## 1 Energy non-conservation in Sean Carroll's paper

Sean carroll shows energy non-conservation by constructing a simple example, with  $\mathcal{H} = \mathcal{H}_s \otimes \mathcal{H}_e$  and  $\mathcal{H}_s = span\{|1\rangle, |2\rangle\}$   $\mathcal{H}_e = span\{|1\rangle, |2\rangle, |3\rangle\}$ . The the Hamiltonian is constructed such that,

$$|\psi(0)\rangle = (\alpha|1\rangle + \beta|2\rangle)|0\rangle \tag{1}$$

$$|\psi(1)\rangle = \alpha|1\rangle|1\rangle + \beta|1\rangle|2\rangle \tag{2}$$

as well as  $|1\rangle|1\rangle$  an eigenstate with energy  $E_1$  and  $|2\rangle|2\rangle$  an eigenstate with energy  $E_2$ .

But Sean Carroll never considers the preparation of the quantum superposition. We see a scenario where we try to prepare a superposition of different energy eigenstates.

## 1.1 The battery pendulum

### 1.1.1 The scenario in words, Schroedinger's battery-pendulum

An isolated system contains a radioactive isotope connected to a detector, closing the circuit between a full battery and an excitable pendulum at rest. After a while, there was 1/2 chance of a decay happening and being detected, thus causing the battery to empty by exciting the pendulum. The other chance is that nothing happens (battery full and pendulum at rest). Considering only the pendulum, we see that we have a system in a superposition of 2 states which have different energies, potentially causing 2 worlds with different energies after measurement. But considering the entangled battery, the energy of both worlds will actually be the same after measurement. World 1 = excited pendulum + empty battery. World 2 = pendulum at rest + full battery

### 1.1.2 More formally

We take our radioactive atom as just a qubit in a superposition of state (we can prepare this via optics). Our pendulum at rest is  $|\downarrow\rangle$  and excited is  $|\uparrow\rangle$ . The full battery is  $|b_f\rangle$ , the empty battery is  $|b_e\rangle$ . The environment outside of the box starts as  $|e\rangle$ .

$$|\psi(t_0)\rangle = (|0\rangle + |1\rangle)|\downarrow\rangle|b_f\rangle|e\rangle \tag{3}$$

$$|\psi(t_1)\rangle = (|0\rangle|\downarrow\rangle|b_f\rangle + |1\rangle|\uparrow\rangle|b_e\rangle)|e\rangle \tag{4}$$

$$|\psi(t_1)\rangle = (|\downarrow\rangle|b_f\rangle + |\uparrow\rangle|b_e\rangle)|e\rangle \tag{5}$$

$$|\psi(t_2)\rangle = |\downarrow\rangle|b_f\rangle|e_1\rangle + |\uparrow\rangle|b_e\rangle|e_2\rangle \tag{6}$$

At  $t_0$ , the qubit is ready in a superposition and the battery+pendulum starts full + at rest. At time  $t_1$ , by interaction, the qubit has places the pendulum+battery in a superposition. We consider the qubit state as part of the battery. At  $t_2$  we opened the box and decoherence caused the wavefucntion to branch, leaving 2 worlds with conserved energy.

**Making it work.** Does adding an intermediary step allows for non-conservation?

$$|\psi(t_0)\rangle = (|0\rangle + |1\rangle)|\downarrow\rangle|b_f\rangle|e\rangle \tag{7}$$

$$|\psi(t_1)\rangle = (|0\rangle|\downarrow\rangle|b_f\rangle + |1\rangle|\uparrow\rangle|b_e\rangle)|e\rangle \tag{8}$$

$$|\psi(t_1)\rangle = (|\downarrow\rangle|b_f\rangle + |\uparrow\rangle|b_e\rangle)|e\rangle \tag{9}$$

$$|\psi(t_2')\rangle = (|\downarrow\rangle + |\uparrow\rangle)|e'\rangle \tag{10}$$

$$|\psi(t_3')\rangle = |\downarrow\rangle|e_1'\rangle + |\uparrow\rangle|e_2'\rangle \tag{11}$$

We change things at  $t'_2$  asking that the battery decoheres without branching the pendulum state such that we recover a separable state with a superposition of pendulum energy state and a unique environment state. Then at  $t'_3$ , we open the box and branch into two worlds with different energies.

For this to work, we need  $e^{-itH_1}|b_f\rangle|e\rangle \approx e^{-itH_2}|b_e\rangle|e\rangle$ , for  $t \geq t_2$  even thought  $\langle b_f|b_e\rangle \approx 0$ . Why is that a tall order?  $|b_f\rangle$  and  $|b_e\rangle$  are macroscopically different objects that we ask evolve to close to the same object at a later time.

## 1.2 Next step/solution

- Working only to creat a tiny energy difference requiring a tiny battery, could make step  $t'_2$  a lot less impossible.
- Show that system size/energy difference acts as a protection to energy non-conservation

in notes it is like this: - Project 2: Energy conservation: I) what initial state does Sean carroll start with and what hamiltonian is necessary to create that state. ii) Study the James cunning model in creating energy superpositions. Is energy conservation respected? Can that state indeed be created? iii) Learn about resource theory of quantum thermodynamics, see if there is a necessary ressource in creating a superposition of energy levels. iv) read the paper that says that you can't create energy superpositions. V) think about tiny battery argument and how size of the system intervenes. Show that system size acts as protector for it not happening. Vi) What about value of the energy difference? Can you find a law that relates energy difference of the superposition and level of fine tuning required?

# 2 Creating a superposition of energies such that branching violates energy conservation

Can we create a system that when measured results in worlds with different energies? Sean Carroll argues yes, but does not explain how to achieve the necessary initial state  $|\psi\rangle = (|\downarrow\rangle + |\downarrow\rangle)|e\rangle$ .

Can we create Sean Carroll's initial state? Making Sean Carroll's energy superpositions is making a system in a superposition of 2 self energy eigenstates which are also pointer states of the environment.

Can we create another system+env such that energy is violated? Multiple ideas/questions:

- The Jaynes-Cummings model
- How to make energy superpositions and what does it mean?
- We don't have to create Sean Carroll's state to get energy violation.
- Can I take the photon in the environment without collapsing the atom state? It needs to not be a pointer state. In the thought experiment, the battery superposition needs to be a pointer state of the environment. (but even then we have energy conservation)
- Small battery idea
- Coherence to make a superposition
- Look at the energy of Schmidt states during branching? ACL model.
- Does the Hamiltonian needed to use coherence to make a superposition allow for energy violation. Both
  my battery example, the JC model and the example used in the catalytic coherence paper have energy
  conservation.

## 2.1 Using the Jaynes-Cummings model

A quantum optics model describing 2 level atom interacting with photon in cavity.

The system has an atom with 2 energy levels with hamiltonian

$$H_A = E_e |e\rangle\langle e| \tag{12}$$

Note that  $E_g = 0$  for simplicity. There is also a photon field with,

$$H_F = \hbar\omega(aa^{\dagger}) \tag{13}$$

Note that we remove the constant scale factor for simplicity. And then there is an interaction Hamiltonian, which after simplification only contains terms relating to de-excitation into a photon and excitation by absorption of a photon,.

$$H_{AF} = \hbar q (\sigma^{+} a + \sigma^{-} a^{\dagger}) \tag{14}$$

We have  $H_{JC} = H_A + H_F + H_{AF}$ , and a bit of maths tells us, is that the energy eigenstates of this system are,

$$|n,+\rangle = \cos(\theta_n/2)|e,n-1\rangle + \sin(\theta_n/2)|g,n\rangle \tag{15}$$

$$|n, -\rangle = \cos(\theta_n/2)|g, n\rangle - \sin(\theta_n/2)|e, n - 1\rangle \tag{16}$$

where n is the number of photon modes, e is the excited atom state, g the atom ground state. So the eigenstates of the system don't have an atom either in the ground state or in the excited state. This means that if the system starts in a state where the atom is just in the excited state, it will evolve to a superposition of excited and ground state atoms. Simply put, The eigenstates of  $H_{JC}$  are superpositions of eigenstates of  $H_A$ . But we cannot create Sean Carroll's initial state, the photon acts as the battery in the thought experiment.

### 2.1.1 A JC model that doesn't conserve energy?

The energy is conserved in branching, because the pointer states of the system have the same energy as the initial state of the system.

If the pointer states of the environment where the energy eigenstates of the cavity, we would have energy non-conservation in branching

To do that we need the interaction between the cavity and the environment to be,

$$H_{JC,env} = |g,n\rangle\langle g,n| \otimes H_e^{(1)} + |e,n-1\rangle\langle e,n-1| \otimes H_e^{(2)}$$

$$\tag{17}$$

as long as the  $H_e$  matrices cause decoherence.

What is the interaction of the environment and an atom? I assume it is such that the pointer states are atoms of definite energy? Why? Is it because the environment is mainly made of other atoms?

Somehow, the JC model evolution causes us to go from 1 pointer state to 2 pointer states, but they have the same energy. Why? the interactions between the electric field and the atom field are energy preserving. so the states that entangle together have exchanged energy, making it constant with entangled states.

-if the interaction hamiltonian doesnt commute with either self interaction hamiltonian, then we get exchange of energy between field 1 and field 2. see thermalization graph for the ACL model –if field 1 commutes with the interaction hamiltonian, then the energy of field 1 is constant in both worlds. Can you get exchange of energy between worlds? see graphs with schmidt states

To make a similar model, just create 2 fields with a self interaction Hamiltonian each, then create an interaction Hamiltonian that doesn't commute with either. what if it commutes with one of them? do i get energy violation then?

## 2.2 Making a superposition (using coherence)

A more quantum info view. 2 papers: "Catalytic coherence" and "Is coherence catalytic?"

Making a superposition unitarly is non-trivial. For example, there is no U unitary such that  $U|00\rangle = \frac{1}{\sqrt{2}}|0\rangle(|0\rangle + |1\rangle)$ . We can see this by attempting to construct U. In the  $\{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$  basis, we can figure out the necessary part of U. (A dot is a un-specified value)

$$U|00\rangle = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 & 0\\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & \cdot & \cdot\\ 0 & \cdot & \cdot & \cdot\\ 0 & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} 1\\0\\0\\0 \end{pmatrix} = \begin{pmatrix} \frac{1}{\sqrt{2}}\\ \frac{1}{\sqrt{2}}\\0\\0 \end{pmatrix}$$
(18)

Next we check if this matrix could be unitary,

What we find is that U cannot be unitary. Note that a unitary matrix is hermitian, which was used to fix  $U_{11}$ 

We can find a U s.t.

$$U\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)|0\rangle = |0\rangle|0\rangle + |1\rangle|1\rangle \tag{20}$$

This is what happens with the battery example, and this is what happens with the block diagonal branching scenario. But this has not created a superposition, it has used a superposition.

Can we find a U s.t,

$$U|0\rangle|0\rangle = \frac{1}{2}(|0\rangle + |1\rangle)(|0\rangle + |1\rangle) \tag{21}$$

This would be like generating coherence from thin air, which I think i impossible... TODO, check

#### 2.2.1 Coherence

Coherence is what is needed to creat superposition in a certain basis. Certain states are better at it than others. Coherence is basis dependent. The maximally mixed state has zero coherence, which makes sense, because how could we make a quantum superposition from a classical state (in the pointer basis)? A measure of coherence are the diagonal terms.

Basic idea: To make a superposition of states you need another superposition of states (check)

TODO: clear definition of coherence. TODO measure of coherence. Issue/question for all non-catlytic ressources: If our world is classical, where is the quantumness? How do we concentrate it to use in quantum computers? We mostly never create a quantum resource, then how "lasers"?) (ZOE question) I guess, a classical state might have coherence wrt to a specific basis? and that basis is used by defining and having access to weird observables? We have sources of these quantum ressources, like T-factories, and lasers for coherence...

Consider,

$$H = a^{\dagger} a \tag{22}$$

with eigenstates  $H|n\rangle n|n\rangle$ . Then a state  $\rho = \sum_{nm} \rho_{nm} |n\rangle \langle m|$  has coherence wrt to the energy eigenbasis, if it has non-zero off diagonal terms. Note that the ground state  $|0\rangle$  has no coherence, but that a coherent states  $a|\alpha\rangle = \alpha |\alpha\rangle does.$  And importantly, that it there is a unitary that goes from one to the other, called the displacement unitary, to  $D(\alpha) = \exp(\alpha a^{\dagger} - \alpha a)(23)$  So clearly, coherence is not conserved under unitary operations.

### 2.2.2 Coherence is non-classical + is related to entanglement

Another intersting note on coherence, is that it is tied to entanglement.

**Distillable coherence** is a measure of how many maximally coherent states can be generated with a state. (maximally coherent state  $\sum_{i} |i\rangle$ , incoherent density matrices:  $\langle p_{i}|i\rangle\langle i|\rangle$ 

Testing our simple battery example: The initial state is  $|\psi(t_0)\rangle \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)|e'\rangle$  with  $|e'\rangle = |\downarrow\rangle |b_f\rangle |e\rangle$ . In the  $\{|0\rangle, |1\rangle\}$  basis this state  $\rho_q = Tr_e[(|0\rangle + |1\rangle)|e'\rangle (\langle 0| + \langle 1|)\langle e'|]$  is the maximally coherent state,

$$C_d(\rho_q) = S(\Delta[\rho_q]) - S(\rho_q) = 1 \tag{24}$$

and evolving it to  $|\psi(t_1)\rangle = \frac{1}{\sqrt{2}}(|0\rangle|e_0'\rangle + |1\rangle|e_1'\rangle$  with  $\langle e_1'|e_0'\rangle = 0$ , we find that  $\rho_q = 1$ , meaning  $C_d(\rho_q) = 0$  and coherence has been used up.

### 2.2.3 How do these papers relate to my problem of studying energy violation in measurements

In the paper, the example where coherence is used to make an energy superposition of a system still leads to energy conservation at the branch level. We then wonder if we can modify the example to lead to energy violation?

TODO quick calculation of why in the paper energy is conserved. + how to modify to lead to energy violation.

### 2.3 Defining the Hamiltonian for it

### 2.3.1 Block diagonal case

In the block diagonal branching case, the energy difference in both worlds is already setup. interestingly it is by difference in total energy eigenstates. in sean carroll's paper its more about a system energy being differente, the environment energy stays the same...

The general block diagonal case is  $H = H_s \otimes \mathbb{1} + H_I + \mathbb{1} \otimes H_e = \sum_i |i\rangle\langle i| \otimes id + \sum_i |i\rangle\langle i| \otimes H_i + \mathbb{1} \otimes H_e$ . We know that if  $[H_s, H_I] = 0$  can cause the environment term and the worlds to have different energies defined by the blocks  $H_i$ . Another way of seeing this is to look at expectation values of the initial and final states,

$$(|0\rangle + |1\rangle)|e\rangle \to |0\rangle|e_0\rangle + |1\rangle|e_1\rangle \tag{25}$$

Then,

$$\langle H(t_0) \rangle = \langle 0|H_s|0 \rangle + \langle 0e|H_I|0e \rangle + \langle e|H_e|e \rangle \tag{26}$$

$$\langle H(t_1) \rangle = \langle 0|H_s|0 \rangle + \langle 0e_1|H_I|0e_1 \rangle + \langle e_1|H_e|e_1 \rangle \tag{27}$$

(28)

Before and after decoherence  $\langle 0|H_s|0\rangle$  is the same, but because of the non-commutativity of  $H_I$  and  $H_e$  we have that  $\langle 0e|H_I|0e\rangle \neq \langle 0e_1|H_I|0e_1\rangle$ . There is exchange of energy between  $\langle H_I\rangle$  and  $\langle H_e\rangle$ , which for a very large environment does not cause any meaningfull change in  $\langle H_e\rangle$ .

There is always conservation of the non-zero weights of the state in the total energy eigenbasis. there doesn't have to be at the level of worlds with the right block diagonal form.

## 2.4 Thinking pointer states

Question: Why does the JC model have energy conservation then? Because of the pointer states of the environment.

Conjecture: energy always seems conserved because pointer states of the environment do not change drastically? If we could have a system in a pointer state, make its Hamiltonian temporarily like the JC model (ie, the pointer state is no longer an energy eigenstate), then it will evolve in a superposition of pointer states of the environment, each with the same energy. Re-activating the interaction between the environment and the system would cause it to appear in a branch as having conserved energy. A way to not have that would be to have a different environment with a different.

## This week's work

- Figure out the general consensus on if coherence is catalytic. To do so, re-read the 2 papers on catalytic coherence...
- Why do we care about coherence? Does it help me solve the issue of understanding why the environment can absorbe the superposition of pulses without branching. Does the catalycity of coherence actually matter at all for our question?
- Does "resource theory of thermodyanmics" contain info that is useful in our quest to understand if sean carroll is right.

## Is it or not?

Aberg: Yes, see my first paper

thoughts: Runnaway coherence. If coherence was catalytic, could i not make the whole universe in a certain state that is not reachable by unitary evolution? ie the near absolutely cohrence state in a choose set of basis. issue with any resource being catalytic, the operation by which you use the ressource catalytically is a free operation, then if we can use it to make more of itself, then we can change the amount of a resource we have for free. Note, you may not necessarly be abel to use a resource catlytically to make more of itself...

resource theory: states we can get for free. free operations. result: state conversion condition, ie from a state rho you know if you can make a state rho' using your free operations. this condition comes in the form of a function of the states following a certain inequality (ie majorization of rho). confusion: Can we ever only do free operations? is it a human limitation or a physical limitation? If it is physical, i dont get how studying resource theories is different from studying unitary evolution... clearing of confusions: the free op, dont necessarly reflect absolut limitations, but arbitrarily choosen ones.(such as just local op.)

Entanglement RT example: Alice and Bob cna actg localy and exchange information classically together (ie do classical stuff). Turns out starting from a certain class of states (classical separble states) you can never get an entangled state like that, so we identified the resource of entanglement. Often operations we want to do usinfg resources (ie teleportation) requier a pure form, so we also study distillation of resources.

## 2.5 So why do we care about a RT of thermodynamics?

I think we might want to consider thermal operations only... Or to see if sean carroll's state is achievable with only TO

"So in conclusion, in general, for TOs it remains an open question as to what are the necessary and sufficient conditions characterizing transitions between arbitrary, energy-coherent states." from the RT of thermo paper, i think this is relevant because we might be attempting to see if sean carroll's state requier maybe some thermodynamic resource to be achieved??? We have this thermodynamic resource, and we have coherence. which one is necessary to achieve SC's state?

"For initial and final states which are block-diagonal, the set of allowed transitions via GPs coincide with thermal operations. However, for general quantum states, GPs may act on energy-incoherent initial states to create energycoherent final states. This is not possible via thermal operations." !!!! So by free thermal operations, we can increase cohjerence in the energy basis? figure it out... GP is one possible RT, a less populare one

Note that b attery is an already loaded term in RT of thermodynamics.

The battery thing is pretty related to my pendulum battery example!