

# ***Chapter 6: Frequency Response of Amplifiers***

**6.1 Basic Current Mirrors**

**6.2 Common-Source Stage**

**6.3 Source Followers**

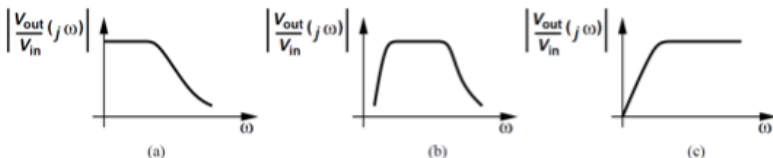
**6.4 Common-Gate Stage**

**6.5 Cascode Stage**

**6.6 Differential Pair**

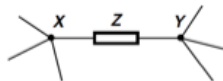
**6.7 Gain-Bandwidth Trade-Offs**

# General Considerations

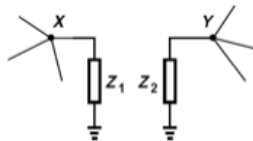


- In this chapter, we are primarily interested in the magnitude of the transfer function.
- The magnitude of a complex number  $a + jb$  is given by  $\sqrt{a^2 + b^2}$
- Zeros and poles are respectively defined as the roots of the numerator and denominator of the transfer function.

# Miller effect



(a)



(b)

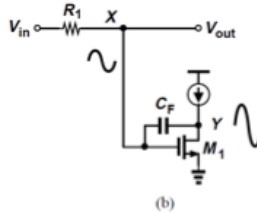
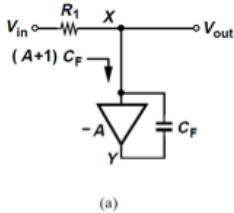
- If the circuit of Fig (a) can be converted to that of Fig (b), then

$$Z_1 = Z / (1 - A_v)$$

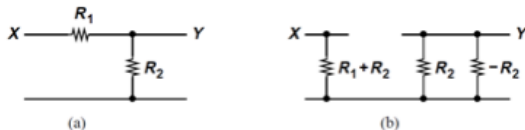
$$Z_2 = Z / (1 - A_v^{-1})$$

$$A_v = V_Y / V_X$$

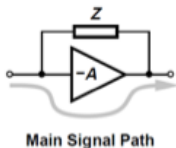
# Example



- A student needs a large Capacitor and decides to utilize the Miller multiplication
- What is the issues in this approach?



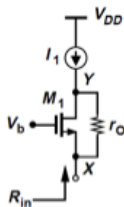
Improper application of Miller's theorem



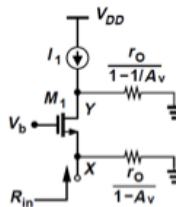
Typical case for valid application of Miller's theorem.

- **Miller's theorem does not stipulate the conditions under which this conversion is valid.**
- **If the impedance  $Z$  forms the only signal path between  $X$  and  $Y$ , then the conversion is often invalid.**

# Example



(a)



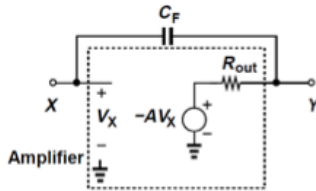
(b)

- Calculate the input resistance of the circuit shown.
- Since  $A_v$  is usually greater than unity,  $r_O/(1 - A_v)$  is a negative resistance.

- $A_v = 1 + (g_m + g_{mb})r_O$

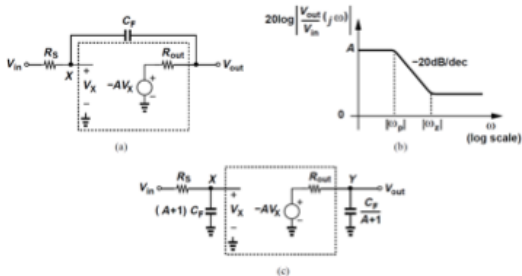
$$\begin{aligned}
 R_{in} &= \frac{r_O}{1 - [1 + (g_m + g_{mb})r_O]} \parallel \frac{1}{g_m + g_{mb}} \\
 &= \frac{-1}{g_m + g_{mb}} \parallel \frac{1}{g_m + g_{mb}} \\
 &= \infty.
 \end{aligned}$$

# Example



- The value of  $A_v = V_Y / V_X$  must be calculated at the frequency of interest.
- In the figure, the equivalent circuit  $V_Y \neq -AV_X$  that at high frequencies.
- In many cases we use the low-frequency value of  $V_Y / V_X$  to gain insight.
- We call this approach “Miller’s approximation.”

# Example



- **Direct Calculation:**

$$\frac{V_{out}}{V_{in}}(s) = \frac{R_{out}C_F s - A}{[(A+1)R_S + R_{out}]C_F s + 1}$$

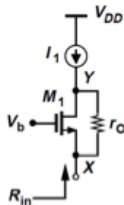
- **Miller Approximation:**

$$\frac{V_{out}}{V_{in}}(s) = \frac{-A}{[(1+A)R_S C_F s + 1] \left( \frac{1}{1+A^{-1}} C_F R_{out} s + 1 \right)}$$

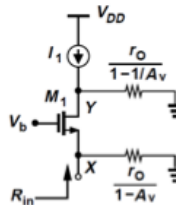
- **Miller's approximation has eliminated the zero and predicted two poles for the circuit!**



# Example



(a)

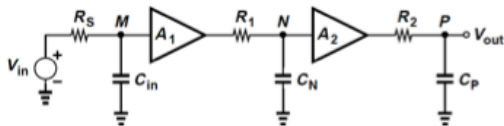


(b)

$$\begin{aligned}
 R_{out} &= \frac{r_O}{1 - 1/A_v} \\
 &= \frac{r_O}{1 - [1 + (g_m + g_{mb})r_O]^{-1}} \\
 &= \frac{1}{g_m + g_{mb}} + r_O,
 \end{aligned}$$

- **Actual Rout = rO**
- **Miller's approximation:**
- **(1) it may eliminate zeros**
- **(2) it may predict additional poles**
- **(3) it does not correctly compute the “output” impedance.**

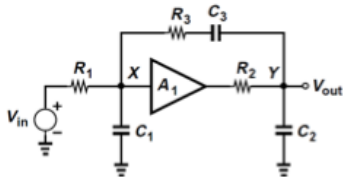
# Association of Poles with Nodes



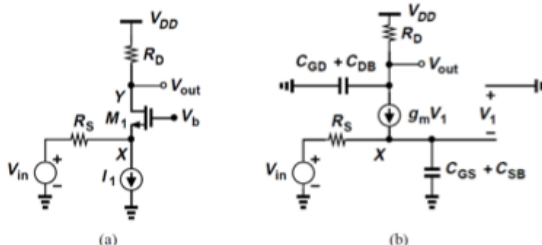
- The overall transfer function can be written as

$$\frac{V_{out}}{V_{in}}(s) = \frac{A_1}{1 + R_S C_{in} s} \cdot \frac{A_2}{1 + R_1 C_N s} \cdot \frac{1}{1 + R_2 C_P s}$$

- Each node in the circuit contributes one pole to the transfer function.
- Not valid in general. Example:



# Example



- At node X:

$$\omega_{in} = \left[ (C_{GS} + C_{SB}) \left( R_S \parallel \frac{1}{g_m + g_{mb}} \right) \right]^{-1}$$

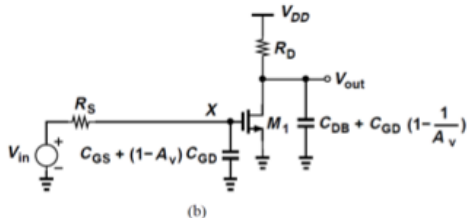
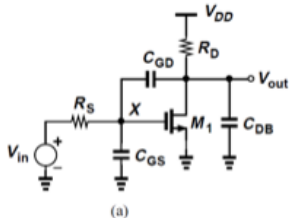
- At node Y:

$$\omega_{out} = [(C_{DG} + C_{DB})R_D]^{-1}$$

- The overall transfer function...

$$\frac{V_{out}}{V_{in}}(s) = \frac{(g_m + g_{mb})R_D}{1 + (g_m + g_{mb})R_S} \cdot \frac{1}{\left(1 + \frac{s}{\omega_{in}}\right) \left(1 + \frac{s}{\omega_{out}}\right)}$$

# Common-Source Stage



- The magnitude of the “input” pole

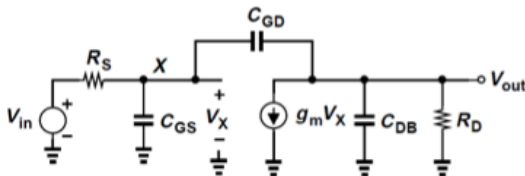
$$\omega_{in} = \frac{1}{R_S[C_{GS} + (1 + g_m R_D)C_{GD}]}$$

- At the output node

$$\omega_{out} = \frac{1}{R_D(C_{DB} + C_{GD})}$$

$$\frac{V_{out}}{V_{in}}(s) = \frac{-g_m R_D}{\left(1 + \frac{s}{\omega_{in}}\right) \left(1 + \frac{s}{\omega_{out}}\right)}$$

# Direct Analysis



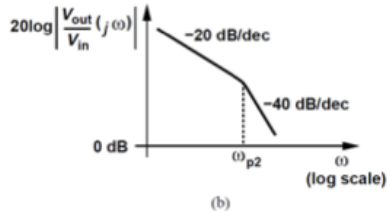
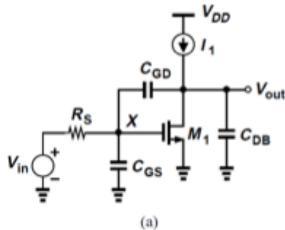
$$\frac{V_{out}}{V_{in}}(s) = \frac{(C_{GD}s - g_m)R_D}{R_S R_D \xi s^2 + [R_S(1 + g_m R_D)C_{GD} + R_S C_{GS} + R_D(C_{GD} + C_{DB})]s + 1}$$

- While the denominator appears rather complicated, it can yield intuitive expressions for the two poles.  $|\omega_{p1}| \ll |\omega_{p2}|$
- “Dominant pole” approximation.

$$\omega_{p1} = \frac{1}{R_S(1 + g_m R_D)C_{GD} + R_S C_{GS} + R_D(C_{GD} + C_{DB})}$$

- The intuitive approach provides a rough estimate with much less effort.

# Example



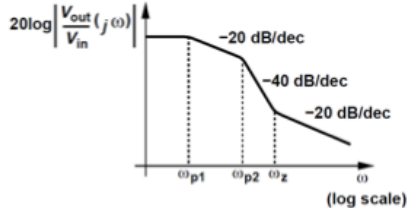
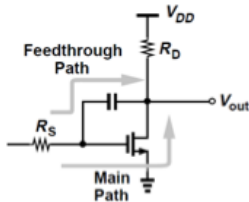
$$\begin{aligned}\frac{V_{out}}{V_{in}}(s) &= \frac{C_{GD}s - g_m}{R_S\xi s^2 + [g_m R_S C_{GD} + (C_{GD} + C_{DB})]s} \\ &= \frac{C_{GD}s - g_m}{s[R_S(C_{GS}C_{GD} + C_{GS}C_{DB} + C_{GD}C_{DB})s + (g_m R_S + 1)C_{GD} + C_{DB}]}\end{aligned}$$

- One pole is at the origin because the dc gain is infinity.
- For a large CDB or load capacitance

$$\omega_{p2} \approx \frac{1}{R_S(C_{GS} + C_{GD})}$$

- No miller multiplication. Why?

# Zero in Transfer Function



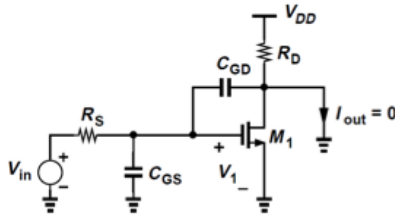
$$\frac{V_{out}(s)}{V_{in}} = \frac{(C_{GD}s - g_m)R_D}{R_S R_D \xi s^2 + [R_S(1 + g_m R_D)C_{GD} + R_S C_{GS} + R_D(C_{GD} + C_{DB})]s + 1}$$

- The transfer function of exhibits a zero given by

$$\omega_z = +g_m/C_{GD}$$

- $C_{GD}$  provides a feedthrough path that conducts the input signal to the output at very high frequencies.

# Calculating zero in a CS stage



- The transfer function  $V_{out}(s)=V_{in}(s)$  must drop to zero for  $s = s_z$ .
- Therefore, the currents through  $C_{GD}$  and  $M1$  are equal and opposite:

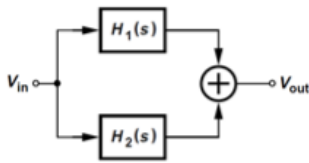
$$V_1 C_{GD} s_z = g_m V_1$$

- That is

$$s_z = +g_m / C_{GE}$$



# Example

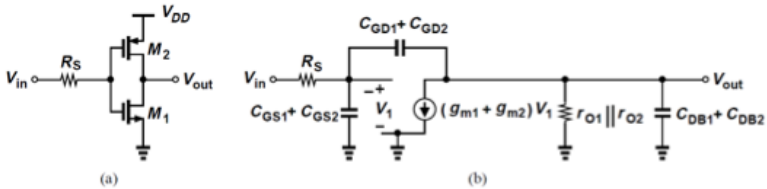


- Can this (the zero) occur if  $H_1(s)$  and  $H_2(s)$  are first-order low-pass circuits?
- $H_1 = A_1 / (1 + s/\omega_{p1})$  and  $H_2 = A_2 / (1 + s/\omega_{p2})$

$$\frac{V_{out}}{V_{in}}(s) = \frac{(\frac{A_1}{\omega_{p2}} + \frac{A_2}{\omega_{p1}})s + A_1 + A_2}{(1 + \frac{s}{\omega_{p1}})(1 + \frac{s}{\omega_{p2}})}$$

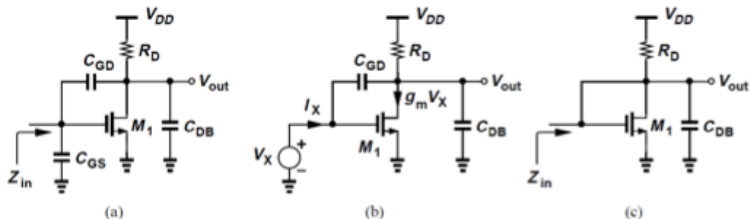
- The overall transfer function contains a zero.

# Example



- Since the corresponding terminals of  $M_1$  and  $M_2$  are shorted to one another in the small-signal model, we merge the two transistors.
- The circuit thus has the same transfer function as the simple CS stage.

# Miller's Approximation



- With the aid of Miller's approximation,

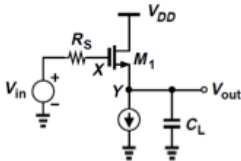
$$Z_{in} = \frac{1}{[C_{GS} + (1 + g_m R_D)C_{GD}]s}$$

- But at high frequencies, the effect of the output node capacitance must be taken into account.
- Ignore  $C_{GS}$

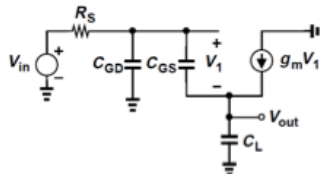
$$\frac{V_X}{I_X} = \frac{1 + R_D(C_{GD} + C_{DB})s}{C_{GD}s(1 + g_m R_D + R_D C_{DB}s)}$$

- if  $C_{GD}$  is large, it provides a low impedance path between the gate and drain of  $M_1$ .

# Source Followers



(a)



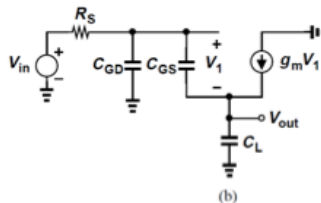
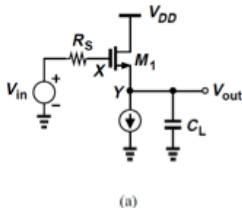
(b)

- The strong interaction between nodes X and Y through  $C_{GS}$  makes it difficult to associate a pole with each node.

$$\frac{V_{out}}{V_{in}}(s) = \frac{g_m + C_{GS}s}{R_S(C_{GS}C_L + C_{GS}C_{GD} + C_{GD}C_L)s^2 + (g_m R_S C_{GD} + C_L + C_{GS})s + g_m}$$

- Contains a zero in the left half plane. Why?

# Example

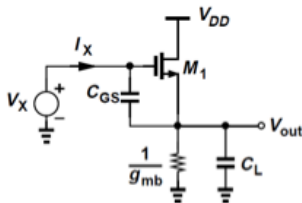


- Transfer function if  $C_L = 0$ ?

$$\begin{aligned}\frac{V_{out}}{V_{in}} &= \frac{g_m + C_{GS}s}{R_S C_{GS} C_{GD} s^2 + (g_m R_S C_{GD} + C_{GS})s + g_m} \\ &= \frac{g_m + C_{GS}s}{(1 + R_S C_{GD}s)(g_m + C_{GS}s)} \\ &= \frac{1}{1 + R_S C_{GD}s}.\end{aligned}$$

- $C_{GS}$  disappear
- In the absence of channel-length modulation and body effect, the voltage gain from the gate to the source is equal to unity.

# Input impedance



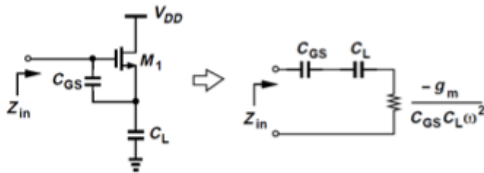
- **$C_{GD}$  simply shunts the input and can be ignored initially.**

$$Z_{in} = \frac{1}{C_{GS}s} + \left(1 + \frac{g_m}{C_{GS}s}\right) \frac{1}{g_{mb} + C_L s}$$

- **If  $g_{mb} = 0$  and  $C_L = 0$ , then  $Z_{in} = \infty$**
- **$C_{GS}$  is entirely bootstrapped by the source follower and draws no current from the input.**
- **At Low frequency the overall input capacitance is equal to  $C_{GD}$  plus a fraction of  $C_{GS}$ .**

$$Z_{in} \approx \frac{1}{C_{GS}s} \left(1 + \frac{g_m}{g_{mb}}\right) + \frac{1}{g_{mb}}$$

# Input Impedance

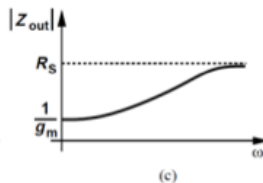
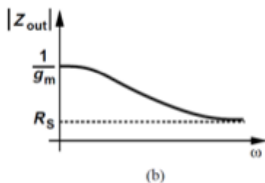
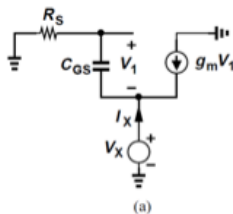


- At high frequencies,  $g_{mb} \ll |C_L s|$

$$Z_{in} \approx \frac{1}{C_{GS}s} + \frac{1}{C_L s} + \frac{g_m}{C_{GS}C_L s^2}$$

- A source follower driving a load capacitance exhibits a negative input resistance, possibly causing instability.

# Output Impedance



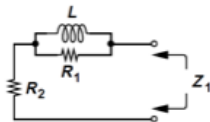
$$Z_{out} = \frac{V_X}{I_X}$$

$$= \frac{R_S C_{GS} s + 1}{g_m + C_{GS} s}$$

- **At low frequency:**  $Z_{out} \approx 1/g_m$
- **At very high frequencies,**  $Z_{out} \approx R_S$
- **Because  $C_{GS}$  shorts the gate and the source.**
- **Which one of these variations is more realistic?**



# Output Impedance

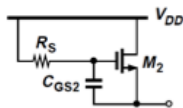


1. Equivalent output impedance of a source follower

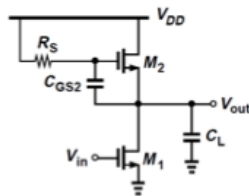
- Since the output impedance increases with frequency, we postulate that it contains an inductive component.

$$Z_{out} - \frac{1}{g_m} = \frac{C_{GS}s \left( R_S - \frac{1}{g_m} \right)}{g_m + C_{GS}s}$$
$$\frac{1}{Z_{out} - \frac{1}{g_m}} = \frac{1}{R_S - \frac{1}{g_m}} + \frac{1}{\frac{C_{GS}s}{g_m} \left( R_S - \frac{1}{g_m} \right)}$$
$$L = \frac{C_{GS}}{g_m} \left( R_S - \frac{1}{g_m} \right)$$

# Example



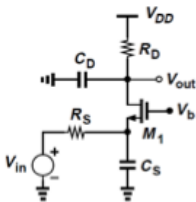
(a)



(b)

- Can we construct a (two-terminal) inductor from a source follower?
- Yes, but non-ideal.
- It also incurs a parallel resistance and a series resistance.
- The inductance can partially cancel the load capacitance,  $C_L$ , at high frequencies, thus extending the bandwidth.

# Common-Gate Stage



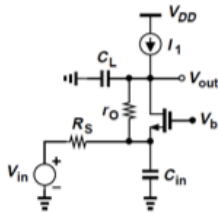
Common-gate stage at high frequencies

- **A transfer function**

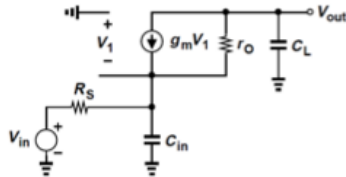
$$\frac{V_{out}}{V_{in}}(s) = \frac{(g_m + g_{mb})R_D}{1 + (g_m + g_{mb})R_S} \frac{1}{\left(1 + \frac{C_S}{g_m + g_{mb} + R_S^{-1}s}\right)(1 + R_D C_D s)}$$

- **No Miller multiplication of capacitances.**
- **$R_D$  is typically maximized, so the dc level of the input signal must be quite low.**
- **As an amplifier in cases where a low input impedance is required**
- **In cascode stages.**

# Example



(a)



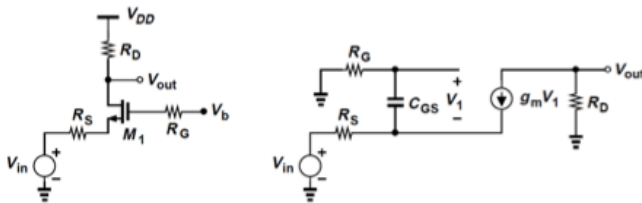
(b)

- Why  $Z_{in}$  becomes independent of  $C_L$  as this capacitance increases?

$$Z_{in} = \frac{1}{g_m + g_{mb}} + \frac{1}{C_L s} \cdot \frac{1}{(g_m + g_{mb})r_O}.$$

- As  $C_L$  or  $s$  increases,  $Z_{in}$  approaches  $1/(g_m + g_{mb})$
- $C_L$  lowers the voltage gain of the circuit, thereby suppressing the effect of the negative resistance introduced by Miller effect through  $r_O$ .

# CG Stage



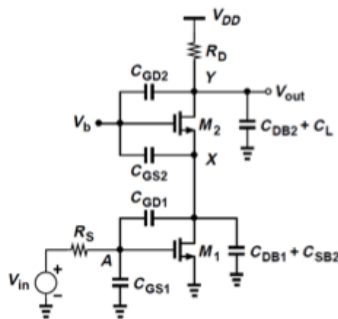
- The bias network providing the gate voltage exhibits a finite impedance.
- Consider only  $C_{GS}$  here.

$$\frac{V_{out}}{V_{in}} = \frac{g_m R_D}{(R_G + R_S)C_{GS}s + 1 + g_m R_S},$$

$$\omega_p = \frac{1 + g_m R_S}{(R_G + R_S)C_{GS}}.$$

- Lowering the pole magnitude.
- Output impedance of the circuit drops at high frequencies.

# Cascode Stage



$$\omega_{p,A} = \frac{1}{R_S \left[ C_{GS1} + \left( 1 + \frac{g_{m1}}{g_{m2} + g_{mb2}} \right) C_{GD1} \right]}$$

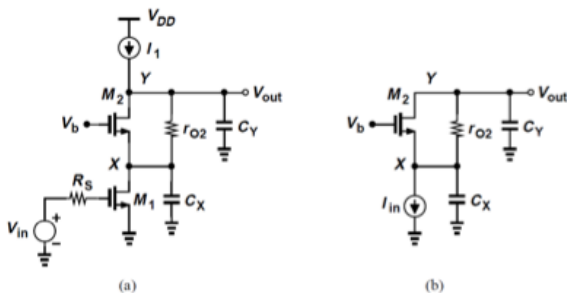
$$\omega_{p,X} = \frac{g_{m2} + g_{mb2}}{2C_{GD1} + C_{DB1} + C_{SB2} + C_{GS2}}$$

$$\omega_{p,Y} = \frac{1}{R_D(C_{DB2} + C_L + C_{GD2})}$$

High-frequency model of a cascode stage.

- **Miller effect is less significant in cascode amplifiers than in common-source stages.**
- **But  $\omega_{p,X}$  is typically quite higher than the other two.**
- **What if  $R_D$  is replaced by a current source?**
- **Pole at node X may be quite lower, but transfer function will not affect much by this. See example.**

# Example



- Compute the transfer function.

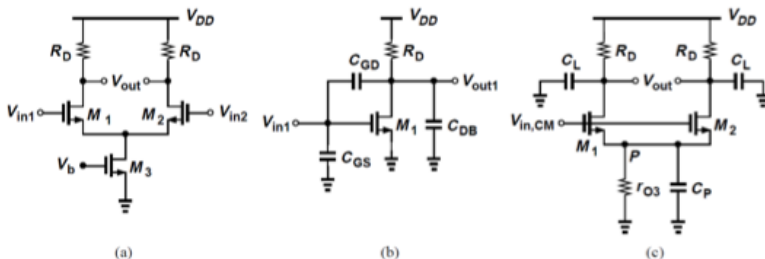
$$\frac{V_{out}}{I_{in}} = -\frac{g_{m2}r_{O2} + 1}{C_X s} \cdot \frac{1}{1 + (1 + g_{m2}r_{O2})\frac{C_Y}{C_X} + C_Y r_{O2} s}$$

- For  $g_{m2}r_{O2} \gg 1$

$$\frac{V_{out}}{I_{in}} = -\frac{g_{m1}g_{m2}}{C_Y C_X s} \frac{1}{g_{m2}/C_X + s}$$

- The magnitude of the pole at node X is still given by  $g_{m2}/C_X$ . Why?

# Differential Pair



(a) Differential pair, (b) half-circuit equivalent, (c) equivalent circuit for common-mode inputs.

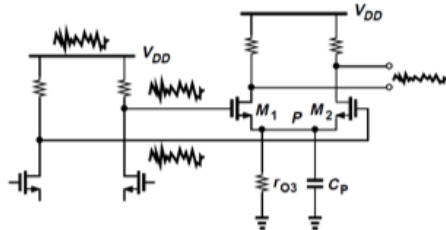
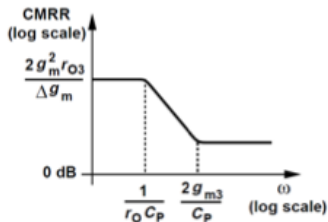
- For differential signals, the response is identical to that of a common-source stage.
- the common-mode rejection of the circuit degrades considerably at high frequencies.

$$A_{v,CM} = -\frac{\Delta g_m \left[ R_D \parallel \left( \frac{1}{C_L s} \right) \right]}{(g_{m1} + g_{m2}) \left[ r_{O3} \parallel \left( \frac{1}{C_P s} \right) \right] + 1}$$

- Channel length modulation, body effect, and other capacitances are neglected.

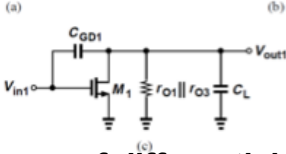


# Differential Pair



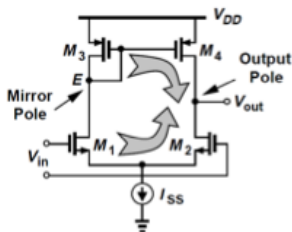
$$\text{CMRR} \approx \frac{g_m}{\Delta g_m} \frac{r_{O3} C_P s + 1 + 2g_m r_{O3}}{r_{O3} C_P s + 1}$$

- This transfer function contains a zero and a pole.
- The magnitude of the zero is much greater than the pole.
- Common-mode disturbance at node P translates to a differential noise component at the output, if the supply voltage contains high-frequency noise and the circuit exhibits mismatches.
- Trade-off between voltage headroom and CMRR.



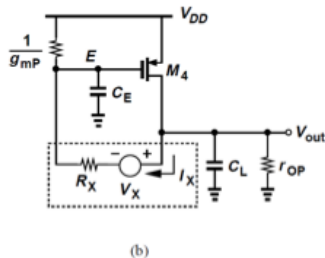
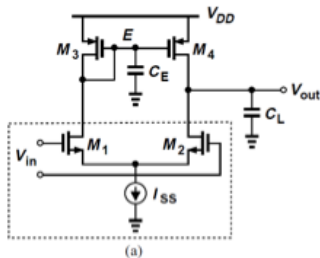
- **More on chapter 10**

# Differential Pair with Active Load



- How many poles does this circuit have?
- The severe trade-off between  $g_m$  and CGS of PMOS devices results in a pole that impacts the performance of the circuit.
- The pole associated with node  $E$  is called a “mirror pole.”

# Active Load



- Replacing  $V_{in}$ ,  $M1$ , and  $M2$  by a Thevenin equivalent.

$$V_X = g_{mN} r_{ON} V_{in}$$

$$R_X = 2r_{ON}$$

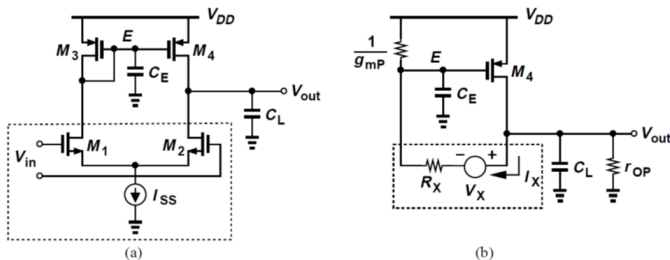
$$\frac{V_{out}}{V_{in}}$$

$$= \frac{g_{mN} r_{ON} (2g_{mP} + C_E s) r_{OP}}{2r_{OP} r_{ON} C_E C_L s^2 + [(2r_{ON} + r_{OP}) C_E + r_{OP} (1 + 2g_{mP} r_{ON}) C_L] s + 2g_{mP} (r_{ON} + r_{OP})}$$

$$\omega_{p1} \approx \frac{1}{(r_{ON} || r_{OP}) C_L}$$

$$\omega_{p2} \approx \frac{g_{mP}}{C_E}$$

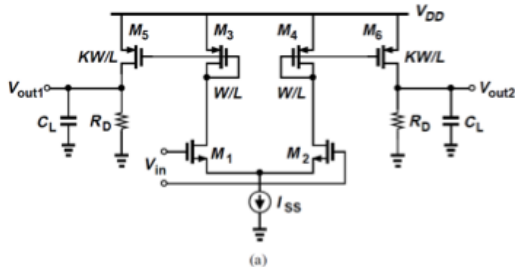
# Active Load



- A zero with a magnitude of  $2g_{mP}/C_E$  in the left half plane.
- The appearance of such a zero can be understood by noting that the circuit consists of a “slow path” (M1, M3 and M4) in parallel with a “fast path” (M1 and M2) by

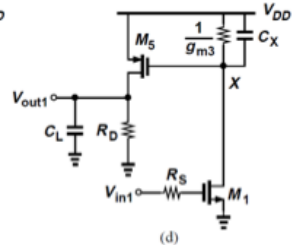
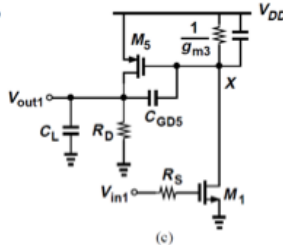
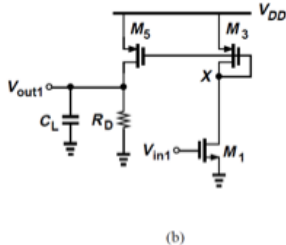
$$\begin{aligned}
 \frac{V_{out}}{V_{in}} &= \frac{A_0}{1 + s/\omega_{p1}} \left( \frac{1}{1 + s/\omega_{p2}} + 1 \right) \\
 &= \frac{A_0(2 + s/\omega_{p2})}{(1 + s/\omega_{p1})(1 + s/\omega_{p2})}
 \end{aligned}$$

# Example



- Not all fully differential circuits are free from mirror poles.
- Estimate the low-frequency gain and the transfer function of this circuit.

# Example



- M5 multiplies the drain current of M3 by K.

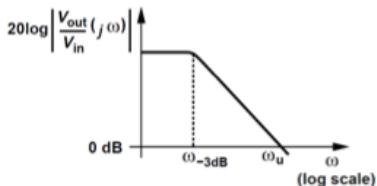
- Assume  $R_{DCL}$  is relatively small so that the Miller multiplication of  $C_{GD5}$  can be approximated as

$$C_{GD5}(1 + g_{m5}R_D)$$

- The overall transfer function is equal to  $V_X/V_{in1}$  multiplied by  $V_{out1}/V_X$ .

$$\frac{V_{out1}}{V_X}(s) = -g_{m5}R_D \frac{1}{1 + R_D C_{Ls}}$$

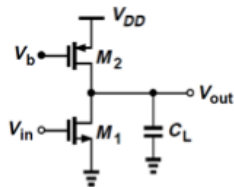
# Gain-Bandwidth Trade-Offs



- We wish to maximize both the gain and the bandwidth of amplifiers.
- we are interested in both the 3-dB bandwidth,  $\omega_{-3dB}$  and the “unity-gain” bandwidth,  $\omega_u$ .



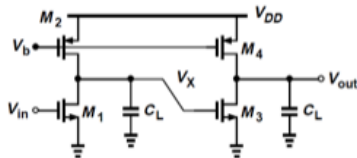
# One pole circuit



$$\begin{aligned} GBW &= A_0 \omega_p \\ &= g_{m1}(r_{O1} || r_{O2}) \frac{1}{2\pi(r_{O1} || r_{O2})C_L} \\ &= \frac{g_{m1}}{2\pi C_L}. \end{aligned}$$

$$\begin{aligned} \omega_u &= \sqrt{A_0^2 - 1} \omega_p \\ &\approx A_0 \omega_p \end{aligned}$$

# Multi-Pole Circuits



- It is possible to increase the GBW product by cascading two or more gain stages.
- Assume the two stages are identical and neglect other capacitances.

$$\frac{V_{out}}{V_{in}} = \frac{A_0^2}{\left(1 + \frac{s}{\omega_p}\right)^2}$$

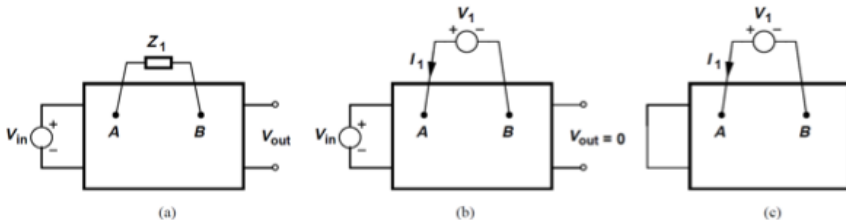
$$\omega_{-3-dB} = \sqrt{\sqrt{2} - 1} \omega_p$$

$$\approx 0.64 \omega_p$$

$$GBW = \sqrt{\sqrt{2} - 1} A_0^2 \omega_p$$

- While raising the GBW product, cascading reduces the bandwidth.

# Appendix A: Extra Element Theorem

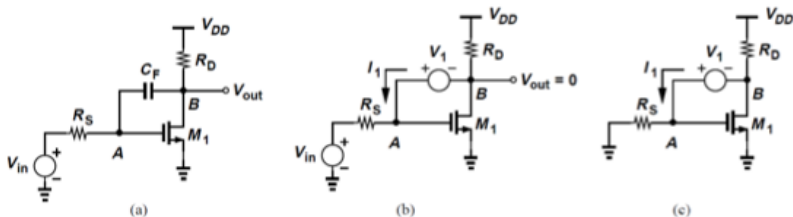


- Suppose the transfer function of a circuit is known and denoted by  $H(s)$ . Add an extra impedance  $Z_1$  between two nodes of the circuit.
- New transfer function:

$$G(s) = H(s) \frac{1 + \frac{Z_{out,0}}{Z_1}}{1 + \frac{Z_{in,0}}{Z_1}}$$

- Particularly useful for frequency response analysis.

# Example 1



- Find the transfer function.

$$H(s) = -g_m(R_D || r_o)$$

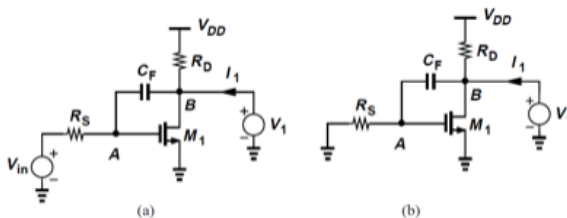
$$Z_{out,0} = -1/g_m$$

- The negative sign of  $Z_{out,0}$  does not imply a negative impedance between A and B, since  $V_{in} \neq 0$

$$Z_{in,0} = (1 + g_m R_S)R_D + R_S = (1 + g_m R_D)R_S + R_D$$

$$G(s) = -g_m(R_D || r_o) \frac{1 - \frac{1}{g_m} C_F s}{1 + [(1 + g_m R_D)R_S + R_D] C_F s}$$

## Example 2



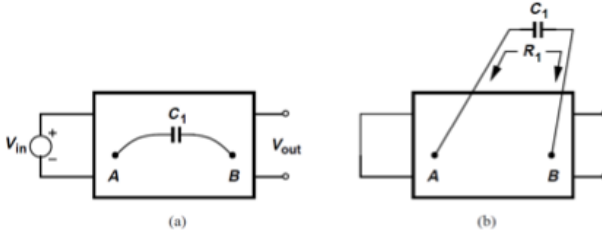
- Include CB, from node B to ground.

$$Z_{out,0} = 0$$

$$Z_{in,0} = \frac{R_D(R_S C_F s + 1)}{R_S(1 + g_m R_D) + R_D C_F s + 1}$$

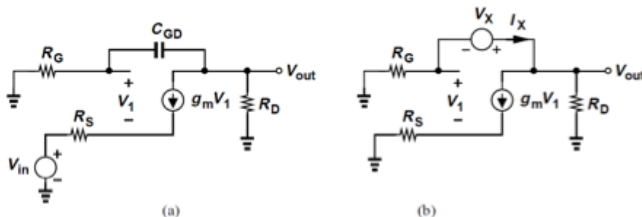
$$G(s) = -g_m(R_D || r_o) \frac{1 - \frac{C_F}{g_m} s}{1 + [(1 + g_m R_D)R_S + R_D]C_F s} \frac{1}{1 + \frac{R_D(R_S C_F s + 1)C_{B s}}{[R_S(1 + g_m R_D) + R_D]C_F s + 1}}$$

# Appendix B: Zero-Value Time Constant



- Suppose a circuit contains one capacitor and no other storage elements.
- Wish to determine the pole of the system.
- Set the input to zero, compute the resistance,  $R_1$ , seen by  $C_1$ , and express the pole as  $1/(R_1C_1)$ .

# Example



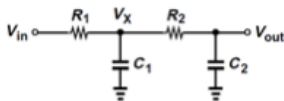
- If only  $C_{GD}$  is considered, determine the pole frequency.

$$g_m V_1 R_S + V_1 = -I_X R_G$$

$$\frac{V_X}{I_X} = R_D + \left( \frac{g_m R_D}{1 + g_m R_S} + 1 \right) R_G = R_{eq}$$

- The pole is given by  $1/(R_{eq} C_{GD})$

# Example



- Writing a KVL around  $V_{in}$ ,  $R_1$ ,  $R_2$ , and  $V_{out}$  gives

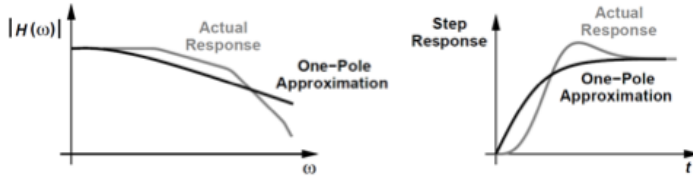
$$\frac{V_{out}}{V_{in}}(s) = \frac{1}{R_1 R_2 C_1 C_2 s^2 + [R_1 C_1 + (R_1 + R_2) C_2] s + 1}$$

- The dominant pole is indeed equal to the inverse of the sum of the zero-value time constants.  
(need to prove)

- The  $B_s$  method proves useful if we wish to estimate the 3-dB bandwidth of a circuit.

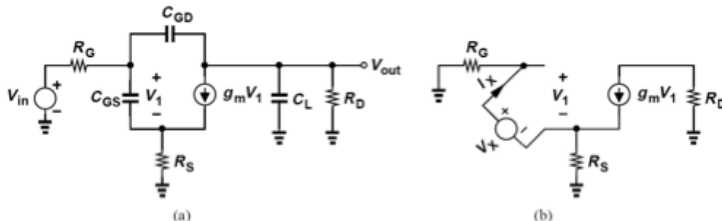


# Example



- **Approximation of the frequency and time responses by one-pole counterparts.**

# Example 1



- Estimate the 3-dB bandwidth.
- Begin with the time constant associated with  $C_{GS}$  and set  $C_{GD}$  and  $C_L$  to zero.

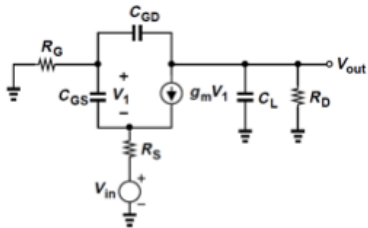
$$R_{CGS} = \frac{R_G + R_S}{1 + g_m R_S}$$

$$R_{CGD} = R_D + \left( \frac{g_m R_D}{1 + g_m R_S} + 1 \right) R_G$$

- The resistance seen by  $C_L$  is simply equal to  $R_D$ .

$$\omega_{-3dB}^{-1} = \frac{R_G + R_S}{1 + g_m R_S} C_{GS} + \left[ R_D + \left( \frac{g_m R_D}{1 + g_m R_S} + 1 \right) R_G \right] C_{GD} + R_D C_L$$

## Example2



- Find 3-dB band width of a common-gate stage containing a gate resistance of  $R_G$  and a source resistance of  $R_S$ .
- The resulting equivalent circuits are identical for CS and CG stages, yielding the same time constants and hence the same bandwidth.
- Does this result contradict our earlier assertion that the CG stage is free from the Miller effect?
- No. Why?