

CLASS

Cosmological Linear Anisotropy Solving System

Markus Mosbech
Institute for Theoretical Particle Physics and Cosmology, RWTH Aachen
University

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Visit <http://class-code.net/> for more info!

class in Les Karellis

What to expect in this *advanced* lecture:

- Theory: What is **class** based upon?
- Coding: Structure of **class**
- Coding: Essential rules and conventions
- Coding: Implementing features (C and python)
- Coding: Using MontePython/Cobaya with **class**

We will learn **the theory behind class** and the fundamental rules of its **code base**.

- 1 Fundamental layout of Einstein-Boltzmann solvers
- 2 Essential steps for each module
- 3 A few details for each of these steps

Fundamental layout of Einstein-Boltzmann solvers

Homogeneous background

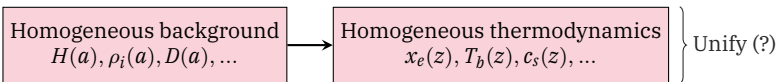
$H(a), \rho_i(a), D(a), \dots$



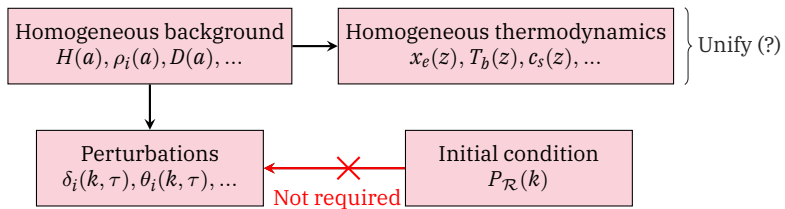
Homogeneous thermodynamics

$x_e(z), T_b(z), c_s(z), \dots$

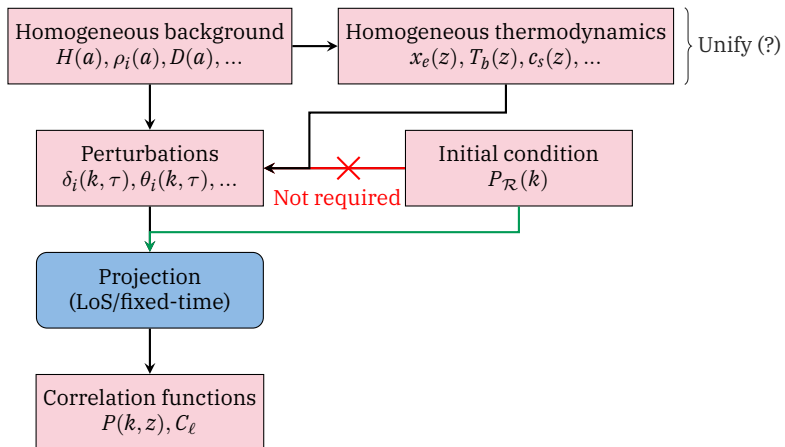
Fundamental layout of Einstein-Boltzmann solvers



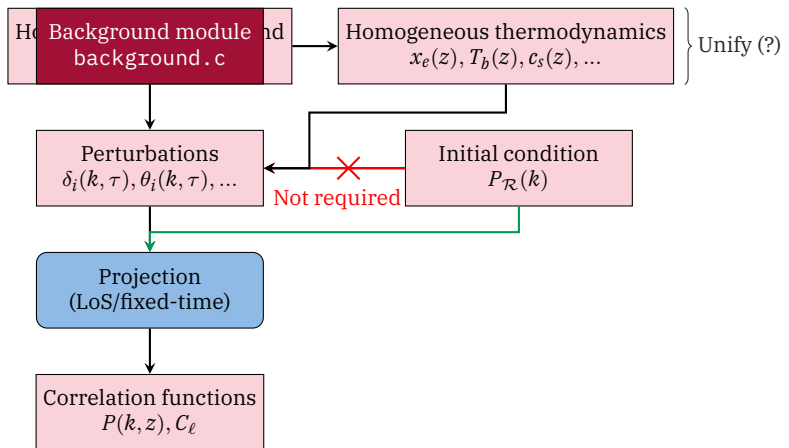
Fundamental layout of Einstein-Boltzmann solvers



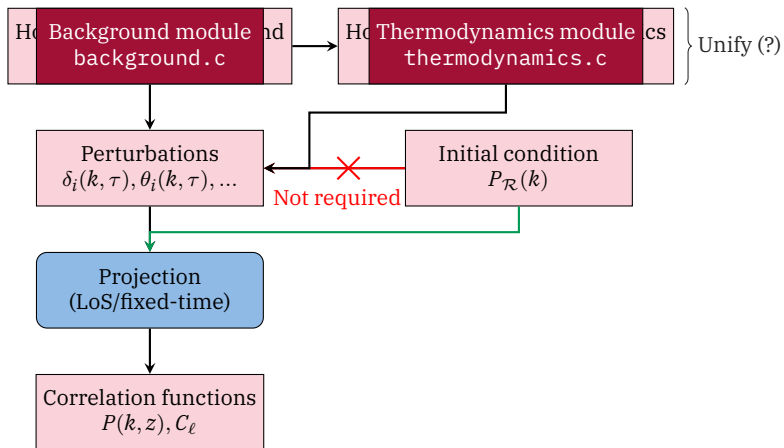
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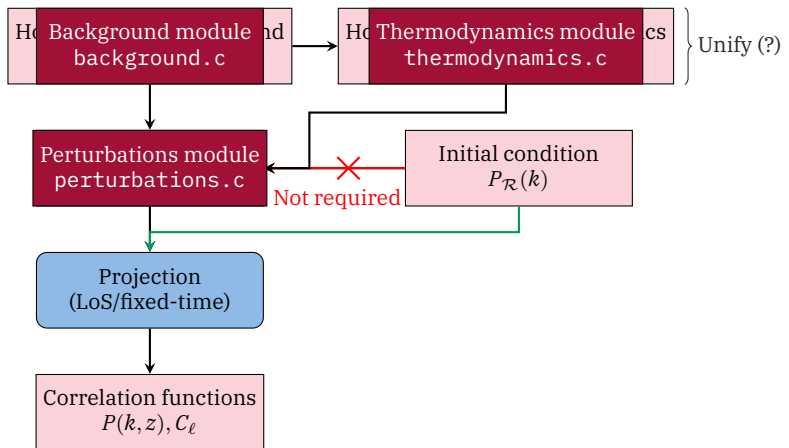
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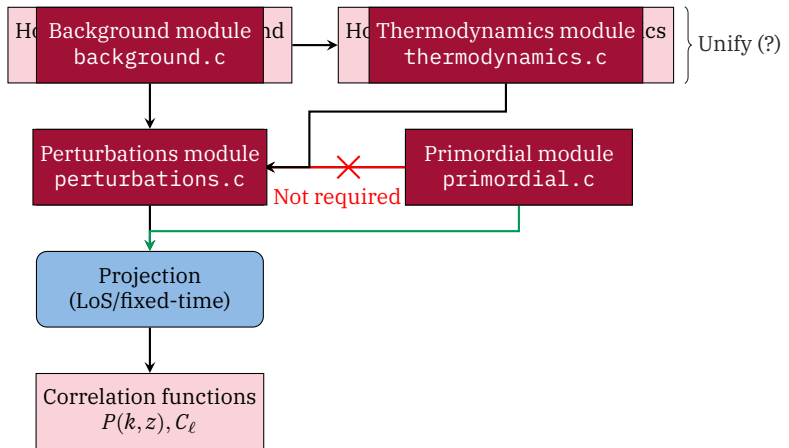
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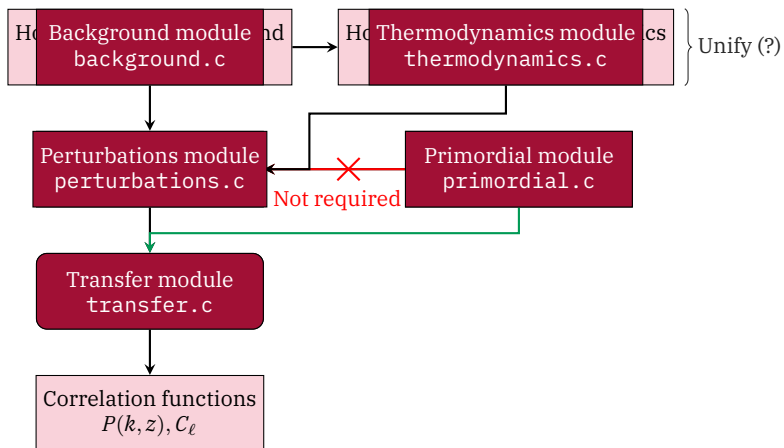
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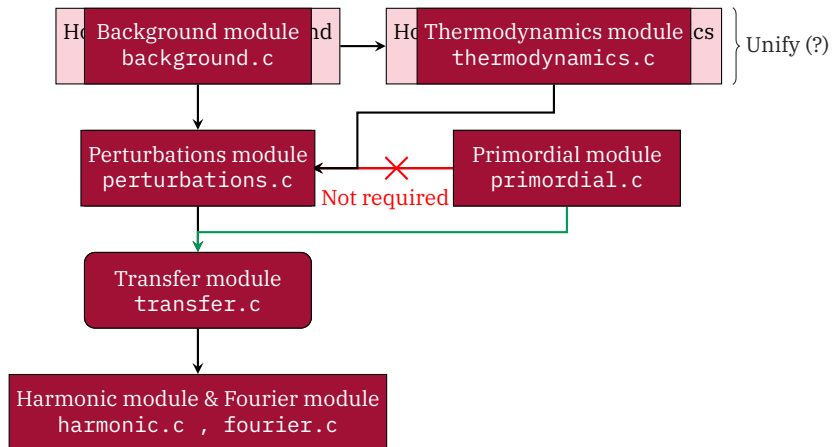
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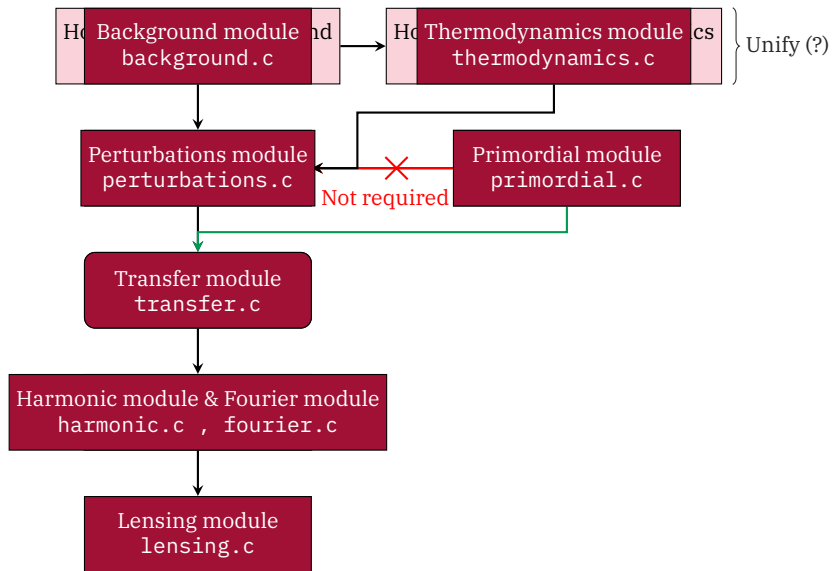
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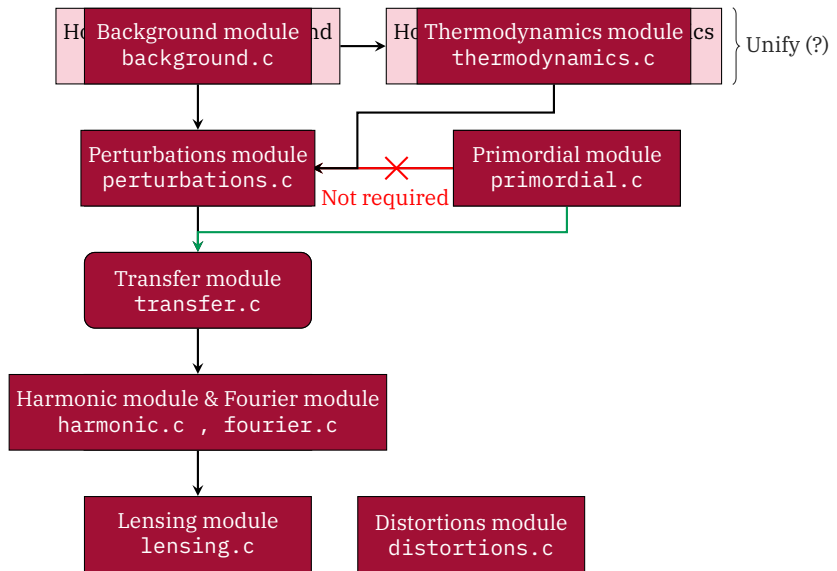
Fundamental layout of Einstein-Boltzmann solvers



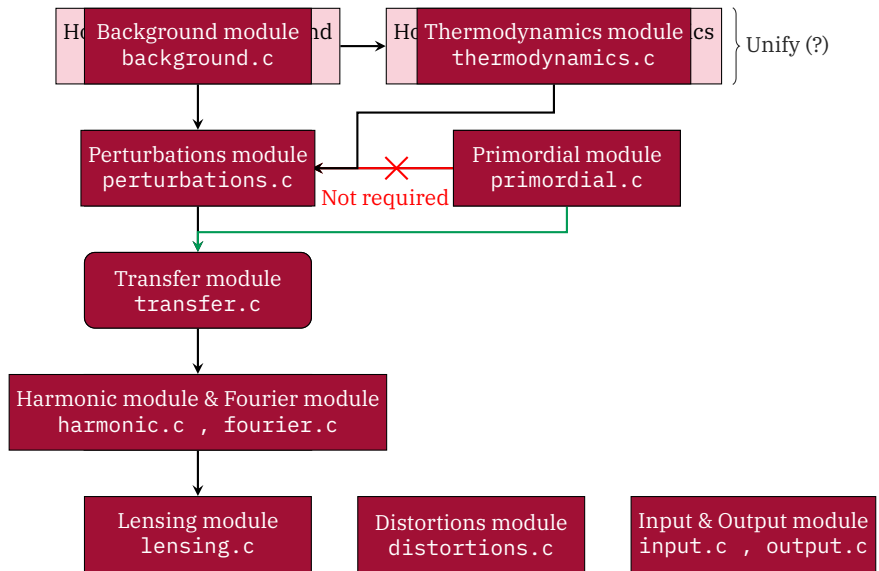
Fundamental layout of Einstein-Boltzmann solvers



Fundamental layout of Einstein-Boltzmann solvers



Fundamental layout of Einstein-Boltzmann solvers



Let's make a journey through each module!

Essential steps in Einstein-Boltzmann solver

Module 1. Input

Read in input files, take care of *shooting*.

```
h = 0.7  
#H0 = 70  
Omega_m = 0.3  
#omega_m = 0.14  
sigma8=0.8
```

Special care for equivalent/unknown parameters

Input management in `class`

For indirect parameters, use `shooting method`

Repeated calls of `input_read_parameters(...)`, `class` executions, from `input_read_from_file(...)` until shooting target is met.

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Try out a few values and narrow down (Example: User wants $100\theta_s = 1.04325$)

h	$100\theta_s$
0.7	1.0522492086422521
0.65	1.0270326366580724
0.68215616173	1.0437999980620178
0.68110138476	1.0432819283581667
0.68103637942	1.0432499363679562
0.68103650871	1.0432499921072458
0.68103652701	1.0432500079710365
...	...

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0.68103637942	1.0432499363679562
0.68103650871	1.0432499921072458
0.68103652701	1.0432500079710365
...	...

In practice, use more sophisticated Ridder's method / Newton's method

Input management in **class**

For **shooting** parameters, establish mapping between *target parameter*, *unknown parameter* and *level*. Currently:

target parameter	unknown parameter	level
$100 \times \theta_s$	h	thermodynamics
Ω_{dcdm}	$\rho_{\text{dcdm}}^{\text{ini}}$	background
σ_8	A_s	spectra
...

... plus a few others (alternative parametrizations of decaying CDM, quintessence parameters).

This is what is used e.g. in models of early dark energy!

If you need to add such parameters: see how it is done e.g. for `100*theta_s` and replicate the structure!

Special exception $\tau_{\text{reio}} \leftrightarrow z_{\text{reio}}$ only concerns reionization and is done independently in `thermodynamics.c`

New Special exception: $\sigma_8 \leftrightarrow A_s$ can be very simply analytically re-scaled (multiplicative property), therefore done independently in `input.c`

Input management in `class`

Budget equation:

$$\sum_X \Omega_X = 1 + \Omega_k$$

To avoid over-constraining the input, one of the last three (`Omega_Lambda`, `Omega_fld`, `Omega_scf`) must be left unspecified and `class` will assign it using budget equation.

Possibly more advanced in the future

- default: `Omega_Lambda` is automatically adjusted
- if you pass `Omega_Lambda`, `Omega_fld` is automatically adjusted
- if you pass `Omega_Lambda` and `Omega_fld`: `Omega_scf` is automatically adjusted (if you allow, by setting to -1)

This allows whatever combination.

E.g. to get Λ plus a DE fluid:

`Omega_Lambda=0.2, Omega_scf=0` or `Omega_fld=0.3, Omega_scf=0`

Helpful output by setting background verbose `>= 2`

Essential steps in Einstein-Boltzmann solver

Module 2. Background

Get all background quantities as function of a scale factor a .

This also gives mapping $a \leftrightarrow z \leftrightarrow t \leftrightarrow \text{conf.time}$

Details of the steps in Einstein-Boltzmann solvers

Let's formalize problem!

Three types of parameters:

- $\{A\}$ are analytical functions of scale factor and $\{B\}$ quantities.
- $\{B\}$ need to be integrated over, and are used to compute $\{A\}$
- $\{C\}$ also need to be integrated over, but are not used to compute $\{A\}$.

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Λ CDM and many simple extensions:

- $\{A\} = \{\rho_i(a), p_i(a), H(a), \dots\}$ with e.g. $H(a) = \left(\sum_X \rho_X(a) - \frac{K}{a^2}\right)^{1/2}$
- $\{B\} = \{\}$ (eliminated since v3.0)
- $\{C\} = \{t, \tau, r_s, D, f\}$ with e.g. $\frac{dt}{da} = 1/H(a)$, $\frac{dr_s}{da} = c_s(a)/(a \cdot H(a))$

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Example of DE/DM/DR fluid:

- $\{A\} = \{\rho_i(a), p_i(a), H(a), \dots, w_{\text{fld}}(a)\}$
- $\{B\} = \{\rho_{\text{fld}}\}$ with $\frac{d\rho_{\text{fld}}}{da} = -3(1 + w_{\text{fld}}(a))\rho_{\text{fld}}$

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Exemple of extended cosmology with quintessence ϕ :

- $\{A\} = \{\rho_i, p_i, H, \dots, V(\phi), \rho_\phi(\phi, \phi')\}$ with e.g. $\rho_\phi(\phi, \phi') = \frac{1}{2}(\phi')^2 + V(\phi)$
- $\{B\} = \{\phi, \phi'\}$ with $\frac{d\phi}{da} = \phi' / [aH(a)]$, $\frac{d\phi'}{da} = -2\phi' - aV(\phi)/H(a)$

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Also Cold Dark Matter decaying into Dark Radiation...

- $\{A\} = \{\rho_i, p_i, H, \dots\}$
- $\{B\} = \{\rho_{\text{cdm}}, \rho_{\text{dr}}\}$ with $\frac{d\rho_{\text{cdm}}}{da} = -3\rho_{\text{cdm}} - \Gamma(a)/H(a) \cdot \rho_{\text{cdm}}$

Details of the steps in Einstein-Boltzmann solvers

Small details:

- Quantities as $D_A(z)$, $D_L(z)$, r_s , t_{age} can be derived after all A,B,C are computed
- Takes care of Λ CDM integration of phase-space distribution
- Useful checks & output
- → Budget equation output at verbosity level 2

Essential steps in Einstein-Boltzmann solver

Module 3. Thermodynamics

Get all thermodynamics quantities as function of a time variable (**class** → redshift z) after integrating differential equations like recombination equations:

$$\frac{dx_e}{dz}, \frac{dT_b}{dz} = \text{excitation, ionization, heating, ...}$$

Then $x_e(z) \rightarrow \kappa'(z)$ (Thomson scattering rate)

- $\kappa(z)$ (Optical depth)
- $\exp(-\kappa(z))$ (factor for Integrated Sachs-Wolfe effect)
- $g(z)$ (visibility function for Sachs-Wolfe effect)
- $g'(z)$ (factor for Doppler effect)

Details of the steps in Einstein-Boltzmann solvers

Simplest model of **recombination** is the **Saha equation**.

It is well known that a non-relativistic ($T \ll m$) species in thermal equilibrium obeys

$$n(\mu, T) \approx g e^{\mu/T} \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T} \quad (1)$$

Thus we find using **complete thermal equilibrium** with $\mu_{\text{ionized}} + \mu_e = \mu_{\text{rec}}$ that

$$\frac{n_e n_{\text{ionized}}}{n_{\text{rec}}} \approx \left(\frac{m_e T}{2\pi} \right)^{3/2} e^{-E_{\text{bind}}/T} \times \underbrace{\left[e^{\mu_{\text{ionized}} + \mu_e - \mu_{\text{rec}}} \left(\frac{g_e g_{\text{ionized}}}{g_{\text{rec}}} \right) \left(\frac{m_{\text{ionized}}}{m_{\text{rec}}} \right)^{3/2} \right]}_{\approx 1}$$

This gives

$$\frac{x_e^2}{1 - x_e} \approx \left(\frac{1.1 \cdot 10^{-10}}{n_{\text{H},0}/T_{\text{cmb},0}^3} \right) \left(\frac{\text{eV}}{T} \right)^{3/2} \exp(39.9 - 13.6 \frac{\text{eV}}{T}) \quad (2)$$

and thus recombination at $T \approx \frac{13.6\text{eV}}{39.9} \approx 0.34\text{eV} \rightarrow z \approx 1400$.

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recombination is a non-equilibrium process

Details of the steps in Einstein-Boltzmann solvers

The **effective multi-level atom** is the basis for recombination codes.

1s 2s 2p 3s 3p 3d ... ionized \rightarrow 1s 2s 2p ionized

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Reason: Intermediate transitions ($4p \rightarrow 3s$) or ($3s \rightarrow 2p$) are comparatively instant. Why? Direct transition $2s \rightarrow 1s$ is forbidden, and $2p \rightarrow 1s$ is immediately reversed by $1s \rightarrow 2p$. The medium is **optically thick** during recombination.

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Instead, focus on $2p \rightarrow 1s$ with subsequent redshifting of photon to escape reabsorption (slow) or $2s \rightarrow 1s$ with two-photon decay (slow).

Peeble's equation

$$\dot{x}_e \approx f_{\text{photo-ion}}(T)x_{\text{rec}} - f_{\text{rec}}(T)x_e x_{\text{ionized}} \quad (3)$$

Solved numerically, basis of **recfast**

Details of the steps in Einstein-Boltzmann solvers

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Fulllest code to date: **CosmoRec** does full numerical computation (iteratively).
Comparatively slow, but highest achievable accuracy

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Further complication: Helium (higher elements don't contribute)

Details of the steps in Einstein-Boltzmann solvers

User can choose to model approximate recombination and get $x_e(z)$, $T_b(z)$ from:

- [RECFAST](#) (Wong, Moss & Scott 2008)
- [HyRec-2](#) (Y. Ali-Haïmoud, N. Lee)
- Possibly soon? [CosmoRec](#) (J. Chluba)

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Recombination needs one more cosmological parameter: the **primordial Helium fraction** Y_{He} .

- Fix it ($Y_{\text{He}} = 0.25$)
- Get it from BBN ($Y_{\text{He}} = \text{BBN}$). **class** has interpolation table pre-pcomputed with a **BBN code** (Parthenope), for each given value of N_{eff} , ω_b (assumes $\mu_{\nu_e} = 0$, easy to generalize).
- BBN interpolation table located in separate directory (in `external/bbn/sBBN_2017.dat`, update inbound)

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For reionization:

- tanh with complicated argument (like **CAMB**)
- multi-tanh
- half tanh
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Mini-shooting to find z_{reio} for given $\tau_{\text{reio}} = \kappa_{\text{reio}}$.

Optical depth $\kappa(z)$ = inverse number of expected interactions $\Rightarrow \kappa'(z) = an_H x_e \sigma_T$

Details of the steps in Einstein-Boltzmann solvers

We also include

- Energy injection (increases ionization, heats T_b)
This can cause changes in scattering $\kappa(z)$ and thus be observable with CMB
- Time-dependent fundamental constants \rightarrow Causes shift in recombination due to fundamental dependencies such as
$$E_{\text{binding}} = \frac{1}{2}\alpha^2 m_e = 13.6\text{eV} (137\alpha)^2 \left(\frac{m_e}{511\text{keV}}\right)$$
We remind ourselves $1 + z_{\text{rec}} = T_{\text{rec}}/T_{\text{cmb}} \approx \frac{E_{\text{binding}}}{12.57\text{meV}}$
- Computation of useful quantities $z_{\text{rec}}, z_{\text{drag}}, z_*, D_A(z_{\text{rec}}), r_s(z_{\text{drag}}), \dots$