Applied Complex Analysis (2021)

1 Lecture 19: Classical orthogonal polynomials

We will also investigate the properties of *classical* OPs. A good reference is Digital Library of Mathematical Functions, Chapter 18.

This lecture we discuss

- 1. Hermite, Laguerre, and Jacobi polynomials
- 2. Legendre, Chebyshev, and ultraspherical polynomials
- 3. Explicit construction for Chebyshev polynomials

1.1 Definition of classical orthogonal polynomials

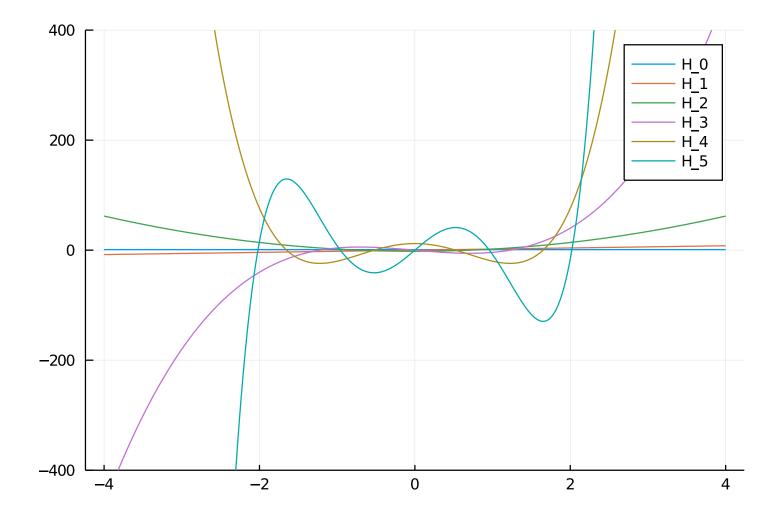
Classical orthogonal polynomials are orthogonal with respect to the following three weights:

Name	(a, b)	w(x)	Notation	$ k_n $
Hermite	$(-\infty,\infty)$	e^{-x^2}	$H_n(x)$	2^n
Laguerre				Table 18.3.1
Jacobi	$\left (-1,1) \right $	$\left (1-x)^{\alpha} (1+x)^{\beta} \right $	$P_n^{(\alpha,\beta)}(x)$	Table 18.3.1

Note out of convention the parameters for Jacobi polynomials are right-to-left order.

We can actually construct these polynomials in Julia, first consider Hermite:

```
using ApproxFun, Plots, LinearAlgebra, ComplexPhasePortrait
H_{-}0 = Fun(Hermite(), [1])
H_{-1} = Fun(Hermite(), [0,1])
H_{-}2 = Fun(Hermite(), [0,0,1])
H_{-3} = Fun(Hermite(), [0,0,0,1])
H_4 = Fun(Hermite(), [0,0,0,0,1])
H_{-}5 = Fun(Hermite(), [0,0,0,0,0,1])
xx = -4:0.01:4
plot(xx, H_0.(xx); label="H_0", ylims=(-400,400))
plot!(xx, H_{-1}(xx); label="H_{-1}", ylims=(-400,400))
plot!(xx, H_2.(xx); label="H_2", ylims=(-400,400))
plot!(xx, H_3.(xx); label="H_3", ylims=(-400,400))
plot!(xx, H_4.(xx); label="H_4", ylims=(-400,400))
plot!(xx, H_5.(xx); label="H_5")
```

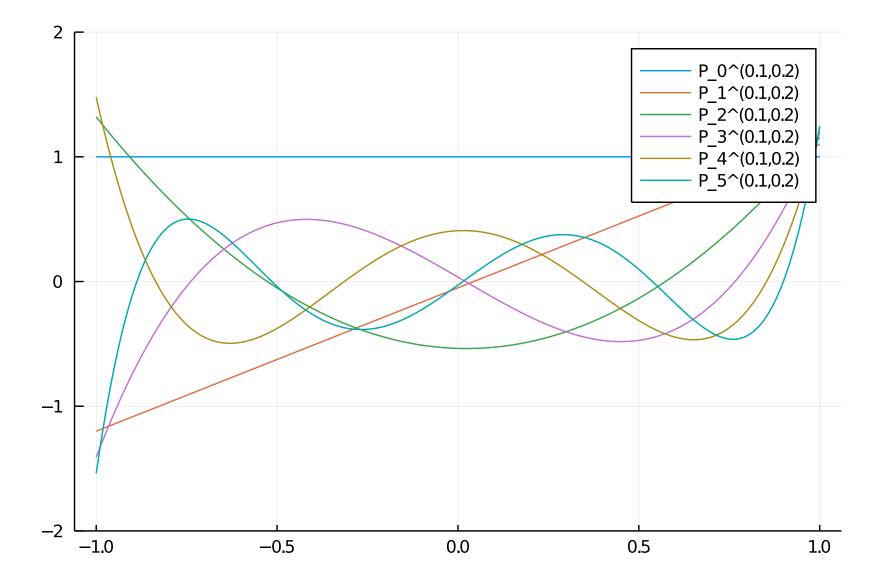


We verify their orthogonality:

```
w = Fun(GaussWeight(), [1.0])
@show sum(H_2*H_5*w) # means integrate
@show sum(H_5*H_5*w);
sum(H_2 * H_5 * w) = 0.0
sum(H_5 * H_5 * w) = 6806.222787477181
```

Now Jacobi:

```
\alpha, \beta = 0.1, 0.2
P_{-0} = Fun(Jacobi(\beta, \alpha), [1])
P_{-1} = Fun(Jacobi(\beta, \alpha), [0, 1])
P_2 = Fun(Jacobi(\beta, \alpha), [0, 0, 1])
P_3 = Fun(Jacobi(\beta, \alpha), [0,0,0,1])
P_{4} = Fun(Jacobi(\beta, \alpha), [0,0,0,0,1])
P_{-5} = Fun(Jacobi(\beta, \alpha), [0,0,0,0,0,1])
xx = -1:0.01:1
plot(xx, P_0.(xx); label="P_0^($\alpha,$\beta)", ylims=(-2,2))
plot!(xx, P_{-1}.(xx); label="P_{-1}^($\alpha,$\beta)")
plot!(xx, P_2.(xx); label="P_2^{(\alpha,\beta)}")
plot!(xx, P_3.(xx); label="P_3^(\alpha,$\beta)")
plot!(xx, P_4.(xx); label="P_4^($\alpha,$\beta)")
plot!(xx, P_5.(xx); label="P_5^(\alpha, $\beta)")
```



1.2 Legendre, Chebyshev, and ultraspherical polynomials

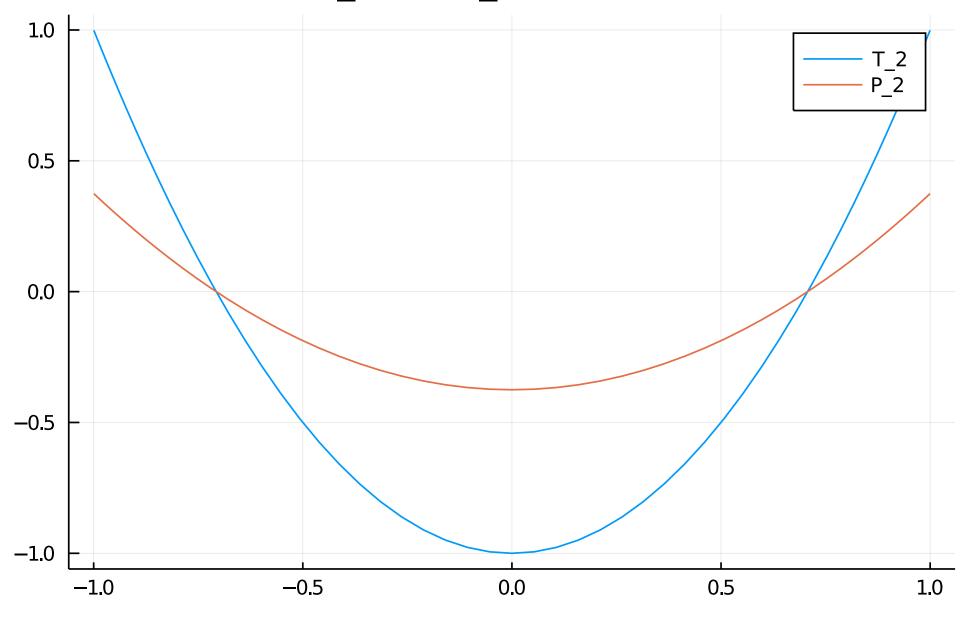
There are special families of Jacobi weights with their own name.

Name	Jacobi parameters	w(x)	Notation	$\mid k_n \mid$
Jacobi	α, β	$(1-x)^{\alpha}(1+x)^{\beta}$	$P_n^{(\alpha,\beta)}(x)$	Table 18.3.1
Legendre	0,0	1	$P_n(x)$	$2^{n}(1/2)_{n}/n!$
Chebyshev (1st)	$-\frac{1}{2}, -\frac{1}{2}$	$\frac{1}{\sqrt{1-x^2}}$	$T_n(x)$	$1(n=0), 2^{n-1}(n \neq 0)$
Chebyshev (2nd)	$\left \frac{1}{2}, \frac{1}{2} \right $	$\sqrt[3]{1-x^2}$	$U_n(x)$	2^n
Ultraspherical	$\lambda - \frac{1}{2}, \lambda - \frac{1}{2}$	$\left (1 - x^2)^{\lambda - 1/2}, \lambda \neq 0 \right $	$C_n^{(\lambda)}(x)$	$2^n(\lambda)_n/n!$

Note that other than Legendre, these polynomials have a different normalization than $P_n^{(\alpha,\beta)}$:

```
T_2 = Fun(Chebyshev(), [0.0,0,1])
P_2 = Fun(Jacobi(-1/2,-1/2), [0.0,0,1])
plot(T_2; label="T_2", title="T_2 is C*P_2 for some C")
plot!(P_2; label="P_2")
```

T_2 is C*P_2 for some C



But because they are orthogonal w.r.t. the same weight, they must be a constant multiple of each-other, as discussed last lecture.

1.2.1 Explicit construction of Chebyshev polynomials (first kind and second kind)

Chebyshev polynomials are pretty much the only OPs with *simple* closed form expressions. **Proposition** (Chebyshev first kind formula) $T_n(x) = \cos n \cos x$ or in other words,

$$T_n(\cos\theta) = \cos n\theta$$

Proof We first show that they are orthogonal w.r.t. $1/\sqrt{1-x^2}$. Too easy: do $x = \cos \theta$, $dx = -\sin \theta$ to get (for $n \neq m$)

$$\int_{-1}^{1} \frac{\cos n a \cos x \cos m a \cos x dx}{\sqrt{1 - x^2}} = -\int_{\pi}^{0} \cos n\theta \cos m\theta d\theta$$
$$= \int_{0}^{\pi} \frac{e^{i(-n-m)\theta} + e^{i(n-m)\theta} + e^{i(m-n)\theta} + e^{i(n+m)\theta}}{4} d\theta = 0$$

We then need to show it has the right highest order term k_n . Note that $k_0=k_1=1$. Using $z=\mathrm{e}^{\mathrm{i}\theta}$ we see that $\cos n\theta$ has a simple recurrence for $n=2,3,\ldots$:

$$\cos n\theta = \frac{z^n + z^{-n}}{2} = 2\frac{z + z^{-1}}{2} \frac{z^{n-1} + z^{1-n}}{2} - \frac{z^{n-2} + z^{2-n}}{2} = 2\cos\theta\cos(n-1)\theta - \cos(n-2)\theta$$

thus

$$\cos n a \cos x = 2x \cos(n-1)a \cos x - \cos(n-2)a \cos x$$

It follows that

$$k_n = 2k_{n-1} = 2^{n-1}k_1 = 2^{n-1}$$

By uniqueness we have $T_n(x) = \cos n \cos x$.

■ Proposition (Chebyshev second kind formula) $U_n(x) = \frac{\sin(n+1)a\cos x}{\sin a\cos x}$ or in other words,

$$U_n(\cos\theta) = \frac{\sin(n+1)\theta}{\sin\theta}$$

Example For the case of Chebyshev polynomials, we have

$$J = \begin{pmatrix} 0 & 1 & & & \\ \frac{1}{2} & 0 & \frac{1}{2} & & & \\ & \frac{1}{2} & 0 & \frac{1}{2} & & & \\ & & \frac{1}{2} & 0 & \ddots & \\ & & \ddots & \ddots & \end{pmatrix}$$

Therefore, the Chebyshev coefficients of xf(x) are given by

$$J^{ op}\mathbf{f} = egin{pmatrix} 0 & rac{1}{2} & & & & \\ 1 & 0 & rac{1}{2} & & & \\ & rac{1}{2} & 0 & rac{1}{2} & & \\ & & rac{1}{2} & 0 & \cdots \end{pmatrix} egin{pmatrix} f_0 \\ f_1 \\ f_2 \\ f_3 \\ dots \end{pmatrix}$$

1.2.2 Demonstration

In the case where f is a degree n-1 polynomial, we can represent J^{\top} as an $n+1\times n$ matrix (this makes sense as xf(x) is one more degree than f):

```
f = Fun(exp, Chebyshev())
n = ncoefficients(f) # number of coefficients
@show n
J = zeros(n,n+1)
```

```
J[1,2] = 1
for k=2:n
    J[k,k-1] = J[k,k+1] = 1/2
end
J^{\top}
n = 14
15\times14 LinearAlgebra.Adjoint(*@{Float64,Array{Float64,2}}:
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```

```
cfs = J'*f.coefficients # coefficients of x*f
xf = Fun(Chebyshev(), cfs)
xf(0.1) - 0.1*f(0.1)
4.163336342344337e-17
We can construct T_0(x), \ldots, T_{n-1}(x) via
                   T_0(x) = 1
                   T_1(x) = xT_0(x)
                 T_{k+1}(x) = 2xT_k(x) - T_{k-1}(x), \qquad 1 \le k \le n-2
Believe it or not, this is much faster than using \cos k a \cos x:
function recurrence_Chebyshev(n,x)
    T = zeros(n)
    T[1] = 1.0
    T[2] = x*T[1]
    for k = 2:n-1
         T[k+1] = 2x*T[k] - T[k-1]
    end
end
trig_Chebyshev(n,x) = [cos(k*acos(x)) for k=0:n-1]
```

```
n = 10
recurrence_Chebyshev(n, 0.1) - trig_Chebyshev(n,0.1) |>norm
1.1102230246251565e-16
n = 10000
@time recurrence_Chebyshev(n, 0.1)
@time trig_Chebyshev(n,0.1);
0.000062 seconds (2 allocations: 78.203 KiB)
    0.000436 seconds (2 allocations: 78.203 KiB)
```

We can also demonstrate Clenshaw's algorithm for evaluating polynomials. To evaluate an expansion in Chebyshev polynomials,

$$\sum_{k=0}^{n-1} f_k T_k(x)$$

we want to solve the system

$$\underbrace{\begin{pmatrix}
1 & -x & \frac{1}{2} \\
1 & -x & \frac{1}{2} \\
& \frac{1}{2} & -x & \cdots \\
& & \frac{1}{2} & \cdots & \frac{1}{2} \\
& & & \ddots & -x \\
& & & \frac{1}{2}
\end{pmatrix}}_{L_{\pi}^{\top}}$$

$$\begin{pmatrix}
\gamma_{0} \\
\vdots \\
\gamma_{n-1}
\end{pmatrix}$$

via

$$\gamma_{n-1} = 2f_{n-1}$$

$$\gamma_{n-2} = 2f_{n-2} + 2x\gamma_{n-1}$$

$$\gamma_{n-3} = 2f_{n-3} + 2x\gamma_{n-2} - \gamma_{n-1}$$

$$\vdots$$

$$\gamma_1 = f_1 + x\gamma_2 - \frac{1}{2}\gamma_3$$

$$\gamma_0 = f_0 + x\gamma_1 - \frac{1}{2}\gamma_2$$

then $f(x)/k_0 = f(x) = \gamma_0$.

```
function clenshaw_Chebyshev(f,x)
    n = length(f)
    \gamma = zeros(n)
    \gamma[n] = 2f[n]
    \gamma[n-1] = 2f[n-1] +2x*f[n]
    for k = n-2:-1:1
         \gamma[k] = 2f[k] + 2x*\gamma[k+1] - \gamma[k+2]
    end
    \gamma[2] = f[2] + x*\gamma[3] - \gamma[4]/2
    \gamma[1] = f[1] + x*\gamma[2] - \gamma[3]/2
    \gamma [1]
end
f = Fun(exp, Chebyshev())
clenshaw_Chebyshev(f.coefficients, 0.1) - exp(0.1)
-1.3322676295501878e-15
```

With some high performance computing tweaks, this can be made more accurate. This is the algorithm used for evaluating functions in ApproxFun:

```
f(0.1) - exp(0.1)
0.0
```

1.3 Approximation with Chebyshev polynomials

Previously, we used the formula, derived via trigonometric manipulations,

$$T_1(x) = xT_0(x),$$
 $T_{n+1}(x) = 2xT_n(x) - T_{n-1}(x)$

Rearranging, this becomes

$$xT_0(x) = T_1(x),$$
 $xT_n(x) = \frac{T_{n-1}(x)}{2} + \frac{T_{n+1}(x)}{2}$

This tells us that we have the three-term recurrence with $a_n = 0$, $b_0 = 1$, $c_n = b_n = \frac{1}{2}$ for n > 0. This can be extended to function approximation. Provided the sum converges absolutely and uniformly in x, we can write

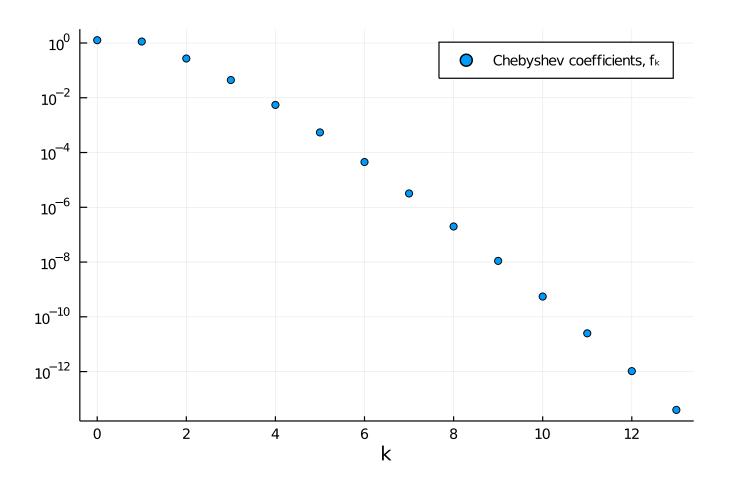
$$f(x) = \sum_{k=0}^{\infty} f_k T_k(x).$$

In practice, we can approximate smooth functions by a finite truncation:

$$f(x) \approx \sum_{k=0}^{n-1} f_k T_k(x)$$

Here we see that e^x can be approximated by a Chebyshev approximation using 14 coefficients and is accurate to 16 digits:

```
f = Fun(x -> exp(x), Chebyshev())
scatter(0:ncoefficients(f)-1,abs.(f.coefficients);yscale=:log10,label="C
coefficients, f_k",xlabel="k")
```



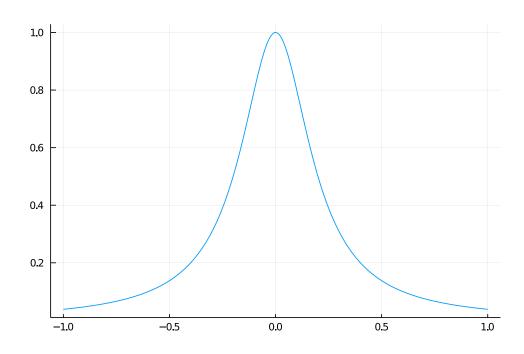
```
Oshow ncoefficients(f)
Oshow f(0.1) # equivalent to f.coefficients'*[cos(k*acos(x)) for
k=0:ncoefficients(f)-1]
Oshow exp(0.1);
ncoefficients(f) = 14
f(0.1) = 1.1051709180756477
exp(0.1) = 1.1051709180756477
```

The accuracy of this approximation is typically dictated by the smoothness of f: the more times we can differentiate, the faster it converges. For analytic functions, it's dictated by the domain of analyticity, just like Laurent/Fourier series. In the case above, e^x is entire hence we get faster than exponential convergence.

Chebyshev expansions work even when Taylor series do not. For example, the following function has poles at $\pm \frac{\mathrm{i}}{5}$, which means the radius of convergence for the Taylor series is $|x| < \frac{1}{5}$, but Chebyshev polynomials continue to work on [-1,1]:

```
f = Fun(x -> 1/(25x^2 + 1), Chebyshev())
@show ncoefficients(f)
plot(f)
```

ncoefficients(f) = 189



This can be explained for Chebyshev expansion by noting that it is the cosine expansion / Fourier expansion of an even function:

$$f(x) = \sum_{k=0}^{\infty} f_k T_k(x) \Leftrightarrow f(\cos \theta) = \sum_{k=0}^{\infty} f_k \cos k\theta$$

1.3.1 Exponential decay of Fourier coefficients of periodic, analytic functions revisited

Before we get to the decay of Chebyshev coefficients, we revisit the proof of the exponential decay of *Fourier* coefficients in Lecture 6. Suppose $f(\theta)$ is 2π -periodic and analytic on $\theta \in [-\pi, \pi)$, then

$$f(\theta) = \sum_{k=-\infty}^{\infty} \hat{f}_k e^{ik\theta}$$

where

$$\hat{f}_k = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(\theta) e^{-ik\theta} d\theta.$$

Recall in Lecture 6 we set $z={\rm e}^{{\rm i}\theta}$ in which case the Fourier series of f becomes a Laurent series of a function g(z):

$$f(\theta) = \sum_{k=-\infty}^{\infty} \hat{f}_k e^{ik\theta} = \sum_{k=-\infty}^{\infty} g_k z^k =: g(z),$$

with $g_k = \hat{f}_k$. We proved that if g(z) is analytic on the closed annulus $A_{r,R} = \{z : r \le |z| \le R\}$, 0 < r < 1, R > 1 then for all $k \in \mathbb{Z}$, $|g_k| \le M \min\left\{\frac{1}{R^k}, \frac{1}{r^k}\right\}$ where $M = \sup_{z \in A_{r,R}} |g(z)|$. This result implies the exponential decay of the Fourier coefficients of f.

An annulus in the z-plane corresponds to a strip of width 2π in the (complex) θ -plane under the transformation $z = e^{i\theta}$, $\Re \theta \in [-\pi, \pi)$:

$$z \in A_{r,R} = \{z : r \le |z| \le R\} \qquad \Longrightarrow_{z=e^{i\theta}}$$

$$\theta \in S_{r,R} = \{\theta : -\pi \le \Re \theta < \pi, -\log(R) \le \Im \theta \le \log(1/r)\}.$$

Suppose $f(\theta)$ is real-valued on $[-\pi,\pi)$, then $\overline{f(\theta)}=f(\overline{\theta})$. Hence if the closest singularity to the real θ -axis is at $\theta=\theta_x+\mathrm{i}\theta_y$, with $\theta_x\in[-\pi,\pi)$ and $\theta_y>0$, then f also has a singularity at $\theta_x-\mathrm{i}\theta_y$. Thus f is analytic in the strip

$$S_{r,R} = \theta \in \{\theta : -\pi \le \Re \theta < \pi, -\log(R) \le \Im \theta \le \log(1/r)\}\$$

with

$$\frac{1}{r} = R < e^{\theta_y}$$

and the Fourier coefficients are bounded by

$$|f_k| = |g_k| \le M \min\left\{\frac{1}{R^k}, \frac{1}{r^k}\right\} = Mr^{|k|} = MR^{-|k|}, \quad k \in \mathbb{Z},$$

where $M=\sup_{z\in A_{r,R}}|g(z)|=\sup_{\theta\in S_{r,R}}|f(\theta)|$. The larger the strip of analyticity, the larger we can make R and the faster the Fourier coefficients of f decay as $|k|\to\infty$ (hence the faster the truncated Fourier expansion $\sum_{k=-n}^n \hat{f}_k \mathrm{e}^{ik\theta}$ of f converges to f as $n\to\infty$).

Example (see also Lecture 6) The function

$$f(\theta) = \frac{1}{2 - \cos \theta},$$

has poles at $\theta = \pm i \log(2 + \sqrt{3})$; it is analytic in the strip $S_{r,R}$ with $R = 1/r < 2 + \sqrt{3}$ and the maximum of $|f(\theta)|$ on $S_{r,R}$ is

$$M = \frac{2}{4 - R^{-1} + R},$$

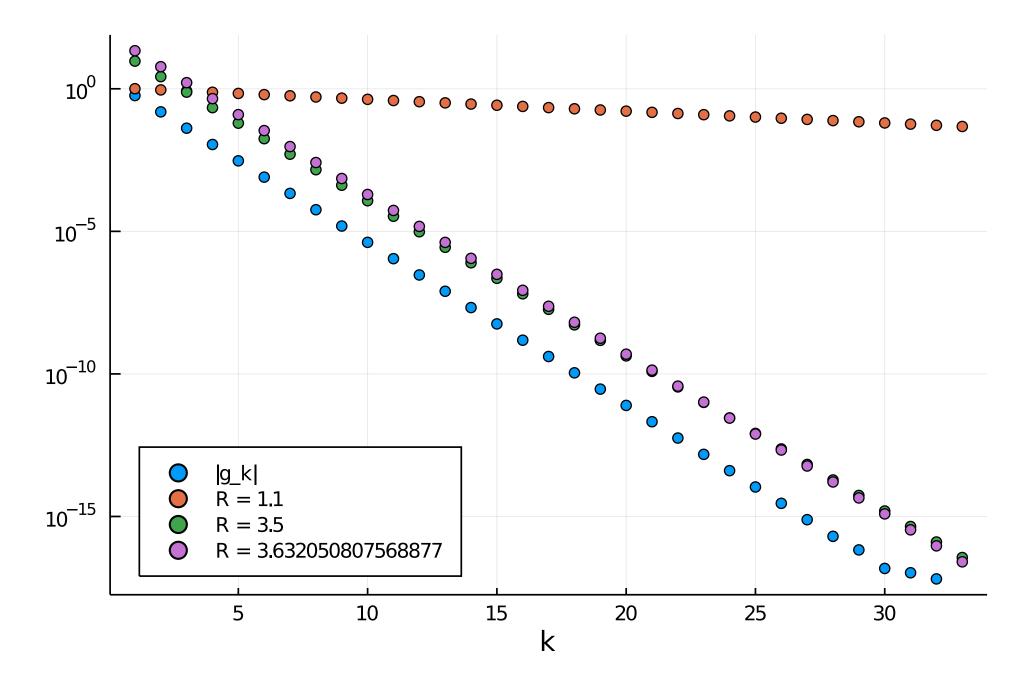
hence

$$|f_k| = |g_k| \le \frac{2}{4 - R - R^{-1}} R^{-|k|}, \qquad k \in \mathbb{Z},$$

for all $R < 2 + \sqrt{3}$.

```
 g = Fun(\theta \rightarrow 1/(2-\cos(\theta)), Laurent(-\pi ... \pi)) 
 g_{-} + = g.coefficients[1:2:end] 
 scatter(abs.(g_{-}+); yscale=:log10, label="|g_{-}k|", 
 legend=:bottomleft,xlabel="k") 
 R = 1.1 
 scatter!(2/(4-R-inv(R))*R.^(-(0:length(g_{-}+))), label = "R = $R") 
 R = 3.5 
 scatter!(2/(4-R-inv(R))*R.^(-(0:length(g_{-}+))), label = "R = $R") 
 R = 2+sqrt(3)-0.1
```

 $scatter!(2/(4-R-inv(R))*R.^(-(0:length(g_+))), label = "R = $R")$



1.3.2 Exponential decay of Chebyshev coefficients of analytic functions

Suppose f(x) is analytic on [-1,1], then

$$f(x) = \sum_{k=0}^{\infty} f_k T_k(x)$$

$$= \sum_{k=0}^{\infty} f_k \cos k\theta \qquad (x = \cos \theta)$$

$$= \sum_{k=0}^{\infty} \frac{f_k}{2} (z^k + z^{-k}) \qquad (z = e^{i\theta})$$

$$=: \sum_{k=-\infty}^{\infty} g_k z^k =: g(z) \qquad (g_0 = f_0, g_k = g_{-k} = f_k/2, k \ge 0)$$

Now we can use the bound on the Laurent coefficients of g(z) to bound the Chebyshev coefficients of f(x). First we need to establish what is the image in the (complex) x-plane of an annulus in the z-plane under the transformation $2x = z + z^{-1}$, which is known as the Joukowsky map.

Let $\rho > 1$ and

$$A_{1,\rho} = \{z : 1 \le |z| \le \rho\}, \qquad A_{\rho^{-1},1} = \{z : \rho^{-1} \le |z| \le 1\}.$$

The Joukowsky transformation maps $A_{1,\rho}$ and $A_{\rho^{-1},1}$ to the following ellipse (known as a Bernstein ellipse) in the x-plane:

$$E_{\rho} = \left\{ x : \frac{(\Re x)^2}{\alpha^2} + \frac{(\Im x)^2}{\beta^2} \le 1, \alpha = \frac{1}{2} \left(\rho + \rho^{-1} \right), \beta = \frac{1}{2} \left(\rho - \rho^{-1} \right) \right\}.$$

We conclude that if f(x) is analytic on E_{ρ} (or g(z) is analytic on $A_{\rho^{-1},\rho}$) and $M=\sup_{x\in E_{\rho}}|f(x)|=\sup_{z\in A_{\rho^{-1},\rho}}|g(z)|$, then

$$|f_k| = 2|g_k| \le 2M\rho^{-k}, \qquad k \ge 1.$$

The larger the Bernstein ellipse on which f(x) is analytic, the faster the decay of the Chebyshev coefficients as $k \to \infty$ (and hence the faster the convergence of the Chebyshev expansion of f). A truncated Chebyshev expansion with n terms of a function f(x) that is analytic on E_{ρ} converges at essentially the same exponential rate as the bound on the Chebyshev coefficients as $n \to \infty$:

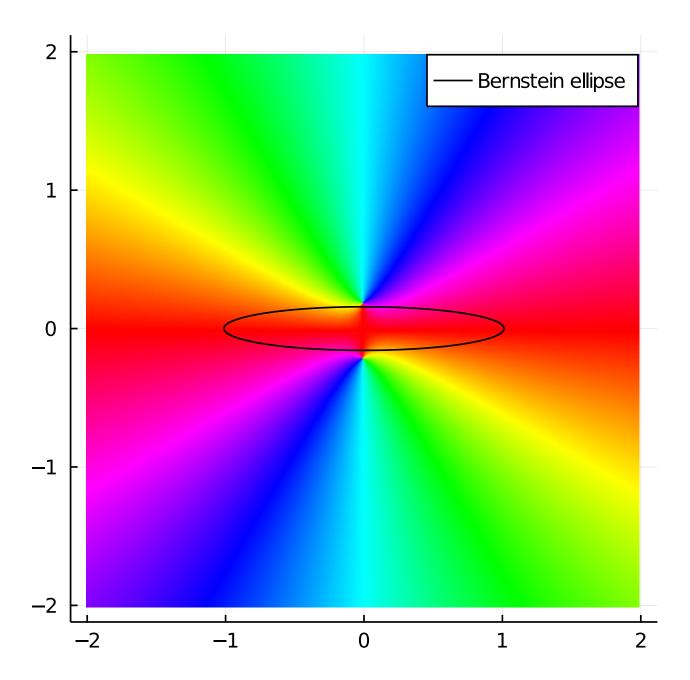
$$\left| f(x) - \sum_{k=0}^{n-1} f_k T_k(x) \right| = \left| \sum_{k=n}^{\infty} f_k T_k(x) \right| \le \sum_{k=n}^{\infty} |f_k| \le 2M \sum_{k=n}^{\infty} \rho^{-k} = 2M \frac{\rho^{-n}}{1 - \rho^{-1}}$$

Example In the case of $f(x) = \frac{1}{25x^2+1}$, setting $2x = z + z^{-1}$, we find that

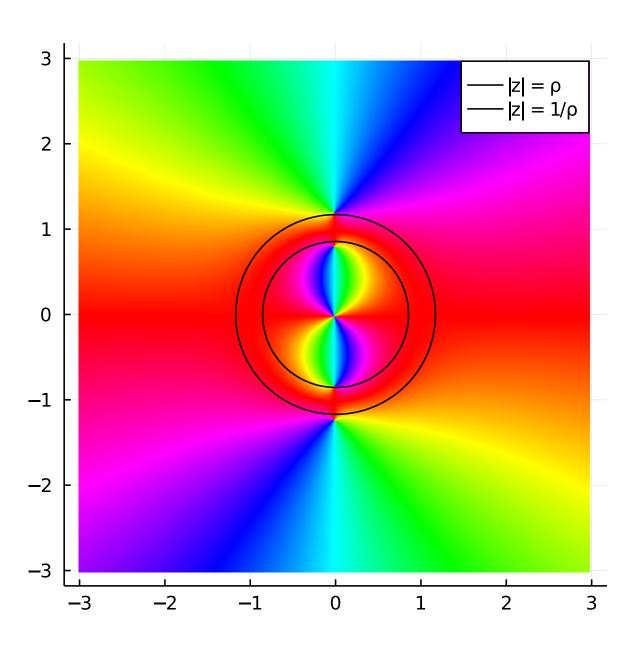
$$f(x) = f\left(\frac{z+z^{-1}}{2}\right) = g(z) = \frac{4z^2}{25 + 54z^2 + 25z^4}.$$

In the complex x-plane, f(x) has poles at $\pm i/5$ and is analytic on E_{ρ} with $\beta=(\rho-\rho^{-1})/2<1/5$, hence $\rho<\frac{1+\sqrt{26}}{5}$. In the z-plane, g(z) has poles at $\pm i\frac{1\pm\sqrt{26}}{5}\approx \pm 0.8198040i, \pm 1.2198i$ and is analytic on the annulus $\rho^{-1}\leq |z|\leq \rho$.

```
\rho = (1 + \text{sqrt}(26))/5-0.05;
\alpha = (\rho + 1/\rho)/2
\beta = (\rho - 1/\rho)/2
\theta = -\pi:0.01:\pi
f = x \rightarrow 1/(25x^2 + 1)
phaseplot(-2..2, -2..2, z \rightarrow f(z))
plot!(\alpha*\cos.(\theta), \beta*\sin.(\theta); linecolor="black", label="Bernstein ellipse")
```



```
phaseplot(-3..3, -3..3, z -> f((z+1/z)/2)) plot!(\rho*\cos.(\theta),\rho*\sin.(\theta),linecolor="black",label="|z| = \rho") plot!(\cos.(\theta)/\rho,\sin.(\theta)/\rho,linecolor="black",label="|z| = 1/\rho")
```



For $\beta = (\rho - \rho^{-1})/2 < 1/5$, we have

$$M = \sup_{x \in E_{\rho}} |f(x)| = \frac{1}{1 - 25\beta^2}$$

hence for $k \geq 1$,

$$|f_k| \le \frac{2}{1 - 25\beta^2} \rho^{-k}, \qquad 1 < \rho < \frac{1 + \sqrt{26}}{5}.$$

Therefore we predict a rate of decay of about 1.2198^{-k} :

```
bound(\beta,k) = 2/(1-25*\beta^2)*(\beta + sqrt(\beta^2 + 1))^(-k)
f = Fun(x \rightarrow 1/(25x^2 + 1), Chebyshev())
scatter(abs.(f.coefficients) .+ eps(); yscale=:log10,
label="Chebyshev coefficients")
plot!(1:ncoefficients(f), bound.(1/5-0.001,1:ncoefficients(f));
label="upper bound")
# Also calculate the error of truncated Chebyshev expansions
with n terms for n = 1, 2, \ldots
xx = -1:0.001:1 \# a fine grid on which to evaluate
Errv = [(maximum(abs.(f.(xx)-Fun(f,Chebyshev(),n).(xx)))) for n =
1:ncoefficients(f)]
scatter!(1:ncoefficients(f),Errv.+eps();label="truncated expansion
error")
```

```
plot!( 1.2198.^(-(0:ncoefficients(f))); label="\rho^(-k)")
```

