

# PHYS 600 : HW 3

## 1. Reviewing the Background

The following are meant to be short problems, just to review key parts of the homogeneous expansion of the Universe. If not specified, assume a cosmology with  $\Omega_{m,0} = 0.3$ ,  $\Omega_{\Lambda} = 0.7$  and  $h = 0.7$ .

### Density Parameters

Calculate  $\Omega_m$  and  $\Omega_{\Lambda}$  at  $z = 0.5$ .

### Luminosity and Angular Diameter distances

Make a plot of the luminosity and angular diameter distances as a function of redshift to  $z = 10$ . Compare these (on the same plot) to the luminosity/angular diameter distance for an  $\Omega_{m,0} = 1$  universe.

Make sure to label your axes. Comment on the non-monotonicity of the angular diameter distance. ### Looking Back

At the colloquium last week, the speaker mentioned that a redshift of 2-3 roughly corresponds to 10 billion years ago. Find the redshift that is 10 billion years ago (three significant figures only).

### A $\Lambda$ -dominated Universe

Consider a Universe with only a cosmological constant. Find  $a(t)$ . What is the age of such a Universe?

## 2. Massive Neutrinos

(This is based on a problem from Baumann.)

Our current understanding of neutrino physics suggests that atleast two of the three neutrino species have a non-zero (although small) mass. Let us work out the cosmological implications of this.

- Assuming that the neutrinos are relativistic at decoupling, show that the energy density at decoupling is given by

$$\rho_\nu = \frac{T_\nu^4}{\pi^2} \int d\xi \frac{\xi^2 \sqrt{\xi^2 + m_\nu^2/T_\nu^2}}{e^\xi + 1}$$

where  $T_\nu$  is the neutrino temperature.

- Consider a series expansion for small  $m_\nu/T_\nu$  and show that

$$\rho_\nu \approx \rho_{\nu,0} \left( 1 + \frac{5}{7\pi^2} \frac{m_\nu^2}{T_\nu^2} \right)$$

where  $\rho_{\nu,0}$  is the energy density of massless neutrinos.

- If  $\rho_\nu$  is “significantly” larger than  $\rho_{\nu,0}$  at the epoch of the CMB ( $z \sim 1000$ ), then the neutrinos can affect the CMB. Estimate the smallest neutrino mass detectable in the CMB. Feel free to make assumptions to do this, but explicitly describe these. You may assume that the CMB temperature today is 0.235 meV.
- Assume that one of the neutrino species has a mass of  $m_\nu$ . Estimate the redshift at which the neutrinos go non-relativistic.
- Estimate the number density of neutrinos today. To do so, you can’t just use the non-relativistic integrals from class, since the neutrinos are not in thermal equilibrium anymore. Instead, calculate the number density at neutrino decoupling, and then dilute the density with the expansion of the Universe.
- Show that these neutrinos have a energy density today of

$$\Omega_{\nu,0} h^2 \approx \frac{m_\nu}{94 \text{ eV}}$$

## 3. Measuring the Expansion History with Standard Candles

You are encouraged to use Mathematica or something similar to simplify/assist with the algebra in this problem, if you want. If/when you need to, assume the standard values of the cosmological parameters from the first problem.

Suppose that the standard candle method directly measures the radial comoving distance

$$\chi = \int dz \frac{c}{H_0 [\Omega_{m,0}(1+z)^3 + \Omega_{\Lambda,0} + (1 - \Omega_{m,0} - \Omega_{\Lambda,0})(1+z)^2]^{1/2}}$$

- Develop a series expansion of  $\chi$  as a function of  $z$  expanding about  $z = 0$ . Work to order  $z^3$  (i.e. include those terms).
- To what redshift is this expansion accurate to 10%?
- Explain the following statement - for very low redshift measurements, the only parameter that can be measured is  $H_0$ .
- As you increase the redshift reach (still low redshifts), you gain sensitivity to the other cosmological parameters. However, it turns out that you are largely sensitive to a combination of  $\Omega_{m,0}$  and  $\Omega_{\Lambda,0}$ , and not both individually. What combination and why?
- Consider a survey that measures  $\chi$  to 1% at  $z = 0.01, 0.1, 0.2, 0.3$ . Forecast the errors on  $H_0$ ,  $\Omega_{m,0}$  and  $\Omega_{\Lambda,0}$ .

# ① Reviewing the Background

Assume  $\Omega_{m,0} = 0.3$ ,  $\Omega_{\Lambda} = 0.7$ ,  $h = 0.7$

## i.) Density Parameters

Find  $\Omega_m$  and  $\Omega_{\Lambda}$  at  $z = 0.5$

$$\Omega_m = \frac{\rho_m}{\rho_{crit}} \leftarrow \text{Definition of } \Omega_m$$

$$\Omega_m = \frac{\rho_{m,0}(1+z)^3}{\rho_{crit,0}} \cdot \frac{\rho_{crit,0}}{\rho_{crit}}$$

$$\Omega_m = \frac{\Omega_{m,0}(1+z)^3}{\Omega_{m,0}(1+z)^3 + \Omega_{\Lambda}} \leftarrow \begin{array}{l} \Omega_{\Lambda} = \Omega_{\Lambda,0} = 0.7 \text{ to be used here.} \\ \Omega_{\Lambda} \text{ has no } a\text{-relation in Friedmann.} \\ H_0^2(\Omega_{m,0}a^{-3} + \Omega_{\Lambda,0}) \end{array}$$

$$\Omega_m = \frac{0.3(1+0.5)^3}{0.3(1+0.5)^3 + 0.7}$$

$$\boxed{\Omega_m = 0.59124} \text{ at } z = 0.5$$

$$\Omega_{\Lambda} = \frac{\Omega_{\Lambda,0}}{\Omega_{\Lambda,0} + \Omega_m(1+z)^3} \leftarrow \text{Actually sort of an inverse of my expression for } \Omega_m.$$

$$\Omega_{\Lambda} = \frac{0.7}{0.7 + 0.3(1+0.5)^3}$$

$$\boxed{\Omega_{\Lambda} = 0.40876} \text{ at } z = 0.5$$

$$\begin{array}{l} \text{At } z = 0.5, \\ \Omega_m = 0.59124 \\ \Omega_{\Lambda} = 0.40876 \end{array}$$

Check that  $\Omega_m$  and  $\Omega_{\Lambda}$  add to 1:  
 $\Omega_m + \Omega_{\Lambda} = 1$   
 $0.59124 + 0.40876 = 1$   
 $1 = 1 \checkmark$

## ii.) Luminosity and Angular Diameter Distances

Mostly done in attached Jupyter Notebook  
Plot is also there.

$$d_A = \frac{d_m}{1+z} \leftarrow \text{angular diameter distance}$$

$$d_L = d_m(1+z) \leftarrow \text{luminosity distance}$$

$$d_L = d_A(1+z)^2$$

$$d_m = \frac{c}{H_0} \int \frac{1}{(1+z)E(z)} dz \leftarrow \text{comoving distance}$$

$$E(z) = \sqrt{\Omega_m(1+z)^3 + \Omega_{\Lambda}}$$

$$H^2 = H_0^2 \left[ \right]$$

$$H = H_0 \sqrt{\quad}$$

See Jupyter Notebook.

$$d_m = \frac{c}{H_0} \int \frac{1}{(1+z)\sqrt{\Omega_m(1+z)^3 + \Omega_{\Lambda}}} dz$$

$$d_L = \frac{c(1+z)}{H_0} \int \frac{1}{(1+z)\sqrt{\Omega_m(1+z)^3 + \Omega_{\Lambda}}} dz$$

$$d_A = \frac{c}{H_0(1+z)} \int \frac{1}{(1+z)\sqrt{\Omega_m(1+z)^3 + \Omega_{\Lambda}}} dz$$

$\leftarrow$  Used these in my code.

## iii.) Looking Back

10 billion years ago = what redshift?

According to Ryden 6.2, (For flat +  $\Omega_{\Lambda,0} > 0$ )

Integrating the Friedmann equation gives:

$$H_0 t = \frac{2}{3\sqrt{1-\Omega_{m,0}}} \ln \left[ \left( \frac{a}{a_{m,1}} \right)^{3/2} + \sqrt{1 + \left( \frac{a}{a_{m,1}} \right)^3} \right] \quad (6.28)$$

$$\text{where } a = \frac{1}{1+z} \text{ and } a_{m,1} = \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{1/3} = \left( \frac{0.3}{0.7} \right)^{1/3} = 0.7559$$

$$t = \frac{2}{3H_0\sqrt{1-\Omega_{m,0}}} \ln \left[ \left( \frac{1}{1+z} \cdot \left( \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} \right)^{3/2} + \sqrt{1 + \left( \frac{1}{1+z} \cdot \left( \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} \right)^3} \right]$$

I was doing it on my own and changed my mind, so...

$$\begin{aligned} t &= \frac{2}{3 \cdot 70 \cdot \sqrt{1-0.3}} \ln \left[ \left( \frac{1.32635}{1+z} \right)^{3/2} + \sqrt{1 + \left( \frac{1.32635}{1+z} \right)^3} \right] \\ 10 \times 10^9 &= 0.011383 \ln \left[ \left( \frac{1.32635}{1+z} \right)^{3/2} + \sqrt{1 + \left( \frac{1.32635}{1+z} \right)^3} \right] \\ 8.78493 \times 10^9 &= \ln \left[ \left( \frac{1.32635}{1+z} \right)^{3/2} + \sqrt{1 + \left( \frac{1.32635}{1+z} \right)^3} \right] \\ e^{8.78493 \times 10^9} &= \left( \frac{1.32635}{1+z} \right)^{3/2} + \sqrt{1 + \left( \frac{1.32635}{1+z} \right)^3} \\ &= \sqrt{2.33333 a^3} + \sqrt{1 + 2.33333 a^3} \end{aligned}$$

$\rightarrow$  Wolfram Alpha (using  $x = \frac{a}{a_{m,1}}$ ,  $b = e^{\frac{3}{2}H_0 t \sqrt{1-\Omega_{m,0}}}$ )

$$\text{Result: } x = \frac{(b^2 - 1)^{2/3}}{2^{2/3} b^{2/3}}$$

Substituting:  $\frac{a}{a_{m,1}} = \frac{(e^{3H_0 t \sqrt{1-\Omega_{m,0}}} - 1)^{2/3}}{2^{2/3} e^{H_0 t \sqrt{1-\Omega_{m,0}}}}$  Unit conversions detailed below.\*

$$\frac{1}{1+z} \cdot \left( \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} = \frac{(e^{(3 \cdot 70 \cdot 10 \times 10^9 \sqrt{1-0.3})} - 1)^{2/3}}{2^{2/3} e^{70 \cdot 10 \times 10^9 \sqrt{1-0.3}}}$$

$$\frac{1}{1+z} \cdot \left( \frac{0.7}{0.3} \right)^{1/3} = \frac{(e^{(3 \cdot 0.390313)} - 1)^{2/3}}{2^{2/3} e^{0.390313}}$$

$$\frac{1}{1+z} (1.32635) = 0.726705$$

Ned Wright calculator:

$$z_{10 \text{ Gyr}} \approx 1.856$$

$$\boxed{z = 1.825}$$

It seems that 10 Gyr ago was slightly less than  $z = 2-3$ .

\* Unit conversion work in exponents:

$$\frac{70 \text{ km}}{\text{s Mpc}} \left| \frac{10 \times 10^9 \text{ yr}}{1 \text{ yr}} \right| \left| \frac{\sqrt{0.3}}{\sqrt{0.3}} \right| \left| \frac{1 \text{ Mpc}}{3.086 \times 10^{19} \text{ km}} \right| = 0.390313$$

## iv.) $\Lambda$ -Dominated Universe

Universe with only  $\Lambda$  constant. Find  $a(t)$  and age.

$$\frac{\dot{a}}{a} = H_0 \sqrt{\Omega_{\Lambda,0}} \leftarrow \Omega_{\Lambda,0} = 1 \text{ because } \sum \Omega = 1$$

$$\frac{\dot{a}}{a} = H_0$$

$$\frac{1}{a} \frac{da}{dt} = H_0$$

$$\int_0^1 \frac{1}{a} da = H_0 \int_t^0 dt$$

$$\ln(1) - \ln(0) = H_0(0-t)$$

$$0 - \text{UND} = -H_0 t$$

$$\boxed{H_0 t = \text{UND}}$$

A value for  $a(t)$  cannot exist for this Universe.  
The age would need to be infinitely old.  
Wild question, cool!



## 2) Massive Neutrinos

$$\rho_\nu = \frac{T_\nu^4}{\pi^2} \int \frac{\xi^2 \sqrt{\xi^2 + \frac{m_\nu^2}{T_\nu^2}}}{e^\xi + 1} d\xi$$

$$\rho_\nu = \frac{T_\nu^4}{\pi^2} \int \frac{\frac{p^2 + m^2}{T^2} \sqrt{\frac{p^2 + m^2}{T^2} + \frac{m^2}{T^2}}}{e^{\frac{p}{T}} + 1} dE$$

$$\xi = \frac{p}{T} = \frac{\sqrt{p^2 + m^2}}{T}$$

$$dP = T d\xi$$

### a) Energy density at decoupling

$$\rho(T) = \frac{g}{(2\pi)^3} \int \frac{1}{e^{\frac{p}{T} + \mu} + 1} E(p) d^3p$$

where  $\mu$  is chemical potential and is relevant only when numbers of particles in system is changing  $\rightarrow \mu = 0$

$$= \frac{g}{(2\pi)^3} \iiint \frac{\sqrt{p^2 + m^2}}{e^{\frac{p}{T} + 1} + 1} d^3p$$

where we treat the  $m$  in the exponent in the denominator to be negligibly small ( $m \ll p$ )

$$= \frac{g}{8\pi^3} \iiint \frac{\sqrt{p^2 + m^2}}{e^{\frac{p}{T}} + 1} d^3p$$

$$= 4\pi p^2 \cdot \frac{g}{8\pi^3} \int \frac{\sqrt{p^2 + m^2}}{e^{\frac{p}{T}} + 1} dp$$

$$= \frac{g p^2}{2\pi^2} \int \frac{\sqrt{p^2 + m^2}}{e^{\frac{p}{T}} + 1} dp$$

where  $g = 2$  degrees of freedom

$$= \frac{2p^2}{2\pi^2} \int \frac{\sqrt{\xi^2 T^2 + m^2}}{e^\xi + 1} T d\xi$$

$$= \frac{p^2}{\pi^2} \int T \frac{\sqrt{\xi^2 + \frac{m^2}{T^2}}}{e^\xi + 1} d\xi$$

$$= \frac{p^2}{\pi^2} \int \frac{\xi^2 \sqrt{\xi^2 + \frac{m^2}{T^2}}}{e^\xi + 1} d\xi$$

$$= \frac{1}{\pi^2} \int \frac{\xi^2 T^2 \sqrt{\xi^2 + \frac{m^2}{T^2}}}{e^\xi + 1} d\xi$$

$$\rho(T) = \frac{T_\nu^4}{\pi^2} \int \frac{\xi^2 \sqrt{\xi^2 + \frac{m_\nu^2}{T_\nu^2}}}{e^\xi + 1} d\xi$$

← Energy density at decoupling, as expected.

### b) Show energy density of neutrinos

$$\rho(T) = \frac{T_\nu^4}{\pi^2} \int \frac{\xi^2 \sqrt{1 + \frac{m_\nu^2}{T_\nu^2 \xi^2}}}{e^\xi + 1} d\xi$$

where  $\frac{m_\nu^2}{T_\nu^2 \xi^2} \ll 1$  so we can say:  $\sqrt{1 + \frac{m_\nu^2}{T_\nu^2 \xi^2}} \approx \sqrt{1 + \text{small}}$

So we use:  $(1+x)^n \approx 1 + nx$   
 $1 + \frac{1}{2} \frac{m_\nu^2}{T_\nu^2 \xi^2}$

$$= \frac{T_\nu^4}{\pi^2} \int \frac{\xi^2 + \frac{m_\nu^2}{2T_\nu^2}}{e^\xi + 1} d\xi$$

$$= \frac{T_\nu^4}{\pi^2} \left( \int \frac{\xi^2}{e^\xi + 1} d\xi + \int \frac{\frac{m_\nu^2}{2T_\nu^2}}{e^\xi + 1} d\xi \right)$$

where we integrate from 0 to  $\infty$ .

$$= \frac{T_\nu^4}{\pi^2} \left( \frac{7\pi^4}{120} + \frac{\pi^2 m_\nu^2}{24 T_\nu^2} \right)$$

$$= \frac{120 T_\nu^4}{7\pi^2} \left( 1 + \frac{120 m_\nu^2}{7 \cdot 24 \pi^2 T_\nu^2} \right)$$

$$= \frac{120 T_\nu^4}{7\pi^2} \left( 1 + \frac{5 m_\nu^2}{7 \pi^2 T_\nu^2} \right)$$

$$= \frac{7\pi^2 T_\nu^4}{120} \left( 1 + \frac{5 m_\nu^2}{7 \pi^2 T_\nu^2} \right)$$

where we can roughly call the coefficient as  $\rho_\nu$  (source: Navya)

$$\rho_\nu(T) \approx \rho_{\nu,0} \left( 1 + \frac{5 m_\nu^2}{7 \pi^2 T_\nu^2} \right)$$

← This is the energy density of neutrinos.

### c) Smallest detectable neutrino mass

$$\rho_\nu \gg \rho_{\nu,0}, \text{ so let's say } \rho_\nu = 10 \rho_{\nu,0} \text{ (on the low end)}$$

From previously-determined equation:

$$10 \rho_{\nu,0} = \rho_{\nu,0} \left( 1 + \frac{5 m_\nu^2}{7 \pi^2 T_\nu^2} \right)$$

$$10 = 1 + \frac{5 m_\nu^2}{7 \pi^2 T_\nu^2}$$

$$9 = \frac{5 m_\nu^2}{7 \pi^2 T_\nu^2}$$

$$m_\nu = \sqrt{\frac{7 \cdot 9}{5} \pi^2 T_\nu^2}$$

$$m_\nu = \sqrt{\frac{63}{5} \pi^2} T_\nu$$

$$m_\nu = \sqrt{\frac{63}{5} \pi^2} \left( \frac{4}{11} \right)^{\frac{1}{3}} T_{\text{CMB},0} (1+z)$$

$$m_\nu = \sqrt{\frac{63}{5} \pi^2} \left( \frac{4}{11} \right)^{\frac{1}{3}} 235 \text{ meV}$$

$$m_\nu = 836.5 \text{ meV}$$

$$m_\nu = 0.837 \text{ eV}$$

← Minimum detectable neutrino mass at  $z \sim 1000$ .

From textbook (for present):  
 CMB<sub>0</sub> → 2.73 K, 0.235 meV  
 Ne<sub>0</sub> → 1.97 K, 0.17 meV

From the scaling between  $T_\nu$  and  $T_\gamma$  in class:

$$T_\nu = \left( \frac{4}{11} \right)^{\frac{1}{3}} T_{\text{CMB}}$$

$$T_{\text{CMB},1000} = T_{\text{CMB},0} (1+z)$$

$$T_{\text{CMB},1000} = 0.235 \text{ meV} (1+1000)$$

$$T_{\text{CMB},1000} = 235 \text{ meV}$$

$$T_{\nu,1000} = \left( \frac{4}{11} \right)^{\frac{1}{3}} 235 \text{ meV}$$

### d) Redshift when neutrinos go non-relativistic

Neutrinos go non-relativistic when  $E_{\text{kin}} = E_{\text{mass}}$

$$m_\nu \propto T_\nu \propto \frac{1}{a} \propto 1+z$$

$$m_\nu = \frac{k_B}{c^2} T_\nu$$

$$k_B = 1.38 \times 10^{-23} \frac{\text{J}}{\text{K}}$$

$$m_\nu = \frac{k_B}{c^2} T_{\nu,0} (1+z)$$

$$1+z = \frac{c^2 m_\nu}{k_B T_{\nu,0}}$$

$$= \frac{(3 \times 10^8)^2 \text{ m}^2}{\text{s}^2} \frac{1.602 \times 10^{-19} \text{ J}}{1.38 \times 10^{-23} \text{ m}^2 \text{ kg} \text{ s}^{-2} \text{ K}^{-1} \cdot 1.97 \text{ K}}$$

$$m_\nu = T_\nu$$

$$m_\nu = T_{\nu,0} (1+z)$$

$$m_\nu = 0.17 \text{ meV} (1+z)$$

$$z+1 = \frac{m_\nu}{0.17 \text{ meV}}$$

$$z = 5.88 m_\nu - 1 \frac{1}{\text{meV}}$$

from textbook value cited in ZC solution

← leaving it in terms of  $m_\nu$  because "m<sub>ν</sub>" was the given mass. Though Bauman does state that  $z_{\text{nr}}$  is when  $m_\nu > 0.24 \text{ meV}$ .  $\rightarrow (z=1411 \text{ with this})$

I don't know how I feel about this answer. It was checked by Navya but I feel like the constants should be there for the calculation. I found a paper that says  $z+1 = \frac{m_\nu}{0.528 \text{ meV}}$ , which doesn't match.

My inclination was to do calculations like this but Navya suggested using  $T=m$  rather than  $T < m$ , so I do that here.

### e) Number density of neutrinos today

Energy density at decoupling  $\propto g T^3$

$T$  falls as  $\frac{1}{a}$

Density falls as  $\frac{1}{a^3}$

- Find relation between number density of photons (known) and number density of neutrinos (unknown).

- We found scaling between  $T_\nu$  and  $n_\nu$  today:

$$n_\nu \propto \frac{3}{4} g T_\nu^3$$

$$n_\gamma \propto g T_\gamma^3$$

$$\frac{n_\nu}{n_\gamma} = \frac{\frac{3}{4} T_\nu^3}{T_\gamma^3} \leftarrow T_\nu = \left( \frac{4}{11} \right)^{\frac{1}{3}} T_\gamma$$

$$n_{\nu,0} = \frac{3}{4} \left( \frac{4}{11} \right)^{\frac{1}{3}} T_{\gamma,0}^3 n_{\gamma,0}$$

$$n_{\gamma,0} = 4.107 \times 10^8 \frac{1}{\text{m}^3} \text{ today}$$

$$n_{\nu,0} = \frac{3}{11} \cdot 4.107 \times 10^8 \frac{1}{\text{m}^3}$$

$$n_{\nu,0} = 1.120 \times 10^8 \frac{1}{\text{m}^3}$$

← A solution for  $n_{\nu,0}$ , number density of neutrinos today.

### f) Energy density of neutrinos today

$\Omega h^2$  vs energy density.

Start with energy density of photons. Units:  $\left[ \frac{\text{h}^2 \text{ eV}}{\text{m}^3} \right]$

↳ In textbook:  $1.053 \times 10^{10} \frac{\text{h}^2 \text{ eV}}{\text{m}^2}$

$$\rho = \Omega_{\nu,0} h^2 = m_\nu n_{\nu,0} \leftarrow \text{Definition of density } \rho$$

$$= m_\nu \frac{3}{11} n_{\gamma,0}$$

$$1.053 \times 10^{10} \frac{\text{h}^2 \text{ eV}}{\text{m}^2} = m_\nu \cdot 1.120 \times 10^8 \frac{1}{\text{m}^3}$$

Using  $n_{\gamma,0}$  found above in Ze.

$$\Omega_{\nu,0} h^2 \approx \frac{m_\nu}{94 \text{ eV}}$$

← Dividing these two quantities returns roughly 94 eV in the denominator.

↳ The energy density of neutrinos today, in terms of neutrino mass. This matches the expected expression.

### ③. Measuring Expansion History

$$\chi = \int dz \frac{c}{H_0 [\Omega_{m,0}(1+z)^3 + \Omega_{\Lambda,0} + (1 - \Omega_{m,0} - \Omega_{\Lambda,0})(1+z)^2]^{\frac{1}{2}}}$$

a.) Series expansion of  $\chi(z)$  about  $z=0$ :

$$\chi(0) + z\chi'(0) + \frac{1}{2}z^2\chi''(0)$$

Used Mathematica to expand. Result:

$$\frac{c}{H_0} + \frac{1}{2}z^2 \left( \frac{3c(3\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))^2}{4H_0} - \frac{c(6\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))}{2H_0} \right) - \frac{cz(3\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))}{2H_0}$$

b.) Used Mathematica to integrate analytically. Result:

$$\frac{cz}{H_0} - \frac{3cz^2\Omega_m}{4H_0} + \frac{cz^3(3\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))^2}{8H_0} - \frac{cz^3(6\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))}{12H_0} - \frac{cz^2(1 - \Omega_m - \Omega_\Lambda)}{2H_0}$$

Used Python to integrate numerically. Result:

[Work is in following pages of code]



But my solution was  $z=1.333$ , as detailed in the Jupyter Notebook and plot.

c.) For very low- $z$  measurements, we can really only measure  $H_0$  because in the expansion:

$$\frac{c}{H_0} + \frac{1}{2}z^2 \left( \frac{3c(3\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))^2}{4H_0} - \frac{c(6\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))}{2H_0} \right) - \frac{cz(3\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))}{2H_0}$$

every term has at least one power of  $z$  except for the very first term,  $\frac{c}{H_0}$ .

Similarly in the result of analytical integration:

$$\frac{cz}{H_0} - \frac{3cz^2\Omega_m}{4H_0} + \frac{cz^3(3\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))^2}{8H_0} - \frac{cz^3(6\Omega_m + 2(1 - \Omega_m - \Omega_\Lambda))}{12H_0} - \frac{cz^2(1 - \Omega_m - \Omega_\Lambda)}{2H_0}$$

every term has multiple  $z$ s multiplied together except for the first term.

We are told that  $z$  is very small and when very small things are multiplied together, the result is extremely small. Because of that, all of the terms after the first will basically go to zero.

In both of these cases, the only other measurable value in the first term is  $H_0$ , so that is the measurable term.

d.) How do  $\Omega_{m,0}$  and  $\Omega_{\Lambda,0}$  impact  $d_L$  when redshifts are not tiny?

$$d_L = \frac{cz}{H_0} \left( 1 + z \frac{1-q_0}{2} \right), \text{ where } q_0 = \Omega_{r,0} + \frac{1}{2}\Omega_{m,0} - \Omega_{\Lambda,0} \leftarrow \text{this comes from Ryden (7.45)}$$

$$d_L = \frac{cz}{H_0} \left( 1 + \frac{z}{2} (1 - (\Omega_{r,0} + \frac{1}{2}\Omega_{m,0} - \Omega_{\Lambda,0})) \right)$$

$$= \frac{cz}{H_0} \left( 1 + \frac{z}{2} (1 - \cancel{\Omega_{r,0}} - \frac{1}{2}\Omega_{m,0} + \Omega_{\Lambda,0}) \right)$$

$$d_L = \frac{cz}{H_0} \left( 1 + \frac{z}{2} (1 - \underbrace{\frac{1}{2}\Omega_{m,0} - \Omega_{\Lambda,0}}) \right)$$

$d_L$  is sensitive to other parameters, but when it comes to  $\Omega_{m,0}$  and  $\Omega_{\Lambda,0}$ , those impact  $d_L$  in the combination of  $\Omega_{\Lambda,0} - \frac{1}{2}\Omega_{m,0}$ .

# ASTR 600 - Cosmology

## HW 3

Iver Warburton

October 2023

```
In [1]: # Our lovely imports

import numpy as np
import matplotlib.pyplot as plt
from scipy.integrate import quad
```

## Problem 1

part ii.)

```
In [2]: # Constants

c = 2.99 * 10**5 # km/s
h = 0.7
H_0 = h*100
#Omega_m0 = 0.3
#Omega_L0 = 0.7
z_max = 10
z = np.linspace(0, z_max, 100)
```

```

# Calculate comoving distance

# Define function to build integrand
def integrand_comdist(z, Omegam0, OmegaL0):
    # Integrand expression for comoving distance that was provided in
    # problem set
    #X = c/(H_0*(Omega_m0*(1+z)**3 + Omega_L0 + (1 - Omega_m0 -
    Omega_L0)*(1+z)**2)**(1/2))
    X_com = (c/H_0) * (1/((1+z) * (Omegam0*(1+z)**3 + OmegaL0)**
    (1/2)))
    return X_com
#integrand_comdist

# Function to calculate comoving distance
def calc_comdist(z_vals, Omegam0, OmegaL0):
    comdist_vals = []
    comdist = [quad(integrand_comdist(z_vals, Omegam0, OmegaL0), 0,
    z_val)[0] for z_val in z]
    comdist_vals.append(comdist)
    return comdist

```

```

# Function to calculate comoving distance
def calc_comdist(z, Omegam0, OmegaL0):
    #X_com_integrand = (c/H_0) * (1/((1+z) * (Omegam0*(1+z)**3 +
    OmegaL0)**(1/2)))
    comdist_vals = []
    for z_val in z:
        comdist = quad(X_com_integrand, 0, z_val, args=(Omegam0,
    OmegaL0))[0]# for z_val in z]
        comdist_vals.append(comdist)
    return comdist_vals

```

```

In [3]: # Calculate comoving distance

# Define function to build integrand
def integrand_comdist(z, Omegam0, OmegaL0):
    # Integrand expression for comoving distance that was provided in
    #X = c/(H_0*(Omega_m0*(1+z)**3 + Omega_L0 + (1 - Omega_m0 - Omega_
    X_com = (c/H_0) * (1/((1+z) * (Omegam0*(1+z)**3 + OmegaL0)**(1/2))
    return X_com
#integrand_comdist

# Function to calculate comoving distance
def calc_comdist(z_vals, Omegam0, OmegaL0):
    #comdist_vals = []
    comdist = np.array([quad(integrand_comdist, 0, z_val, args=(Omegam
    #comdist_vals.append(comdist)
    return comdist

```



```

# Function to calculate luminosity distance
def calc_lumdist(z_vals, Omegam0, OmegaL0): #, comdist): #z_vals, Omegam0, OmegaL0):
    comdist = calc_comdist(z_vals, Omegam0, OmegaL0)
    lumdist_vals = []
    for z in z_vals:
        lumdist = comdist * (1+z)
        lumdist_vals.append(lumdist)
    return lumdist

```

```

In [4]: # Function to calculate luminosity distance
def calc_lumdist(z_vals, Omegam0, OmegaL0): #, comdist): #z_vals, Omegam0, OmegaL0):
    comdist = calc_comdist(z_vals, Omegam0, OmegaL0)
    #lumdist_vals = []
    #for z in z_vals:
    lumdist = comdist * (1+z_vals)
    #lumdist_vals.append(lumdist)
    return lumdist

```

```

# Function to calculate angular distance
def calc_angdist(z_vals, Omegam0, OmegaL0): #comdist): #z_vals, Omegam0, OmegaL0):
    comdist = calc_comdist(z_vals, Omegam0, OmegaL0)
    angdist_vals = []
    for z in z_vals:
        angdist = comdist / (1+z)
        angdist_vals.append(angdist)
    return angdist_vals

```

```

In [5]: # Function to calculate angular distance
def calc_angdist(z_vals, Omegam0, OmegaL0): #comdist): #z_vals, Omegam0, OmegaL0):
    comdist = calc_comdist(z_vals, Omegam0, OmegaL0)
    #angdist_vals = []
    #for z in z_vals:
    angdist = comdist / (1+z_vals)
    #angdist_vals.append(angdist)
    return angdist

```

```

In [6]: # Plot results
def dists_plot(z_vals):

    ## SETUP ##
    # Set figure
    fig, ax = plt.subplots(figsize=(10,8)) #(8,6))

    Omegam0_03 = 0.3
    Omegam0_1 = 1.0
    OmegaL0_07 = 0.7
    OmegaL0_0 = 0

    # Calculate distances

```

```

lumdist_vals03 = calc_lumdist(z_vals, Omegam0_03, OmegaL0_07)
lumdist_vals1 = calc_lumdist(z_vals, Omegam0_1, OmegaL0_0)

comdist_vals03 = calc_comdist(z_vals, Omegam0_03, OmegaL0_07)
comdist_vals1 = calc_comdist(z_vals, Omegam0_1, OmegaL0_0)

angdist_vals03 = calc_angdist(z_vals, Omegam0_03, OmegaL0_07)
angdist_vals1 = calc_angdist(z_vals, Omegam0_1, OmegaL0_0)

## PLOTTING ##
# Plot distances
plt.plot(z_vals, np.log10(lumdist_vals03), label='$d_L$ for $\Omega_{m,0} = 0.3$', ls='-', color='lightblue', lw=2)
plt.plot(z_vals, np.log10(lumdist_vals1), label='$d_L$ for $\Omega_{m,0} = 1$', ls=':', color='blue', lw=2)

plt.plot(z_vals, np.log10(comdist_vals03), label='$d_m$ for $\Omega_{m,0} = 0.3$', ls='-', color='lightblue', lw=2)
plt.plot(z_vals, np.log10(comdist_vals1), label='$d_m$ for $\Omega_{m,0} = 1$', ls=':', color='blue', lw=2)

plt.plot(z_vals, np.log10(angdist_vals03), label='$d_A$ for $\Omega_{m,0} = 0.3$', ls='-', color='lightblue', lw=2)
plt.plot(z_vals, np.log10(angdist_vals1), label='$d_A$ for $\Omega_{m,0} = 1$', ls=':', color='blue', lw=2)

## LABELLING ##
# Label axes
plt.xlabel('Redshift z', fontsize=16)
plt.ylabel('Log Distances [Mpc]', fontsize=16)

# Increase axis numbering text size
ax.tick_params(axis='both', which='major', labelsize=10)
ax.tick_params(axis='both', which='minor', labelsize=10)

# Plot features
plt.xlim(0, 10)
plt.legend(fontsize=10) #14)
plt.grid(True, color='lightgrey', ls='-.')
plt.title('Various Distance Measures', fontsize=18)

## SAVING ##
# Save and show
plt.savefig("HW3Q1DistancesPlot.pdf", format="pdf", bbox_inches="tight")
plt.show()

```

In [7]: dists\_plot(z)

```

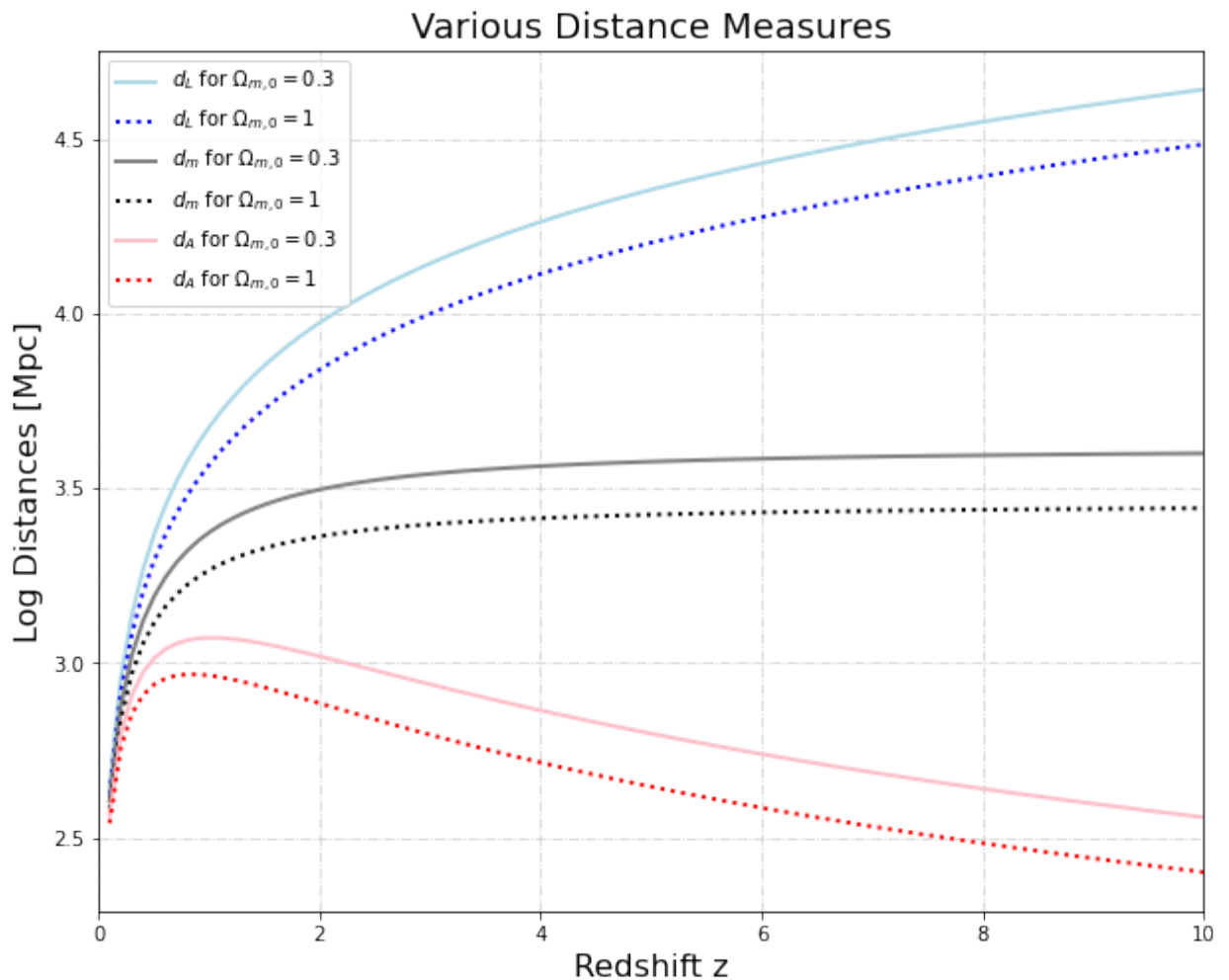
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:26: RuntimeWarning: divide by zero encountered in log10
  plt.plot(z_vals, np.log10(lumdist_vals03), label='$d_L$ for $\Omega_{m,0} = 0.3$', ls='-', color='lightblue', lw=2)
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:27: RuntimeWarning: divide by zero encountered in log10
  plt.plot(z_vals, np.log10(lumdist_vals1), label='$d_L$ for $\Omega_{m,0} = 1$', ls=':', color='blue', lw=2)

```

```

/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:29: RuntimeWarning: divide by zero encountered in log10
  plt.plot(z_vals, np.log10(comdist_vals03), label='$d_m$ for $\Omega_{m,0} = 0.3$', ls='-', color='grey', lw=2)
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:30: RuntimeWarning: divide by zero encountered in log10
  plt.plot(z_vals, np.log10(comdist_vals1), label='$d_m$ for $\Omega_{m,0} = 1$', ls=':', color='black', lw=2)
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:32: RuntimeWarning: divide by zero encountered in log10
  plt.plot(z_vals, np.log10(angdist_vals03), label='$d_A$ for $\Omega_{m,0} = 0.3$', ls='-', color='pink', lw=2)
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:33: RuntimeWarning: divide by zero encountered in log10
  plt.plot(z_vals, np.log10(angdist_vals1), label='$d_A$ for $\Omega_{m,0} = 1$', ls=':', color='red', lw=2)
/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel_32956/1624
047874.py:54: MatplotlibDeprecationWarning: savefig() got unexpected
keyword argument "overwrite" which is no longer supported as of 3.3 a
nd will become an error in 3.6
  plt.savefig("HW3Q1DistancesPlot.pdf", format="pdf", bbox_inches="ti
ght", overwrite=True)

```



This plot very closely matches the prediction plot I drew before coding this notebook.

- Luminosity distance is simply the comoving distance increased by a factor of  $(1 + z)$ , so this distance increases a little bit more quickly as  $z$  increases.
- Angular diameter distance is certainly the most curious result of these, as it increases with  $z$  and then suddenly slopes back down, but it's entirely possible to come up with this concept before plotting. This behaviour has to do with the way that angular sizes behave in an expanding Universe. Earlier in the process of expansion, distant objects were much closer to us and therefore covered a larger angular area on the sky. The images we see of these objects are from light that left the objects when they were much closer and larger on our sky, so this turnover point is closely related to the expansion rate of the Universe.

## Problem 3

```
In [8]: # Constants
c = 2.9979 * 10**8 # m/s
h = 0.7
H_0 = 100 * h # units of km s-1 Mpc-1
z_max = 2
z = np.linspace(0, z_max, 100)
Omega_m0 = 0.3
Omega_L0 = 0.7
```

## Analytical integration

Completed in Mathematica and result is rewritten here to be plotted.

```
In [9]: # Expression for result of analytical integration
X_analint = ((c * z)/(H_0)) - ((3 * c * z**2 * Omega_m0)/(4 * H_0)) +
```

## Numerical integration

```
In [10]: # Define function to build integrand
def integrand(z):
    # Integrand expression for comoving distance that was provided in
    X = c/(H_0*(Omega_m0*(1+z)**3 + Omega_L0 + (1 - Omega_m0 - Omega_L0)))
    return X
```



In [11]:

```
# Calculate the integral for comoving distance numerically
X_numint = [quad(integrand, 0, z_val)[0] for z_val in z]
#print(np.shape(X_numint))
```

In [12]: # Plot results

```
def comdist_plot(z, results_numint, results_analint):

    ## PLOT ##
    # Set figure
    fig, ax = plt.subplots(figsize=(8,6))

    # Plot distances
    #plt.plot(z, results_numint, label='Numerical', ls='-', color='lightblue')
    #plt.plot(z, results_analint, label='Analytical', ls='-', color='black')
    # Plot distances normalised to numerical result
    plt.plot(z, results_numint/np.max(results_numint), label='Numerical', color='black')
    plt.plot(z, results_analint/np.max(results_numint), label='Analytical', color='black')

    # Plot difference between the calculations
    difference = results_numint/np.max(results_numint) - results_analint/np.max(results_numint)
    plt.plot(z, difference, label='Difference', ls='--', color='black')

    ## POINTS ##
    # Add point
    index_diff10 = np.where(difference >= 0.1)[0][0]
    diff10 = difference[index_diff10]
    z_diff10 = z[index_diff10]
    print("z at 10% difference:", z_diff10.round(3))
    plt.scatter(z[index_diff10], difference[index_diff10], color='red')

    # Add pointer lines
    plt.axhline(0.1, ls=':', c='red', lw=1)
    plt.axvline(z_diff10, ls=':', c='red', lw=1)

    ## LABELLING ##
    # Label axes
    plt.xlabel('Redshift z', fontsize=16)
    plt.ylabel('Comoving Distance', fontsize=16)

    # Increase axis numbering text size
    ax.tick_params(axis='both', which='major', labelsize=10)
    ax.tick_params(axis='both', which='minor', labelsize=10)

    # Plot features
    plt.xlim(0, 2)
```

```
plt.ylim(0, 1)
plt.legend(fontsize=14)
plt.grid(True, color='lightgrey', ls='-.')
plt.title('Comoving Distance Calculations', fontsize=18)

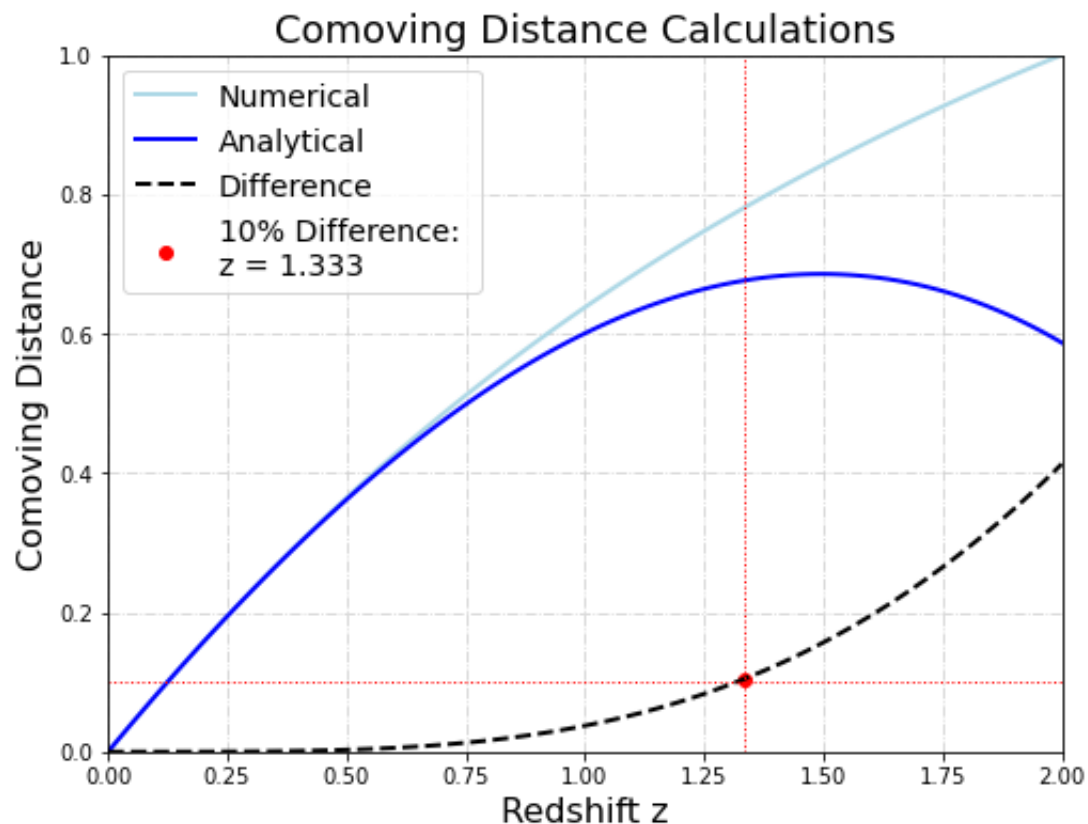
## SAVE ##
# Save and show
plt.savefig("HW3Q3Plot.pdf", format="pdf", bbox_inches="tight", ov
plt.show()
```

In [13]: *# Create the plot*  
comdist\_plot(z, X\_numint, X\_analint)

z at 10% difference: 1.333

/var/folders/1w/ktxtfrr91bj5bztz50dqm0fr0000gn/T/ipykernel\_32956/3647817420.py:52: MatplotlibDeprecationWarning: savefig() got unexpected keyword argument "overwrite" which is no longer supported as of 3.3 and will become an error in 3.6

```
plt.savefig("HW3Q3Plot.pdf", format="pdf", bbox_inches="tight", over
write=True)
```



The analytical expansion is only accurate to 10% out to redshift  $z = 1.333$ . After that point, the expansion result curves sharply downward and becomes extremely unreliable. It's an excellent match out to roughly  $z = 0.6$ , so this is the realm in which it could be extremely useful to calculate comoving distance.

Oops

