

# Tidal Deformation of Neutron Stars

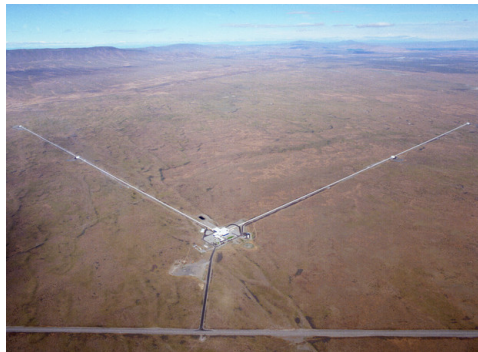
Marco Zenari

Supervisor: Prof. Albino Perego

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# Introduction

- GW150914: first detection of GWs;
- GWs are ripples in the curvature of space-time;
- GW170817: GWs from the inspiral phase of BNSs;
- Our aim: main aspects of NSs, BNSs, tidal effects and Love number  $k_2$  calculation.



**Figure:** LIGO detector site at Hanford (Washington). Credits: LIGO.

# Neutron stars

- Remnant left behind SNs can give birth to NSs;
- First detection of a pulsar: Cavendish Laboratory in 1967;
- Pulsars are highly magnetized, rotating NSs;
- Magnetic dipole model:
  - ▶  $\dot{T} \propto B^2$ ;
  - ▶  $B \propto \sqrt{T \dot{T}}$ .

# Tolman-Oppenheimer-Volkoff equations

- The structure of stars is determined by *hydrostatic equilibrium*;
- Gravity inside and at the surface of NSs is so strong that to describe their structure GR is needed;
- In GR all form of energy, including pressure, contribute to gravity.

## TOV equations

$$\frac{dm}{dr} = \frac{4\pi r^2 \varepsilon}{c^2}$$
$$\frac{dP}{dr} = -\frac{G\varepsilon m(r)}{c^2 r^2} \left(1 + \frac{P}{\varepsilon}\right) \left[1 + \frac{4\pi r^3 P}{m(r)c^2}\right] \left[1 - \frac{2Gm(r)}{c^2 r}\right]^{-1}$$

# Equation of state of neutron stars

- The EOS is the information on the matter composition together with the relationships between:

- ▶ matter pressure  $P$ ;
- ▶ energy density  $\varepsilon$ ;
- ▶ baryonic density  $n$ ;
- ▶ temperature  $T$ .

- Non-interacting fermions:

## Polytrope EOS

$$P = K\rho^\Gamma$$

## Polytrope exponent

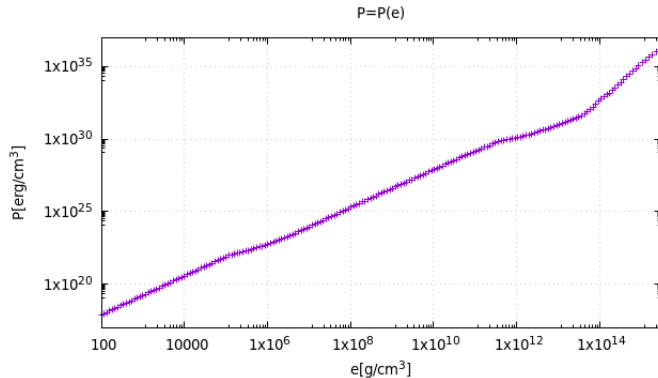
Non-relativistic limit  $\Gamma = 5/3$

Ultra-relativistic limit  $\Gamma = 4/3$

- Interactions in a NS cannot be ignored;
- Some neutrons will decay via  $n \rightarrow p + e + \bar{\nu}_e$ .

# Bombaci-Logoteta EOS<sup>1</sup>

- The EOS can be used in tabulated form in numerical simulations.
- Plot in logarithmic scale of the pressure  $P$  in  $\text{erg}/\text{cm}^3$  as function of the energy density  $\varepsilon$  in  $\text{g}/\text{cm}^3$  for the BL equation of state.



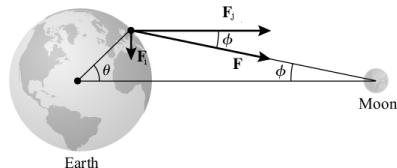
<sup>1</sup>Astronomy & Astrophysics 646 (2021), A55

# Compact binary systems and tidal deformation

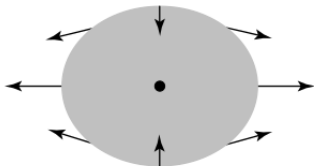
- Binary systems in which at least one star is a compact object (WDs, NSs, BHs);
- Hulse-Taylor binary (PSR 1913+16), discovered in 1974;
- GR predicts the emission of GWs that carry away energy and momentum;
- The two objects become closer and closer, until they merge;
- In the final phase of the inspiral tidal forces between the NSs tend to stretch the compact objects.

# Physics of tides: Earth-Moon system

- The force on an object due to his non-zero size is known as a tidal force;
- This phenomenon happens in every binary system;



**Figure:** The geometry of tidal forces acting on Earth due to the Moon.



**Figure:** The differential gravitational force on Earth, relative to its center.

## Differential gravitational force

$$\Delta \mathbf{F} \simeq \frac{GMmR}{r^3} (2 \cos \theta \hat{\mathbf{i}} - \sin \theta \hat{\mathbf{j}})$$

The result is that Earth is elongated.



# General relativity tidal theory in Newtonian approximation

- The inspiral phase of BNSs can be studied with the PPN formalism;
- Tidal effects depend on the EOS of the NSs.

In PPN  $U_{\text{ext}}$  generates a quadrupolar tidal field  $\mathcal{E}_{ij}$ . In GR  $\mathcal{E}_{ij}$  depends on the Riemann Tensor.

Quasi-static approximation and expansion at the linear order (for some constant  $\lambda$ ):

We can define the dimensionless Love number  $k_2$  as:

## Quadrupolar tidal field

$$\mathcal{E}_{ij} = -\partial_i \partial_j U_{\text{ext}}$$

## Induced quadrupole

$$Q_{ij}(t) = -\lambda \mathcal{E}_{ij}(t).$$

## Love number $k_2$

$$k_2 = \frac{3}{2} \frac{G\lambda}{R^5}.$$

## Love number $k_2$

- Typical value of  $k_2$  for NSs is between 0.05 and 0.15;
- To calculate  $k_2$  we need to integrate the following equations<sup>2</sup>:

### Love number $k_2$

$$k_2 = \frac{8C^5}{5}(1-2C)^2[2+2C(y_R-1)-y_R] \left\{ 2C[6-3y_R+3C(5y_R-8)] + 4C^3[12-11y_R+C(3y_R-2)+2C^2(1+y_R)] + 3(1-2C)^2[2-y_R+2C(y_R-1)] \times \ln(1-2C) \right\}^{-1}$$

### Compactness

$$C = \frac{GM}{c^2 R},$$

### y function

$$y_R = y(R)$$

<sup>2</sup>Physical Review C 98.3 (2018), p. 035804; Phys. Rev. D80 (2009), p. 084035; The Astrophysical Journal 677.2 (2008), p. 1216.

# y function

## Differential equation

$$r \frac{dy(r)}{dr} + y(r)^2 + y(r)F(r) + r^2 Q(r) = 0$$

## Initial condition

$$y(0) = 2$$

## F(r)

$$F(r) = \frac{r - \frac{4\pi G}{c^4} r^3 [\varepsilon(r) - P(r)]}{r - \frac{2m(r)G}{c^2}}$$

## Q(r)

$$Q(r) = \frac{\frac{4\pi G}{c^4} r \left[ 5\varepsilon(r) + 9P(r) + \frac{\varepsilon(r) + P(r)}{\partial P(r)/\partial \varepsilon(r)} \right] - \frac{6}{4\pi r^2}}{r - \frac{2Gm(r)}{c^2}} - 4 \left[ \frac{\frac{Gm(r)}{c^2} + 4\pi \frac{Gr^3 P(r)}{c^4}}{r^2 \left[ 1 - \frac{2Gm(r)}{c^2 r} \right]} \right]^2$$

# Overview of the code

- Code written in C;
- Dimensionless units in order to avoid numerical errors;
- We have integrated the TOV equations together with the equations for the Love number  $k_2$  and the information from the BL EOS;
- Integrator: Runge-Kutta IV;
- Linear interpolation of the EOS.

## Initial condition

$$P(0) = P_c$$

$$m(0) = 0$$

$$y(0) = 2$$

## End of the integration

$$P < \delta$$

with  $\delta$  arbitrarily chosen

# Results

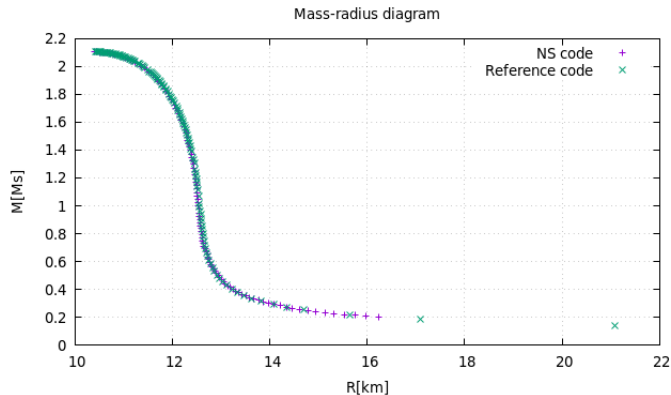
- The output of our code consists in a sequence of possible value of mass, radius, compactness and Love number  $k_2$  that correspond to a given central pressure, for a NS with BL EOS;
- The results are compatible with the ones obtained with a public Matlab code <sup>3</sup>.

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<sup>3</sup>Phys. Rev. D78 (2008), p. 024024; Phys. Rev. D80 (2009), p. 084035

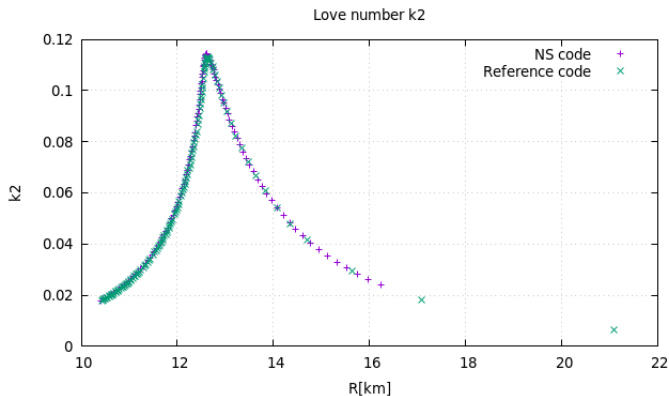
## Results: mass-radius diagram

- The radius decrease with the increasing of the mass;
- Typical radius for a NS is between 12 and 13 km;
- The mass has a maximum value of about  $2.102 M_{\odot}$ .



## Results: Love number $k_2$

- The value of the obtained  $k_2$  are in the range expected (0.05 – 0.15);
- There is a cusp for a radius that is between 12.5 and 13 km. This is a typical trend for low masses NSs ( $M < 1M_\odot$ ).



# Conclusions

- The detection of GWs had a profound impact in our understanding of gravity and the physics of dense matter;
- GWs from the inspiral phase of BNSs can be used to measure the tidal deformability of the NSs, setting stringent constraints on their EOS;
- The derivation of an EOS suitable for their description requires a tremendous theoretical and computational effort;
- Continuing the detection of GWs due to inspiral of binary systems could provide a large amount of data that could help in the understanding of matter under such extreme conditions.