

**Lecture August
25, 2023,
FYS4480**

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Computational basis-

- Example : two-level system

$$|\psi_0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \quad |\psi_1\rangle = \begin{bmatrix} 0 \\ -1 \end{bmatrix}$$

ONS : $\langle \psi_1 | \psi_0 \rangle = 0$

$$\hat{P} = |\psi_0\rangle \langle \psi_0| = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \end{bmatrix}$$

$$\hat{P} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \quad \hat{Q} = |\psi_1\rangle \langle \psi_1|$$
$$= \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\hat{P} \cdot \hat{Q} = 0$$

$$\hat{P} + \hat{Q} = \hat{I} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

in general

$$\hat{P} = \sum_{i=0}^{d-1} \underbrace{|i\rangle\langle i|}_{\text{ONB}} \quad \hat{Q} = \sum_{i=d}^{\infty} |i\rangle\langle i|$$

$$\hat{P}^2 = \hat{P} \quad | \quad \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$

$$\hat{Q}^2 = \hat{Q} \quad | \quad \begin{bmatrix} \hat{P} & \hat{Q} \end{bmatrix}$$

Computational SP-basis (single-particle)

Example : H₀ in 2dime

$$n_x, n_y \quad -x^2 d$$

$$|n_x\rangle = H_{n_x}(x) e^{-x^2 d}$$

$$|n_y\rangle = H_{n_y}(y) e^{-y^2 d}$$

$$n_x, n_y = 0, 1, 2, \dots$$

$$|n_x\rangle \otimes |n_y\rangle = |n_x n_y\rangle$$

adding spin $s=1/2 \quad m_s=\pm 1/2$

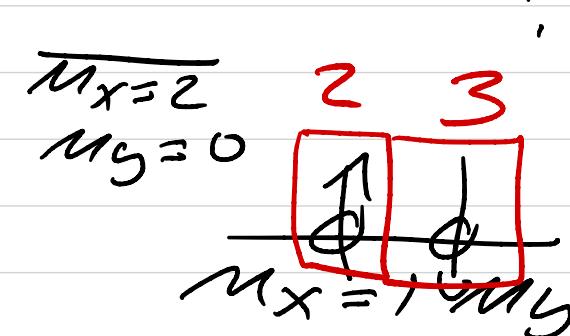
$$|n_x\rangle \otimes |n_y\rangle \otimes |s m_s\rangle$$

$$= |m_x m_y s m_s\rangle = |\alpha\rangle$$

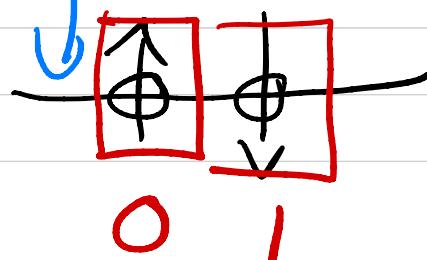
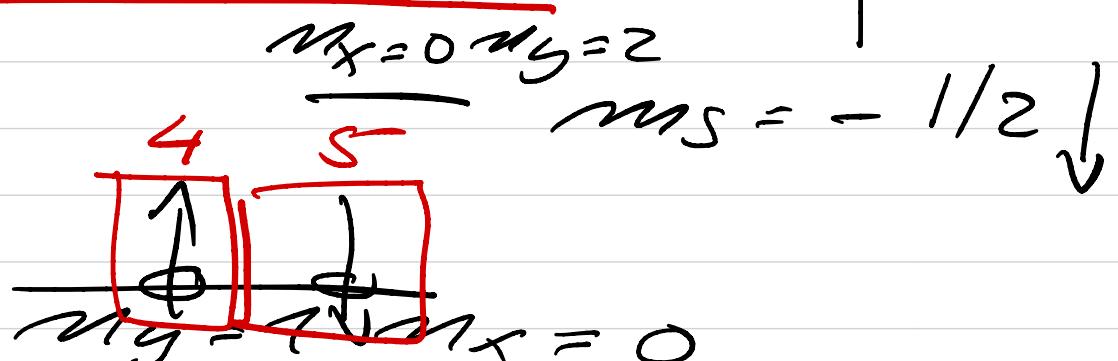
$$\alpha = \{ m_x m_y s m_s \}$$

$$\alpha = 0, 1, 2, -d-1 \quad \infty$$

$$E_{m_x m_y} = \hbar w (m_x + m_y + 1) \quad m_s = +1/2$$



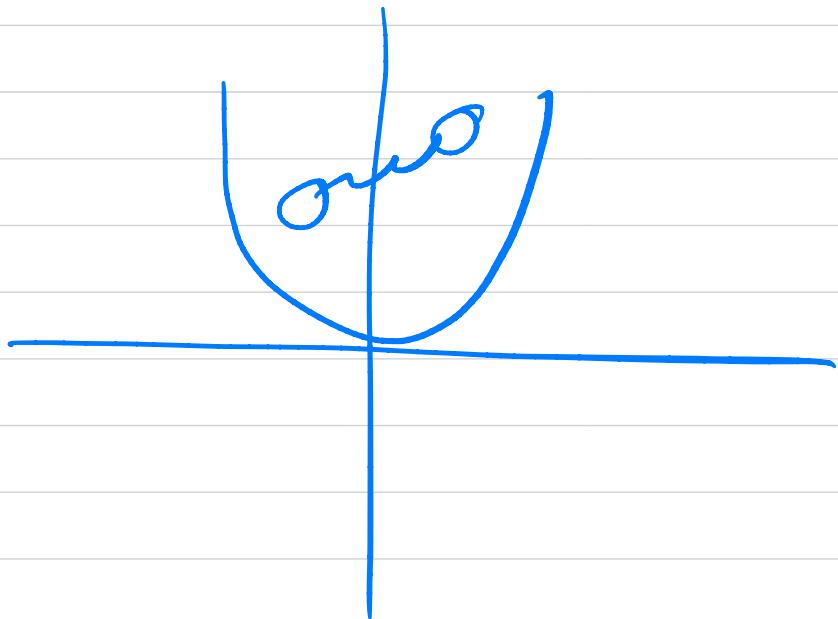
$$m_x = m_y = 1$$



can expand a new state

$$|N\rangle = \sum_{i=0}^{\infty} c_i |i\rangle$$

2-fermions in a HO potential



Hamiltonian: $\hat{H} = \hat{H}_0 + \hat{H}_I$ interaction

$$\hat{H}_0 = \sum_{i=1}^N \hat{h}_0(\vec{x}_i) \quad N = \text{Number of particles}$$

$$\hat{h}_0(\vec{x}_i) = -\frac{\hbar^2}{2m} D_i^2 + U_{ext}(\vec{x}_i)$$

$$\hat{h}_0 |i\rangle = \epsilon_i |i\rangle \quad \langle i|j\rangle = \delta_{ij}$$

assumption made.

$$\hat{H}_I = \sum_{i < j} v(\vec{x}_i, \vec{x}_j)$$

$$v(\vec{x}_i, \vec{x}_j) = v(|\vec{x}_i - \vec{x}_j|)$$

$$|\vec{x}_i - \vec{x}_j| = \sqrt{(x_i - x_j)^2 + (y_i - y_j)^2}$$

ansatz for 2-body state function

$$|ab\rangle = |a\rangle \otimes |b\rangle$$

$$= \frac{1}{\sqrt{2}} \begin{vmatrix} \varphi_a(\vec{x}_1) & \varphi_a(\vec{x}_2) \\ \varphi_b(\vec{x}_1) & \varphi_b(\vec{x}_2) \end{vmatrix}$$

$$\varphi_a(\vec{x}_i) = \langle \vec{x}_i | a \rangle$$

$$\hat{h}_0 \varphi_a(\vec{x}_i) = \epsilon_a \varphi_a(\vec{x}_i)$$

Example $|a\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$ $|b\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$

$$|ab\rangle = |a\rangle \otimes |b\rangle = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix}$$

$$\langle ab | \hat{H} | ab \rangle$$

$$|ab\rangle = \frac{1}{\sqrt{2}} (\varphi_a(x_1)\varphi_b(x_2) - \varphi_a(x_2)\varphi_b(x_1))$$

1) $\langle ab | \hat{H}_0 | ab \rangle =$

$$\begin{aligned} \frac{1}{2} \iint d\vec{x}_1 d\vec{x}_2 & [\varphi_a^*(x_1) \varphi_b^*(x_2) - \varphi_a^*(x_2) \varphi_b^*(x_1)] \\ & \times [\hat{h}_0(\vec{x}_1) + \hat{h}_0(\vec{x}_2)] [\varphi_a(x_1) \varphi_b(x_2) \\ & - \varphi_a(x_2) \varphi_b(x_1)] \end{aligned}$$

$$\langle H \rangle = E[H] = \frac{\int d\vec{r} |4^*(\vec{r})|^2 \psi(\vec{r})}{\int d\vec{r} |\psi(\vec{r})|^2}$$

$$d\vec{r} = dx, dy, dx_2 dy_2$$

First term:

$$\begin{aligned} & \frac{1}{2} \int d\vec{x}_1 \int d\vec{x}_2 \psi_a^*(x_1) \psi_b^*(x_2) [h_o(x_1) + h_o(x_2)] \\ & \quad \times [\psi_a(x_1) \psi_b(x_2)] \\ & = \frac{1}{2} (\varepsilon_a + \varepsilon_b) \end{aligned}$$

Second term

$$\frac{1}{2} \int \int \vec{dx_1} \vec{dx_2} \varphi_a^*(x_2) \varphi_b^*(x_1) \left[\hat{h}_0(x_1) + \hat{h}_0(x_2) \right]$$

$$+ \varphi_a(x_1) \varphi_b(x_2) = 0$$

same with the third
last

$$\frac{1}{2} \int \int \vec{dx_1} \vec{dx_2} \varphi_a^*(x_2) \varphi_b^*(x_1) \left(\hat{h}_0(x_1) + \hat{h}_0(x_2) \right)$$

$$+ \varphi_a(x_2) \varphi_b(x_1)$$

$$= \frac{1}{2} (\varepsilon_a + \varepsilon_b) \Rightarrow$$

$$\Rightarrow \langle ab | \hat{H}_0 | ab \rangle = \epsilon_q + \epsilon_e$$

$$\hat{H}_0 |ab\rangle = (\epsilon_q + \epsilon_e) |ab\rangle$$

$|ab\rangle$ is an eigenvector of
 \hat{H}_0

$$\hat{A} |ab\rangle = (\hat{H}_0 + \hat{H}_1) |ab\rangle$$

$$\neq E |ab\rangle$$

$|\Phi_{ab}\rangle = |ab\rangle$ $ab \rightarrow i$
configuration

configuration : way to distribute $1 \text{ or more particles}$
on all possible computational basis states (n)

Fermions with n states
and N particles \Rightarrow

$$\# \text{ possible config} = \binom{n}{N}$$

$$= \frac{n!}{(n-N)! N!} \quad n \geq N$$

$$|\psi_{ab}\rangle = \sum_{i=0}^{d-1} c_i |\underbrace{i}_{ab}\rangle$$

Example $n = 100$

$$N = 2$$

$$\# \text{ configs} = \frac{100!}{98! 2!} = 49 \cdot 25$$

$$\langle ab | H_1 | ab \rangle =$$

$$= \frac{1}{2} \iint d\vec{x}_1 d\vec{x}_2 [\varphi_a^*(x_1) \varphi_b^*(x_2) - \varphi_a^*(x_2) \varphi_b^*(x_1)] v(\vec{x}_1, \vec{x}_2) [v(x_1, x_2) = v(x_2, x_1) - \varphi_a(x_2) \varphi_b(x_1)]$$

$$1) \frac{1}{2} \iint d\vec{x}_1 d\vec{x}_2 \varphi_a^*(x_1) \varphi_b^*(x_2) v(x_1, x_2)$$

$$\varphi_a(x_1) \varphi_b(x_2)$$

$$= \frac{1}{2} \langle ab | v | ab \rangle$$

4) Last term

$$\frac{1}{2} \langle \bar{c}a | v | c\bar{a} \rangle$$

$$= \frac{1}{2} \langle ab | \sigma | ab \rangle$$

$$2+3) - \frac{1}{2} \langle ab/r/a \rangle$$

- $\sum_{n=1}^{\infty} \langle ab/r/ba \rangle \Rightarrow$

$$\langle \tilde{a}^\dagger | H_I | \tilde{a} \rangle =$$

$\langle ab|v|ab \rangle - \langle ab|v|ba \rangle$