

**Lecture FYS4480,
October 27, 2023**

Electron-gas in Solid

$$\langle H \rangle = \langle T \rangle + \langle H_I \rangle$$

$$\langle T \rangle = \frac{\sqrt{4\pi^2 k_F^5}}{10m}$$

$$a_0 = \frac{e^2}{2me^2}$$

$$1 \text{ Ry} = \frac{me^4}{2t^2}$$
$$= 13.6 \text{ eV}$$

$$f = \frac{N}{V} = \frac{k_F}{3\pi^2}$$

$$= \frac{3}{4\pi r_s^3}$$

$$\langle T \rangle = N \cdot 2.21 \left(\frac{a_0}{r_s} \right)^2 \cdot 1 \text{ Ry}$$

$$\langle +I_I \rangle = \dots$$

$$8\pi \int_0^{k_F} dP^3 \int_0^P dk \sum_L \left(\frac{k}{P} \right)^{2L+2} \frac{1}{2L+1}$$

$$= 8\pi^2 k_F^4 \sum_L \frac{1}{(2L+1)(2L+3)}$$

$$\sum_L \frac{1}{(2L+1)(2L+3)} = \frac{1}{2} \sum_L \left(\frac{1}{2L+1} - \frac{1}{2L+3} \right)$$

$$= \frac{1}{2} + \frac{1}{2} \sum_{L=0}^{\infty} \left(\frac{1}{2L+3} - \frac{1}{2L+5} \right) = 1/2$$

\Rightarrow

$$\langle H_I \rangle = - \frac{e^2 V K_F^4}{4\pi^3}$$

using a_0 and a_s

$$\langle H_I \rangle = - 0.916 \cdot a_0/h_s \cdot N \cdot 1Ry$$

$$\frac{\langle H \rangle}{N} = \left[\frac{2.21}{(a_s/a_0)^2} - \frac{0.916}{(a_s/a_0)} \right] 1Ry$$

$$\frac{\langle H_I \rangle}{N} = - e^2 g^{-1} \frac{K_F^4}{4\pi^3} \leq \alpha g^{4/5}$$

S = N/V $= - \text{const } g^{1/3}$

→ local density approx for
the exchange energy

$$P = - \left(\frac{\partial E_0^{\text{HF}}}{\partial V} \right)_N$$

$$E_0^{\text{HF}} = \langle \tau \rangle - \text{const} g^{1/3}$$

$$B = -V \left(\frac{\partial P}{\partial V} \right)_N$$

Density functional theory.

$$\begin{aligned} & [t(\vec{r}) + v_{\text{ext}}(\vec{r}) + \int d\vec{r}' \rho(\vec{r}') \\ & \times v(\vec{r}, \vec{r}') + V_{\text{xc}}[S](\vec{r})] \psi_k(\vec{r}) \end{aligned}$$

$$\begin{aligned} & \text{KS = Kohn-sham} \\ & = \epsilon_K \psi_K^{\text{KS}}(\vec{r}) \end{aligned}$$

Hartree - Fock theory

$$[t(\vec{r}) + V_{\text{ext}}(\vec{r}) + u^{\text{HARTREE}}(\vec{r})]$$

$$\times \psi_K(\vec{r}) + \int d\vec{r}' n^{\text{FOCK}}(\vec{r}') \phi_K^{\text{HF}}(\vec{r}')$$

$$= \epsilon_K^{\text{HF}} \psi_K^{\text{HF}}(\vec{r})$$

$$u^{\text{HARTREE}}(\vec{r}) =$$

$$\sum_{i \in F} \int d\vec{r}' |\psi_i(\vec{r}')|^2 / v(\vec{r}, \vec{r}')$$

$$g(\vec{r}) = \sum_{i=1}^n |\psi_i(\vec{r})|^2 = \boxed{\int d\vec{r}' f(\vec{r}') v(\vec{r}, \vec{r}')}}$$

$$U^{FOCK}(\vec{r}) = - \sum_{n \in F} \psi_n^*(\vec{r}') v(\vec{r}, \vec{r}') \times \psi_n(\vec{r})$$

Exchange term.

Cannot write it as something proportional to $\delta(\vec{r})$

$$\left[\langle \Phi_0 | \rho(\vec{r}) | \Phi_0 \rangle \right] = \sum_{i=1}^N |\psi_i(\vec{r})|^2$$

Total energy with an SD

$$\Psi_0 = \frac{1}{\sqrt{N!}} \cdot a_c^+ |0\rangle$$

$c=1$

$$\Psi_0 = \frac{1}{\sqrt{N!}} \begin{vmatrix} \psi_1(\vec{r}_1) & \psi_1(\vec{r}_2) & \dots & \psi_1(\vec{r}_N) \\ \psi_2(\vec{r}_1) & ; & - & - \\ \vdots & ; & - & - \\ \psi_N(\vec{r}_1) & ; & \psi_N(\vec{r}_N) \end{vmatrix}$$

$$\langle \Psi_0 | H | \Psi_0 \rangle = E_0^{HF}$$

$$H | \Psi_0 \rangle$$

$$\neq E_0 | \Psi_0 \rangle$$

$$= \sum_{c \leq F} \langle i | b_{ci} | i \rangle$$

$$+ \frac{1}{2} \sum_{ij} \left\{ \langle i j | v | i j \rangle - \langle i j | h | i j \rangle \right\}$$

$$\frac{1}{Z} \sum_{ij} \langle \psi_j | \omega | \psi_i \rangle =$$

$$\frac{1}{Z} \sum_{ij} \int d\vec{z} \int d\vec{z}' \psi_i^*(\vec{z}) \psi_j^*(\vec{z}') \sigma(\vec{z}, \vec{z}') \\ \times \psi_i(\vec{z}') \psi_j(\vec{z}')$$

$$= \frac{1}{Z} \int d\vec{z} \int d\vec{z}' \rho(\vec{z}) \rho(\vec{z}') \sigma(\vec{z}, \vec{z}')$$

$$\frac{1}{Z} \sum_{ij} \langle \psi_j | \omega | \psi_j \rangle =$$

$$\frac{1}{Z} \sum_{ij} \int d\vec{z} \int d\vec{z}' \psi_i^*(\vec{z}) \psi_j^*(\vec{z}') \sigma(\vec{z}, \vec{z}') \\ \times \psi_i(\vec{z}') \psi_j(\vec{z})$$

HF-energy:

$$\bar{E}_0^{\text{HF}} = \sum_{n \leq F} \varepsilon_n^{\text{HF}} - \frac{1}{2} \sum_{ij} \langle ij | v | ij \rangle$$

$+ \frac{1}{2} \sum_{ij} \langle ij | v | ji \rangle$

$$= \sum_{n \leq F} \varepsilon_n^{\text{HR}} - \frac{1}{2} \int d\vec{r} \int d\vec{r}' f(\vec{r}) f(\vec{r}') \\ \times v(\vec{r}, \vec{r}')$$

$+ \boxed{\bar{E}_{xc}^{\text{HF}}}$

$$\bar{E}_0^{\text{KS}} = \sum_{n \leq F} \varepsilon_n^{\text{KR}} - \frac{1}{2} \int d\vec{r} \int d\vec{r}' f(\vec{r}) f(\vec{r}') \\ \times v(\vec{r}, \vec{r}')$$

$$+ \bar{E}_{xc}[\rho] - \int d\vec{r} V_{xc}[\rho](\vec{r})$$

E_{xc} = the exchange
correlation energy which
replaces $\frac{1}{2} \sum_{ij} \langle i j | v_{ij} | i j \rangle$
(the exchange term in HF
theory)

$$V_{xc}[\rho](\vec{r}) = \frac{\delta \bar{E}_{xc}[\rho]}{\delta \rho}$$

instead of

$$E[\phi_0] = \arg \min_{\phi_0} \langle \phi_0 | H | \phi_0 \rangle$$

(HF)

$$E[\bar{\rho}] = \arg \min_S \langle \bar{\rho}_c | H | \bar{\rho}_c \rangle$$

defined by

in the non-interacting case

$$H = T + V_{ext}$$

$$T = \sum_i -\frac{D_i^2}{2} \quad V_{ext} = \sum_i V_{ext}^{(i)}$$

$$E[\rho] = T[\rho] + \int d\vec{z} V_{ext}(\vec{z}) \rho(\vec{z})$$

variation w.r.t ρ

$$\frac{\delta T}{\delta \rho} + V_{ext}(\vec{z}) = \lambda \rho$$

Lagrange
multiplier

$$\text{no } \psi_i^*(\vec{z}) = \varepsilon_i \psi_i(\vec{z})$$

$$H_0 |\underline{\Phi}_0\rangle = E_0 |\underline{\Phi}_0\rangle$$

$$\langle \underline{\Phi}_0 | T | \underline{\Phi}_0 \rangle \neq \langle \Psi_0 | T | \Psi_0 \rangle$$

$$H |\underline{\Phi}_0\rangle \neq E |\underline{\Phi}_0\rangle$$

$$H |\Psi_0\rangle = E |\Psi_0\rangle$$

$$H = T + V_{ext} + H_I$$

$$E[\bar{S}] = T[\bar{S}] + \int d\vec{z} p(\vec{z}) V_{ext}(\vec{z}) \\ + \frac{1}{2} \int d\vec{z} \int d\vec{z}' p(\vec{z}) p(\vec{z}') \omega(\vec{z}, \vec{z}')$$

$$+ \text{Exc } [\bar{S}]$$

vary with \bar{S}

$$\frac{\delta T}{\delta S} + \frac{\delta \text{Exc}}{\delta S} + \int d\vec{z}' p(\vec{z}') \omega(\vec{z}, \vec{z}')$$

$$+ \frac{\delta \text{Exc}}{\delta S} = \lambda p(\vec{z})$$

$$N_{\text{eff}}(\vec{r}) = N_{\text{ext}}(\vec{r}) + \frac{\delta E_{\text{xc}}}{\delta \delta} = V_{\text{xc}}$$

$$+ \int d\vec{r}' \rho(\vec{r}') v(\vec{r}, \vec{r}')$$

can define an effective SP-equation

$$[\bar{t}(\vec{r}) + N_{\text{eff}}(\vec{r})] \psi_{\text{xc}}^{<5}(\vec{r})$$

$$= \sum_k \psi_k^{<5}$$

$$V_{\text{xc}}[\delta] = \frac{\delta E_{\text{xc}}}{\delta \delta} = V_{\text{xc}}[\delta(\vec{r}), D\delta(\vec{r})]$$

$$\hat{\partial}(\hat{\partial} \rho(\vec{r})) \dots]$$

Local density approx

$$V_{xc}[\bar{s}](\vec{r}) = V_{xc}(s(\vec{r}))$$

$$\bar{E}_{xc} \doteq \int d\vec{r} E_{xc}[\bar{s}(\vec{r})] s(\vec{r})$$

$$E_{xc}[s] = -\text{const } s^{1/3}(\vec{r})$$

(From homogeneous el gas)

$$E_{xc} = -\text{const} \int d\vec{r} s^{4/3}$$

$$V_{xc} = \frac{\partial E_{xc}}{\partial S} = - 4/5 \cdot \text{const} \times S^{1/3}$$

$$= - \text{const} S^{1/3}$$

$$\begin{aligned} & \left[t(\vec{r}) + V_{\text{ext}}(\vec{r}) + \int d\vec{r}' \rho(\vec{r}') \right. \\ & \quad \times \nu(\vec{r}, \vec{r}') \\ & \quad \left. - \text{const} S^{1/3}(\vec{r}) \right] \psi_K(\vec{r}) \\ & = \varepsilon_K^{\text{KS}} \psi_K(\vec{r}) \end{aligned}$$