

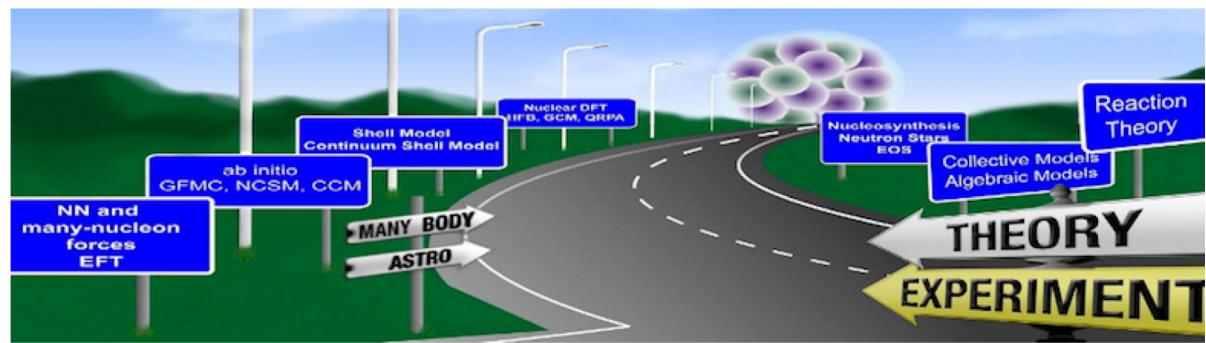
Coupled-Cluster for open-shell nuclei

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Overview of close-shell coupled-cluster(by T.P)

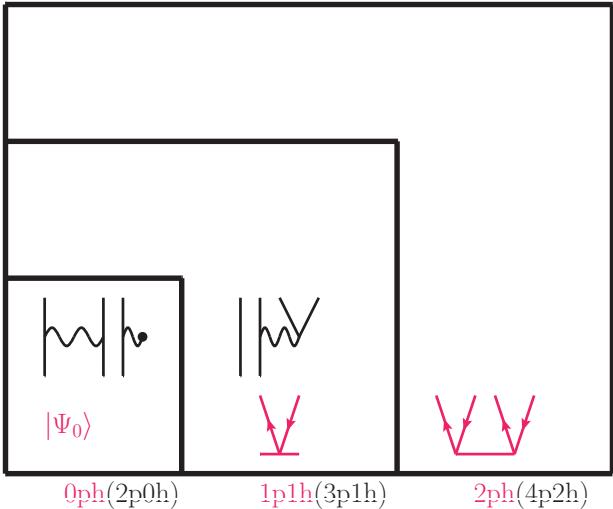
- Starting from a Hamiltonian normal ordered to a reference state(new definition of vacuum)

$$\mathcal{H} = H - E_{\text{ref}} = \sum_{pq} f_{pq} \{p^\dagger q\} + \frac{1}{4} \sum_{pqrs} V_{pqrs} \{p^\dagger q^\dagger s r\}.$$

- Coupled-cluster(CC) is similarity transformation(result in a non-Hermitian Hamiltonian)

$$\begin{aligned}\overline{\mathcal{H}} &\equiv \exp(-T)\mathcal{H}\exp(T) \\ &= (\mathcal{H}\exp(T))_C \\ &= \Delta E + \sum_{pq} \overline{\mathcal{H}}_{pq} \{p^\dagger q\} + \sum_{pqrs} \overline{\mathcal{H}}_{pqrs} \{p^\dagger q^\dagger s r\} \quad \langle \Psi_i | \mathcal{H} | \Psi_j \rangle \neq \langle \Psi_j | \mathcal{H} | \Psi_i \rangle \\ &\quad + \sum_{pqrsut} \overline{\mathcal{H}}_{pqrsut} \{p^\dagger q^\dagger r^\dagger u t s\} + \dots\end{aligned}$$

- One-body , two-body , etc. excitations decoupled from the system. Ground state and any excited states are **not linked** after the transformation by definition of CC.



$$\Psi = \Phi_{\text{SCF}} + \sum_{i,a} C_i^a \Phi_i^a + \sum_{i < j, a < b} C_{ij}^{ab} \Phi_{ij}^{ab} + \dots \quad (\text{up to } N \text{ excitations}),$$

$$\hat{C}_4 = \hat{T}_4 + \frac{1}{2}\hat{T}_2^2 + \frac{1}{4!}\hat{T}_1^4 + \hat{T}_1\hat{T}_3 + \frac{1}{2}\hat{T}_1^2\hat{T}_2.$$

- CC is extensive(FCI, CC, MBPT, etc.)
computational less expensive compared with CI

Approximation of CC

Truncation of high-rank ops.

$$\hat{T} = \sum_{m=1}^N \hat{T}_m .$$

$$\begin{aligned}\overline{\mathcal{H}} &\equiv \exp(-T)\mathcal{H}\exp(T) \\ &= (\mathcal{H}\exp(T))_C \\ &= \Delta E + \sum_{pq} \overline{\mathcal{H}}_{pq}\{p^\dagger q\} + \sum_{pqrs} \overline{\mathcal{H}}_{pqrs}\{p^\dagger q^\dagger sr\} \\ &\quad + \sum_{pqrstu} \overline{\mathcal{H}}_{pqrstu}\{p^\dagger q^\dagger r^\dagger uts\} + \dots.\end{aligned}$$

$$\begin{aligned}\hat{T}_m &= \frac{1}{(m!)^2} \sum_{\substack{i_1 i_2 \dots i_m \\ a_1 a_2 \dots a_m}} \langle a_1 a_2 \dots a_m | \hat{t}_m | i_1 i_2 \dots i_m \rangle_A \{ \hat{a}_1^\dagger \hat{a}_2^\dagger \dots \hat{a}_m^\dagger \hat{i}_m \dots \hat{i}_2 \hat{i}_1 \} \\ &= \frac{1}{(m!)^2} \sum_{\substack{i_1 i_2 \dots i_m \\ a_1 a_2 \dots a_m}} t_{i_1 i_2 \dots i_m}^{a_1 a_2 \dots a_m} \{ \hat{a}_1^\dagger \hat{i}_1 \hat{a}_2^\dagger \hat{i}_2 \dots \hat{a}_m^\dagger \hat{i}_m \dots \},\end{aligned}\tag{9.82}$$

$$T = T_1 \qquad \qquad \qquad \text{HF}$$

$$T = T_1 + T_2 \qquad \qquad \qquad \text{CCSD}$$

$$T = T_1 + T_2 + T_3 \qquad \qquad \qquad \text{CCSDT}$$

Ground state decoupling, Coupled-cluster effective Hamiltonian

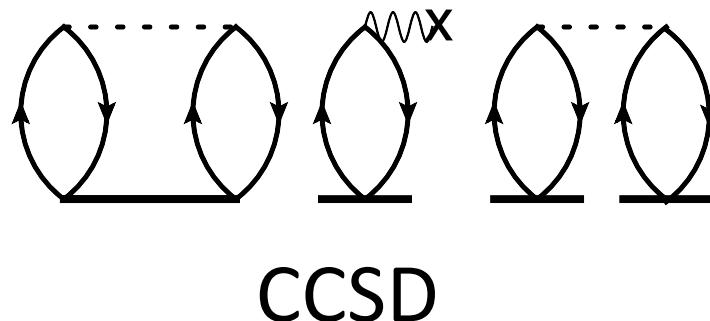
$$|\Psi_0\rangle = e^{\hat{T}}|\Phi_0\rangle \quad \quad \hat{T} = \sum_i T_i$$

$$\hat{T}_1 = \sum_{ia} t_i \alpha_a^\dagger \alpha_i, \quad \hat{T}_2 = \frac{1}{2} \sum_{i < j, a < b} t_{ij}^{ab} \alpha_a^\dagger \alpha_b^\dagger \alpha_i \alpha_j$$

$$\bar{H} = e^{-T} H e^T = H + [H, T] + \frac{1}{2!} [[H, T], T] + \dots = (H e^T)_C$$

$$\langle \Psi_0 | e^{-T} H e^T | \Psi_0 \rangle = E$$

$$\langle \Psi_{ij\dots}^{ab\dots} | e^{-T} H e^T | \Psi_0 \rangle = 0$$



Expectation of observables

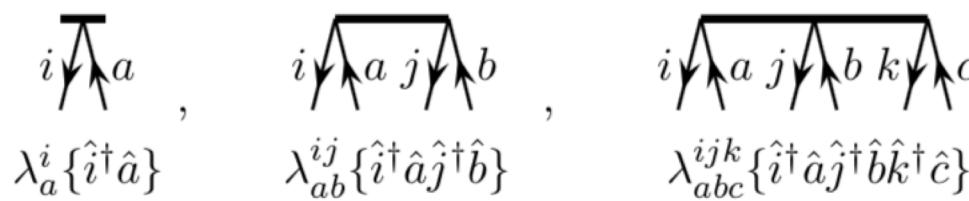
The expectation value in quantum mechanics

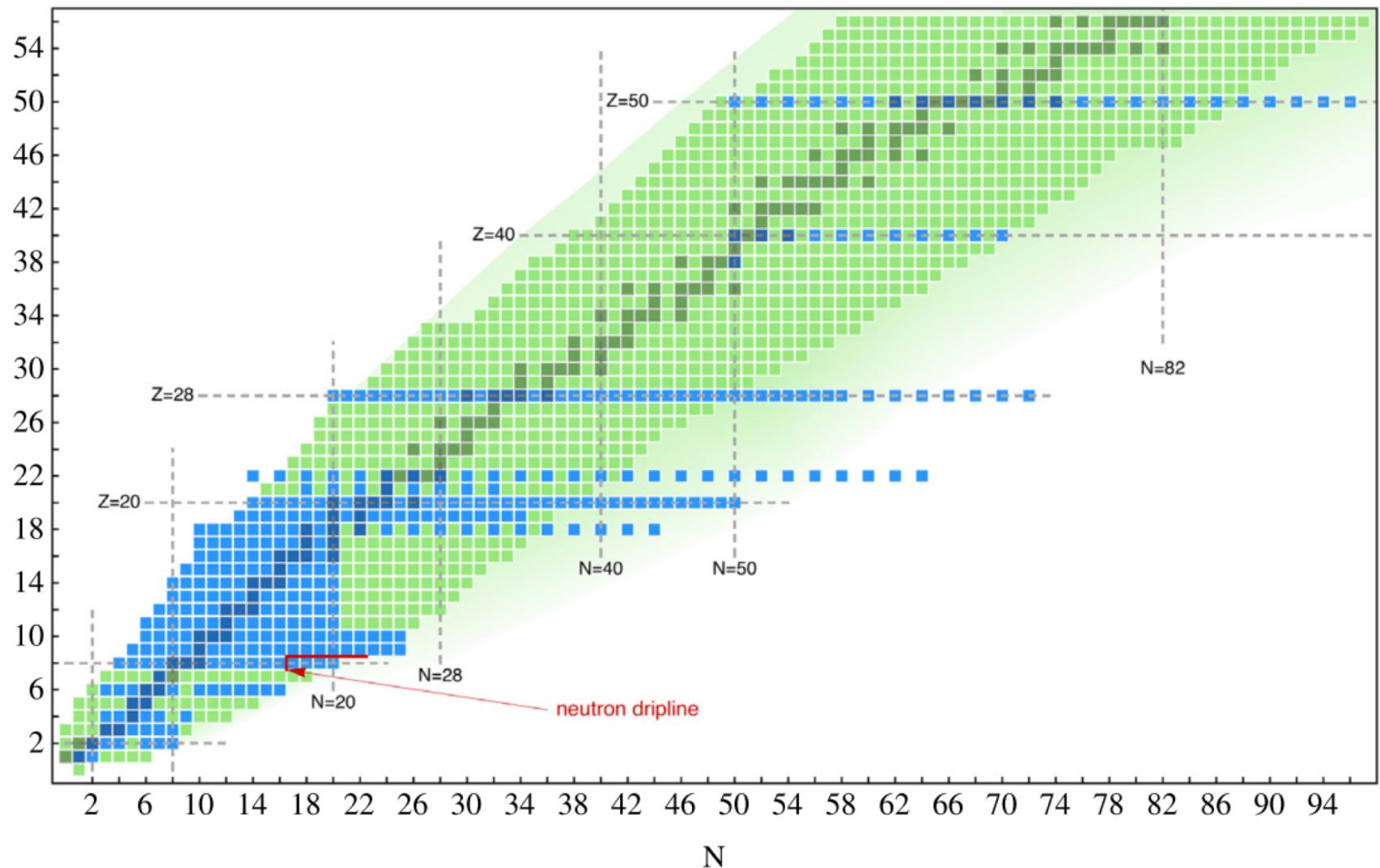
$$\bar{O} = \frac{\langle \Psi | O | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

$$\bar{O} = \frac{\langle 0 | e^{T^\dagger} O e^T | 0 \rangle}{\langle 0 | e^{T^\dagger} e^T | 0 \rangle} = \frac{\langle 0 | e^{T^\dagger} e^T e^{-T} O e^T | 0 \rangle}{\langle 0 | e^{T^\dagger} e^T | 0 \rangle} = \frac{\langle 0 | e^{T^\dagger} e^T}{\langle 0 | e^{T^\dagger} e^T | 0 \rangle} \bar{O} | 0 \rangle$$

For CC:

$$\langle O \rangle = P(1 + \Lambda)\bar{O}P \quad P(1 + \Lambda Q)(\bar{H} - \delta E)Q = 0$$





Equation of motion CC

- Coupled-cluster effective Hamiltonian

$$\chi_{ia} \{\hat{a}^\dagger \hat{a}\} \equiv i \begin{array}{c} \nearrow \\ \nwarrow \\ \text{---} \end{array} a \sim = i \begin{array}{c} \nearrow \\ \nwarrow \\ \text{---} \end{array} a^{-\times} + i \begin{array}{c} \nearrow \\ \nwarrow \\ \text{---} \end{array} a \underline{\text{---}} ,$$

$$\begin{aligned} \chi_{ab} \{\hat{a}^\dagger \hat{b}\} &\equiv \begin{array}{c} \nearrow \\ \uparrow \\ b \\ \downarrow \\ a \end{array} \sim = \begin{array}{c} \nearrow \\ \uparrow \\ b \\ \downarrow \\ a \end{array}^{-\times} + \begin{array}{c} \nearrow \\ \uparrow \\ b \\ \downarrow \\ a \end{array} \underline{\text{---}} + a \begin{array}{c} \nearrow \\ \downarrow \\ \text{---} \\ \nearrow \\ b \end{array}^{-\times} + a \begin{array}{c} \nearrow \\ \downarrow \\ \text{---} \\ \nearrow \\ b \end{array} \underline{\text{---}} + a \begin{array}{c} \nearrow \\ \downarrow \\ \text{---} \\ \nearrow \\ b \end{array} \underline{\text{---}} \\ &= \begin{array}{c} \nearrow \\ \uparrow \\ b \\ \downarrow \\ a \end{array}^{-\times} + \begin{array}{c} \nearrow \\ \uparrow \\ b \\ \downarrow \\ a \end{array} \underline{\text{---}} + a \begin{array}{c} \nearrow \\ \downarrow \\ \text{---} \\ \nearrow \\ b \end{array} \sim + a \begin{array}{c} \nearrow \\ \downarrow \\ \text{---} \\ \nearrow \\ b \end{array} \underline{\text{---}} , \\ \chi_{ij} \{\hat{i}^\dagger \hat{j}\} &\equiv \begin{array}{c} \downarrow \\ \uparrow \\ j \\ \downarrow \\ i \end{array} \sim = \begin{array}{c} \downarrow \\ \uparrow \\ j \\ \downarrow \\ i \end{array}^{-\times} + \begin{array}{c} \downarrow \\ \uparrow \\ j \\ \downarrow \\ i \end{array} \underline{\text{---}} + j \begin{array}{c} \downarrow \\ \nearrow \\ \text{---} \\ \downarrow \\ i \end{array}^{-\times} + j \begin{array}{c} \downarrow \\ \nearrow \\ \text{---} \\ \downarrow \\ i \end{array} \underline{\text{---}} + j \begin{array}{c} \downarrow \\ \nearrow \\ \text{---} \\ \downarrow \\ i \end{array} \underline{\text{---}} \\ &= \begin{array}{c} \downarrow \\ \uparrow \\ j \\ \downarrow \\ i \end{array}^{-\times} + \begin{array}{c} \downarrow \\ \uparrow \\ j \\ \downarrow \\ i \end{array} \underline{\text{---}} + j \begin{array}{c} \downarrow \\ \nearrow \\ \text{---} \\ \downarrow \\ i \end{array} \sim + j \begin{array}{c} \downarrow \\ \nearrow \\ \text{---} \\ \downarrow \\ i \end{array} \underline{\text{---}} , \end{aligned}$$

$$\begin{aligned} \overline{\mathcal{H}} &\equiv \exp(-T) \mathcal{H} \exp(T) \\ &= (\mathcal{H} \exp(T))_C \\ &= \Delta E + \sum_{pq} \overline{\mathcal{H}}_{pq} \{p^\dagger q\} + \sum_{pqrs} \overline{\mathcal{H}}_{pqrs} \{p^\dagger q^\dagger s r\} \\ &\quad + \sum_{pqrstu} \overline{\mathcal{H}}_{pqrstu} \{p^\dagger q^\dagger r^\dagger u t s\} + \dots . \end{aligned}$$

Configuration interaction with CC Hamiltonian

- The coupled cluster transformation is a similarity transformation even if the T operator is truncated.
- The CC Hamiltonian is fully renormalized. Truncations in CI is a good approximation

reference state: $\hat{H}\Psi_0 = E_0\Psi_0$ $|\Psi_k\rangle = Q_k|\Psi_0\rangle$

excited state: $\hat{H}\Psi_k = E_k\Psi_k$

EOM

$$\begin{aligned} H|\Psi_k\rangle &= HQ_k|\Psi_0\rangle = E_kQ_k|\Psi_0\rangle \\ &= [H, Q_k]|\Psi_0\rangle + Q_kH|\Psi_0\rangle \\ &= [H, Q_k]|\Psi_0\rangle + E_0Q_k|\Psi_0\rangle \end{aligned}$$

$$\begin{aligned} [H, Q_k]|\Psi_0\rangle &= (E_k - E_0)Q_k|\Psi_0\rangle \\ &= \omega_k Q_k|\Psi_0\rangle \end{aligned}$$

EOM-CC

Ground state:

$$|\Psi_0\rangle = e^T |HF\rangle$$

k-th excited state:

$$|\Psi_k\rangle = \hat{R}_k |\Psi_0\rangle$$

Close shell nuclei

$$\hat{R}_k = r_0 + \sum_{i,a} r_i^a \{\hat{a}^\dagger i\} + \sum_{i < j, a < b} r_{ij}^{ab} \{\hat{a}^\dagger i \hat{b}^\dagger j\} + \dots$$

One particle attached(core + 1 system)

$$\hat{R} = \sum_a r^a \hat{a}^\dagger + \sum_{a > b, j} r_j^{ba} \hat{b}^\dagger j \hat{a}^\dagger + \dots$$

$$[T_k, R_k] = 0$$

One particle removed(core -1 system)

$$\hat{R} = \sum_i r_i \hat{i} + \sum_{b, j > i} r_{ji}^b \hat{b}^\dagger j \hat{i} + \sum_{b > c, j > k > i} r_{jki}^{bc} \hat{b}^\dagger j \hat{c}^\dagger k \hat{i} + \dots$$

Two particle attached/removed(core +2/-2 system)

EOM-CC

$$H_N R_k e^T |HF\rangle = \Delta E_k R_k e^T |HF\rangle$$

R_k Commutes with T_k

$$\mathcal{H} R_k |HF\rangle = \Delta E_k R_k |HF\rangle$$

$$[\mathcal{H}, R_k] |HF\rangle = (\Delta E_k - \Delta E_0) R_k |HF\rangle$$

H can only contract with R_k from right side(same as T_k) as it contains only particle Or hole generation operator.

$$(\mathcal{H} R_k)_C |HF\rangle = \omega_k R_k |HF\rangle$$

Good Eigenvalue problem

$$(\mathcal{H}R_k)_C\big|HF\rangle=\omega_k R_k \big|HF\rangle$$

$$\big(\hat P \mathcal{H} \hat Q \hat R_k \hat P\big)_{\mathrm C} = \omega_k r_0 \hat P\,,$$

$$\big(\hat Q \mathcal{H} \hat Q \hat R_k \hat P\big)_{\mathrm C} = \omega_k \hat Q \hat R_k \hat P\,.$$

$$\sum_d \chi_{ad}r_i^d - \sum_l \chi_{li}r_l^a + \sum_{dl} \chi_{ladi}r_l^d + \sum_{dl} \chi_{ld}r_{il}^{ad} + \tfrac{1}{2}\sum_{del} \chi_{alde}r_{il}^{de} \\ - \tfrac{1}{2}\sum_{dlm} \chi_{lmid}r_{lm}^{ad} + \tfrac{1}{4}\sum_{delm} \langle lm \| de \rangle r_{ilm}^{ade} = \omega r_i^a \qquad (\text{for all } i,a),$$

$$\hat{P}(ij)\sum_d \chi_{abej}r_i^d - \hat{P}(ab)\sum_l \chi_{lbij}r_l^a + \sum_{dl} \chi_{labdij}r_l^d + \hat{P}(ab)\sum_d \chi_{bd}r_{ij}^{ad} \\ - \hat{P}(ij)\sum_l \chi_{lj}r_{il}^{ab} + \tfrac{1}{2}\sum_{de} \chi_{abde}r_{ij}^{de} + \tfrac{1}{2}\sum_{lm} \chi_{lmi}r_{lm}^{ab} \\ + \hat{P}(ab|ij)\sum_{dl} \chi_{lbdj}r_{il}^{ad} + \tfrac{1}{2}\hat{P}(ij)\sum_{del} \chi_{albdej}r_{il}^{de} - \tfrac{1}{2}\hat{P}(ab)\sum_{dlm} \chi_{lmbidj}r_{lm}^{ae} \\ + \sum_{dl} \chi_{ld}r_{ijl}^{abd} + \tfrac{1}{2}\hat{P}(ab)\sum_{del} \chi_{alde}r_{ilj}^{deb} - \tfrac{1}{2}\hat{P}(ij)\sum_{dlm} \chi_{lmid}r_{lmj}^{adb} \\ = \omega r_{ij}^{ab} \qquad (\text{for all } i > j, \; a > b),$$

$$\omega r_0 = \sum_{dl} \chi_{ld}r_l^d + \tfrac{1}{4}\sum_{delm} \langle lm \| de \rangle r_{lm}^{de} \, .$$

$$\omega r_0 = \text{Diagram A} + \text{Diagram B} \, .$$

$$\text{Diagram A} = \text{Diagram A}_1 + \text{Diagram A}_2 + \text{Diagram A}_3 + \text{Diagram A}_4 + \text{Diagram A}_5 \\ + \text{Diagram A}_6 + \text{Diagram A}_7$$

$$\text{Diagram B} = \text{Diagram B}_1 + \text{Diagram B}_2 + \text{Diagram B}_3 + \text{Diagram B}_4 + \text{Diagram B}_5 + \text{Diagram B}_6 + \text{Diagram B}_7 + \text{Diagram B}_8 + \text{Diagram B}_9 + \text{Diagram B}_{10}$$

$$\text{Diagram C} = \text{Diagram C}_1 + \text{Diagram C}_2 + \text{Diagram C}_3 + \text{Diagram C}_4 + \text{Diagram C}_5 + \text{Diagram C}_6 + \text{Diagram C}_7 + \text{Diagram C}_8 + \text{Diagram C}_9 + \text{Diagram C}_{10} \\ + \text{Diagram C}_{11} + \text{Diagram C}_{12} + \text{Diagram C}_{13} + \text{Diagram C}_{14} + \text{Diagram C}_{15} + \text{Diagram C}_{16} + \text{Diagram C}_{17} + \text{Diagram C}_{18} + \text{Diagram C}_{19} + \text{Diagram C}_{20}$$

Operator expectation of EOM-CC

Left wavefunction:

$$\langle 0 | \hat{L}_k \mathcal{H} = \langle 0 | \hat{L}_k \Delta E_k .$$

$$\hat{L}_k = l_0 + \sum_{i,a} l_a^i \{ \hat{i}^\dagger \hat{a} \} + \sum_{i < j, a < b} l_{ab}^{ij} \{ \hat{i}^\dagger \hat{a} \hat{j}^\dagger \hat{b} \} + \dots$$

Neither left wf or right wf are normalized to one:

$$\langle 0 | \hat{L}_k \hat{R}_l | 0 \rangle = \delta_{kl} .$$

$$\overline{O} = \frac{\langle \Psi | \hat{O} | \Psi \rangle}{\langle \Psi | \Psi \rangle} .$$

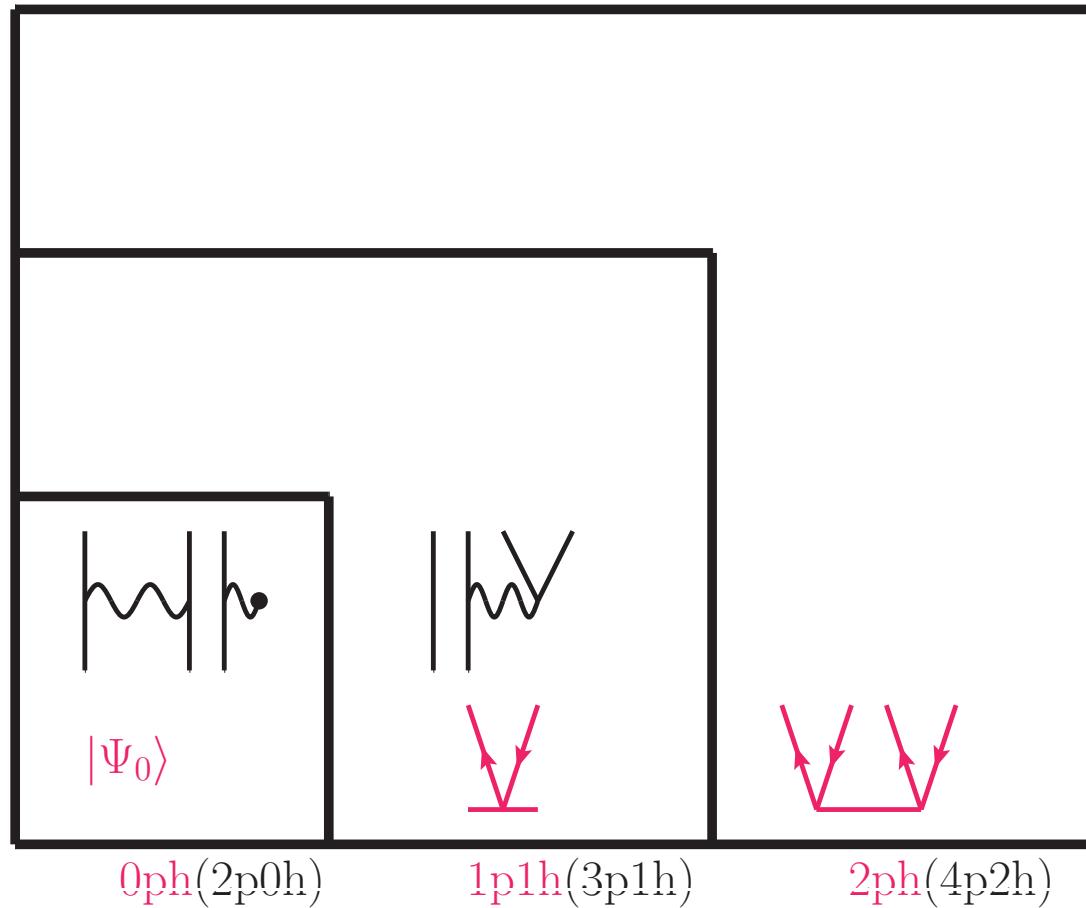
For energy:

$$\langle 0 | \hat{L}_k \mathcal{H} \hat{R}_l | 0 \rangle = E_l \delta_{lk}$$

For other observables:

$$\langle 0 | \hat{L}_k \mathcal{O} \hat{R}_l | 0 \rangle = \langle \mathcal{O} \rangle$$

Other method: SMCC



Thanks

Non-Hermitian Hamiltonian

- Expectation value of coupled cluster

$$\overline{O} = \frac{\langle \Psi | \hat{O} | \Psi \rangle}{\langle \Psi | \Psi \rangle}.$$

$$\begin{aligned}\overline{O} &= \frac{\langle 0 | e^{\hat{T}^\dagger} \hat{O} e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} \\ &= \frac{\langle 0 | [1 + T^\dagger + \frac{1}{2!}(T^\dagger)^2 + \frac{1}{3!}(T^\dagger)^3 + \dots] \hat{O} [1 + T + \frac{1}{2!}T^2 + \frac{1}{3!}T^3 + \dots] | 0 \rangle}{\langle 0 | [1 + \hat{T}^\dagger + \frac{1}{2!}(\hat{T}^\dagger)^2 + \frac{1}{3!}(\hat{T}^\dagger)^3 + \dots] [1 + \hat{T} + \frac{1}{2!}\hat{T}^2 + \frac{1}{3!}\hat{T}^3 + \dots] | 0 \rangle}.\end{aligned}$$

$$\begin{aligned}\Delta E &= \frac{\langle 0 | e^{\hat{T}^\dagger} \hat{H}_N e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} = \frac{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} (\hat{P} + \hat{Q}) e^{-\hat{T}} \hat{H}_N e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} \\ &= \frac{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle \langle 0 | e^{-\hat{T}} \hat{H}_N e^{\hat{T}} | 0 \rangle}{\langle 0 | e^{\hat{T}^\dagger} e^{\hat{T}} | 0 \rangle} = \langle 0 | e^{-\hat{T}} \hat{H}_N e^{\hat{T}} | 0 \rangle = \langle 0 | \mathcal{H} | 0 \rangle.\end{aligned}$$

- CC decoupled the Hamiltonian **only** for energy

Failure of Feynman-Hellman

Solution to Left wave function of CC

$$\hat{P}(1 + \Lambda \hat{Q})(\mathcal{H} - \Delta E) = 0$$

$$\Lambda = \Lambda_1 + \Lambda_2 + \Lambda_3 + \cdots,$$

where

$$\begin{aligned}\Lambda_1 &= \sum_{ia} \lambda_a^i \{i^\dagger a\}, \\ \Lambda_2 &= \frac{1}{4} \sum_{ijab} \lambda_{ab}^{ij} \{i^\dagger a j^\dagger b\}, \\ \Lambda_n &= \frac{1}{(n!)^2} \sum_{\substack{ij\dots \\ ab\dots}} \lambda_{ab\dots}^{ij\dots} \{i^\dagger a j^\dagger b \dots\},\end{aligned}$$

$$\hat{P}\mathcal{H}\hat{Q} + \hat{P}(\Lambda\mathcal{H})_C\hat{Q} + \hat{P}\mathcal{H}\hat{Q}\Lambda\hat{Q} = 0$$

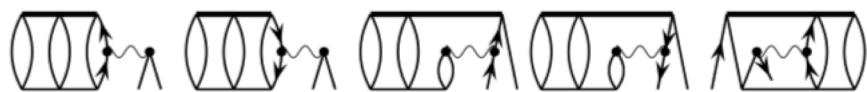
$$\begin{aligned}&\langle 0 | \hat{H}_N e^{\hat{T}} | \Phi_{ij\dots}^{ab\dots} \rangle_C + \langle 0 | \Lambda(\hat{H}_N e^{\hat{T}})_C | \Phi_{ij\dots}^{ab\dots} \rangle_C \\ &+ \sum_{\substack{k < l < \dots \\ c < d < \dots}} \langle 0 | \hat{H}_N e^{\hat{T}} | \Phi_{kl\dots}^{cd\dots} \rangle_C \langle \Phi_{kl\dots}^{cd\dots} | \Lambda | \Phi_{ij\dots}^{ab\dots} \rangle = 0.\end{aligned}$$



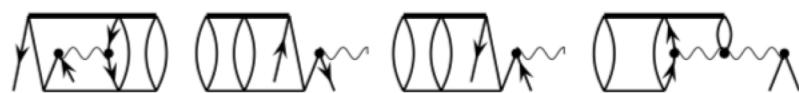
1 3 4 5 18 19 26 27



48 49 50 51 52



53 54 55 56 57



58 59 60 75



1 2 4 5 8 9 10



11 12 23 24 33 34



45 46 47 48 55



56 57 58 59