

DEEPER INSIGHTS  
UNVEILING DOPANT PROFILES WITH GIXRD

by

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DEPARTMENT APPROVAL

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This thesis has been reviewed by the research advisor, research coordinator, and department chair and has been found to be satisfactory.

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## ABSTRACT

## DEEPER INSIGHTS

### UNVEILING DOPANT PROFILES WITH GIXRD

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In this work, grazing incidence X-ray diffraction (GIXRD) is utilized to investigate the depth profile of dopants in silicon. By systematically varying the X-ray incidence angle, one can precisely control the depth penetrated into doped silicon samples. This approach enables a non-destructive analysis of how dopant concentration evolves with depth. The findings provide crucial insights into the distribution and behavior of dopants, which is essential for optimizing semiconductor device performance and fabrication processes.

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# Chapter 1

## Background

Solids can be divided into two broad types: crystalline, and amorphous. Crystalline solids have repeating structure that is predictable, and is consistent throughout the entire solid. Amorphous solids have no repeating structures, and have no preferred orientation.

X-Ray diffraction (XRD) utilizes x-rays to non-destructively probe the inner atomic layers of materials. In this work it is used to probe different depths of doped silicon to see dopant levels as a function of depth. As the x-ray interacts with the sample it will reflect, refract, and diffract. Diffraction is the process of note, and the only one that will be investigated for this work.

When XRD is used on crystalline solids a diffraction pattern will emerge. The spacing between the atoms acts as the slit size that allows for the diffraction interference pattern, in accordance with Bragg's Law ( $2d \sin \theta = n\lambda$ ). There also exist different spacings between atoms that are not nearest neighbors. Each of these different spacings are relegated into planes that act like a diffraction grating, existing for every possible repeating spacing within the atom. All of these in superposition produce a graph that shows sharp peaks correlating directly to the atomic plane's

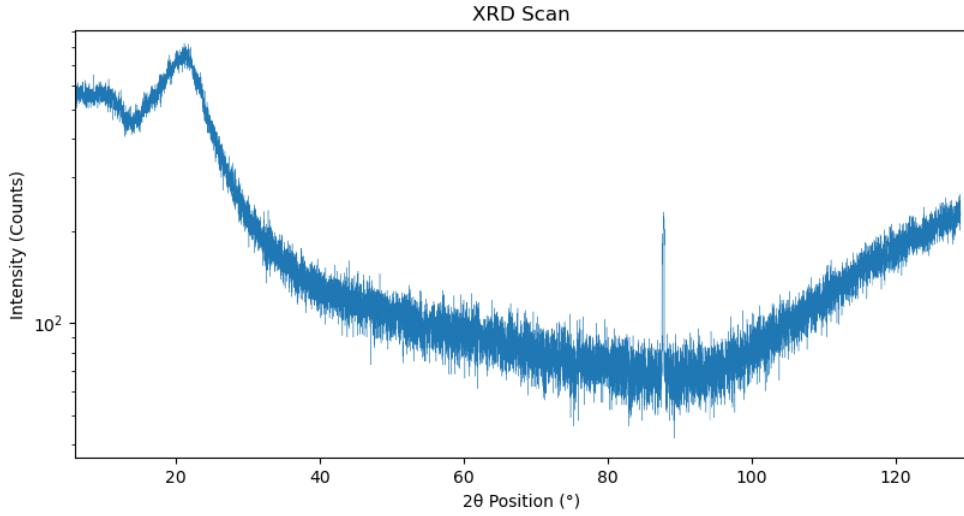
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spacing (shown in figure 1.1).



**Figure 1.1** Diffractogram of pure crystalline solid

Amorphous solids under the x-ray beam behave in a similar manner. The difference being because every atom is arranged randomly, there are an infinite number of planes with their own spacing. As shown in figure 1.2, this correlates to a moderately low angle hump (called an amorphous hump) with all spacings that would correlate to higher angled peaks being nonexistent because of deconstructive interference[1–3].



**Figure 1.2** Diffractogram of pure amorphous solid

The materials analyzed in this work are silicon doped with phosphorous, silicon doped with boron, and pure silicon. The first two materials will be best thought of as a gradient from pure phosphorous or boron to pure silicon from top to bottom. Silicon is a crystalline material, while phosphorous and boron are amorphous. This means that the diffractogram will be a superposition of an amorphous diffractogram on a crystalline diffractogram. The silicon sample was only used to verify that the amorphous hump was present on the diffractograms.

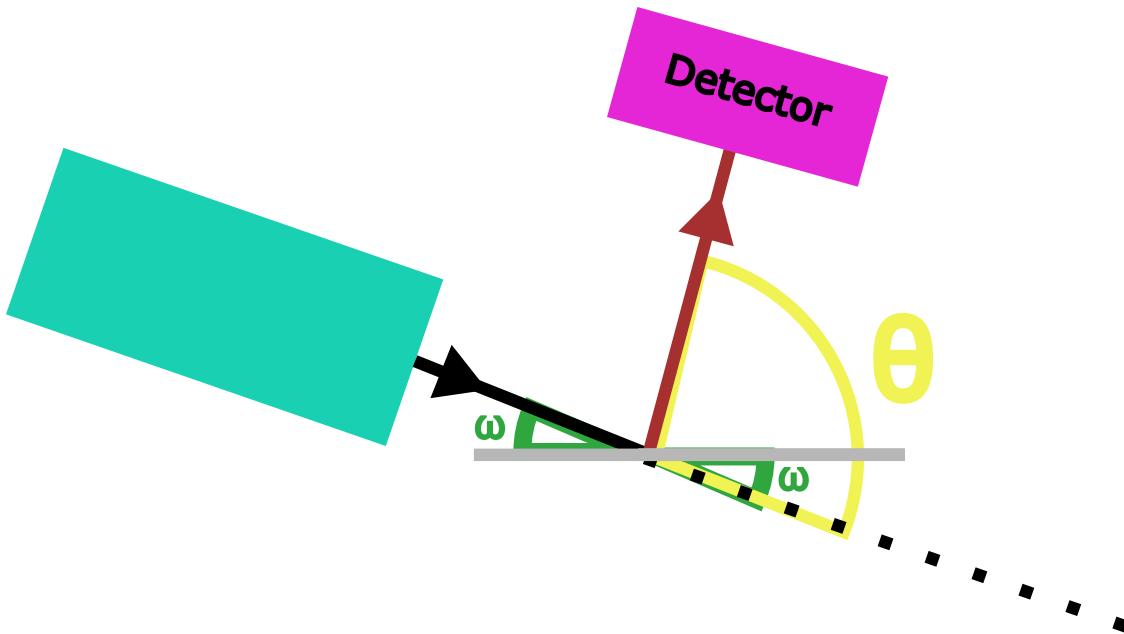
The different incident angles probe different depths due to the different path lengths they provide the x-rays as they interact with the samples. This is due to the exponential attenuation of intensity as the beam travels given by this equation:  $I(y) = I_0 e^{-\alpha y}$  with  $\alpha = 2\omega n_I / c$  [4, 5]. With the complex index of refraction ( $n_I$ ) strongly dependent on wavelength. The x-ray source is copper, producing  $K_\alpha$ , and  $K_\beta$  in most abundance. Since one wavelength provides the most uniformity on the diffractogram, the incident beam optics include a filter to limit the unnecessary wavelength.  $K_\alpha$  is what is most common, so the  $K_\beta$  gets mostly filtered out by a nickel filter. This

results in a x-ray beam that has a wavelength of 1.54Å(uniformity in the wavelength is also necessary for Bragg's law).

# Chapter 2

## Methods

The process by which the XRD is used in this work to probe increasingly deeper depths is called grazing incidence x-ray diffraction (GIXRD; also called glancing incidence x-ray diffraction). GIXRD holds an incident angle relative to the surface of the sample (called omega ( $\omega$ )), while the detector on the diffracted beam side sweeps over a range of angles (figure 2.1). The detector is used in 0D mode in conjunction with a parallel plate collimator; This means it simply counts all the x-rays that interact with it. The counts get added to bins corresponding to the current detector angle according to this equation:  $2(\omega + \theta)$ , with  $\theta$  being the current detector angle. After the detector has swept through all programmed angles, a new scan is started with a different  $\omega$  corresponding to a different depth.



**Figure 2.1** Diagram of XRD utilizing GIXRD

Three different methods for determining relative percentages of the dopant in the substrate are used. The methods are named: 1. Intensity, 2. Curve Fitting, and 3. Summation. Before discussing the intricacies of these methods, it is important to first know what shared simplifying assumptions on which they are predicated. The first of these assumptions is that the mixture of the probed depth is uniform. Uniformity of the mixture cannot be achieved in the sample due to the diffusion process starting at the top of the sample working its way down. However, this assumption is needed as the basis for future work to use as a starting point. The next assumption is that the intensity of the x-ray has been attenuated by 90%. The thirdly, that the sample is put in its preferred orientation that allows for the silicon structure to be aligned with the x-ray beam. The last assumption is that the index of refraction is the same as bulk samples[6] as those values will be utilized.

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### Method 1: Intensity

This method takes the highest count of the amorphous hump, and the highest count of the crystalline peaks as input parameters, to take the ratio of them. This ratio is then the relative percentages of the dopant and substrate[7]. The following formula is used for the calculation:

$$m = \frac{I_a}{I_a + I_c}$$

$I_a$  is the highest count of the amorphous hump,  $I_c$  is the highest count of the crystalline peaks, and  $m$  is the mixture percent.

### Method 2: Curve Fitting

The amorphous hump, and the crystalline peaks can be approximated by Gaussian curves. As such, it is possible to fit a curve to the hump, and the desired peak. With this fit, analytic integration is performed. The values from the integrations are then taken as a ratio as in the intensity method. This method also gives a value to the background noise that can then be subtracted out.

### Method 3: Summation

This method is similar to the curve fitting method, but instead of fitting the hump, and peak to a Gaussian, a numerical integration is performed. This numerical integration sums up the heights in the region of the relevant area to take their ratio. As stated, only the heights are summed; This is due to the spacing between all points being identical. Below is the proof:

$$m = \frac{\sum_{\theta_i}^{\theta_f} y_a dx_a}{\sum_{\theta_i}^{\theta_f} y_a dx_a + \sum_{\theta_i}^{\theta_f} y_c dx_c}$$

$$dx_a = dx_c = Constant$$

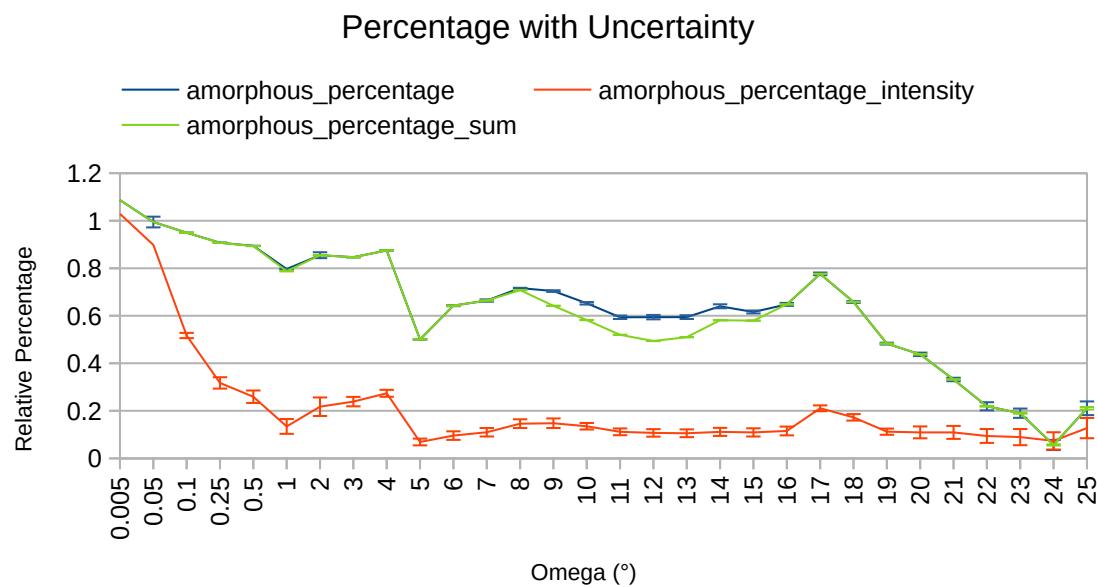
$$\begin{aligned}
m &= \frac{dx_a \sum_{\theta_i}^{\theta_f} y_a}{dx_a \sum_{\theta_i}^{\theta_f} y_a + dx_c \sum_{\theta_i}^{\theta_f} y_c} \\
m &= \frac{dx_a \sum_{\theta_i}^{\theta_f} y_a}{dx_a \left( \sum_{\theta_i}^{\theta_f} y_a + \sum_{\theta_i}^{\theta_f} y_c \right)} \\
m &= \frac{\sum_{\theta_i}^{\theta_f} y_a}{\sum_{\theta_i}^{\theta_f} y_a + \sum_{\theta_i}^{\theta_f} y_c}
\end{aligned}$$

$y_a$ , and  $y_c$  is the count of the current *theta* position for the amorphous and crystalline parts respectively.  $dx_a$ , and  $dx_c$  is the spacing between the *theta* positions for the amorphous and crystalline parts respectively.  $m$  is the mixture amount, and  $\theta_{tf}$ , and  $\theta_{ti}$  are the starting and ending *theta* positions for the amorphous hump or crystalline peak respectively.

A similar proof can also be performed to show the counts do not need to be normalized.

# Chapter 3

## Analysis

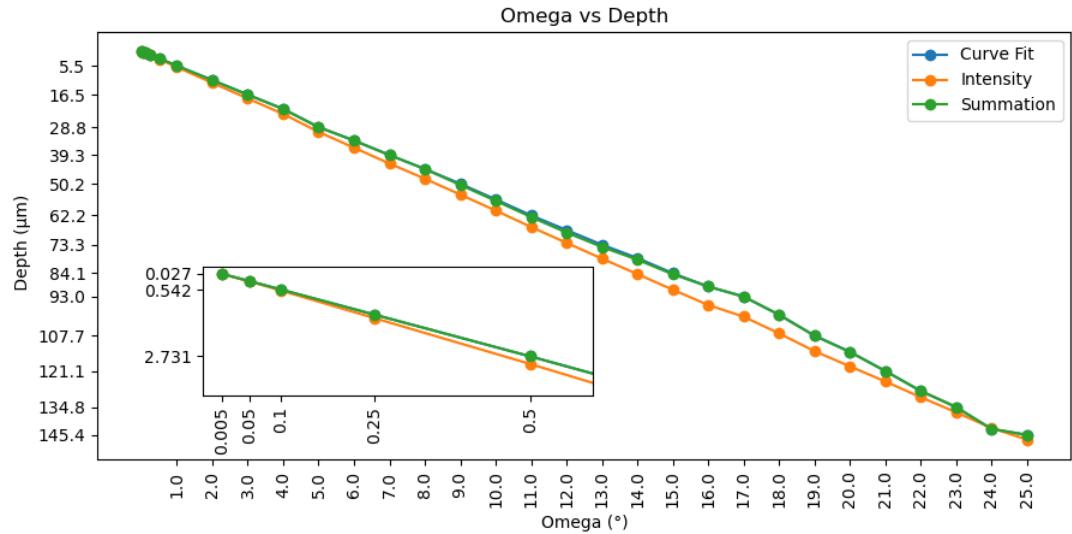


**Figure 3.1** Percentage of amorphous vs omega for the three methods with fractional uncertainty

omega	Intensity	Intensity Fractional Uncertainty	Summation	Summation Fractional Uncertainty
0.005	1.0291		1.0869	
0.05	0.8983	0.0110	0.9946	0.0001
0.1	0.5168	0.0239	0.9495	0.0003
0.25	0.3174	0.0263	0.9070	0.0004
0.5	0.2593	0.0313	0.8936	0.0004
1	0.1344	0.0389	0.7877	0.0008
2	0.2172	0.0197	0.8552	0.0003
3	0.2387	0.0146	0.8467	0.0002
4	0.2737	0.0142	0.8752	0.0002
5	0.0690	0.0181	0.5004	0.0005
6	0.0958	0.0179	0.6426	0.0004
7	0.1098	0.0183	0.6639	0.0004
8	0.1459	0.0203	0.7082	0.0004
9	0.1478	0.0138	0.6417	0.0003
10	0.1351	0.0142	0.5820	0.0004
11	0.1115	0.0157	0.5196	0.0004
12	0.1071	0.0165	0.4938	0.0005
13	0.1056	0.0169	0.5099	0.0005
14	0.1113	0.0173	0.5812	0.0004
15	0.1091	0.0185	0.5794	0.0004
16	0.1153	0.0122	0.6478	0.0003
17	0.2106	0.0136	0.7765	0.0002
18	0.1726	0.0131	0.6580	0.0003
19	0.1123	0.0249	0.4827	0.0008
20	0.1093	0.0273	0.4377	0.0009
21	0.1093	0.0292	0.3320	0.0014
22	0.0942	0.0341	0.2191	0.0021
23	0.0894	0.0364	0.1900	0.0024
24	0.0735	0.0427	0.0564	0.0053
25	0.1273	0.0333	0.2107	0.0025

**Figure 3.2** Table for percentage of amorphous vs omega for the three methods with fractional uncertainty

From figure 3.1 we can see that the intensity method is not in agreement with the curve fit or summation methods. The curve fit, and summation methods are expected to be in agreement with each other since the methods are both integral methods. From the table in figure 3.2 we can see, with more precision, that the uncertainty is lower as well for the summation method over the intensity method. The 0.005° shows that that depth is composed of a value greater than unity for the sample; This is physically impossible. The greater than unity error comes about from having no crystalline peak from the silicon, so the amorphous hump from the phosphorous is the only thing present on the diffractogram (this is the diffractogram in figure 1.2).



**Figure 3.3** Depth vs omega for the three methods

In Figure 3.3, we can see the depth that correlates to the three methods. The three methods have close agreement at lower depths before splitting between the intensity method and the curve fit and summation methods. The curve fit and summation method in particular are right on top of each other.

# **Chapter 4**

## **Conclusion**

Using GIXRD to determine relative composition of samples led to three analysis methods. Two of the methods utilized integration, and their results are more closely aligned. One method used peak intensity ratios, and provides results that noticeably differ. All three methods provide reasonable agreement as to the depth that each omega probed. Using GIXRD to see relative mixture amount dependent on depth has been shown to be useful, keeping in mind the limitations given by the assumptions.

### **4.1 Future Work**

Performing this analysis on a known sample would be able to provide greater validity to the results of this analysis. Creating an analysis to correct for the real diffusion process by not assuming uniform layers. This would produce the greatest impact on the real world application of these results.

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