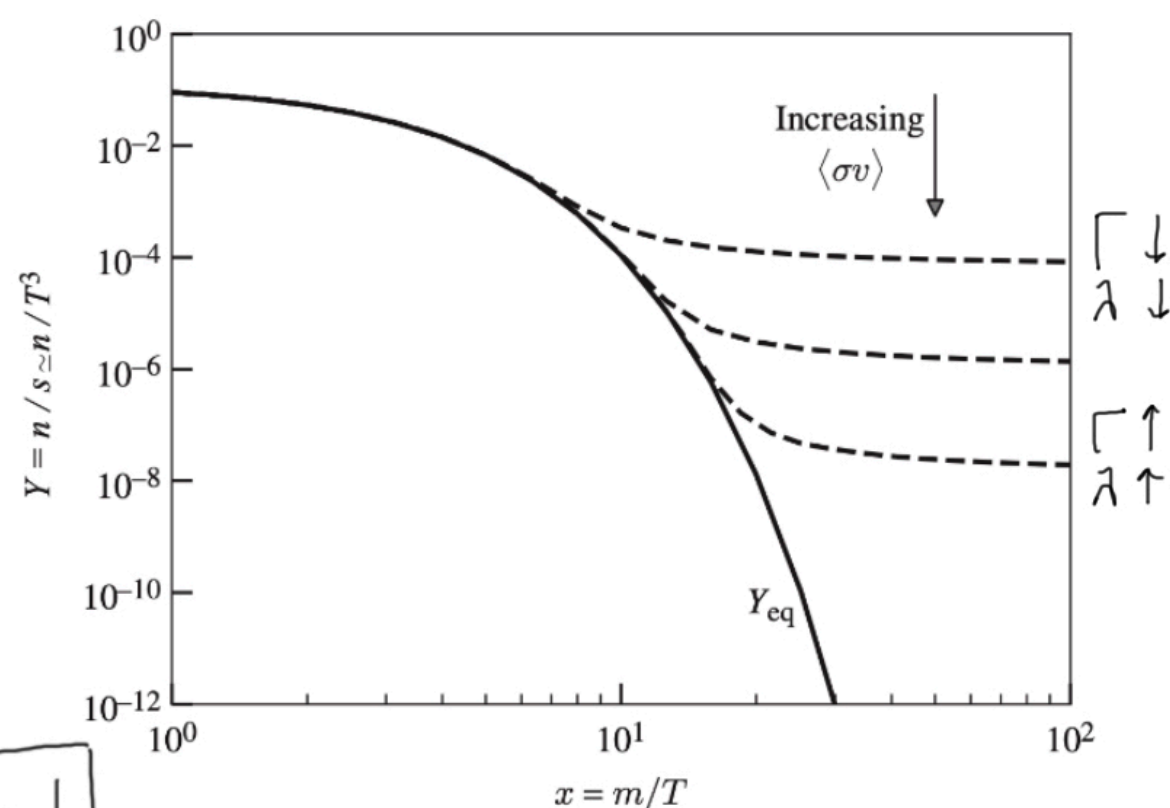


FIG 6.1

Freezeout of particle species. The  $y$ -axis is the comoving number-density-to-entropy ratio  $Y \equiv n/s$ , which is proportional to the comoving number density of particles ( $na^3$ ). The  $x$ -axis is the mass of the particle divided by the temperature of the universe,  $x = m/T$ ; basically, time flows to the right. The freezeout abundance depends on the annihilation cross-section of the particle multiplied by its velocity; the more efficient annihilation is (larger  $\langle \sigma v \rangle$ ), the lower the final abundance. The three dashed curves correspond respectively to cases  $\lambda = 10^5, 10^7, 10^9$ ; see text for details.

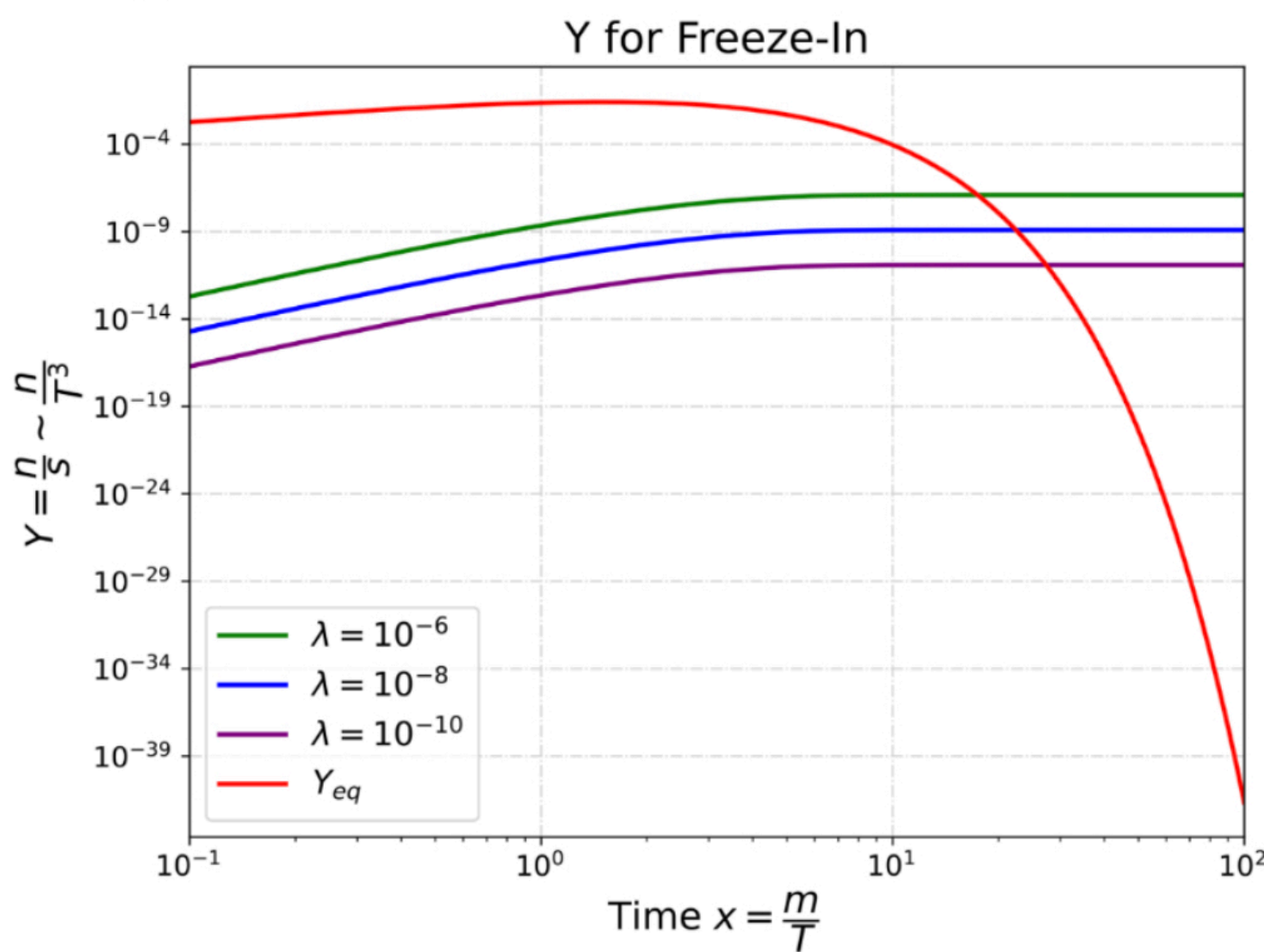
#### b.) Plot $Y(x)$

$$\frac{dY}{dx} = \lambda_1 x h(x) Y_{eq}(x), \text{ where } Y_{eq} \equiv \frac{n_{eq}}{T^3} = \frac{n_{eq} x^3}{M^3}$$

$$\frac{dY}{dx} = \lambda_1 x \left( \frac{x}{x+2} \right) \frac{n_{eq} x^3}{M^3} \quad \left[ \begin{array}{l} x \ll 1 \quad Y_{eq} \approx 0.1 \\ x \gg 1 \quad Y_{eq} \sim \left( \frac{x}{2\pi} \right)^{\frac{3}{2}} e^{-x} \end{array} \right]$$

$$\int_{10^{-20}}^{100} dY = \lambda_1 \int_{0.01}^{100} \left( \frac{x^2}{x+2} \right) \left( \frac{x}{2\pi} \right)^{\frac{3}{2}} e^{-x} dx \quad \leftarrow$$

I integrated and plotted this in my code file. Here's the result:



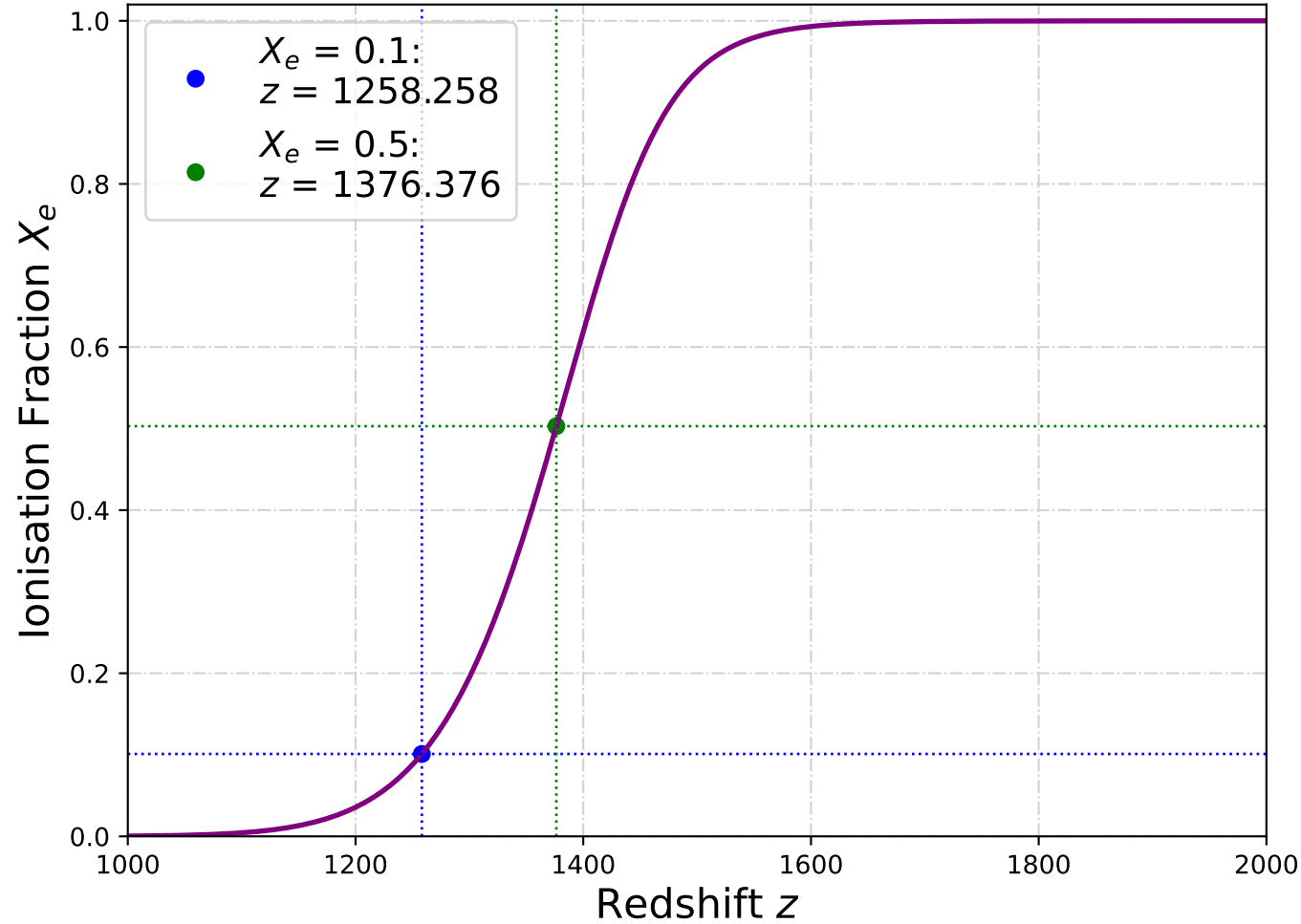
#### c.) Difference between freeze-out and freeze-in:

In terms of time dependence, we see that for freeze-out the  $Y_\chi$  curves decrease before levelling out, whereas for freeze-in the  $Y_\chi$  curves increase before levelling out.

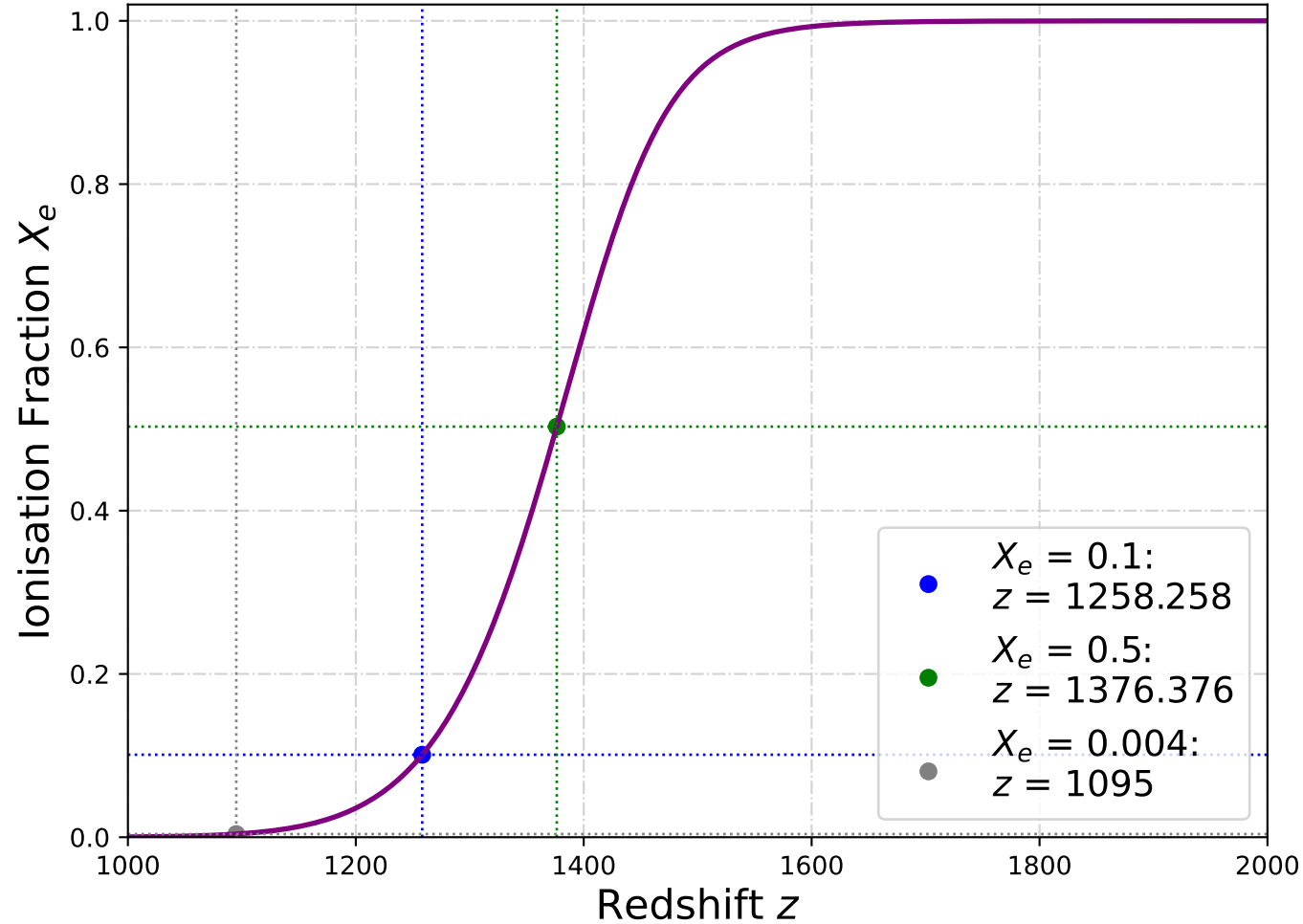
Additionally, the  $Y_{eq}$  curve drops off much earlier in time for freeze-in than freeze-out.

Regarding the impact of the reaction rate  $\Gamma$  as this question asks, in Huterer 6.1 to the left for freeze-out, the greater reaction rates are lower on the plot, meaning that annihilation efficiency correlates with lower final abundances. The opposite is actually true of the freeze-in scenario, as our higher- $\lambda$  curves (higher reaction rate  $\Gamma$ ) are actually at higher  $Y_\chi$  values.

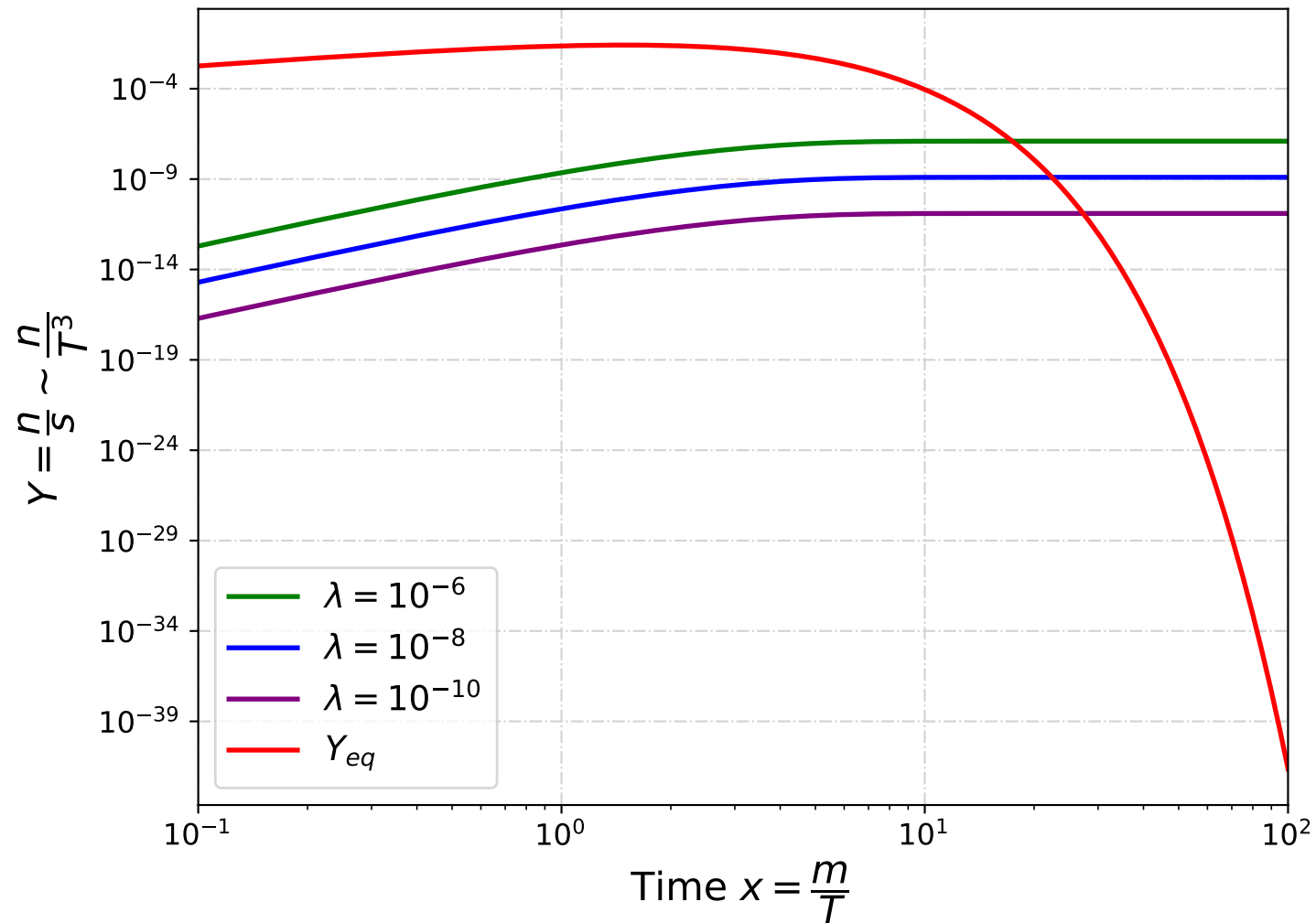
# Ionisation Fraction over Redshift



# Ionisation Fraction over Redshift



# Y for Freeze-In



# ASTR 600 - Cosmology

## HW 4

Iver Warburton

October 2023

```
In [1]: # Our lovely imports

import numpy as np
import matplotlib.pyplot as plt
from scipy.integrate import quad
from scipy.special import zeta
```

## Problem 1

part ii.)

```
In [2]: # Constants

# Masses
m_p = 938.272 # MeV/c^2
m_n = 940.6 # MeV/c^2
m_e = 0.511 # MeV/c^2
m_phot = 0
m_H = 938.8 #939.0 # MeV/c^2

# Number densities
#n_p =
#n_n =
#n_e =
#n_phot =

#n_H = n_p - n_n
#n_He = (1/2)*n_n
#n_B = n_p + n_H

# Defined
E_I = 0.0000136 # MeV, so 13.6 eV #m_p + m_e - m_H
print(E_I)
eta = 6 * 10**-10 # ____ #n_B/n_gamma
```

1.36e-05

```
In [3]: z = np.linspace(1000, 2000, 1000)
```

```
T_o = 2.3525 * 10**-10 #MeV #2.73 K ###0.8 MeV  
T = T_o * (1 + z)
```

```
In [4]: # Build the solution function
```

```
def build_Xe_solution(T, eta, E_I, m_e):  
    f = ((2*zeta(3))/(np.pi**2)) * eta * (((2*np.pi*T)/m_e)**(3/2)) *  
    X_e = (-1 + np.sqrt(1 + (4*f))) / (2*f)  
    return X_e
```

```
In [5]: # Plot results
```

```
def Xe_solution_plot(z, X_e):  
  
    ## PLOT ##  
    # Set figure  
    fig, ax = plt.subplots(figsize=(8,6))  
  
    # Plot  
    plt.plot(z, X_e, ls='-', color='purple', lw=2)  
  
    ## POINTS ##  
    ## For X_e = 0.1 ##  
    # Add point  
    index_Xe01 = np.where(X_e >= 0.1)[0][0]  
    Xe01 = X_e[index_Xe01]  
    z_Xe01 = z[index_Xe01]  
    print("z at X_e = 0.1:", z_Xe01.round(3))  
    plt.scatter(z[index_Xe01], X_e[index_Xe01], color='blue', label='$X_e = 0.1$')  
  
    # Add pointer lines  
    plt.axhline(Xe01, ls=':', c='blue', lw=1)  
    plt.axvline(z_Xe01, ls=':', c='blue', lw=1)  
  
    ## For X_e = 0.5 ##  
    # Add point  
    index_Xe05 = np.where(X_e >= 0.5)[0][0]  
    Xe05 = X_e[index_Xe05]  
    z_Xe05 = z[index_Xe05]  
    print("z at X_e = 0.5:", z_Xe05.round(3))  
    plt.scatter(z[index_Xe05], X_e[index_Xe05], color='green', label='$X_e = 0.5$')  
  
    # Add pointer lines  
    plt.axhline(Xe05, ls=':', c='green', lw=1)  
    plt.axvline(z_Xe05, ls=':', c='green', lw=1)  
  
    ## Part 1e ##  
    ## For z=1095 ##  
    # Add point  
    z_1095 = 1095  
    X_e_1095 = X_e[z_1095]  
    print("X_e at z=1095:", X_e_1095.round(3))  
    plt.scatter(z_1095, X_e_1095, color='red', label='$z = 1095$')  
    plt.axvline(z_1095, ls=':', c='red', lw=1)
```



```

index_z1095 = np.where(z >= 1095)[0][0]
z_1095 = z[index_z1095]
Xe_z1095 = X_e[index_z1095]
print("X_e at z = 1095:", Xe_z1095.round(3))
plt.scatter(z[index_z1095], X_e[index_z1095], color='grey', label='

# Add pointer lines
plt.axhline(Xe_z1095, ls=':', c='grey', lw=1)
plt.axvline(z_1095, ls=':', c='grey', lw=1)
## End of part 1e addition ##

## LABELLING ##
# Label axes
plt.xlabel('Redshift $z$', fontsize=16)
plt.ylabel('Ionisation Fraction $X_e$', fontsize=16)

# Increase axis numbering text size
ax.tick_params(axis='both', which='major', labelsize=10)
ax.tick_params(axis='both', which='minor', labelsize=10)

# Plot features
plt.xlim(1000, 2000)
plt.ylim(0, 1.02)
plt.legend(fontsize=14)
plt.grid(True, color='lightgrey', ls='-.')
plt.title('Ionisation Fraction over Redshift', fontsize=18)

## SAVE ##
# Save and show
plt.savefig("HW4Q1ePlot.pdf", format="pdf", bbox_inches="tight", ov
plt.show()

```

```
In [6]: Xe_solution_plot(z, build_Xe_solution(T, eta, E_I, m_e))
```

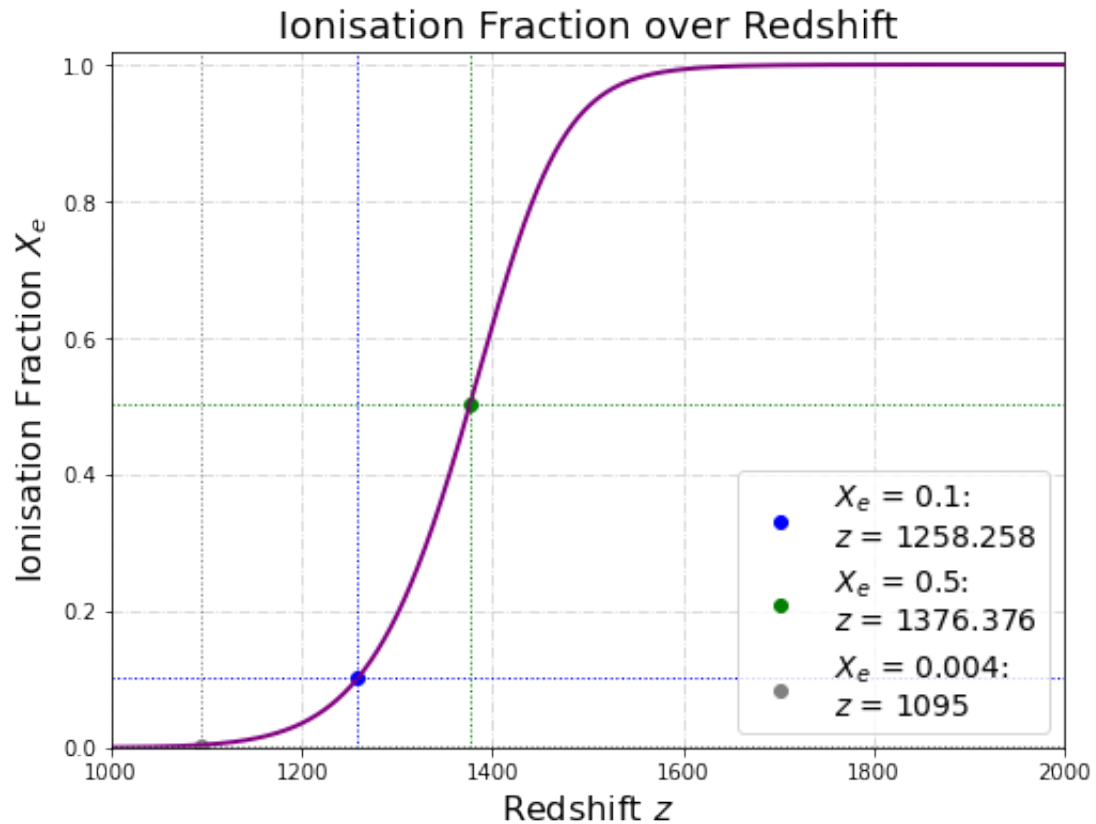
z at  $X_e = 0.1$ : 1258.258

z at  $X_e = 0.5$ : 1376.376

$X_e$  at  $z = 1095$ : 0.004

/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel\_12415/3203391629.py:74: MatplotlibDeprecationWarning: savefig() got unexpected keyword argument "overwrite" which is no longer supported as of 3.3 and will become an error in 3.6

plt.savefig("HW4Q1ePlot.pdf", format="pdf", bbox\_inches="tight", overwrite=True)



```
build_Xe_solution(T, eta, E_I, m_e)
```

```
In [7]: zeta(3)
```

```
Out[7]: 1.2020569031595942
```

```
In [ ]:
```

```
In [ ]:
```

```
In [ ]:
```

## Problem 3

part ii.)

### Numerical integration

```
In [8]: # Define function to build integrand
def integrand(x):
    # Integrand expression for comoving distance that was provided in
    dY = ((x**2)/(x+2)) * (x/(2*np.pi))**(3/2) * np.exp(-x)
    return dY
```

```
lambdas = [10**-6, 10**-8, 10**-10]
x = np.linspace(0.01, 100, 1000)
Y_numint_list = {}

for lambda_val in lambdas:
    # Calculate the integral for comoving distance numerically
    Y_numint = lambda_val * quad(integrand, 0.1, 100)[0]
    #Y_numint = lambda_val * Y_numint
    Y_numint_list.append(Y_numint)
    print(np.shape(Y_numint))
```

```
In [30]: lambda1 = 10**-6
lambda2 = 10**-8
lambda3 = 10**-10

x = np.linspace(0.01, 100, 1000)
#Y_numint_list = {}

# Calculate the integral numerically
Y1 = [lambda1 * quad(integrand, 0, x_val)[0] for x_val in x]
Y2 = [lambda2 * quad(integrand, 0, x_val)[0] for x_val in x]
Y3 = [lambda3 * quad(integrand, 0, x_val)[0] for x_val in x]

Yeq = (x/(2*np.pi))**(3/2) * np.exp(-x)

#Y_numint_list.append(Y_numint)
#print(np.shape(Y1))
#print(Y1)
```



```

In [43]: # Plot results
def Y_plot(x, Y1, Y2, Y3):

    ## PLOT ##
    # Set figure
    fig, ax = plt.subplots(figsize=(8,6))

    # Plot
    plt.loglog(x, Y1, ls='-', color='green', lw=2, label='$\lambda = 1$')
    plt.loglog(x, Y2, ls='-', color='blue', lw=2, label='$\lambda = 10$')
    plt.loglog(x, Y3, ls='-', color='purple', lw=2, label='$\lambda = 100$')
    plt.loglog(x, Yeq, ls='-', color='red', lw=2, label='$Y_{eq}$')
    #for Y in Y_list:
        #plt.plot(x, Y, ls='-', color='purple', lw=2)

    ## LABELLING ##
    # Label axes
    plt.xlabel('Time $x = \dfrac{m}{T}$', fontsize=16)
    plt.ylabel('$Y = \dfrac{n}{s} \sim \dfrac{n}{T^3}$', fontsize=16)

    # Increase axis numbering text size
    ax.tick_params(axis='both', which='major', labelsize=12)
    ax.tick_params(axis='both', which='minor', labelsize=12)

    # Plot features
    plt.xlim(0.1, 100)
    #plt.ylim(0, 1.02)
    plt.legend(fontsize=14)
    plt.grid(True, color='lightgrey', ls='-.')
    plt.title('Y for Freeze-In', fontsize=18)

    ## SAVE ##
    # Save and show
    plt.savefig("HW4Q3Plot.pdf", format="pdf", bbox_inches="tight", ov
    plt.show()

```

```
In [44]: Y_plot(x, Y1, Y2, Y3)
```

```
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_12415/629557688.py:36: MatplotlibDeprecationWarning: savefig() got unexpected keyword argument "overwrite" which is no longer supported as of 3.3 and will become an error in 3.6
```

```
plt.savefig("HW4Q3Plot.pdf", format="pdf", bbox_inches="tight", overwrite=True)
```

