

# Exercises FYS4480, week 37, September 11-15, 2023

## Exercise 1

We will study a schematic model (the Lipkin model, Nucl. Phys. **62** (1965) 188) for the interaction among 4 fermions that can occupy two different energy levels. Each level has degeneration  $d = 4$ . The two levels have quantum numbers  $\sigma = \pm 1$ , with the upper level having  $\sigma = +1$  and energy  $\varepsilon_1 = \varepsilon/2$ . The lower level has  $\sigma = -1$  and energy  $\varepsilon_2 = -\varepsilon/2$ . In addition, the substates of each level are characterized by the quantum numbers  $p = 1, 2, 3, 4$ .

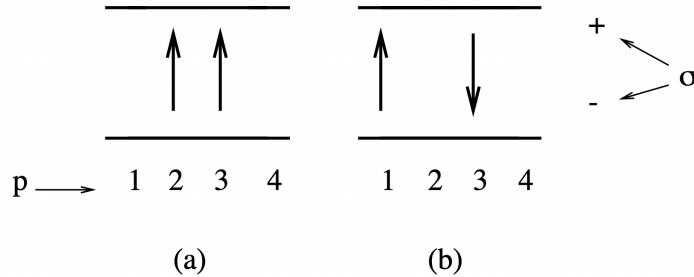
We define the single-particle states

$$|u_{\sigma=-1,p}\rangle = a_{-p}^\dagger |0\rangle \quad |u_{\sigma=1,p}\rangle = a_{+p}^\dagger |0\rangle.$$

The single-particle states span an orthonormal basis. The Hamiltonian of the system is given by

$$\begin{aligned} \hat{H} &= \hat{H}_0 + \hat{H}_1 + \hat{H}_2 \\ \hat{H}_0 &= \frac{1}{2}\varepsilon \sum_{\sigma,p} \sigma a_{\sigma,p}^\dagger a_{\sigma,p} \\ \hat{H}_1 &= \frac{1}{2}V \sum_{\sigma,p,p'} a_{\sigma,p}^\dagger a_{\sigma,p'}^\dagger a_{-\sigma,p'} a_{-\sigma,p} \\ \hat{H}_2 &= \frac{1}{2}W \sum_{\sigma,p,p'} a_{\sigma,p}^\dagger a_{-\sigma,p'}^\dagger a_{\sigma,p'} a_{-\sigma,p} \end{aligned}$$

where  $V$  and  $W$  are constants. The operator  $H_1$  can move pairs of fermions as shown in the left part of the figure (a). while  $H_2$  is a spin-exchange term. As shown in (b),  $H_2$  moves a pair of fermions from a state  $(p\sigma, p' - \sigma)$  to a state  $(p - \sigma, p'\sigma)$ .



We will encounter this model again in our analysis of mean field methods like the Hartree-fock method and full configuration interaction theory. It is a model which has been used widely in many-body physics and recently also in quantum computing, see for example <https://journals.aps.org/prc/abstract/10.1103/PhysRevC.104.024305>.

a. *Quasispin operators* Introduce the quasispin operators

$$\begin{aligned} \hat{J}_+ &= \sum_p a_{p+}^\dagger a_{p-} \\ \hat{J}_- &= \sum_p a_{p-}^\dagger a_{p+} \\ \hat{J}_z &= \frac{1}{2} \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} \\ \hat{J}^2 &= J_+ J_- + J_z^2 - J_z \end{aligned}$$

Show that these operators obey the commutation relations for angular momentum.

*b. Number operator* Express  $\hat{H}$  in terms of the above quasispin operators and the number operator

$$\hat{N} = \sum_{p\sigma} a_{p\sigma}^\dagger a_{p\sigma}.$$

*c. Commutation relations* Show that  $\hat{H}$  commutes with  $J^2$ , viz.,  $J$  is a good quantum number. Does it commute with  $J_z$ ?

*d. Wick's theorem* Consider thereafter a state with all four fermions in the lowest level (see the above figure). We can write this state as

$$|\Phi_0\rangle = |\Phi_{J_z=-2}\rangle = a_{1-}^\dagger a_{2-}^\dagger a_{3-}^\dagger a_{4-}^\dagger |0\rangle.$$

This state has  $J_z = -2$  (convince yourself about this) and belongs to the set of possible projections of  $J = 2$ . We introduce the shorthand notation  $|J, J_z\rangle$  for states with different values of spin  $J$  and its projection  $J_z$ . We can think of this as our computational basis for  $J = 2$  and all five projections  $J_z$ . We will also assume that the state  $\Phi_0$  can be considered as an ansatz for the ground state of the system.

Use Wick's theorem to calculate the expectation values of

$$\langle \Phi_0 | \hat{N} | \Phi_0 \rangle,$$

and

$$\langle \Phi_0 | \hat{H} | \Phi_0 \rangle.$$

Comment your results.

*e. Using quasispin operators* Show that you can obtain the same result for

$$\langle \Phi_0 | \hat{H} | \Phi_0 \rangle.$$

using the quasispin representation of the Hamiltonian (plus the number operator). Comment your results.