F45 4980, NOV 11, 2022 Many-body pest theory SE (3) = < \$\_1 H, / \$\_0 > - < \$0 (V/\$0) in particle famalism <\$.|

1<u>\$</u>c> in particle-lide formalism くずしずの>= イナーーー か < i1 i2/v/ú, 12> - < i, 1/2/v/i2i2> = in --- iz +  $\Delta E^{(3)} = 2 \sum_{i \in \mathcal{I}} \left( \partial_{i}^{i} - \partial_{j}^{i} + \partial_{i}^{i} \right)$ = 0---0 + == SE = < \$1/5 v(\$0)  $= \sum_{m \neq 0} \langle \underline{\sharp}_{0} | v | \underline{\sharp}_{m} \rangle \langle \underline{\sharp}_{m} | v | \underline{\sharp}_{0} \rangle$  $\vec{Q} = \sum_{m \neq 0} |\vec{\Phi}_m\rangle \langle \vec{\Phi}_m|$ 

e = Wo-Ho Holemin = Wm / Im> Schematic plet in particle jamalisme (1,) (1z) 1 m) ( Jm) no  $\sum_{ij}$  ae  $\sum_{i,j}$   $\sum_{i,$ 

 $\frac{\sum_{i} \langle i | i \rangle | \omega t \rangle \langle \omega t | v | i v | i \rangle}{\sum_{i} t + \sum_{i} t - \sum_{i} t - \sum_{i} t}$ in particle-hole formatism in fa of is ano the contribation particle forma les un to the fine to the second ph- Janmahsm

Final contribation

contribution (R5) Zake of the state Unlinked diagrams + Pauli' violating diagram -  $\Delta E^{(3)} = \langle \mathcal{F}_0 | v \frac{\mathcal{F}_0}{w_0 - \mathcal{H}_0} v \frac{\mathcal{F}_0}{w_0 - \mathcal{H}_0} v / \mathcal{F}_0 \rangle$ - < \$c/ Ulac > < \$c/ U \(\frac{\partial}{w\_0 - H\_0} \) U \(\frac{\partial}{\partial} \) in ph- Jamalism 0--0 × () unliked dagran 10m) 6 (50)

SE(3) = < \$125 v \$ v 1 \$c> e= Wo-Ho - (\$0| v(\$) v/\$0>  $x < \mathcal{F}_{c} / v / \mathcal{F}_{c} >$ the distributed of the signal this will be canceled by 11 1/2 13 14 1/2 13 14 1-1 1 x 1 1 a f-1 = 2 i, iz is ig < 4,10/4,> Example: Exacise 1, weekgs H = \( \sum\_{\alpha = 1} \in \alpha \

$$-E_{2} \qquad H_{0}|\underline{\mathfrak{H}}\rangle = \mathcal{E}_{1}|\underline{\mathfrak{H}}\rangle$$

$$-E_{1} \qquad H_{0}|\underline{\mathfrak{H}}\rangle = \mathcal{E}_{2}|\underline{\mathfrak{H}}\rangle$$

$$+E_{1} \qquad H_{0}|\underline{\mathfrak{H}}\rangle = \mathcal{E}_{2}|\underline{\mathfrak{H}}\rangle$$

$$+E_{1} \qquad +C|\underline{\mathfrak{H}}\rangle = 0$$

$$+E_{1} \qquad +C|\underline{\mathfrak{$$

BW: 
$$XE_1 = E_1 - E_1$$

$$= \frac{\chi^2}{E_1 - E_2} = 7$$

$$E_1 = E_1 + \frac{\chi^2}{E_1 - E_2} \left( (E_1 - E_1)(E_1 - E_2) - \chi^2 = 0 \right)$$

$$SE_1(BW) = \langle \Phi_0 | H_1 + H_2(E_1 - H_1) \Phi_0 \rangle$$
Diagrammatically in the particle fermalism

$$\frac{1}{2} \frac{1}{1 - \chi} \times \chi \longrightarrow \frac{1}{\sqrt{2}} \left( \frac{1}{2} \frac{\chi}{1 - \chi} \right)$$

$$\lim_{z \to 1} ph \text{ faminalism}$$

$$|\Phi_1| = q_2^{\dagger} q_1 |\Phi_1|$$

$$|\Phi_2| = q_2^{\dagger} |\Phi_1|$$

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$$E_{1} \rightarrow E_{1} \text{ in the expansion}$$

$$\Delta E_{1} = \langle \bar{J}_{1}|H_{1}|\bar{J}_{1}\rangle + \langle \bar{J}_{2}|H_{1}|\Phi_{1}\rangle \langle \bar{J}_{2}|H_{1}|\Phi_{2}\rangle \langle \bar{J}_{2}|H_{1}|\Phi_{1}\rangle \langle \bar{J}_{2}|H_{1}|\Phi_{2}\rangle \langle \bar{J}_{2}|H_{1}|\Phi_{2}\rangle$$

$$E = \frac{1}{2} \left[ (\varepsilon_{1} + \varepsilon_{2}) \pm ((\varepsilon_{1} - \varepsilon_{2})^{2} + 4 \lambda^{2} \right]$$

$$= \frac{1}{2} \left[ (\varepsilon_{1} + \varepsilon_{2} - (\varepsilon_{2} - \varepsilon_{1})) + \frac{4 \lambda^{2}}{(\varepsilon_{2} - \varepsilon_{1})^{2}} \right]$$

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$$= \frac{1}{2} \left[ (\varepsilon_{1} - \varepsilon_{2}) + \frac{\lambda^{2}}{(\varepsilon_{2} - \varepsilon_{1})^{2}} \right]$$

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$$= \frac{\lambda^{2}}{\varepsilon_{1} - \varepsilon_{2}} \left[ (\varepsilon_{1} - \varepsilon_{2})^{2} + \cdots \right]$$

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$$\Delta E_{1}^{(4)}(es) = \langle H_{1} \frac{Q}{e} H_{1} \frac{Q}{e} H_{1} \frac{Q}{e} H_{1} \rangle$$

$$e = \varepsilon_{1} - H_{0}$$

$$- \langle H_{1} \frac{Q}{e} + H_{2} \frac{Q}{e} + H_{1} \rangle$$

$$- \langle H_{1} \frac{Q}{e} + H_{2} \frac{Q}{e} + H_{1} \rangle$$

$$+ \langle H_{1} \frac{Q}{e} + H_{2} \frac{Q}{e} + H_{2} \rangle$$

$$- \langle H_{1} \frac{Q}{e} + H_{2} \frac{Q}{e} + H_{2} \rangle$$

$$- \langle H_{1} \frac{Q}{e} + H_{2} \frac{Q}{e} + H_{3} \rangle$$

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$$- \langle H_{1} \frac{Q}{e}$$

$$\frac{\chi^{2}}{\epsilon_{1}-\epsilon_{2}} < \underline{\Phi}_{1} | \underline{H}_{\underline{e}} | \underline{\Phi}_{1} | \underline{\Phi}_{1} \rangle$$

$$= \frac{\chi^{2}}{\epsilon_{1}-\epsilon_{2}} < \underline{\Phi}_{1} | \underline{H}_{1} | \underline{\Phi}_{2} \rangle < \underline{\Phi}_{2} | \underline{H}_{1} | \underline{\Phi}_{1} \rangle$$

$$= \frac{\chi^{2}}{(\epsilon_{1}-\epsilon_{2})^{2}} = \frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{3}} = -\frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{3}}$$

$$\Delta E_{1}^{(2)} + \Delta E_{1}^{(4)} = \frac{\lambda^{2}}{(\epsilon_{1}-\epsilon_{2})^{2}} - \frac{\lambda^{2}}{(\epsilon_{2}-\epsilon_{1})^{2}} = \frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{2}} - \frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{2}} = \frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{2}} = \frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{2}} - \frac{\chi^{2}}{(\epsilon_{2}-\epsilon_{1})^{2}} = \frac{\chi^{2}}{($$

$$= \langle \underline{\mathcal{F}}_{1} | \mathcal{H}_{1} | \underline{\mathcal{F}}_{2} \rangle \langle \underline{\mathcal{F}}_{1} | \mathcal{H}_{1} | \underline{\mathcal{F}}_{1} \rangle \langle \underline{\mathcal{F}}_{2} | \mathcal{F}_{2} | \mathcal{H}_{1} \rangle \langle \underline{\mathcal{F}}_{2} | \mathcal{F}_{2} | \mathcal{F}_{2} \rangle \langle \underline{\mathcal{F}_{2} | \mathcal{F}_{2} \rangle$$

 $\frac{1}{2} \int_{0}^{2\pi} \frac{d^{2}}{dt} = \frac{1}{2} \int_{0}^{2\pi} \frac{1}{2} \int_{0}^{$