

Projectes d'Enginyeria Física - 2

Modelling and designing a Paul ion trap

V2024.2

Juan M. Rius

Universitat Politècnica de Catalunya (UPC)

Barcelona, Spain

- Introduction to the Paul ion trap
- Introduction to Computational Electromagnetics
- Modeling the problem with Integral Equation methods
- Discrete equations with the Method of Moments
- Implementation in 2D: validation with capacitor
- Implementation in 3D: validation with capacitor
- Analyze and design a Paul Trap:
 - Static potential, one ion
 - AC potential, one ion
 - AC potential, multiple ions
 - Design the trap for optimum ion confinement

The Paul ion trap

- **Pioneers of ion trapping:
Hans Dehmelt and Wolfgang Paul (Nobel price 1989)**

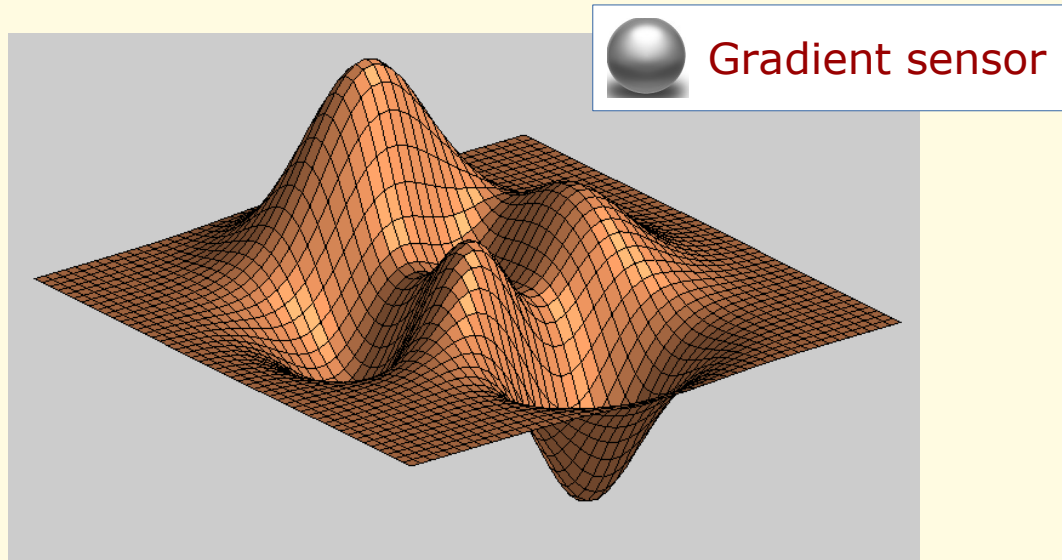


大



The Paul ion trap

The force on ions follows the direction of field: $\vec{E} = -\nabla V$



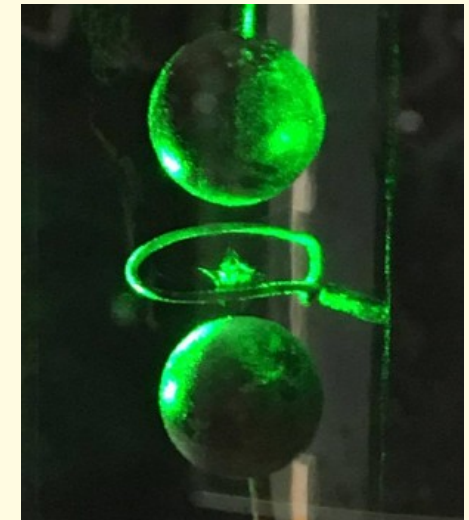
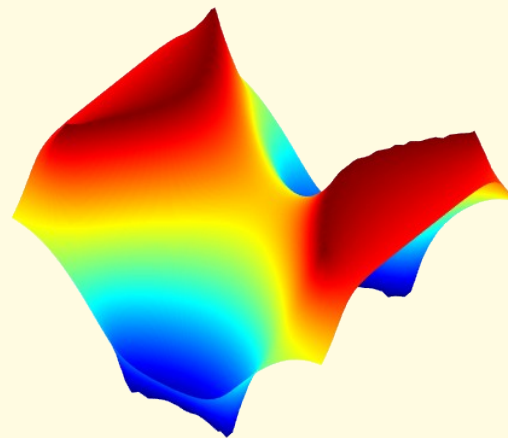
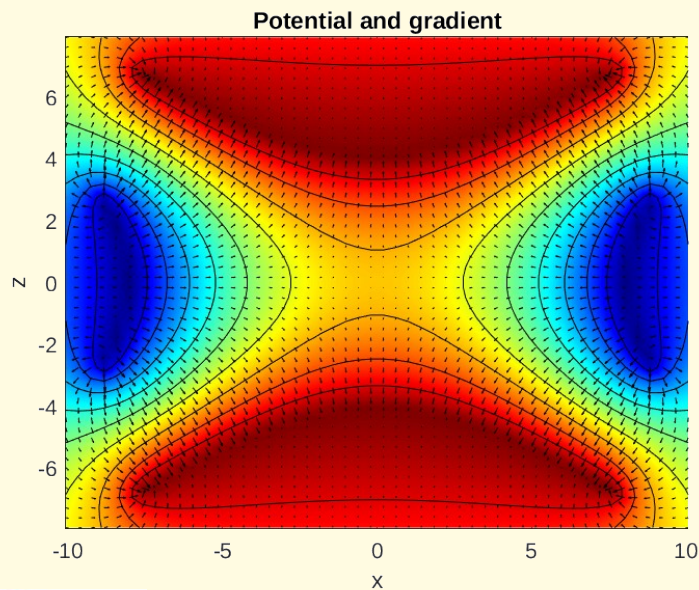
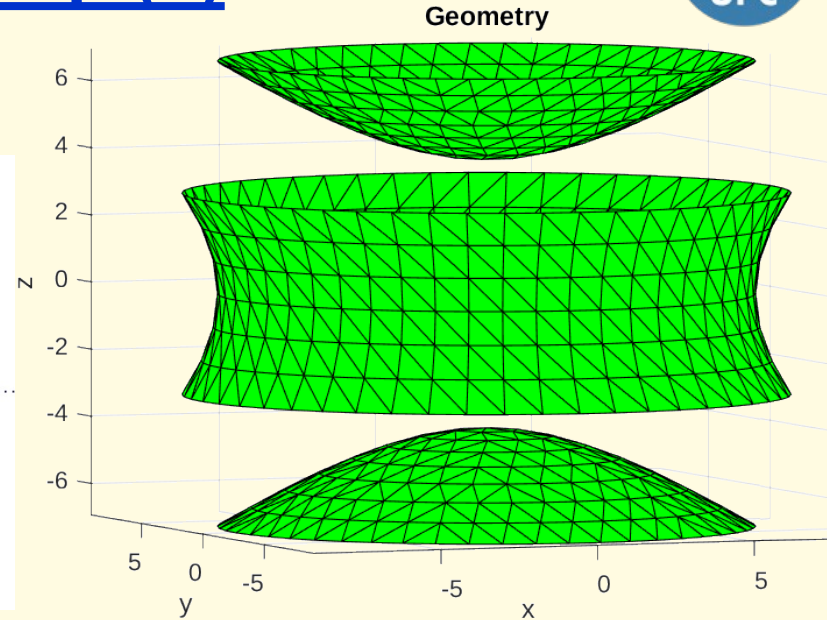
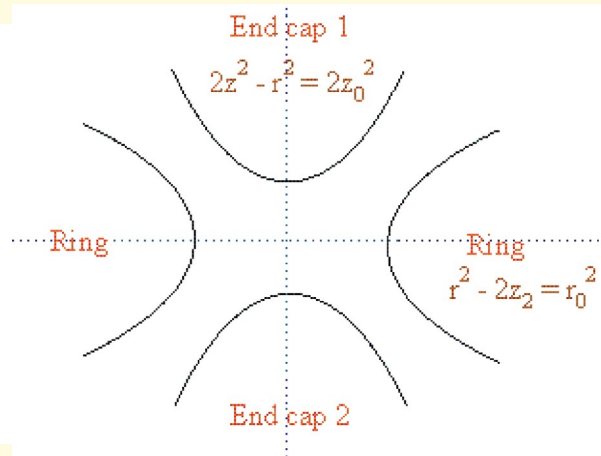
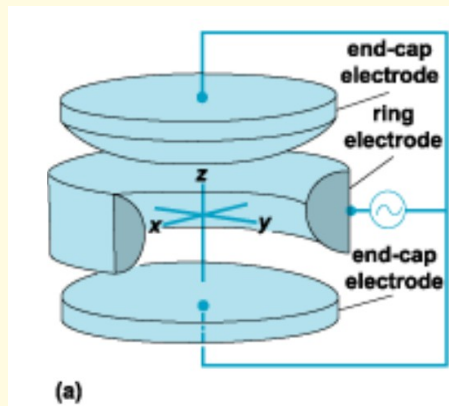
- It is not possible to have a 3D minimum or maximum of $V(\mathbf{r})$, since

$$\nabla^2 V(\vec{r}) = \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) V(\vec{r}) = 0$$

- Anyway, it would work only for positive or negative ions

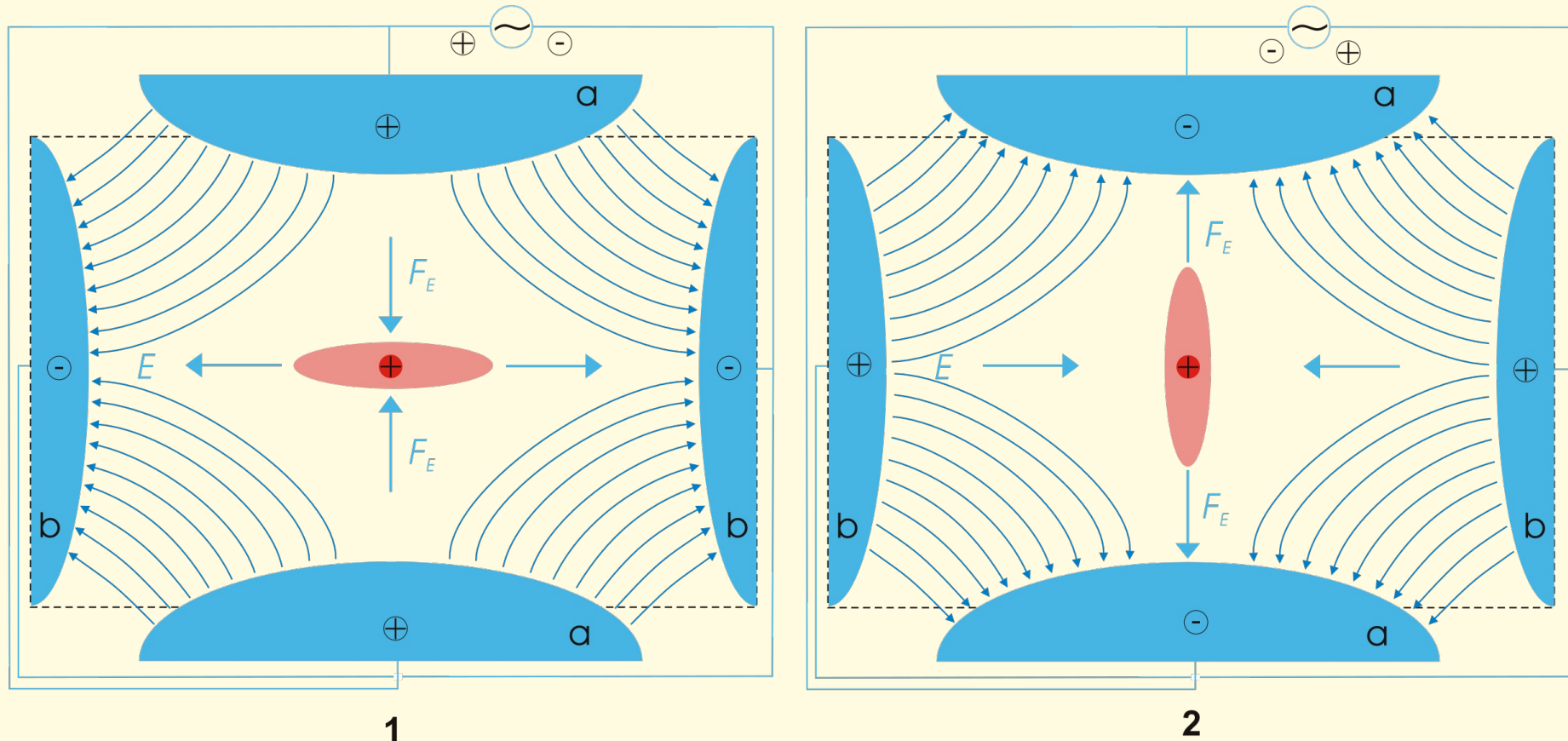
The Paul ion trap (2)

■ Quadrupole:



The Paul ion trap (3)

■ Quadrupole with AC potential variation (low freq.)

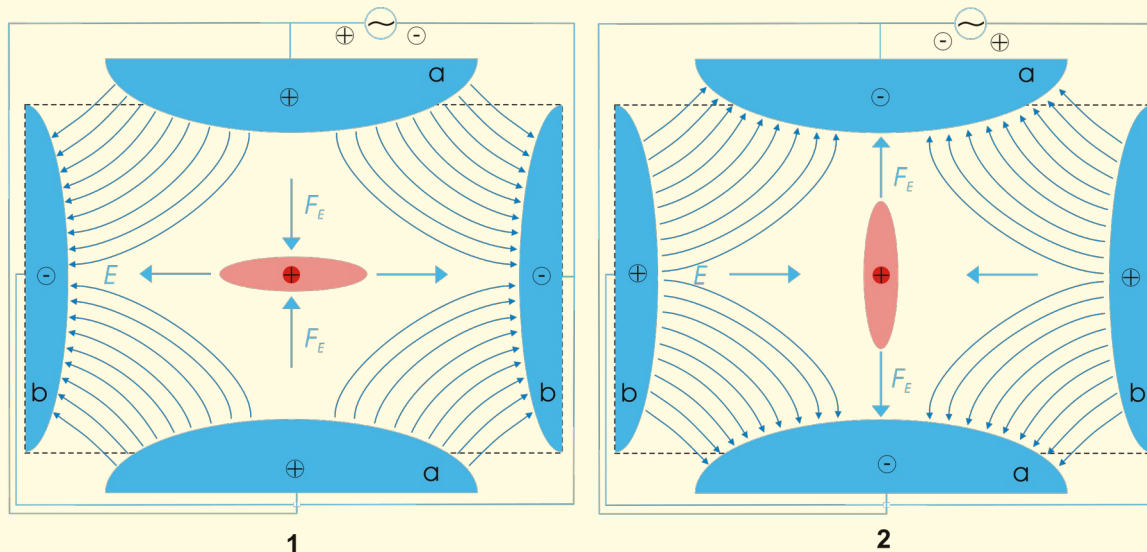


- Now, ions rotate movement direction, following AC changes, and become confined at the center of the trap

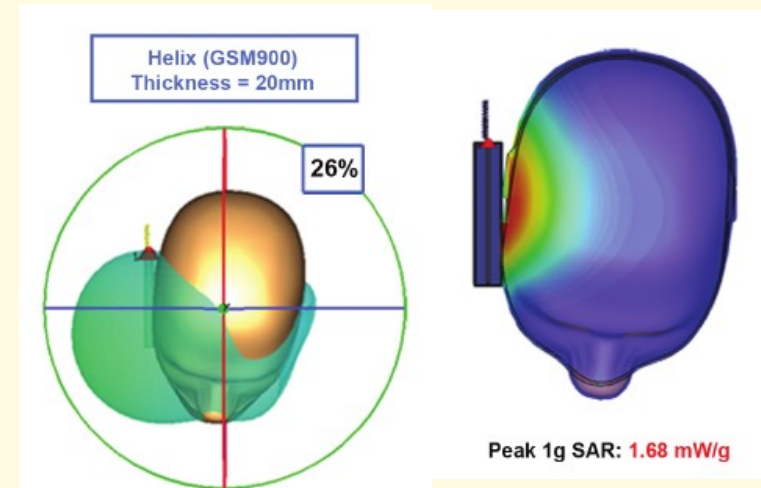
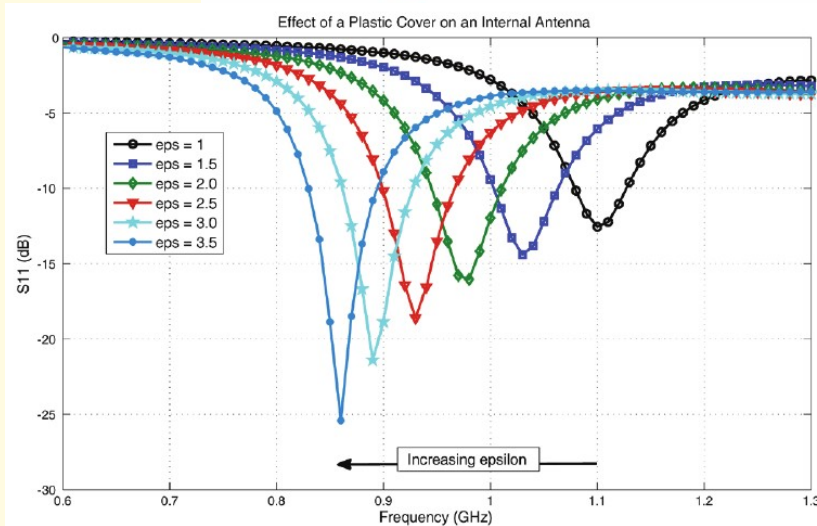
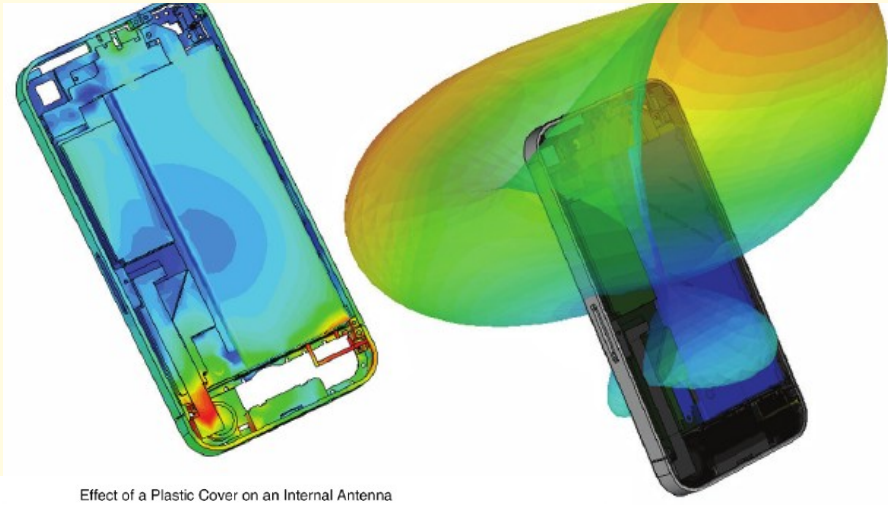
The Paul ion trap (4)

■ Objective:

- Develop computer code for computing the potential due to hyperbolic quadrupole.
- Compute ion trajectories within the trap.
- Account for AC variation of potential → ion trajectories
- Design size of the trap, AC frequency, etc... for ion confinement with weights and charge.

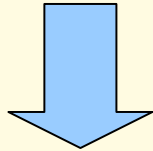


- Antenna analysis and design using numerical simulation software:



$$\hat{n} \times \vec{E}^i(\vec{r}) = \hat{n} \times \iint_S \left[jk\eta G(\vec{r}' - \vec{r}) \vec{J}_s(\vec{r}') + \frac{\eta}{jk} \nabla' G(\vec{r}' - \vec{r}) \nabla' \cdot \vec{J}_s(\vec{r}') \right] ds$$

Diff. or integ. equations



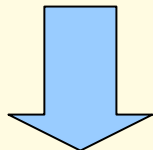
Discretization

$$G = \frac{e^{-jk|\vec{r} - \vec{r}'|}}{4\pi|\vec{r} - \vec{r}'|}$$

Singular!

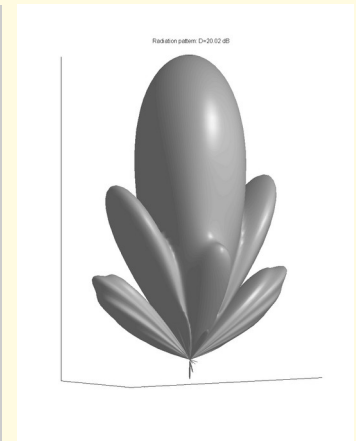
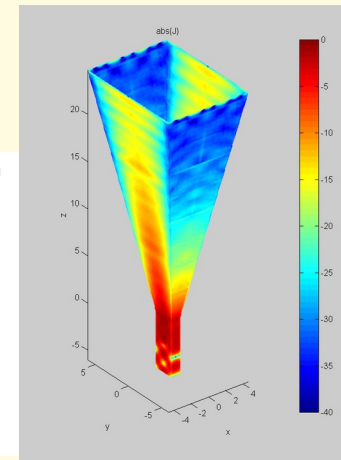
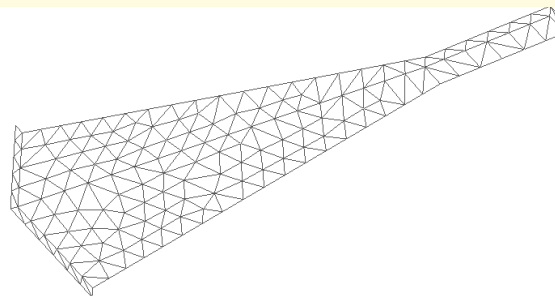
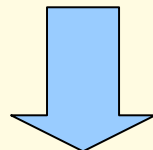
Linear system of equations

Full matrix with
millions of unknowns!



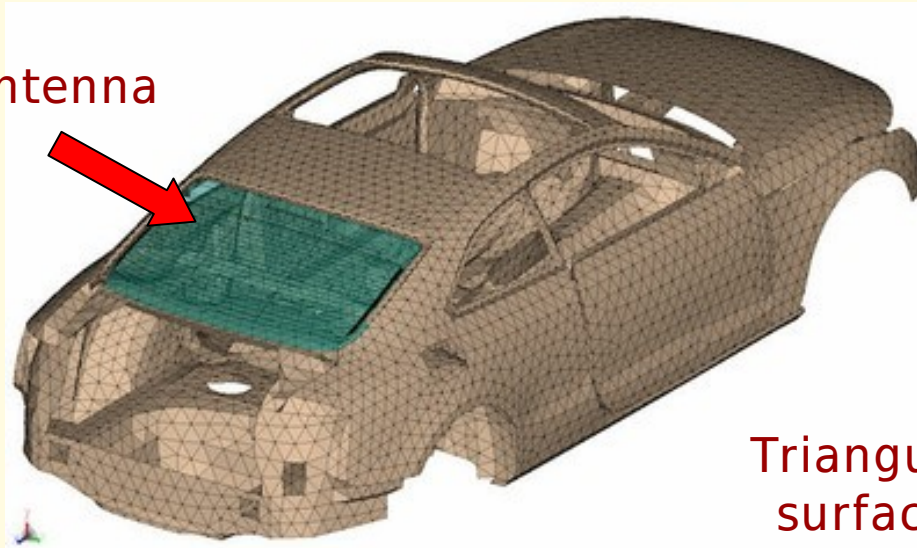
Fast Solver

Solution $\vec{J}(\vec{r}')$

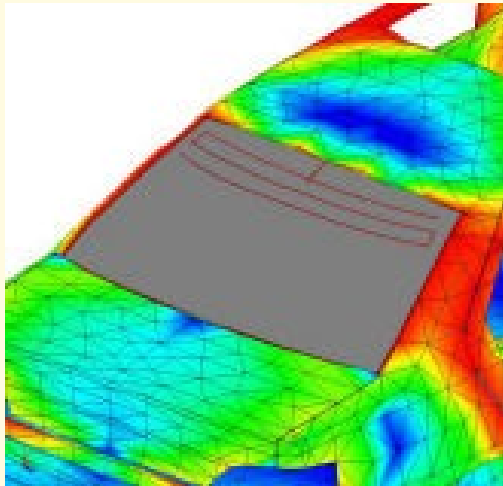


Antenna parameters: input impedance, radiation pattern, etc.

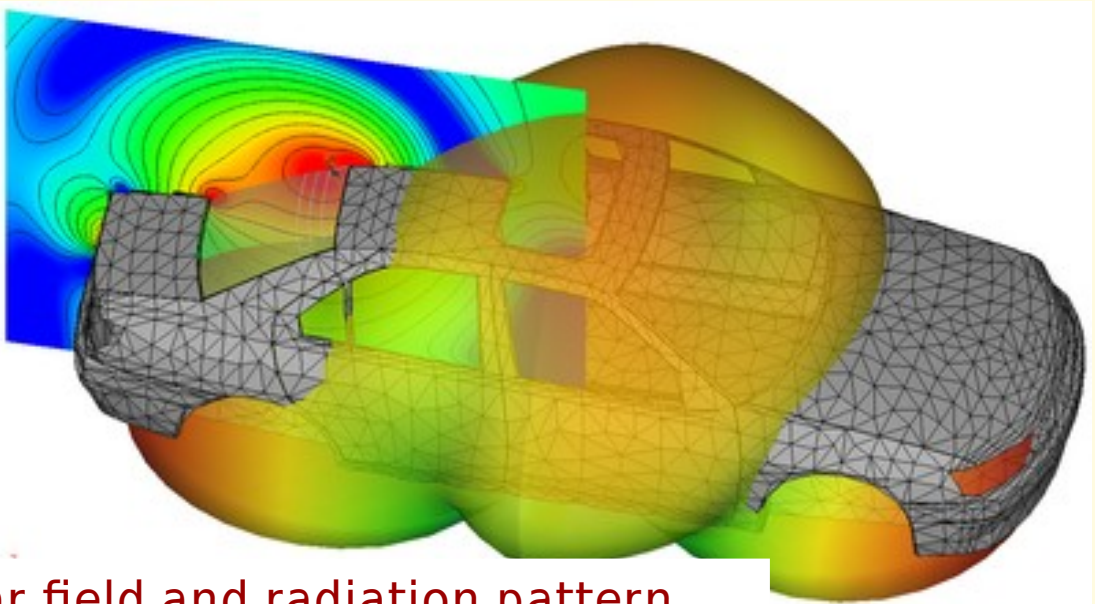
Antenna



Triangular mesh
surface model



Induced currents



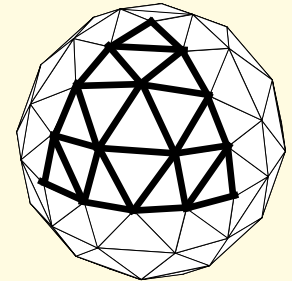
Near field and radiation pattern

■ Formulation of (linear) differential or integral equations

- Time or frequency domain
- Examples: Maxwell's or wave equations, EFIE, MFIE,....

■ Discretization of equations into a linear system

- Method of moments, finite element method, etc..

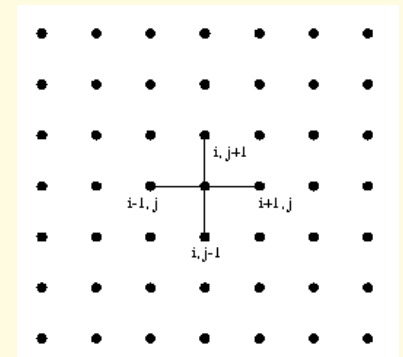


■ Discretization of boundary conditions

- Discretize surfaces into triangle meshes....
- Set numerical boundary condition (samples, min sq error, ...)

■ Set mesh truncation conditions

- Only for differential equations, not necessary for IE
- Finite mesh: a different equation must be set at mesh truncation nodes



- **Simplifications and approximations in the equations:**
 - Problem modeling
 - Material parameters
- **Discretization:**
 - Equations into linear system (projection into finite-dimensional space)
 - Numerical integral computation (often with singularities)
 - Boundary conditions surfaces into meshes
 - Boundary conditions into numbers
 - Truncation mesh conditions are always approximate
- **Linear system solution**
 - Truncation of real numbers into a finite-length word
 - Iterative methods
 - Fast solvers for huge linear systems: always involve approximations

There are many things than can go wrong!!

Integral equation methods

■ Transform DE into IE with Green's function:

$\phi(\vec{r}) = \mathcal{L}f(\vec{r})$ \mathcal{L} =linear operator, $\phi(\vec{r})$ =field/potential, $f(\vec{r})$ = sources (invariant)

$G(\vec{r}) = \mathcal{L}\delta(\vec{r})$ G is the impulse response or **Green's function**

$$\phi(\vec{r}) = \int_S f(\vec{r}')G(\vec{r} - \vec{r}')d\vec{r}'$$

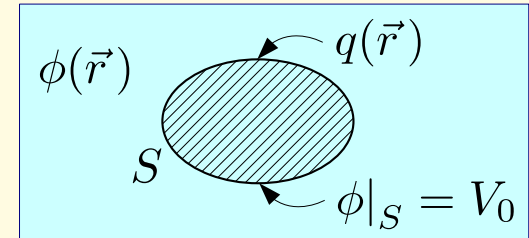
■ Set boundary condition: $\phi(\vec{r})|_S = \phi_0$

$$\left. \int_S f(\vec{r}')G(\vec{r} - \vec{r}')d\vec{r}' \right|_S = \phi_0$$

Solution domain on S instead of whole volume
 \Rightarrow **discretize only on S** \Rightarrow much less unknowns

■ Electrostatic potential (3D):

– Poisson's equation: $\nabla^2 \phi(\vec{r}) = -\frac{q(\vec{r})}{\varepsilon}$



– Green's function: $\nabla^2 G(\vec{r}) = -\delta(\vec{r})$ $G = \frac{1}{4\pi\varepsilon|\vec{r}|}$

$$\phi(\vec{r}) = \int_S q(\vec{r}') \frac{1}{4\pi\varepsilon|\vec{r} - \vec{r}'|} d\vec{r}'$$

– Boundary condition:

$$\phi(\vec{r})|_S = V_0$$

– Integral equation:

$$\int_S q(\vec{r}') \frac{1}{4\pi\varepsilon|\vec{r} - \vec{r}'|} d\vec{r}' \Big|_S = V_0$$

■ Potential Green's function in 2D:

- Green's function: Potential due to a **line of charge**
- First we compute E field due to a point source $q_0=1$,
 E must be $\vec{E}(\rho, \phi) = E_\rho(\rho)\hat{\rho}$ and according to Gauss' law:

Per unit length: $\frac{q_0}{\varepsilon} = \oint \vec{E} \cdot \hat{n} dl = 2\pi\rho E_\rho \Rightarrow \vec{E} = \frac{q_0}{2\pi\varepsilon\rho}\hat{\rho}$

- The potential is such that $\vec{E} = -\nabla\phi = -\frac{\partial}{\partial\rho}\phi(\rho)\hat{\rho} = \frac{q_0}{2\pi\varepsilon\rho}\hat{\rho}$

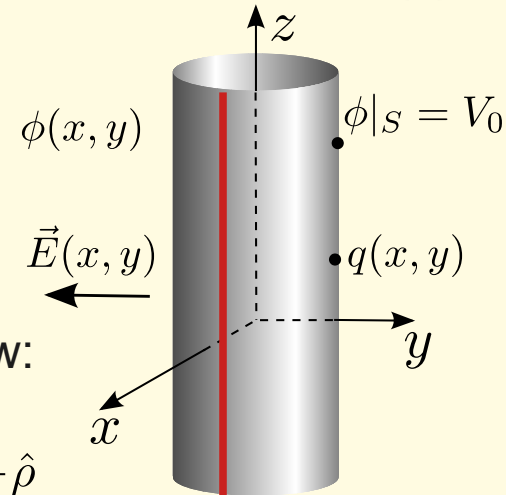
so, the potential $\phi = \frac{-q_0}{2\pi\varepsilon} \ln \rho$ is the Green's function in 2D ($q_0 = 1$)

$$G = \frac{-1}{2\pi\varepsilon} \ln \rho$$

$$\phi(\vec{\rho})|_C = V_0$$

$$\left. \frac{-1}{2\pi\varepsilon} \int_C q(\vec{\rho}') \ln(|\vec{\rho} - \vec{\rho}'|) d\ell' \right|_C = V_0$$

Electrostatics integral equation in 2D



Method of Moments (MoM) (1)

- The most commonly used method to discretize electromagnetic integral equations
- Valid for any linear equation (also differential eq.)

$$\mathcal{L}X(\vec{r}) = Y(\vec{r})$$

\mathcal{L} is a **linear** operator

X is the **unknown** function

Y is a **known** function

it will be discretized as a linear system $[Z][a] = [b]$

$$\frac{-1}{2\pi\epsilon} \int_C q(\vec{\rho}') \ln(|\vec{\rho} - \vec{\rho}'|) d\ell' \Big|_C = V_0$$

Electrostatics integral equation in 2D

$$\int_S q(\vec{r}') \frac{1}{4\pi\epsilon|\vec{r} - \vec{r}'|} d\vec{r}' \Big|_S = V_0$$

Electrostatics integral equation in 3D

Method of Moments (MoM) (2)

- The unknown charge is discretized as a linear combination of **basis functions**:

$$q_N(\vec{r}) = \sum_{n=1}^N q_n x_n(\vec{r})$$

$$\mathcal{L}q(\vec{r}) = V_0(\vec{r}) \quad \mathcal{L}q_N(\vec{r}) = \sum_{n=1}^N q_n \mathcal{L}x_n(\vec{r}) \approx V_0(\vec{r})$$

The **unknowns are now the coefficients** q_n

- But the equation is still a functional equation, and we need a linear system to solve with a computer.
- We want a very small (negligible) residual error:

$$R = V_0(\vec{r}) - \sum_{n=1}^N q_n \mathcal{L}x_n(\vec{r}) \approx 0$$

Method of Moments (MoM) (3)

- We set to zero inner products of the residual error

$$R(\vec{r}) = V_0(\vec{r}) - \sum_{n=1}^N q_n \mathcal{L}x_n(\vec{r})$$

with a set of **weighting functions** $w_m(\vec{r})$

where the inner product is: $\langle w(\vec{r}), f(\vec{r}) \rangle = \int w^*(\vec{r}) f(\vec{r}) d\vec{r}$

$$\langle w_m(\vec{r}), R(\vec{r}) \rangle = \langle w_m(\vec{r}), V_0(\vec{r}) \rangle - \sum_{n=1}^N q_n \langle w_m(\vec{r}), \mathcal{L}x_n(\vec{r}) \rangle = 0$$

$$w_m(\vec{r}) = \delta(\vec{r} - \vec{r}_m) \Rightarrow R(\vec{r}_m) = V_0(\vec{r}_m) - \sum_{n=1}^N q_n \mathcal{L}x_n(\vec{r})|_{\vec{r}=\vec{r}_m} = 0$$

“Point matching”

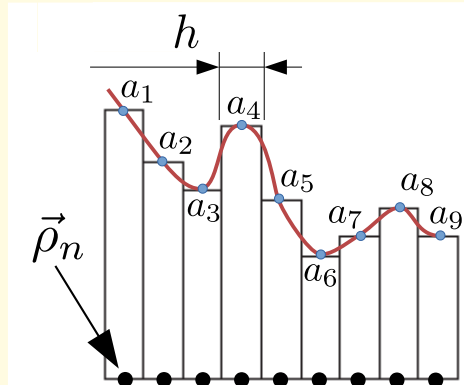
- **Linear system:**

$$[Z][q] = [b]$$

$$Z_{mn} = \mathcal{L}x_n(\vec{r})|_{\vec{r}=\vec{r}_m}$$

$$b_m = V_0(\vec{r}_m)$$

■ Pulse basis functions and point matching:



Unit pulse basis functions

$$\frac{-1}{2\pi\epsilon} \int_S q(\vec{\rho}') \ln(|\vec{\rho} - \vec{\rho}'|) d\vec{\rho}' \Big|_S = V_0$$

$$[Z][q] = [b] \quad \begin{aligned} Z_{mn} &= \mathcal{L}x_n(\vec{\rho})|_{\vec{\rho}=\vec{\rho}_m} \\ b_m &= V_0(\vec{\rho}_m) \end{aligned}$$

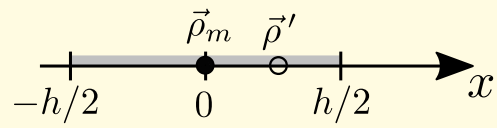
$$\mathcal{L}x_n(\vec{\rho}) = \frac{-1}{2\pi\epsilon} \int_S \Pi\left(\frac{\vec{\rho}' - \vec{\rho}_n}{h_n}\right) \ln(|\vec{\rho} - \vec{\rho}'|) d\vec{\rho}'$$

- With 1-point integration: $\mathcal{L}x_n(\vec{\rho}) \approx \frac{-1}{2\pi\epsilon} h_n \ln(|\vec{\rho} - \vec{\rho}_n|)$
- Point matching at $\vec{\rho}_m$: $Z_{mn} = \frac{-h_n}{2\pi\epsilon} \ln(|\vec{\rho}_m - \vec{\rho}_n|), \quad b_m = V_0(\vec{\rho}_m)$

$$Z_{mn} = \frac{-h_n}{2\pi\epsilon} \ln(|\vec{\rho}_m - \vec{\rho}_n|), \quad b_m = V_0(\vec{\rho}_m)$$

■ **Self-interaction** (diagonal terms): when $\vec{\rho}_m = \vec{\rho}_n$, $Z_{mm} = \infty$

For $m = n$ we have to do the source integral analytically



$$Z_{mm} = \mathcal{L}x_m(\vec{\rho}_m) = \frac{-1}{2\pi\epsilon} \int_S \Pi\left(\frac{\vec{\rho}' - \vec{\rho}_m}{h_m}\right) \ln(|\vec{\rho}_m - \vec{\rho}'|) d\vec{\rho}'$$

$$= \frac{-1}{2\pi\epsilon} \int_{x_m} \ln(|\vec{\rho}_m - \vec{\rho}'|) d\vec{\rho}' = \frac{-1}{2\pi\epsilon} \int_{-h_m/2}^{h_m/2} \ln|x| dx$$

$$= \frac{-1}{2\pi\epsilon} 2 \int_0^{h_m/2} \ln x dx = \frac{-1}{\pi\epsilon} [x(\ln x - 1)]_0^{h_m/2} = \boxed{\frac{-h_m}{2\pi\epsilon} \left[\ln\left(\frac{h_m}{2}\right) - 1 \right] = Z_{mm}}$$

■ Discretization of the geometry:

- Use complex numbers

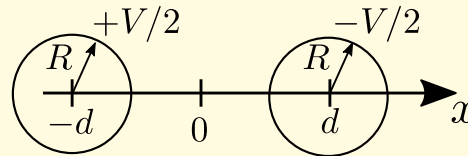
$$\begin{aligned}\vec{\rho} &: r = x + j y; \\ \vec{\rho}_1 - \vec{\rho}_2 &: r1-r2; \\ |\vec{\rho}_1 - \vec{\rho}_2| &: \text{abs}(r1-r2)\end{aligned}$$

- Write a function that returns a vector `rn` with the center point of the basis functions and another vector `hn` with the lengths

■ 2D plot:

- Use meshgrid: $x = \text{linspace}(\dots), y = \text{linspace}(\dots)$
 $[xx,yy] = \text{meshgrid}(x,y); rr = xx + 1j*yy;$
- Loop BF: $V=\text{zeros}(\dots); \text{for } n=1:N, R = \text{abs}(rr-rn(n)); V=V+\dots; \text{end}$
- Potential and field: use $\text{surf}(), \text{pcolor}(), \text{contour}(), \text{gradient}(), \text{quiver}()$

■ Error check:



- Exact capacitance of two infinite cylinders (per u.L.): $C = \frac{\pi \epsilon}{\text{arccosh}\left(\frac{d}{R}\right)}$

- Computation $C = \frac{Q}{V_0} = \frac{1}{V_0} \sum_{n=1}^N h_n q_n = \text{sum}(hn.*qn)/V$

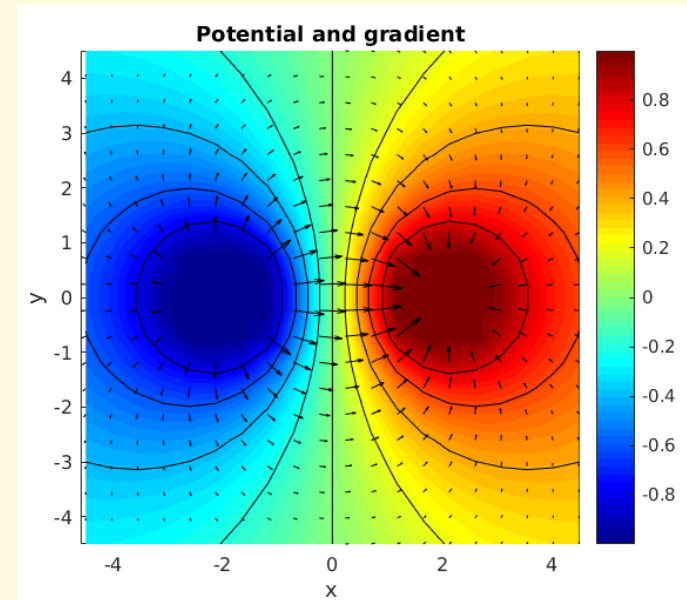
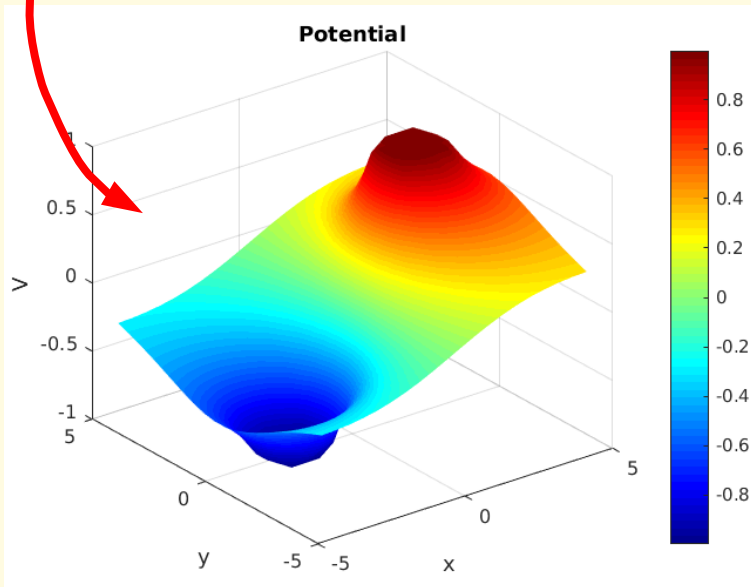
Programming tips in 2D (2)

■ Post-processing:

– 2D plots:

```
x = linspace(...), y = linspace(...);
[xx,yy] = meshgrid(x,y); rr = xx + 1j*yy;

for ..., V = ... ; end; % Computation of potential at xx,yy
[Ex, Ey] = gradient(-V);
surf(x,y,V);
pcolor(x,y,V); quiver(xx,yy,Ex,Ey); contour(xx,yy,V):
colormap jet; shading interp; colorbar
```



■ Object geometry:

- Create functions that return `rn` and `hn` for planar and circular objects.
- Create an object consisting of two (planar or circular) plates.
- Plot the geometry.

■ Linear System:

- Compute linear system elements.
- Compute independent term: set one plate to $+V/2$ and the other to $-V/2$.
- Solve linear system and compute capacitance. Check with reference.
- Compute potential and field.

■ Plot potential and E field:

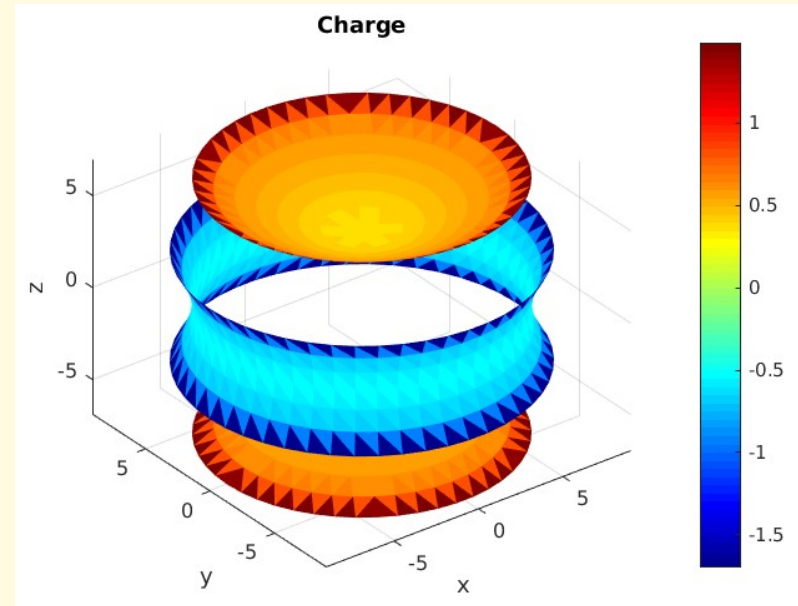
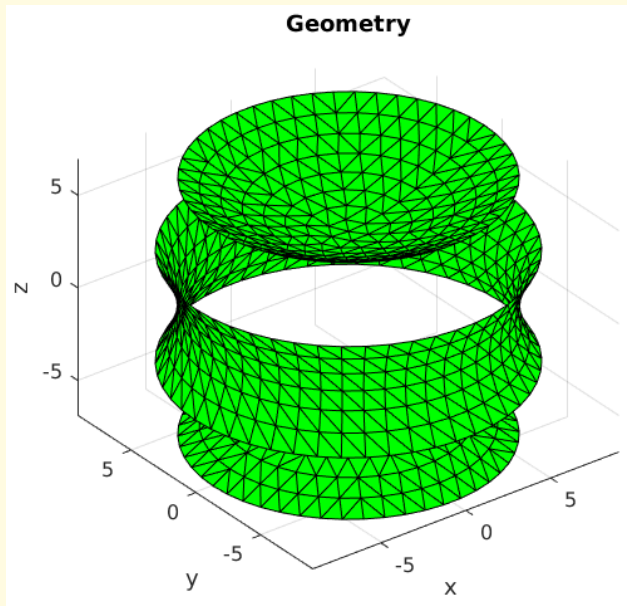
- Obtain `surf()`, `pcolor()`, `contour()` and `quiver()` plots of potential and field.

MoM Electrostatics 3D

$$\int_S q(\vec{r}') \frac{1}{4\pi\epsilon|\vec{r} - \vec{r}'|} d\vec{r}' \Big|_S = V_0$$

■ Triangular basis functions and point matching:

- Planar triangles mesh
- Constant charge on each triangle



■ MoM matrix elements:

- The integral of $\mathcal{L}x_n(\vec{r})$ can be computed analytically

(see Rao et al., “A simple numerical solution procedure for statics problems involving arbitrarily shaped surfaces” IEEE Trans AP, Vol. 27, No. 5, sept 1979)

- Matlab function by E. Úbeda

Computation of integral and point matching

Returns **a row** of [Z]: single observation point and all source triangles.

$$\int_{S_{T_n}} \frac{1}{|\vec{r}_m - \vec{r}'|} d\vec{r}'$$

```
int_S_1divR(rf, r_1, r_2, r_3, un_S, cent_S);
```

Output: intS

Input: rf: field point [3x1]

r_1 : first vertex of triangles [3xNt]

r_2 : second vertex of triangles [3xNt]

r_3 : third vertex of triangles [3xNt]

un_S: normal vectors to the set of triangles <r1-r2-r3> [3xNt]

< **IMPORTANT**: r1-r2-r3 must turn in the same sense as un_S >

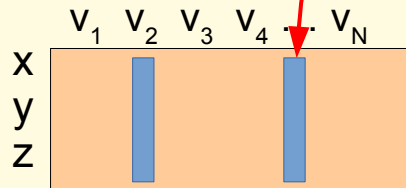
cent_S: set of centroids of the triangles <r1-r2-r3> [3xNt]

	T ₁	T ₂	T ₃	T ₄	...	T _N
x						
y						
z						

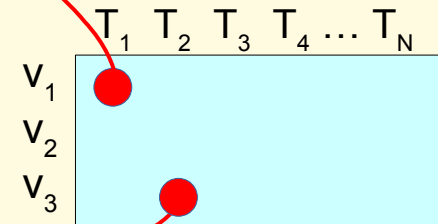
■ Discretization of the geometry

- Use a vertex matrix and a topology matrix

Vertex matrix:
coordinates
of vertices



Topology matrix:
vertices v_1, v_2, v_3
of triangles



- Write a function that returns a structure with 'vertex' and 'topol' fields:
obj.vertex and obj.topol = vertex and topology matrices
 - Create a cloud of surface points, possibly with `x = linspace(...)`,
`y = linspace(...)`, `[xx,yy] = meshgrid(x,y)`, `zz = fun(xx,yy)`
 - `obj.vertex = [xx(:) yy(:) zz(:)]'`;
 - Use `obj.topol = delaunay(xx(:), yy(:))'`; Matlab function
- We also provide a function that reads a FEM triangular mesh from a file created with GiD mesher (www.gidhome.com) and returns the obj struct

■ Geometry management: `int_S_ldivR()` function parameters

- **v1, v2, v3 of all triangles:**

```
v1 = obj.vertex(:,obj.topol(1,:));
v2 = obj.vertex(:,obj.topol(2,:));
v3 = obj.vertex(:,obj.topol(3,:));
```

- **Centroid of all triangles:**

```
obj.cent = (v1+v2+v3)/3;
```

- **Unit normal and area of triangles:**

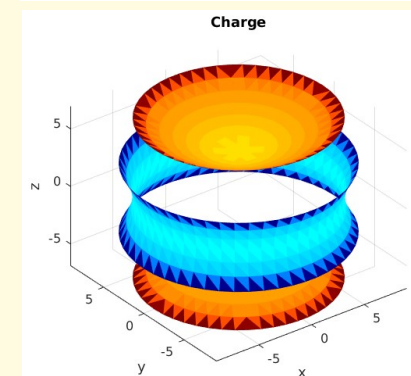
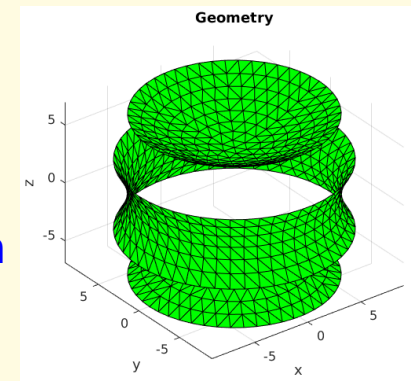
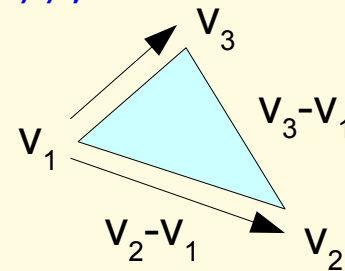
```
c = cross(v2-v1, v3-v1);
obj.ds = sqrt(sum(c.^2))/2;
obj.un = c./repmat(2*obj.ds,3,1); % Normalization
```

- **View geometry:**

```
patch('Faces',obj.topol,'Vertices',obj.vertex',
'FaceColor','g','EdgeColor','k');
axis equal;
```

- **With charge:**

```
patch('Faces',obj.topol,'Vertices',obj.vertex',
'CData',q,'FaceColor','flat','EdgeColor','none');
axis equal;
```

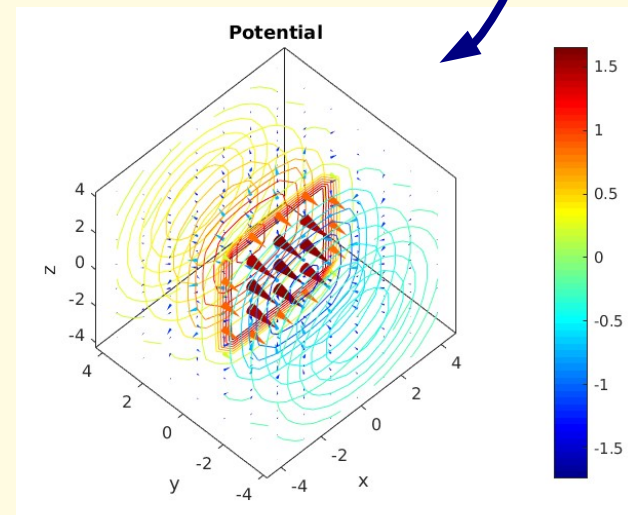
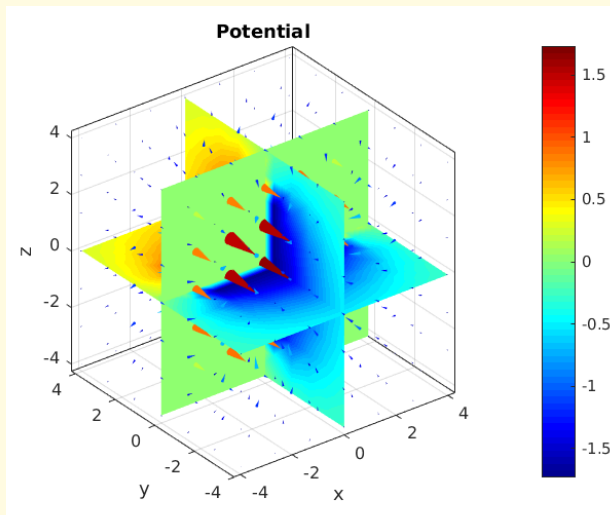


Programming tips in 3D (3)

■ Post-processing:

– 3D plots:

```
x = linspace(...), y = linspace(...); z = linspace(...);
[xx,yy,zz] = meshgrid(x,y,z);
for ..., V = ... ; end; % Computation of potential at xx,yy,zz
[Ex, Ey, Ez] = gradient(-V);
slice(xx,yy,zz, V, xs,ys,zs);
contourslice(xx,yy,zz, V, xs,ys,zs);
coneplot(xx,yy,zz, Ex,Ey,Ez, cxx,cyy,czz, absE);
colormap jet; shading interp; colorbar
```



■ Object geometry:

- Create function that returns `obj.vertex` and `obj.topol` for planar plate.
- Create an object consisting of two planar plates. Plot the geometry.

■ Linear System:

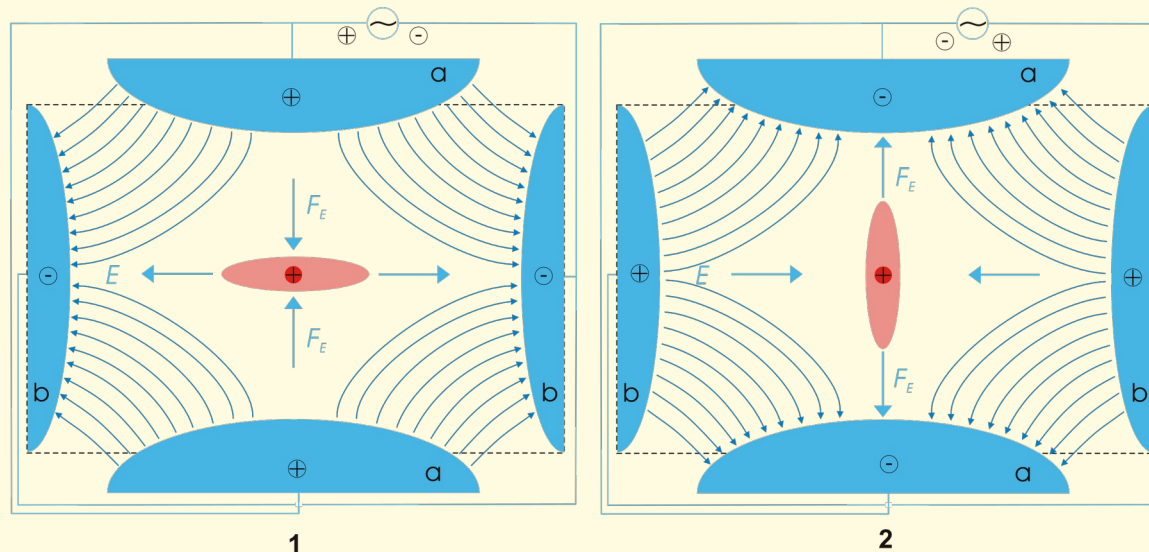
- Compute the arguments of `int_S_1divR()` function.
- Compute the linear system elements.
- Compute independent term: set one plate to $+V/2$ and the other to $-V/2$.
- Solve linear system and compute capacitance. Check with reference.
- Compute potential and field.

■ Plot potential and E field:

- Obtain `slice()`, `contourslice()` and `coneplot()` plots of potential and field.

Methodology:

- Develop computer code for computing the potential due to hyperbolic quadrupole
- Compute ion trajectories within the trap.
- Account for AC variation of potential → ion trajectories
- Design size of the trap, AC frequency, etc... for ion confinement with weights and charge.

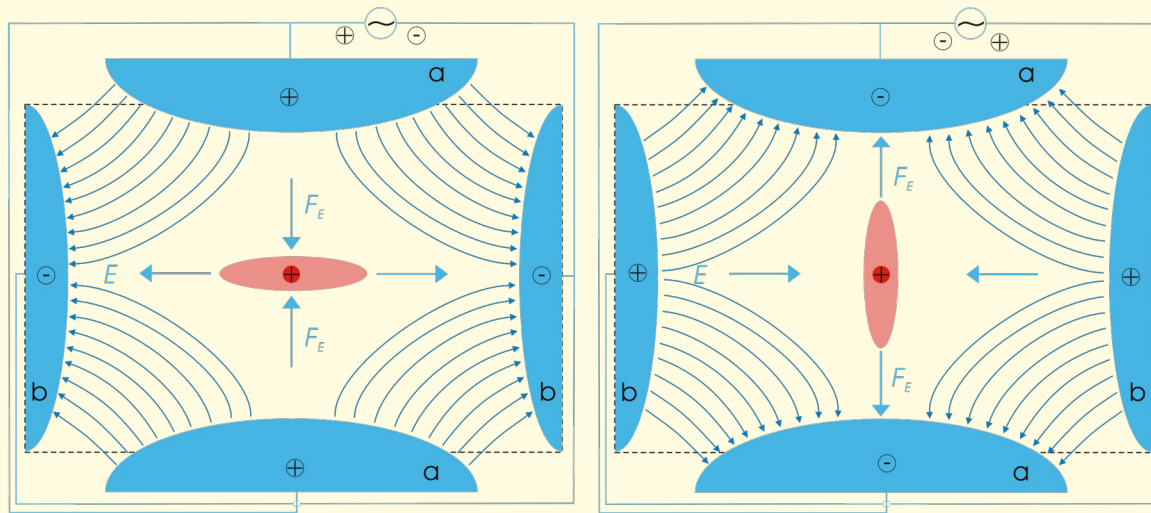


■ Methodology (cont):

With only one ion in the trap (initially static), we just have to lower the voltage to ensure that it does not escape. For $V_0=0$ it just remains static.

Set **two (or more) equal ions** in the trap, at random positions: repulsion will make them escape if V_0 is too low.

- Compute the two ion trajectories **accounting for the repulsion force** between them. Extend to more than two ions, if possible.
- Determine the range for V_0 and AC frequency that confine the ions.



- Week 1: Implementation in 2D: validation with capacitor
- Week 2: Implementation in 3D: validation with capacitor
- Week 3: Paul Trap: Hiperbolic electrodes, static potential
- Week 4: Trajectory of (i) one ion and (ii) two ions with repulsion
- Week 5: AC potential and trajectories of multiple ions
- Week 6: Design the trap for optimum ion confinement
 - Size of the trap, voltage V_0 , AC frequency, etc...