

# Applied Complex Analysis (2021)

## 1 Lecture 12: Cauchy transforms and Plemelj's theorem

This lecture concerns the properties of the Cauchy transform:

**Definition (Cauchy transform)** For a contour  $\gamma$  and  $f : \gamma \rightarrow \mathbb{C}$  define the *Cauchy transform* as

$$\mathcal{C}_\gamma f(z) := \frac{1}{2\pi i} \int_\gamma \frac{f(\zeta)}{\zeta - z} d\zeta$$

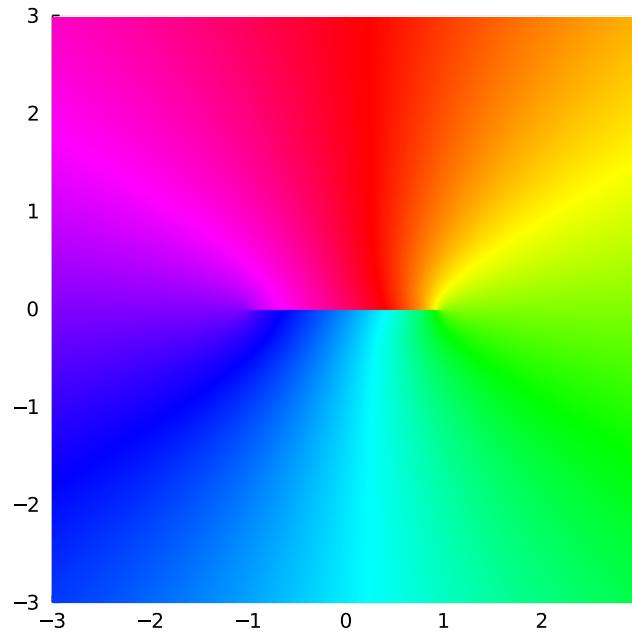
Unlike in Cauchy's integral formula,  $f$  need not be analytic and  $\gamma$  need not be closed.

We focus on the case of an interval  $[a, b]$ :

$$\mathcal{C}_{[a,b]} f(z) := \frac{1}{2\pi i} \int_a^b \frac{f(x)}{x - z} dx$$

Here is a phase portrait of the Cauchy transform of a simple function:

```
using ApproxFun, SingularIntegralEquations, ComplexPhasePortrait,  
Plots  
x = Fun(-1 .. 1)  
f = exp(x)*sqrt(1-x^2)  
phaseplot(-3..3, -3..3, z -> cauchy(f,z))
```



What's evident here is that it has a jump on the contour. It turns out that the Cauchy transform has a very simple subtractive jump. Here we denote

$$\mathcal{C}_{[a,b]}^+ f(x) = \lim_{\epsilon \rightarrow 0^+} \mathcal{C}_{[a,b]} f(x + i\epsilon), \quad \mathcal{C}_{[a,b]}^- f(x) = \lim_{\epsilon \rightarrow 0^+} \mathcal{C}_{[a,b]} f(x - i\epsilon)$$

that is the limit from above and below. For more complicated contours these would denote the limit from the left/right or in the case of a simple closed contour, interior/exterior.

**Theorem (Plemelj on the interval I)** Suppose  $(b-x)^\alpha(x-a)^\beta f(x)$  is differentiable on  $[a, b]$ , for  $\alpha, \beta < 1$ . Then the Cauchy transform has the following properties:

1. *Analyticity*:  $\mathcal{C}_{[a,b]}f(z)$  is analytic in  $\bar{\mathbb{C}} \setminus [a, b]$

2. *Decay*:  $\mathcal{C}_{[a,b]}f(\infty) = 0$

3. *Jump*: It has the subtractive jump:

$$\mathcal{C}_{[a,b]}^+ f(x) - \mathcal{C}_{[a,b]}^- f(x) = f(x) \quad \text{for} \quad a < x < b$$

4. *Regularity*:  $\mathcal{C}_{[a,b]}f(z)$  has weaker than pole singularities at  $a$  and  $b$

**Sketch of Proof** We show the proof for  $[-1, 1]$ .

1. From the dominated convergence theorem, we know that  $\mathcal{C}f(z)$  is complex-differentiable off  $[-1, 1]$ :

$$\frac{d}{dz}\mathcal{C}f(z) = \frac{1}{2\pi i} \int_{-1}^1 \frac{d}{dz} \frac{f(x)}{x-z} dx = \frac{1}{2\pi i} \int_{-1}^1 \frac{f(x)}{(x-z)^2} dx$$

We know it is analytic at  $\infty$  because

$$\mathcal{C}f(z^{-1}) = z \frac{1}{2\pi i} \int_{-1}^1 \frac{f(x)}{zx-1} dx$$

is differentiable at zero.

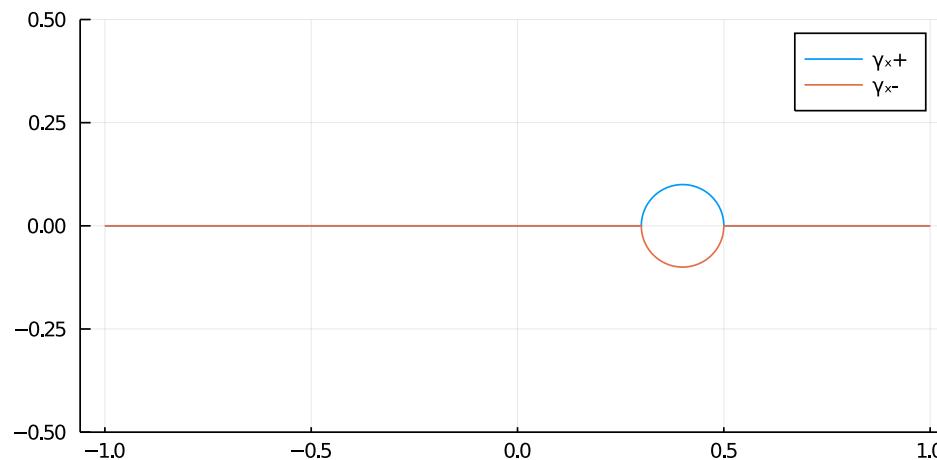
2.

$$\mathcal{C}f(\infty) = 0$$

follows from uniform convergence of  $\frac{1}{z-x}$  to zero as  $z \rightarrow \infty$ .

3. For the constant function, which is analytic, this follows by considering a contour  $\gamma_x^+$  perturbed above  $x$  and  $\gamma_x^-$  perturbed below  $x$ , see the plot below:

```
x = 0.4; r = 0.1
tt = range(pi, 0.; length=100)
plot([-1.; x .+ r*cos.(tt); 1.], [0.; r*sin.(tt);
0.]; ylims=(-0.5,0.5),label="γ_x+", ratio=1.0)
plot!([-1.; x .+ r*cos.(tt); 1.], [0.; -r*sin.(tt);
0.]; ylims=(-0.5,0.5),label="γ_x-")
```



Therefore, by the Cauchy integral formula we have

$$\mathcal{C}^+ 1(x) - \mathcal{C}^- 1(x) = \frac{1}{2\pi i} \int_{\gamma_x^-} \frac{1}{x-z} dx - \frac{1}{2\pi i} \int_{\gamma_x^+} \frac{1}{x-z} dx = \frac{1}{2\pi i} \oint \frac{1}{x-z} dx = 1.$$

For other functions, we consider, for  $z = x + i\epsilon$ ,

$$\mathcal{C}f(z) = \frac{1}{2\pi i} \int_{-1}^1 \frac{f(t) - f(x)}{t - z} dt + f(x)\mathcal{C}1(z)$$

For  $\epsilon = 0$ , the first integral exists because the singularity at  $t = x$  is removable:

$$\lim_{t \rightarrow x} \frac{f(t) - f(x)}{t - x} = f'(x)$$

We leave it as an exercise (or see [Trogdon & Olver 2015, *Riemann–Hilbert Problems, Their Numerical Solution, and the Computation of Nonlinear Special Functions*, Lemma 2.7]) to show that  $\int_{-1}^1 \frac{f(t) - f(x)}{t - z} dt$  converges to  $\int_{-1}^1 \frac{f(t) - f(x)}{t - x} dt$  as  $z \rightarrow x$ . It follows that

$$\mathcal{C}^\pm f(x) = \frac{1}{2\pi i} \int_{-1}^1 \frac{f(t) - f(x)}{t - x} dt + f(x)\mathcal{C}^\pm 1(x)$$

and in particular

$$\mathcal{C}^+ f(x) - \mathcal{C}^- f(x) = f(x)(\mathcal{C}^+ 1(x) - \mathcal{C}^- 1(x)) = f(x)$$

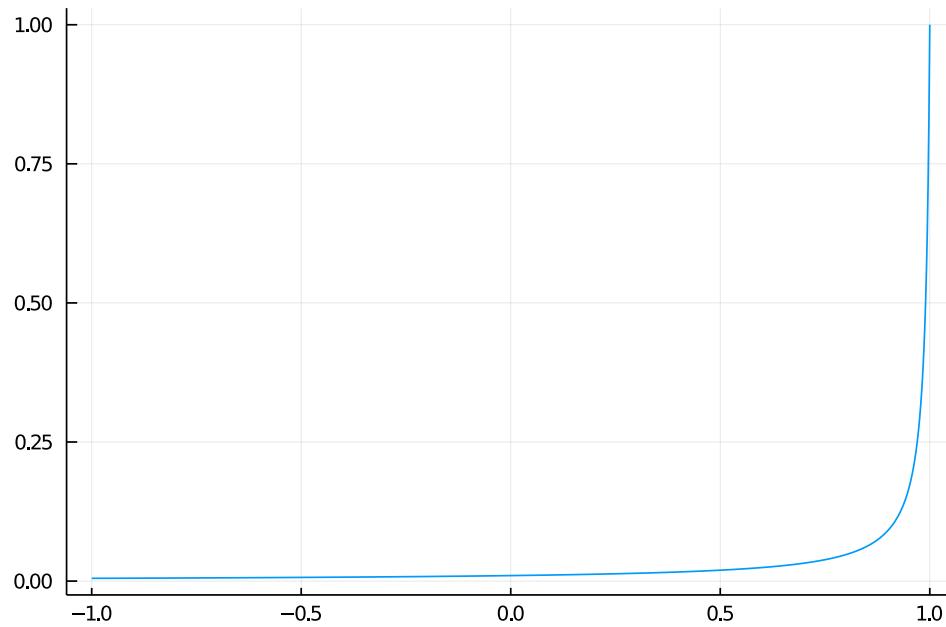
4. We show that it has a weaker than pole singularity at  $+1$ , with  $-1$  following by the same argument. First note that  $f$  is absolutely integrable.

If we assume we approach  $1$  at an angle of  $-\pi + \delta \leq \theta \leq \pi - \delta$ , the uniform convergence of  $(z - 1)\mathcal{C}f(z)$  to zero follows from observing that  $\frac{z-1}{z-t}$  can be made arbitrarily small in a larger and larger interval. This is easiest to see for real  $x > 1$ , where for  $1 \leq x \leq 1 + \epsilon^2$  we have

$$\left| \frac{x-1}{x-t} \right| \leq \epsilon$$

for all  $t \leq 1 + \epsilon^2 - \epsilon$ , or more generously,  $t \leq 1 - \epsilon$ . Here is a plot of  $\frac{x-1}{x-t}$  showing that it is small on an increasing portion of the interval as  $x \rightarrow 1$  from the right:

```
x = 1 + 0.01
tt = range(-1., 1.; length=1000)
plot(tt, abs((x - 1) ./ (x .- tt)); legend=false)
```



Therefore,

$$|(x-1)\mathcal{C}f(x)| \leq \frac{1}{2\pi} \int_{-1}^{1-\epsilon} |f(t)| \left| \frac{x-1}{x-t} \right| dt + \int_{1-\epsilon}^1 |f(t)| dt \leq \epsilon \int_{-1}^1 |f(t)| dt + \int_{1-\epsilon}^1 |f(t)| dt$$

Both terms tends to zero as  $\epsilon \rightarrow 0$ , hence so does  $|(x-1)\mathcal{C}f(x)|$ . To extend this to the interval itself (that is,  $\delta = 0$ ), we use the stronger requirement that  $(1-x)^\alpha(1+x)^\beta f(x)$  is differentiable. For  $\alpha = \beta = 0$ , this follows from the expression in condition (3) and the fact that (found via direct integration)

$$\mathcal{C}1(z) = \frac{\log(z-1) - \log(z+1)}{2\pi i}$$

has only logarithmic singularities, and  $f(x)$  is bounded. ■

*Remark:* The singularities of  $\mathcal{C}_{[a,b]}f(z)$  at  $z = a, b$  are analysed in detail in [Lemma 7.2.2](#), [Ablowitz & Fokas, Complex Variables: Introduction and Applications](#). Roughly speaking, if  $f$  is smooth on  $[a, b]$ , then  $\mathcal{C}_{[a,b]}f(z)$  will have logarithmic singularities at  $z = a, b$  (e.g.,  $\mathcal{C}1(z)$ ). If  $f$  has an algebraic branch point at  $a$  or  $b$  of order  $\alpha$ , then so has  $\mathcal{C}_{[a,b]}f(z)$ . For example, if  $f(x) = g(x)/\sqrt{1 - x^2}$  and  $g$  is smooth on  $[-1, 1]$ , then  $f$  has branch points of order  $-1/2$  at  $\pm 1$  and

$$\mathcal{C}_{[a,b]}f(z) = O\left((1 - z)^{-1/2}\right), \quad z \rightarrow 1,$$

and

$$\mathcal{C}_{[a,b]}f(z) = O\left((1 + z)^{-1/2}\right), \quad z \rightarrow -1.$$

We can use the previous results, combined with Liouville's theorem, to show that the function satisfying (1)-(4) is in fact unique:

**Theorem (Liouville)** If  $f$  is entire and bounded in  $\mathbb{C}$ , then  $f$  must be constant.

**Theorem (Plemelj on the interval II)** Suppose  $\phi(z)$  satisfies the following properties:

1. *Analyticity*:  $\phi(z)$  is analytic in  $\bar{\mathbb{C}} \setminus [a, b]$
2. *Decay*:  $\phi(\infty) = 0$
3. *Jump*: It has the subtractive jump:

$$\phi^+(x) - \phi^-(x) = f(x) \quad \text{for} \quad a < x < b$$

where  $(b - x)^\alpha (x - a)^\beta f(x)$  is differentiable in  $[a, b]$  for  $\alpha, \beta < 1$ .

4. *Regularity*:  $\phi(z)$  has weaker than pole singularities at  $a$  and  $b$

Then  $\phi(z) = \mathcal{C}_{[a,b]} f(z)$ .

## Sketch of Proof Consider

$$A(z) = \phi(z) - \mathcal{C}_{[a,b]} f(z)$$

This is continuous (hence analytic) on  $(a, b)$  as

$$A^+(x) - A^-(x) = \phi^+(x) - \phi^-(x) - \mathcal{C}_{[a,b]}^+ f(x) + \mathcal{C}_{[a,b]}^- f(x) = f(x) - f(x) = 0$$

Also,  $A$  has weaker than pole singularities at  $a$  and  $b$ , hence is analytic there as well: it's entire. Only entire functions that are bounded are constant, since it vanishes at  $\infty$  the constant must be zero.



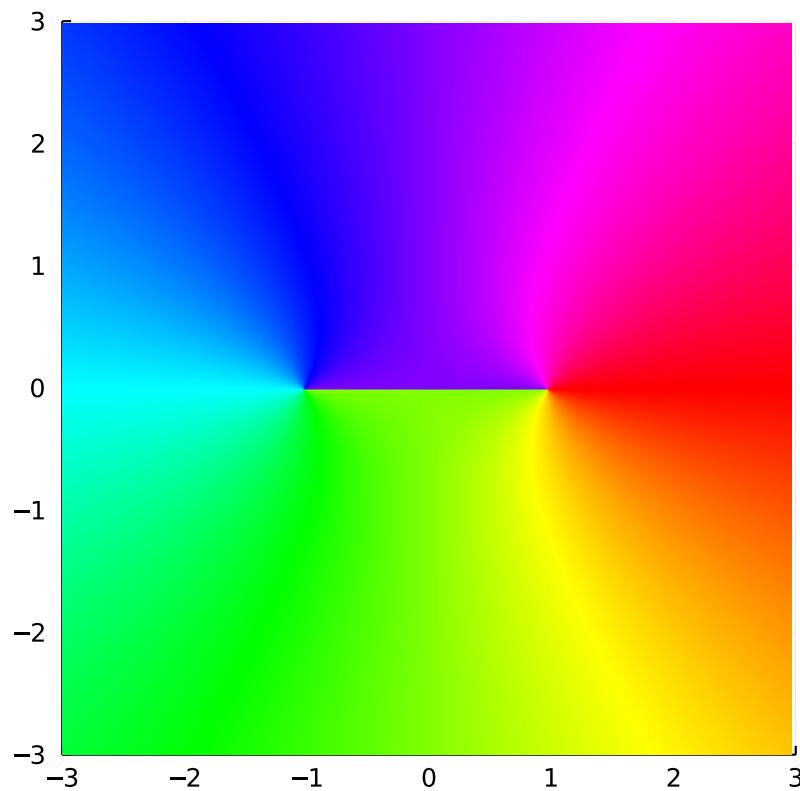
*Example 1* We can use this theorem to prove the following relationships (using  $\diamond$  for the dummy variable):

$$\frac{1}{\sqrt{z-1}\sqrt{z+1}} = -2i\mathcal{C} \left[ \frac{1}{\sqrt{1-\diamond^2}} \right] (z) = -\frac{1}{\pi} \int_{-1}^1 \frac{dx}{\sqrt{1-x^2}(x-z)}$$

(1) follows because the jumps cancel. (2 and 4) are immediate. (3) follows from a simple calculation.

**Demonstration** From the phase plot we see it has a branch cut on  $[-1, 1]$ :

```
 $\kappa = z \rightarrow 1/(\text{sqrt}(z-1)\text{sqrt}(z+1))$ 
phaseplot(-3..3, -3..3,  $\kappa$ )
```



On the branch there is the expected jump:

$$x = 0.1$$

$$\kappa(x + 0.0im) - \kappa(x - 0.0im), \quad -2im/\sqrt{1-x^2}$$

$$(0.0 - 2.010075630518424im, 0.0 - 2.010075630518424im)$$

For  $x < -1$  the branch cut is removable: we have continuity and therefore analyticity:

$$x = -2.3$$

$$\kappa(x + 0.0im) - \kappa(x - 0.0im)$$

$$0.0 - 0.0im$$

*Example 2* Now consider a problem of reducing

$$\phi(z) = \sqrt{z-1}\sqrt{z+1}$$

to its behaviour near its singularities. It has two singularities: it blows up at  $\infty$  and has a branch cut on  $[-1, 1]$

We can subtract out the singularity at infinity first to determine

$$\phi(z) = z + 2i\mathcal{C}[\sqrt{1-\diamond^2}](z)$$

Note this works because, as  $z \rightarrow \infty$ , we have

$$\phi(z) = z(\sqrt{1-1/z}\sqrt{1+1/z}) = z(1+O(1/z))(1+O(1/z)) = z + O(1/z)$$

hence  $\phi(z) - z$  vanishes at  $\infty$ . This is an example of summing over the behaviour at each singularity to recover the function (in this case,  $\phi$  has a singularity along the cut  $[-1, 1]$  and polynomial growth at  $\infty$ ).

Because  $\phi(z) - z$  decays, we can now deploy Plemelj II to determine:

$$\phi(z) - z = \mathcal{C}[\phi_+ - \phi_-](z)$$

where

$$\phi_+(x) - \phi_-(x) = 2i\sqrt{1-x^2}$$

*Example 3* Finally, we have the following (also verifiable using indefinite integration):

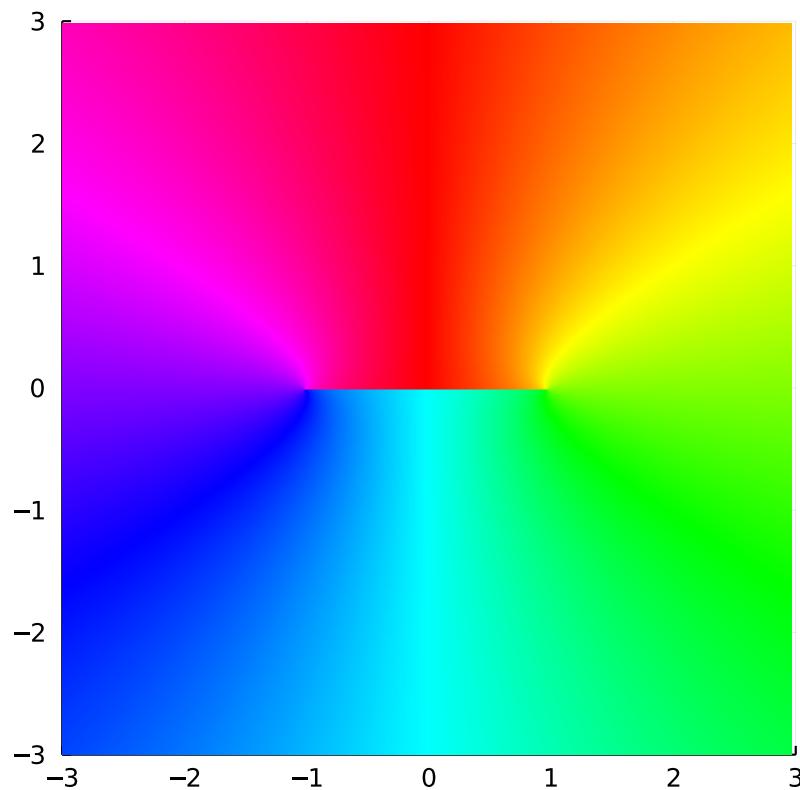
$$\frac{\log(z-1) - \log(z+1)}{2\pi i} = \mathcal{C}[1](z) = \frac{1}{2\pi i} \int_{-1}^1 \frac{dx}{x-z}$$

3.  $\mu(z) = \frac{\log(z-1) - \log(z+1)}{2\pi i}$  is the unique function analytic in  $\mathbb{C} \setminus [-1, 1]$  with weaker than pole singularities at  $\pm 1$  satisfying  $\mu(\infty) = 0$  and

$$\mu_+(x) - \mu_-(x) = 1 \quad \text{for} \quad -1 < x < 1.$$

*Demonstration* Here we see from the phase plot of  $\mu$  that it has a branch cut on  $[-1, 1]$ :

```
μ = z → (log(z-1) - log(z+1))/(2π*im)
phaseplot(-3..3, -3..3, μ)
```



For  $-1 < x < 1$  we have the jump 1:

$$x = 0.3$$

$$\mu(x + 0.0im) - \mu(x - 0.0im)$$

$$1.0 + 0.0im$$

For  $x < -1$  we see that the branch cuts cancel and we have continuity:

$$x = -4.3$$

$$\mu(x + 0.0im) - \mu(x - 0.0im)$$

$$0.0 + 0.0im$$

**Remark** As an aside, these integrals are computationally difficult because of the singularity in the integrand, hence standard integration methods become slow as  $z$  approaches the interval. There are other specialised routines (as implemented in `cauchy(f,z)`) that are much more efficient:

```
using BenchmarkTools
z = 0.1 +0.0001im
x = Fun()
f = exp(x)
#  $\mu s$  is micro seconds (1E-6 seconds) while  $ms$  is milliseconds
# (1E-3 seconds)
@btime cauchy(f, z) # specialised routine
@btime sum(f/(x-z))/(2π*im) # standard quadrature
10.799 μ@*(s (18 allocations: 1.64 KiB)344.853 ms (305 allocations:
207.54 MiB)0.5525638794334992 - 0.3181012137711561im
```

We can evaluate the limit from above and below using the specialised routine, where standard quadrature breaks down. Here we see numerically that we recover  $f$  from taking the difference:

```
cauchy(f, 0.1+0.0im)-cauchy(f, 0.1-0.0im) , f(0.1)
(1.1051709180756475 + 0.0im, 1.1051709180756475)
```