

FYS4480/9480, lecture
November 14, 2025

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FCI

$$|4_0\rangle = (1 + \hat{C}) |\Phi_0\rangle$$

$$\hat{C} = \sum_{p_H > 0} c_H^p |\Phi_H^p\rangle$$

$$= \sum_{p_H > 0} c_H^p \underbrace{A_H^p} |\Phi_0\rangle$$

example

$$\sum_{a_i} c_i^a a_a^\dagger a_i |\Phi_0\rangle$$

$$|\psi_0\rangle_{\text{(MBPT(2))}} = \sum_{k=0}^{\infty} \left\{ \left[\hat{R} \hat{H}_I \right]^k |\Phi_0\rangle \right\}_L$$

$$\hat{R} = \sum_{M>0} \frac{|\Phi_M\rangle \langle \Phi_M|}{\epsilon_0 - \epsilon_M}$$

CC-Theory

$$\hat{T} = \underbrace{\hat{T}_1}_{1p1h} + \underbrace{\hat{T}_2}_{2p2h} + \dots + \hat{T}_{NpNh}$$

$$|\psi_0\rangle = \exp\{\hat{T}\} |\Phi_0\rangle$$

$$\begin{aligned}
 T = T_1 &= \sum_{a_i} c_i^a a_a^\dagger a_i \\
 &= \sum_{a_i} t_i^a a_a^\dagger a_{i'}
 \end{aligned}$$

$$|4_0\rangle = \exp\{T_1\} |\Phi_0\rangle$$

$$= \exp\left\{\sum_{a_i} t_i^a a_a^\dagger a_{i'}\right\} |\Phi_0\rangle$$

$$\begin{aligned}
 &= \prod_{i \in F} \left\{ 1 + \sum_{a \in F} t_i^a a_a^\dagger a_{i'} + \right. \\
 &\quad \left. \frac{1}{2!} \left(\sum_a t_i^a a_a^\dagger a_{i'} \right)^2 + \dots \right\} |\Phi_0\rangle \\
 &\quad \dots a_i a_{i'} \dots |\Phi_0\rangle
 \end{aligned}$$

$$|\psi_0\rangle = \prod_{i \in F} \left(1 + \sum_a t_i^a a_i^\dagger a_{i'} \right) |\Phi_0\rangle$$

$$T_2 : 2p = 4$$

$$a_a^\dagger a_b^\dagger a_j a_{i'} |\Phi_0\rangle = |\Phi_{ij}^{ab}\rangle$$

$$T_2 = \frac{1}{4} \sum_{\substack{ab \\ i'j'}} t_{ij}^{ab} \underbrace{a_a^\dagger a_b^\dagger a_j a_{i'}}_{A_{ij}^{ab} \rightarrow A_H^p}$$

General expansion to 9C

$$|\psi_0\rangle_{cc} = \prod_H \left(1 + \sum_p t_H^p A_H^p \right) |\Phi_0\rangle$$

$$(A_H^p)^2 |\Phi_0\rangle = 0$$

$$(A_H^p)^n |\Phi_0\rangle = 0 \quad n > 1$$

$$|4_0\rangle_{cc} = \exp(T) |\Phi_0\rangle$$

$$T_2 = \frac{1}{4} \sum_{\substack{ab \\ ij}} t_{ij}^{ab} a_a^\dagger a_b^\dagger a_j a_i$$

$$\frac{1}{(2!)^2}$$

$$t_{ij}^{ab} = -t_{ji}^{ab} = -t_{ij}^{ba}$$

$$= t_{ji}^{ba}$$

$$\langle ab|v|ij\rangle_{AS} = -\langle ab|v|ji\rangle_{AS} \dots$$

$$\text{MBPT}(2) \\ \text{2p2h}$$

$$t_{ij}^{ab} = \frac{\langle ab|v|ij\rangle_{AS}}{\epsilon_i + \epsilon_j - \epsilon_a - \epsilon_b}$$

$$\overline{T}_3 = \frac{1}{(3!)^2} \sum_{abc, ijk} t_{ijk}^{abc} (a_a^\dagger a_k^\dagger a_c^\dagger a_k a_j a_i)$$

CC

FCI

$$e^T |\Phi_0\rangle$$

$$\sum_{PH} C_H^P |\Phi_0\rangle$$

$$= \sum_{\lambda=0}^{\infty} C_{\lambda} |\Phi_{\lambda}\rangle$$

$$T = T_1 + T_2 + \dots + T_{N_P N_H}$$

$$\underline{C_0} = \underline{1} = C_{0H}^{0P}$$

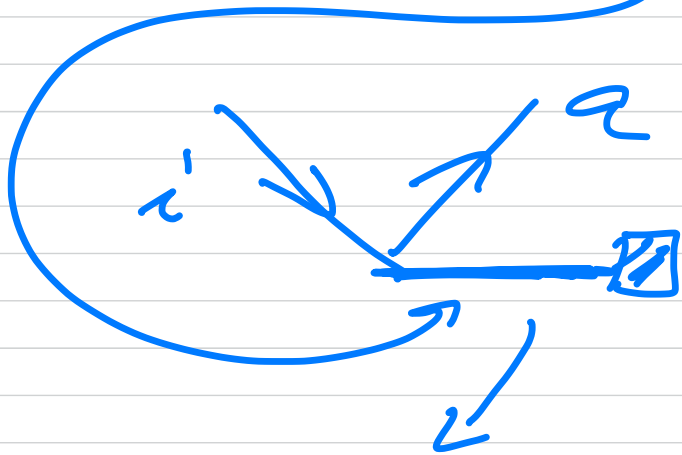
$$\underline{T_1} = C_{1H}^{1P} = C_{\underline{1}}$$

$$\underline{T_2} + \frac{1}{2} \underline{T_1}^2 = C_2 (= C_{2H}^{2P})$$

$$C_3 = T_3 + T_1 T_2 + \frac{1}{6} T_1^3 \text{ (3p3h)}$$

$$(\exp(T))$$

$$T_1 = \sum_{a, a'} t_a^a q_a^\dagger q_{a'}$$

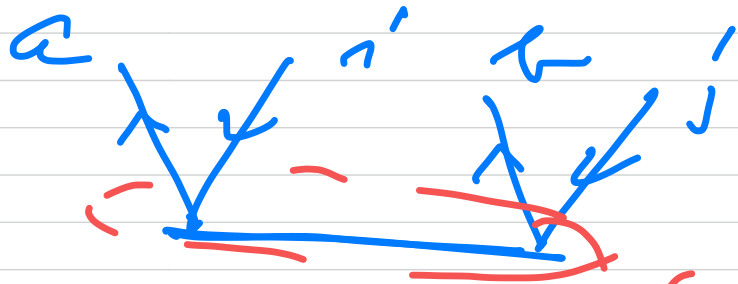


$$\left(\begin{array}{c} \text{diagram} + \\ \text{diagram} + \dots \end{array} \right)$$

$$\text{MBPT}(2) : \frac{\langle a | f | i' \rangle}{\epsilon_{i'} - \epsilon_a}$$

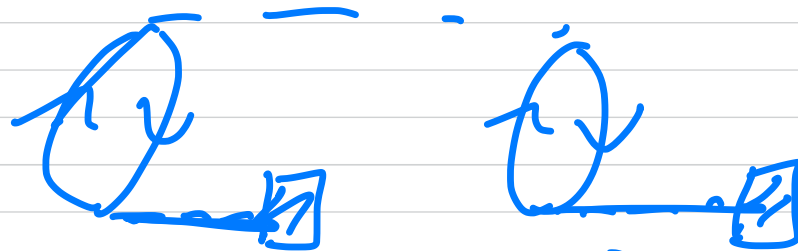
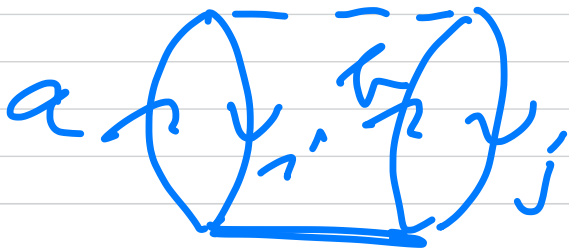
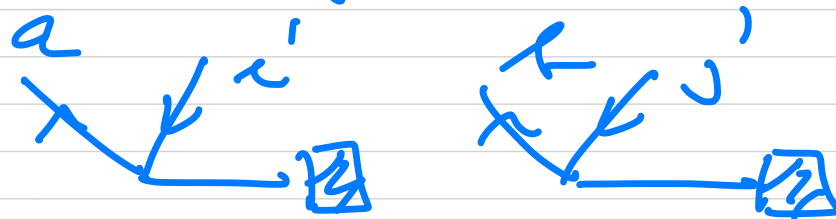
$$C_2 = \frac{1}{2} + \frac{1}{2} \frac{1}{1}^2$$

$$\frac{1}{2} =$$



$$= \frac{1}{4} \sum_{\substack{a,b \\ r,s}} (t_{ij})^{ab} q_a^+ q_b^+ q_s q_{s'}$$

$$\frac{1}{2} \rightarrow 2$$



$\chi_{\chi} - G$

$$|\psi_0\rangle_{cc} = \exp\{T\} |\Phi_0\rangle = e^T |\Phi_0\rangle$$

$$\hat{H} |\psi_0\rangle_{cc} = E_0 |\psi_0\rangle_{cc}$$

$$H e^T |\Phi_0\rangle = E_0 e^T |\Phi_0\rangle$$

$$E_{min} = \arg \min_{t_H} \frac{\langle \Phi_0 | e^{-T} H e^T | \Phi_0 \rangle}{\underbrace{\langle \psi_0 | \psi_0 \rangle}_{\langle \Phi_0 | \Phi_0 \rangle}}$$

non-linear dependence on t_H ; in principle this leads to an intractable set of

non-linear equations,

Equations (projection approach)

$$\hat{H} e^T |\Phi_0\rangle = E_0 e^T |\Phi_0\rangle$$

$$\langle \Phi_0 | \hat{H} e^T |\Phi_0\rangle = E_0 \underbrace{e^T |\Phi_0\rangle}_{1}$$

$$\langle \Phi_0 | \psi_0 \rangle = 1$$

Think to FCI: First row

$$\langle \Phi_0 | \hat{H} | \sum_{PH} C_H^P |\Phi_H^P\rangle = E_0$$

$$\Delta E_0 = E_0 - E_0^{REF} = \sum_{PH>0} \langle \Phi_0 | \hat{H} | \Phi_H^P \rangle \times C_H^P$$

$$\Delta E_0 = \sum_{a i'} c_{i'}^a \langle a | f | i' \rangle + \sum_{\substack{ab \\ rj}} c_{ij}^{ab} \langle ab | r | i' j \rangle$$

$$\langle \Phi_n^a | \hat{H} e^T | \Phi_0 \rangle = E_0 \langle \Phi_n^a | e^T | \Phi_0 \rangle$$

For a truncated $T \approx \underbrace{T_1 + T_2}_{\text{CCSD}}$

$$(i) \quad \langle \Phi_0 | H e^T | \Phi_0 \rangle = E_0$$

$$(ii) \quad \langle \Phi_n^a | H e^T | \Phi_0 \rangle =$$

$$E_0 \langle \Phi_n^a | e^T | \Phi_0 \rangle$$

$$(iii) \quad \langle \Phi_{ij}^{\alpha\beta} | H e^T | \Phi_0 \rangle$$

$$= E_0 \langle \Phi_{ij}^{\alpha\beta} | e^T | \Phi_0 \rangle$$

$$\hat{H} = E_0^{\text{Ref}} + \underbrace{\hat{F}_N + \hat{V}_N}_{\hat{H}_N}$$

$$e^{-T} H e^T |\Phi_0\rangle = E_0 |\Phi_0\rangle$$

$\underbrace{(e^T |\Phi_0\rangle)}_{|\Psi_0\rangle}$

$$\Delta E_0 = E_0 - \underbrace{E_0^{\text{ref}}}$$

(i)

$$= \langle \Phi_0 | e^{-T} H_N e^T | \Phi_0 \rangle$$

(i')

$$\langle \Phi_n^a | e^{-T} H_N e^T | \Phi_0 \rangle = 0$$

(i'')

$$\langle \Phi_{ij}^{ab} | e^{-T} H_N e^T | \Phi_0 \rangle = 0$$

FCi'

<opch | fl | ipm>

opoh

ipih

zpz4

- - -

npnh

opch

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0

0

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qpi4

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zpz4

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1

0

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1

1

1

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1

1

1

1

1

1

npnh

$D=4$ 0p0L
 $D=3$

2p2L
~~00~~

2p2L
~~00~~

$D=2$ ~~00~~
 $D=1$ ~~00~~

$|\Phi_0\rangle$

~~00~~

$|\Phi_1\rangle$

~~00~~

$|\Phi_2\rangle$

2p2L

2p2L
~~00~~

4p4L

~~00~~
~~00~~

$|\Phi_3\rangle$

$|\Phi_4\rangle$

$|\Phi_5\rangle$

\Rightarrow Hamiltonian matrix
 6×6

$$|4_0\rangle = c_0 |\Phi_0\rangle + c_1 |\Phi_1\rangle + \dots + c_r |\Phi_r\rangle$$

$$|4_0^{(1)}\rangle = \sum_{a_i} \cancel{\frac{|\Phi_i^a\rangle \langle a|g|i\rangle}{\epsilon_i - \epsilon_a}} \quad \text{1p1h}$$

$$+ \frac{1}{4} \sum_{a_i, j} |\Phi_{ij}^{a_i}\rangle \frac{\langle a_i | h | i, j \rangle}{\epsilon_i + \epsilon_j - \epsilon_a - \epsilon_b} \quad \text{2p2h}$$