

Lecture fys4480,
September 1,
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ansatz

$$\underline{\phi}_0(\vec{x}_1 \vec{x}_2, \dots \vec{x}_N; d_0 \dots d_{N-1})$$
$$= \frac{1}{\sqrt{N!}} \begin{vmatrix} \varphi_{d_0}(\vec{x}_1) & \varphi_{d_0}(\vec{x}_2) & \dots & \varphi_{d_0}(\vec{x}_N) \\ \varphi_{d_1}(\vec{x}_1) & , & & \\ ; & & , & \\ ; & & & \\ \varphi_{d_{N-1}}(\vec{x}_1) & & & \varphi_{d_{N-1}}(\vec{x}_N) \end{vmatrix}$$

$$= \frac{1}{\sqrt{N!}} \sum_P (-)^P \hat{P} \underbrace{\varphi_{d_0}(\vec{x}_1) \varphi_{d_1}(\vec{x}_2) \dots \varphi_{d_{N-1}}(\vec{x}_N)}_{\phi_H}$$

$$\hat{A} = \frac{1}{N!} \sum_P (-1)^P \hat{P}$$

$$N=2 \quad \hat{A}_2 = \frac{1}{2} (1 - \hat{P}_{12})$$

$$\langle \Phi_0 | H_0 | \Phi_0 \rangle = \sum_{\alpha_i=0}^{N-1} \sum_{i=1}^N$$

$$\hat{h}_0(\vec{x}_i) \varphi_\alpha(\vec{x}_i) = \varepsilon_\alpha \varphi_\alpha(\vec{x}_i)$$

↓

$$\hat{\sigma}_1 \quad \varepsilon_{\alpha i} \rightarrow \langle \varphi_{\alpha j} | \hat{\sigma}_1 | \varphi_{\alpha i} \rangle$$

$$\langle \Phi_i | H_0 | \Phi_0 \rangle$$

Φ_i : differs by one sp-state

$$\langle \vec{J}_0 | H_1 | \vec{J}_0 \rangle$$

$$= \sum_{i < j} \int d\vec{x}_1 d\vec{x}_2 \dots d\vec{x}_i \dots d\vec{x}_j \dots d\vec{x}_N$$

$$\times \varphi_{d_0}^*(\vec{x}_1) \dots \varphi_{d_{N-1}}^*(\vec{x}_N) v(x_i x_j)$$

$$\sum_P (-1)^P \hat{P} \cdot (\varphi_{d_0}(\vec{x}_1) \dots \varphi_{d_{N-1}}^*(\vec{x}_N))$$

$$P = 0$$

$$\langle d_i d_j | v(x_i x_j) \rangle$$

$$\mathcal{P} = 1$$

$$\sum_{i,j} \langle \alpha_i^\dagger \alpha_j^\dagger \nu | \alpha_j \alpha_i \rangle$$

$$\mathcal{P} \geq 2$$

$$i=1 \quad j=2$$

$$P_{12} P_{34}$$

$$- \int dx_1 dx_2 \varphi_{\alpha_0}^*(\vec{x}_1) \varphi_{\alpha_1}^*(\vec{x}_2) \nu(x_1 x_2)$$

$$\times \varphi_{\alpha_0}(\vec{x}_2) \varphi_{\alpha_1}(\vec{x}_1) \int dx_3 dx_4$$

$$\varphi_{\alpha_2}^*(\vec{x}_3) \varphi_{\alpha_3}^*(\vec{x}_4) \varphi_{\alpha_2}(\vec{x}_4) \varphi_{\alpha_3}(\vec{x}_3)$$

$$= 0$$

$$\langle \Phi_0 | H_1 | \Phi_0 \rangle =$$

$$\sum_{i < j} \left\{ \langle d_i d_j | v | d_i d_j \rangle - \langle d_i d_j | v | d_j d_i \rangle \right\}$$

$$= \frac{1}{2} \sum_{i < j} \left\{ \quad \quad \quad \quad \right\}$$

$$= \frac{1}{2} \sum_{i < j} \langle d_i d_j | v | d_i d_j \rangle_{AS}$$

$$E_0 = \sum_{i=0}^{N-1} \sum_{i=0}^N \langle d_i | \hat{H} | d_i \rangle + \frac{1}{2} \sum_{i,j} \langle d_i d_j | v | d_i d_j \rangle_{AS}$$

$\hat{H} \Phi_0 \neq \bar{\Phi}_{exact} \Phi_0$

$$\langle \Phi_0 | H_I | \Phi_i \rangle \propto \sum_k \langle \delta_{ek} / v_{k\alpha} \rangle$$

↑
differs by one
single-particle state
only,
 $\varphi_{ek}^{(x_i)} \neq \varphi_{ek}^{(x_j)}$

$$\langle \Phi_0 | H_I | \Phi_i \rangle$$

differ by two
single particle
states

$$\psi_{d_K}(x_i) \wedge \psi_{d_E}(r_j)$$

Ket side

Bra - side

$$\varphi_{d_{K'}}(x_i) \wedge \varphi_{d_{J'}}(x_j)$$

$$\langle \underline{\psi}_C | H_I | \underline{\psi}_N \rangle = \langle \varphi_i \varphi_j | \text{overlap} \rangle_N$$

Coulomb-Slater rule

if we have more than two
sp-states which differ and

if we have at most a two-body interaction

$$\langle \Psi_0 | H_I | \Psi_i \rangle = 0$$

Encode SD as bit strings

— α_9

— .

— .

— .

— .

— .

ϕ .

ϕ α_2

ϕ α_1

ϕ α_0

$N=4$

$|11110\cdots0\rangle$
 $\alpha_0 \alpha_1 \alpha_2 \alpha_3 \cdots \alpha_9$
1111000000

— α_9
 — α_8
 \circ α_7
 \circ α_6
 \circ α_5
 \circ α_4
 — α_3
 — α_2
 — α_1
 — α_0

$$|000011100\rangle = \Phi_1$$

$$\langle \Phi_0 | H_1 | \Phi_1 \rangle = 0$$

compare

$|\Phi_0\rangle$ and $|\Phi_1\rangle$

Second quantization, basic definitions -

$$N=4, \quad d_0 d_1, \dots d_9$$

$|\underline{\Phi}_0\rangle = |d_0 d_1 d_2 d_3\rangle$ compact representation of a SD.

We define a creation operator

$$d^\dagger |0\rangle = |d\rangle$$

↑ vacuum

2 particles

$$|d_B\rangle = d^\dagger_0 d^\dagger_3 |0\rangle$$

$$|\underline{\Phi}_0\rangle = d^\dagger_0 d^\dagger_1 d^\dagger_2 d^\dagger_3 |0\rangle$$

$$\underbrace{|\alpha_0 \alpha_1 \alpha_2 \alpha_3\rangle}_{\text{standard ordering}} = - |\alpha_0 \alpha_2 \alpha_1 \alpha_3\rangle$$

4-particle state

$$|\alpha_0 \alpha_1 \alpha_2 \alpha_3\rangle = a_{\alpha_0}^+ a_{\alpha_1}^+ a_{\alpha_2}^+ a_{\alpha_3}^+ |0\rangle$$

$$a_{\alpha}^+ |\alpha_0 \alpha_1 \alpha_2 \alpha_3\rangle = a_{\alpha}^+ a_{\alpha_0}^+ a_{\alpha_1}^+ a_{\alpha_2}^+ a_{\alpha_3}^+ |0\rangle$$

$$\alpha = \{\alpha_0 \alpha_1 \alpha_2 \alpha_3\}$$



$$a_{\alpha_0}^+ a_{\alpha_0}^+ |0\rangle = 0$$

$$\alpha \neq \{d_0 d_1, d_2 d_3\}$$

$$|\alpha d_0 d_1, d_2 d_3\rangle$$

$$d_0 < d_1 < d_2 < d_3 \dots \langle dd_{-1}$$

$$\alpha = d_5$$

$$|\alpha_0, \alpha_1, \alpha_2, \alpha_3, \alpha_5\rangle$$

- 5-particle

$$|\alpha d_0 d_1, d_2 d_5\rangle = a_\alpha^+ a_{d_0}^\dagger a_{d_1}^+ a_{d_2}^+ a_{d_5}^\dagger |0\rangle$$

$$= - |\alpha d d_1, d_2 d_3\rangle = - a_{d_0}^+ a_\alpha^\dagger a_{d_1}^+ a_{d_2}^+ a_{d_3}^\dagger$$

$$a_{\alpha}^+ a_{\alpha_0}^+ + a_{\alpha_0}^+ a_{\alpha}^+ = 0$$

$$\Rightarrow [a_{\alpha}^+, a_{\alpha_0}^+]_+ = \{a_{\alpha}^+, a_{\alpha_0}^+\}$$

$$= 0 \Rightarrow \{a_{\alpha}^+, a_{\beta}^+\} = 0$$

$$a_{\alpha}^+ |\alpha\rangle = 0$$

Define annihilation operator

$$a_{\alpha} |0\rangle = 0$$

$$a_{\alpha} |\alpha\rangle = |0\rangle$$

$$a_\alpha |\alpha_0 \alpha_1 \alpha_2 \alpha_3 \rangle = ?$$

$$\alpha \in \{\alpha_0 \alpha_1 \alpha_2 \alpha_3\}$$

$$\alpha = \alpha_0$$

$$a_\alpha |\alpha_0 \alpha_1 \alpha_2 \alpha_3 \rangle = |\alpha_1 \alpha_2 \alpha_3 \rangle$$

$$N=4$$

$$N=3$$

$$a_\alpha = (a_\alpha^+)^+$$

hermitian
conjugate

$$\Rightarrow \{a_\alpha, a_\beta\} = 0 \text{ of } a_\alpha^+$$

2 commutation relations

$$\{ a_{\alpha}^{\dagger}, a_{\beta}^{+} \} = \{ a_{\alpha}, a_{\beta} \}$$

$$= 0$$

↑

$$a_{\alpha} a_{\beta} + a_{\beta} a_{\alpha}$$

we need an additional rule

$$\{ a_{\alpha}^{+}, a_{\beta} \} = S_{\alpha \beta}$$

Example

$$| \alpha \rangle = a_{\alpha}^{+} | 0 \rangle$$

$$| \beta \rangle = a_{\beta}^{+} | 0 \rangle$$

mmm

$$\langle \beta | \alpha \rangle = \langle 0 | \alpha_\beta \alpha_\alpha^\dagger | 0 \rangle$$

$$\alpha_\beta \alpha_\alpha^\dagger + \alpha_\alpha^\dagger \alpha_\beta = S_{\alpha\beta}$$

$$\alpha_\beta \alpha_\alpha^\dagger = S_{\alpha\beta} - \alpha_\alpha^\dagger \alpha_\beta$$

$$\langle 0 | \alpha_\beta \alpha_\alpha^\dagger | 0 \rangle = S_{\alpha\beta} \underbrace{\langle 0 | 0 \rangle}_{=1}$$

$$- \langle 0 | \alpha_\alpha^\dagger \alpha_\beta | 0 \rangle$$

$$= S_{\alpha\beta}$$