

## Lecture 8:

### Monte Carlo Methods



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**Bayesian Data Analysis and  
Machine Learning for Physical  
Sciences**



## Course Map

Module 1	Maximum Entropy and Information, Bayes Theorem
Module 2	Naive Bayes, Bayesian Parameter Estimation, MAP
Module 3	MLE, Lin Regression
Module 4	Model selection I: Comparing Distributions
Module 5	Model Selection II: Bayesian Signal Detection
Module 6	Variational Bayes, Expectation Maximization
Module 7	Hidden Markov Models, Stochastic Processes
<b>Module 8</b>	<b>Monte Carlo Methods</b>
Module 9	Machine Learning Overview, Supervised Methods
Module 10	Unsupervised Methods
Module 11	ANN: Perceptron, Backpropagation
Module 12	ANN: Basic Architecture, Regression vs Classification, Backpropagation again
Module 13	Convolution and Image Classification and Segmentation
Module 14	Graphs and GNNs
Module 15	RNNs and LSTMs
Module 16	Transformer and LLMs



## CHECKING IN



## Outline

**Basic Idea & Finding Pi**

**Mapping Distributions & Gibbs Sampling**

**Gillespie Algorithm**

**Metropolis (- Hastings) Algorithm**

**Bootstrapping**





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**Basic Idea & Finding Pi**

Mapping Distributions & Gibbs Sampling

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Metropolis (- Hastings) Algorithm

Bootstrapping



**idea:** generating a set of values randomly  
i.e. repeated random sampling → **Monte Carlo** method

**pros:** for many sample repetitions → the actual probability density function emerges  
pretty **simple** set up & easy to implement, **easy to parallelize**

**cons:** not directed like e.g. gradient descent (see later)

**applications:** - numerical evaluation of complicated integrals

$$\rho = \frac{P(M_A|D, I)}{P(M_B|D, I)} = \frac{P(M_A)}{P(M_B)} \cdot \frac{\int P(D|\{\alpha\}_A, M_A, I) d\alpha_{Aj}}{\int P(D|\{\alpha\}_B, M_B, I) d\alpha_{Bj}} \cdot \frac{\prod_j \alpha_{jB}(max) - \alpha_{jB}(min)}{\prod_j \alpha_{jA}(max) - \alpha_{jA}(min)}$$

- estimating posteriors: 
$$P(q|D) = \frac{\binom{n}{k} q^k (1-q)^{n-k}}{P(D)} P(q) \sim q^k (1-q)^{n-k} P(q)$$

$$q_i(Z_i|D) = \frac{1}{Z} \exp \left( \langle E(Z_i, \{Z_{j \neq i}\}, D) \rangle_{\{j \neq i\}} \right)$$



**applications:**

- numerical evaluation of complicated integrals

$$\rho = \frac{P(M_A|D, I)}{P(M_B|D, I)} = \frac{P(M_A)}{P(M_B)} \cdot \frac{\int P(D|\{\alpha\}_A, M_A, I) d\alpha_{Aj}}{\int P(D|\{\alpha\}_B, M_B, I) d\alpha_{Bj}} \cdot \frac{\prod_j \alpha_{jB}(max) - \alpha_{jB}(min)}{\prod_j \alpha_{jA}(max) - \alpha_{jA}(min)}$$

- estimating posteriors:  $P(q|D) = \frac{\binom{n}{k} q^k (1-q)^{n-k}}{P(D)} P(q) \sim q^k (1-q)^{n-k} P(q)$

$$q_i(Z_i|D) = \frac{1}{Z} \exp \left( \langle E(Z_i, \{Z_{j \neq i}\}, D) \rangle_{\{j \neq i\}} \right)$$

- modelling stochastic processes: Gillespie, Metropolis

- optimization: simulated annealing

- bootstrapping: estimating confidence intervals, random forest





let's start simple:  $A_{circ} = \pi r^2 = \pi$  for  $r = 1$

**example:** finding  $\pi$  via Monte Carlo

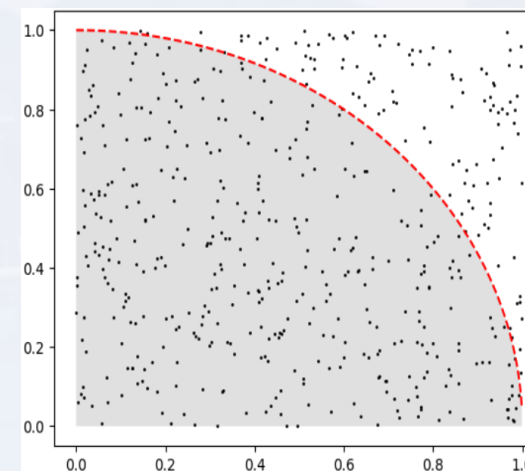


$$A_{square} = 1$$

$$A_{section} = \frac{\pi}{4}$$

$$\pi = 4 \frac{A_{section}}{A_{square}}$$

picking  $N_{tot}$  random values  $[0,1] \times [0,1]$

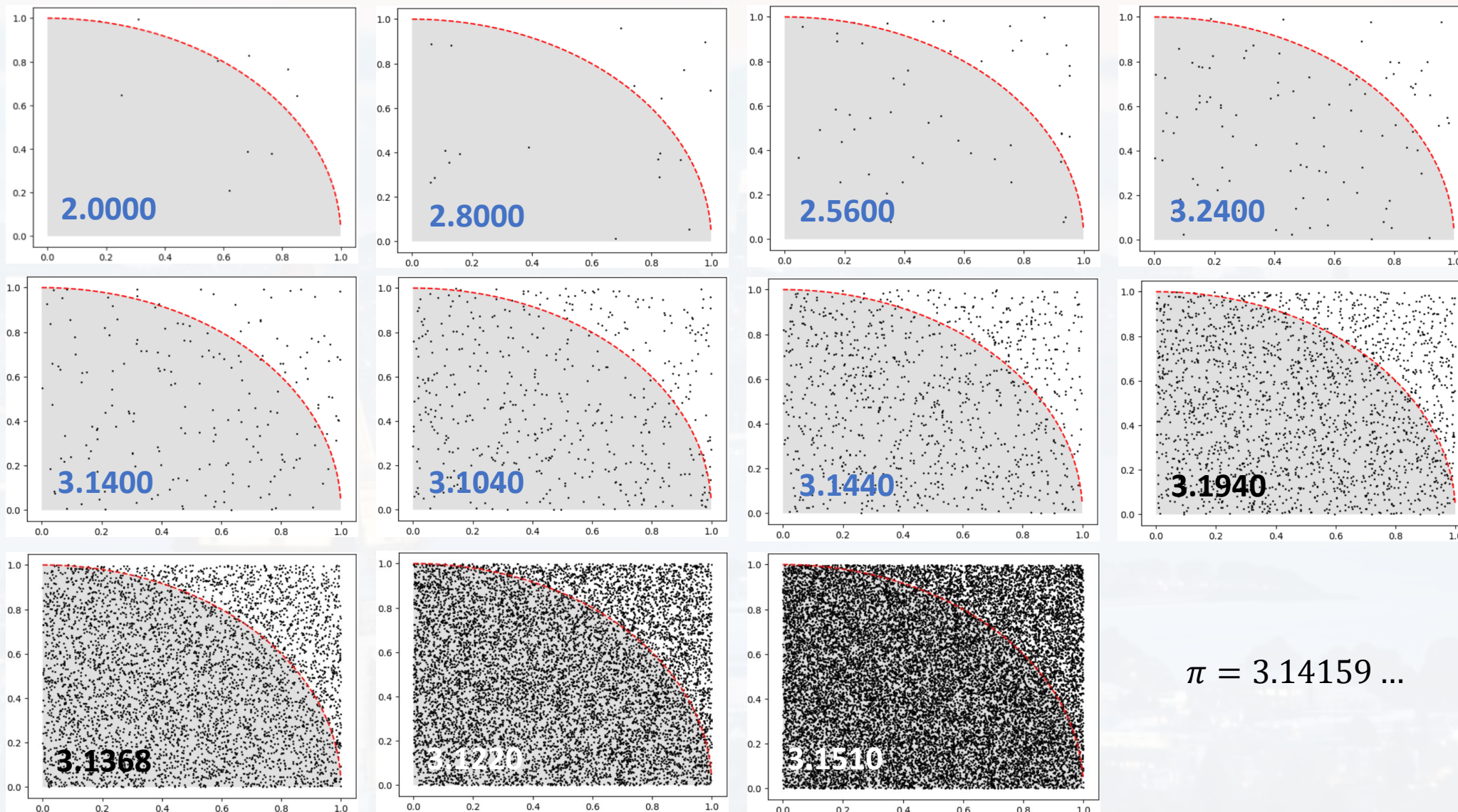


$$\pi \approx 4 \frac{N_{section}}{N_{tot}}$$



idea: generating a set of values randomly  
i.e. repeated random sampling

→ Monte Carlo method



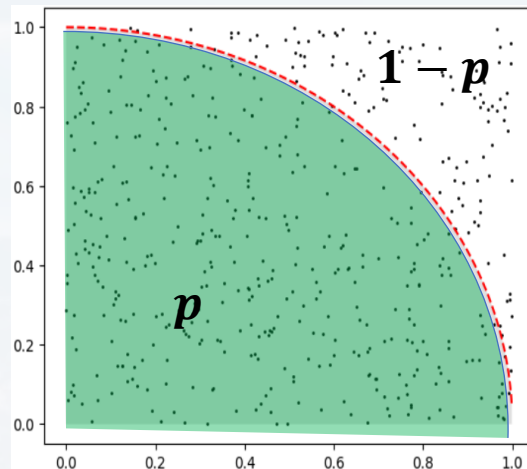
$$\pi = 3.14159 \dots$$





**example:** finding  $\pi$  via Monte Carlo

How does the accuracy of  $\pi$  depend on  $N_{tot}$ ?



$$\pi \approx 4 \frac{N_{section}}{N_{tot}}$$

We know that a point within the section is drawn with the probability  $p$

That is a **binomial problem!**

In practice,  $k$  points fall into the section with probability  $p$  and  $N_{tot} - k$  don't, with a probability of  $1 - p$

Thus, we can tell the mean and the variance of **one simulation for one specific  $N_{tot}$**

$$\sigma(k)^2 = N_{tot} p(1 - p)$$

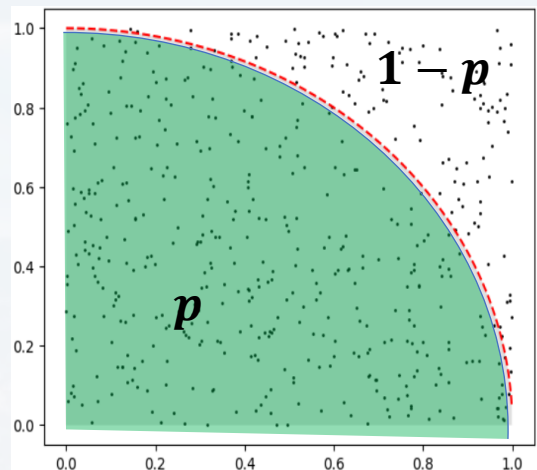
$$\mu(k) = pN_{tot}$$

(see Physics 77/88)



**example:** finding  $\pi$  via Monte Carlo

How does the accuracy of  $\pi$  depend on  $N_{tot}$ ?



$$\pi \approx 4 \frac{N_{section}}{N_{tot}}$$

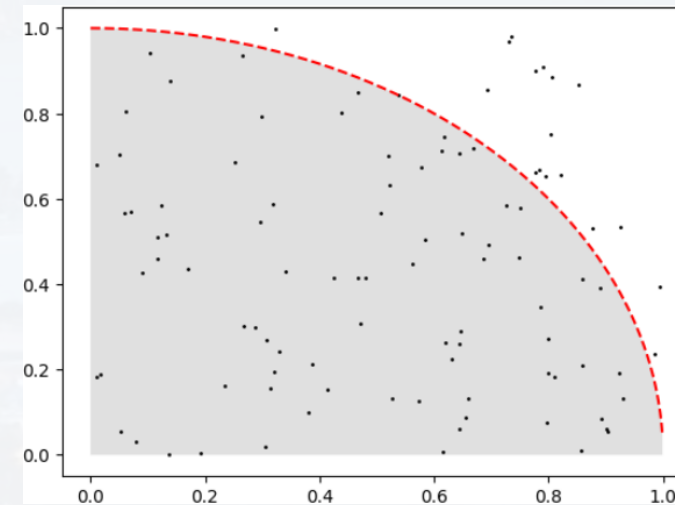
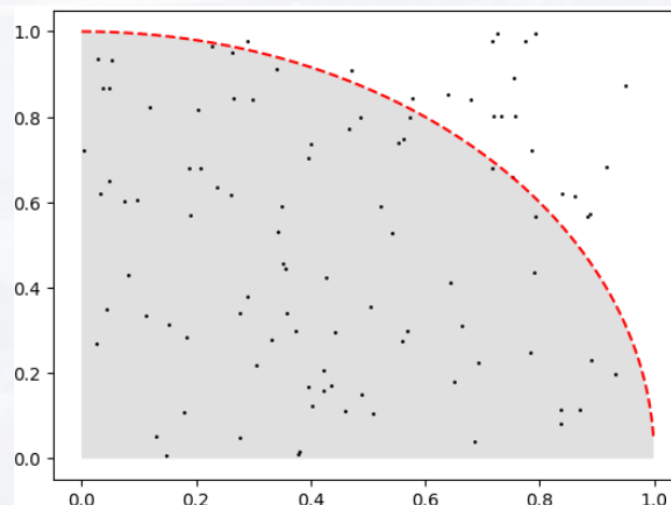
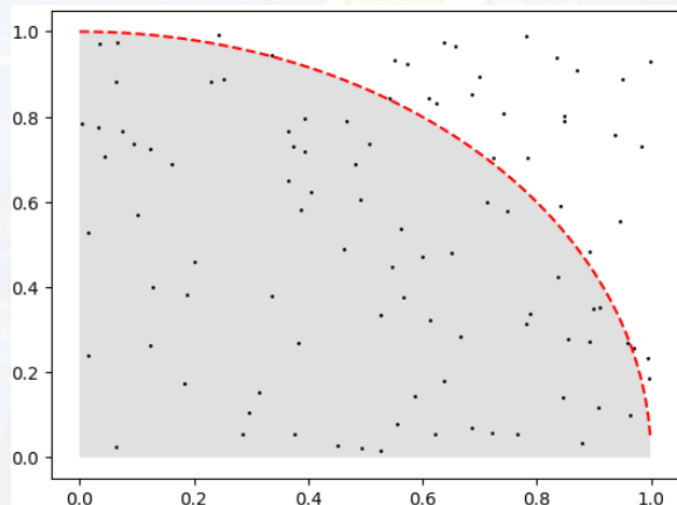
$$\sigma(k)^2 = N_{tot} p(1 - p) \quad \mu(k) = pN_{tot}$$

$$\text{error of } \pi: 4\sigma\left(\frac{N_{section}}{N_{tot}}\right) = 4\sigma\left(\frac{k}{N_{tot}}\right)$$

standard deviation  $\sigma$  of the ratio  $\frac{N_{section}}{N_{tot}}$

Say we run the simulation for a **specific**  $N_{tot}$  many times  $\rightarrow Var(k)$

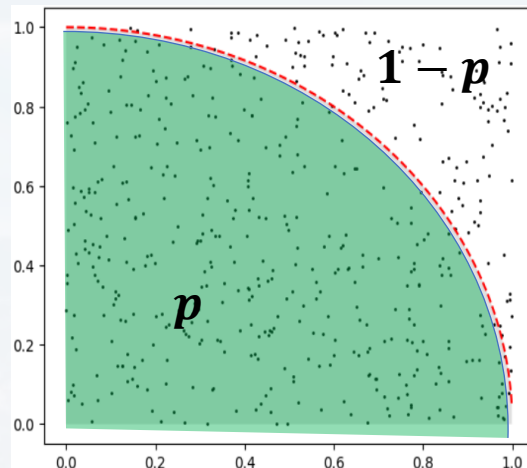
$N_{tot} = 100$





**example:** finding  $\pi$  via Monte Carlo

How does the accuracy of  $\pi$  depend on  $N_{tot}$ ?



$$\pi \approx 4 \frac{N_{section}}{N_{tot}}$$

$$\sigma(k)^2 = N_{tot} p(1 - p) \quad \mu(k) = pN_{tot}$$

$$\text{error of } \pi: 4\sigma\left(\frac{N_{section}}{N_{tot}}\right) = 4\sigma\left(\frac{k}{N_{tot}}\right)$$

Say we run the simulation for a **specific**  $N_{tot}$  many times  $\rightarrow Var(k)$

$$\text{error of } \pi: 4\sigma\left(\frac{N_{section}}{N_{tot}}\right) = 4\sigma\left(\frac{k}{N_{tot}}\right) = 4\sqrt{Var\left(\frac{k}{N_{tot}}\right)} = 4\sqrt{\frac{1}{N_{tot}^2} Var(k)} = 4\sqrt{\frac{1}{N_{tot}^2} N_{tot} p(1 - p)}$$

We know:  $Var(x) = \langle x^2 \rangle - \langle x \rangle^2$

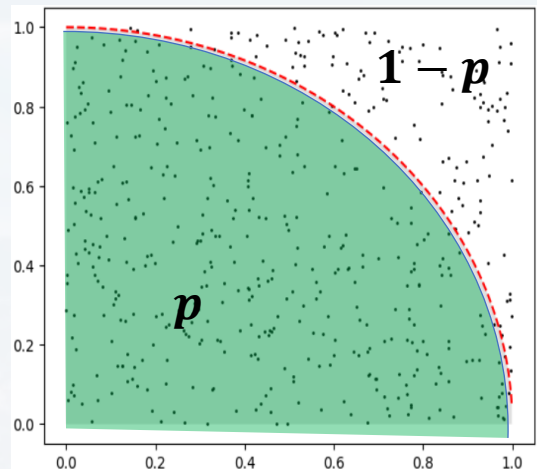
$$\mathbf{a = const} \quad x \rightarrow ax \quad Var(ax) = \langle (ax)^2 \rangle - \langle ax \rangle^2 = a^2(\langle x^2 \rangle - \langle x \rangle^2) = a^2 Var(x)$$





**example:** finding  $\pi$  via Monte Carlo

How does the accuracy of  $\pi$  depend on  $N_{tot}$ ?



$$\pi \approx 4 \frac{N_{section}}{N_{tot}}$$

$$\text{error of } \pi: 4\sigma\left(\frac{N_{section}}{N_{tot}}\right) = 4\sqrt{\frac{1}{N_{tot}^2} N_{tot} p(1-p)} = 4\sqrt{\frac{1}{N_{tot}} p(1-p)}$$

of course we **don't know**  $p = \frac{\pi}{4} = \frac{A_{section}}{A_{square}}$

because we wanted find  $\pi$  in the first place

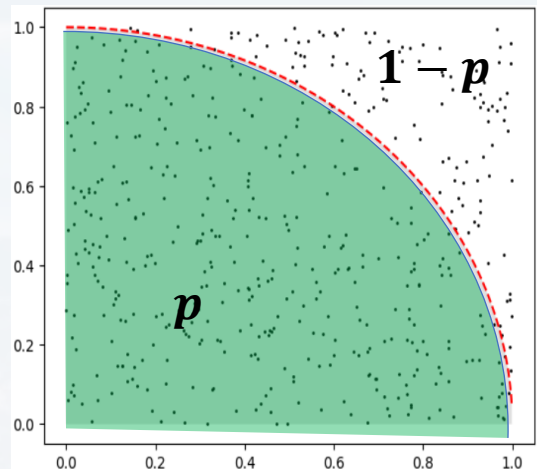
but we can **estimate** it during the simulation

$$4\sqrt{\frac{1}{N_{tot}} p(1-p)} \approx 4 \cdot 0.4 \sqrt{\frac{1}{N_{tot}}}$$



**example:** finding  $\pi$  via Monte Carlo

How does the accuracy of  $\pi$  depend on  $N_{tot}$ ?



$$\pi \approx 4 \frac{N_{section}}{N_{tot}}$$

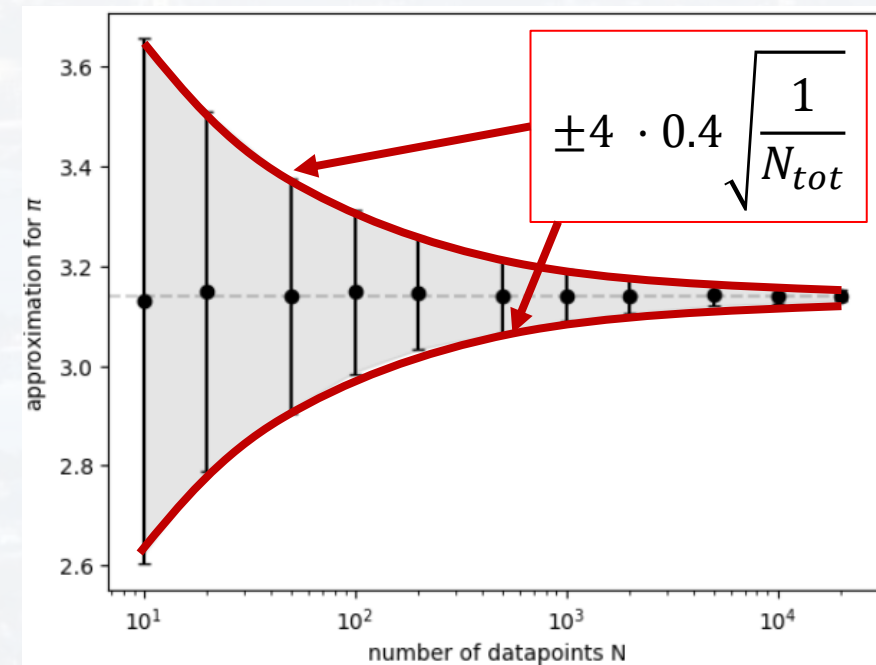
$$\text{error of } \pi: 4\sigma\left(\frac{N_{section}}{N_{tot}}\right) = 4 \sqrt{\frac{1}{N_{tot}}} p(1-p) \approx 4 \cdot 0.4 \sqrt{\frac{1}{N_{tot}}}$$

running 100 simulations **for each**  $N_{tot}$

→ calculating standard deviation of  $\pi$

→ comparing to  $4 \cdot 0.4 \sqrt{\frac{1}{N_{tot}}}$

see `Monte_Carlo_Simulation_PI.ipynb`





## CHECKING IN



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Mapping Distributions & Gibbs Sampling

Gillespie Algorithm

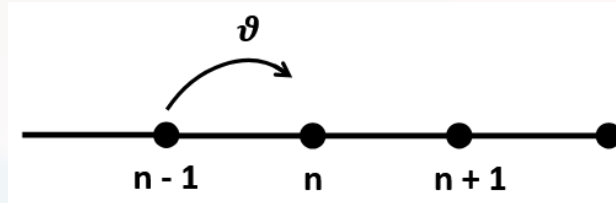
Metropolis (- Hastings) Algorithm

Bootstrapping





last time:

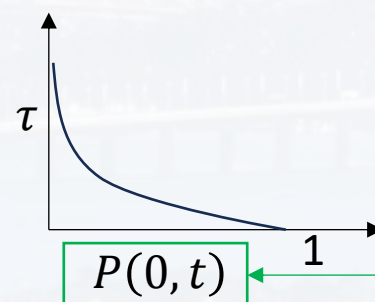
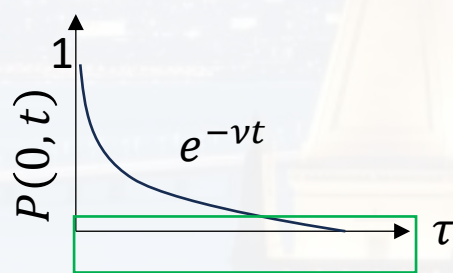


$n$ : different states  
 $\vartheta$ : hopping rate (*probability/time*)

calculating the **waiting time** (time  $\tau$  between two events)

$$P(0, t) = \frac{(\vartheta t)^0}{0!} e^{-\vartheta t} \quad \tau = -\frac{1}{\vartheta} \ln[P(0, t)]$$

We want to understand the process a bit better (i. e. find the **waiting time distribution**  $w(\tau)$ )



random number  
 $\rho \in [a, b]$   
 $P(\rho) = \text{const} = 1/(b - a)$

waiting time  $\tau$ :  
 $\tau \in [0, \infty)$   
 $w(\tau) \neq \text{const}$

$$\int_a^b P(\rho) d\rho = 1 = \int_0^\infty w(\tau) d\tau$$



$$\tau = -\frac{1}{\vartheta} \ln[P(0, t)] = -\frac{1}{\vartheta} \ln[\rho]$$

**$n$ :** different states  
 **$\vartheta$ :** hopping rate (*probability/time*)  
 **$\rho$ :** uniformly dist random number

$$\int_a^b f(x) dx = \int_{y(a)}^{y(b)} f[x(y)] \frac{dx}{dy} dy \quad \text{(substitution)}$$

$$= \int_{y(a)}^{y(b)} g(y) dy \quad g(y) := f[x(y)] \frac{dx}{dy}$$

note: we actually need  $\left| \frac{dx}{dy} \right|$

inverse of  $y$ :  $x(y) = y^{-1}$

**our situation:**  $\int_{\rho=0}^{\rho=1} P(\rho) d\rho = \int_{\tau(0)}^{\tau(1)} P[\rho(\tau)] \left| \frac{d\rho}{d\tau} \right| d\tau = \int_{\tau(0)}^{\tau(1)} w(\tau) d\tau$

$$w(\tau) = P[\rho(\tau)] \left| \frac{d\rho}{d\tau} \right| \quad x(y) = \rho(\tau) = e^{-\vartheta\tau}$$

$$w(\tau) = 1 \cdot \vartheta e^{-\vartheta\tau}$$

mean waiting time:  $t^* = \int_{\tau=0}^{\tau=\infty} \tau w(\tau) d\tau = \frac{1}{\vartheta}$



actual Gibbs Sampling:

we know from module 6:  $q_i(Z_i|D) = \frac{1}{Z} \exp \left( \langle E(Z_i, \{Z_{j \neq i}\}, D) \rangle_{\{j \neq i\}} \right)$

which led to circular dependencies

$D$	: data set of size $K$
$Z$	: set of $n$ (latent) parameter
$\sigma^2$	: variance
$\mu$	: mean
$\frac{1}{\sigma^2} = \lambda$	: precision

$$q_\mu(\mu|D) \sim \mathcal{N}(\mu|\mu_K, \lambda_K^{-1})$$

where

$$\mu_K = \frac{\tau_0 \mu_0 + K \bar{x}}{\tau_0 + K}$$

$$\lambda_K = (\tau_0 + K) \langle \tau \rangle_\tau$$

$$\bar{x} = \frac{1}{K} \sum_{k=1}^K x_k$$

$$q_\tau(\tau|D) \sim \Gamma(\tau|a_K, b_K)$$

where

$$a_K = a + \frac{K+1}{2}$$

$$b_K = b + \frac{1}{2} \langle \sum_k (x_k - \mu)^2 + \tau_0 (\mu - \mu_0)^2 \rangle_\mu$$

$$\langle \tau \rangle_\tau = \int \tau q_\tau(\tau|D) d\tau = \frac{a_K}{b_K}$$

set  $\tau_0$ ,  $\mu_0$ ,  $a$  and  $b$  to **small positive values** (largest ignorance)





**actual Gibbs Sampling:**

we know from [module 6](#):  $q_i(Z_i|D) = \frac{1}{Z} \exp \left( \langle E(Z_i, \{Z_{j \neq i}\}, D) \rangle_{\{j \neq i\}} \right)$

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$q_i(Z_i|\{Z_{j \neq i}\})$  sampling for all  $i$  in a particular order

say we have **three** parameters

- randomly (or MLE guess from data) initialize  $Z_1, Z_2, Z_3$

iteration  $t$

draw  $Z_1$   $q(Z_1(t+1)|Z_2(t), Z_3(t))$

draw  $Z_2$   $q(Z_2(t+1)|Z_1(t+1), Z_3(t))$

draw  $Z_3$   $q(Z_3(t+1)|Z_1(t+1), Z_2(t+1))$

iteration  $t+2$





**actual Gibbs Sampling:**

we know from [module 6](#):  $q_i(Z_i|D) = \frac{1}{Z} \exp \left( \langle E(Z_i, \{Z_{j \neq i}\}, D) \rangle_{\{j \neq i\}} \right)$

$q_i(Z_i|\{Z_{j \neq i}\})$  sampling for all  $i$  in a particular order

$D$	: data set of size $K$
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$\frac{1}{\sigma^2} = \lambda$	: precision

for the example from [module 6](#)

- now: randomly (or MLE guess from data) initialize  $\mu, \lambda$

$$D|\mu, \lambda \sim \mathcal{N}(\mu, \lambda^{-1})$$

$$\mu \sim \mathcal{N}(\mu_0, \lambda_0^{-1})$$

$$\lambda \sim \Gamma(a_0, b_0)$$

instead for calculating the means, we now sample:

$$\mu|\lambda, D \sim \mathcal{N}(M_\lambda, L_\lambda^{-1}) \quad L_\lambda = (\lambda_0 + K)\lambda \quad M_\lambda = \frac{\mu_0 \lambda_0 + \sum_{k=1}^K x_k}{\lambda_0 + K}$$

$$\lambda|\mu, D \sim \Gamma(a_K, b_K) \quad a_K = a_0 + \frac{K}{2} \quad b_K = b_0 + \frac{1}{2} \sum_{k=1}^K (x_k - \mu)^2$$

iteration  $t \rightarrow t + 1$



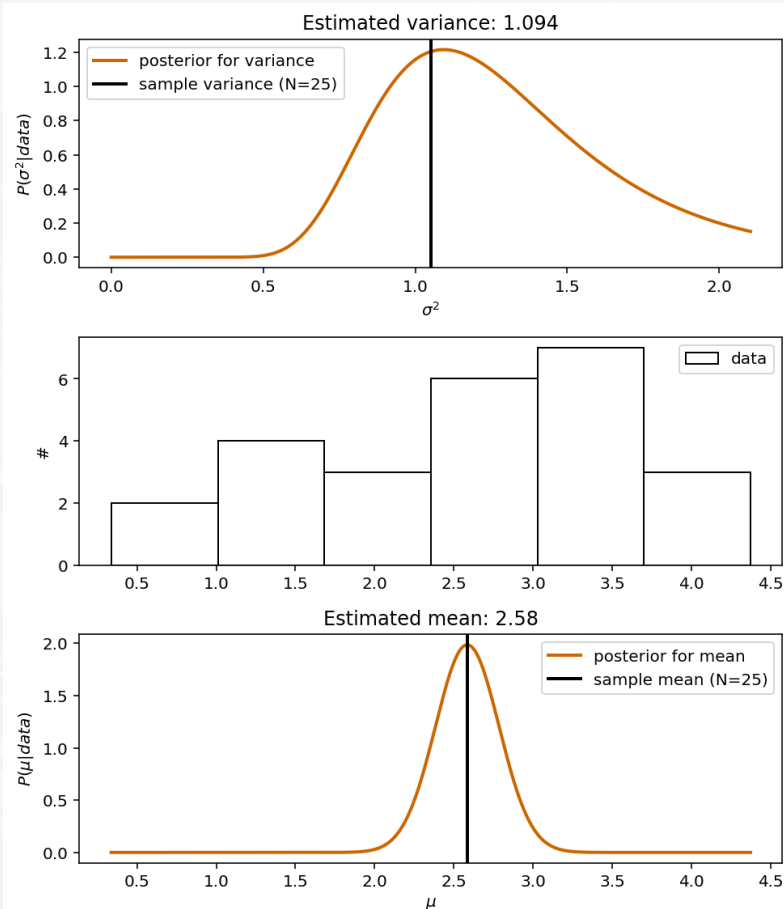
**actual Gibbs Sampling:**

we know from [module 6](#):  $q_i(Z_i|D) = \frac{1}{Z} \exp \left( \langle E(Z_i, \{Z_{j \neq i}\}, D) \rangle_{\{j \neq i\}} \right)$

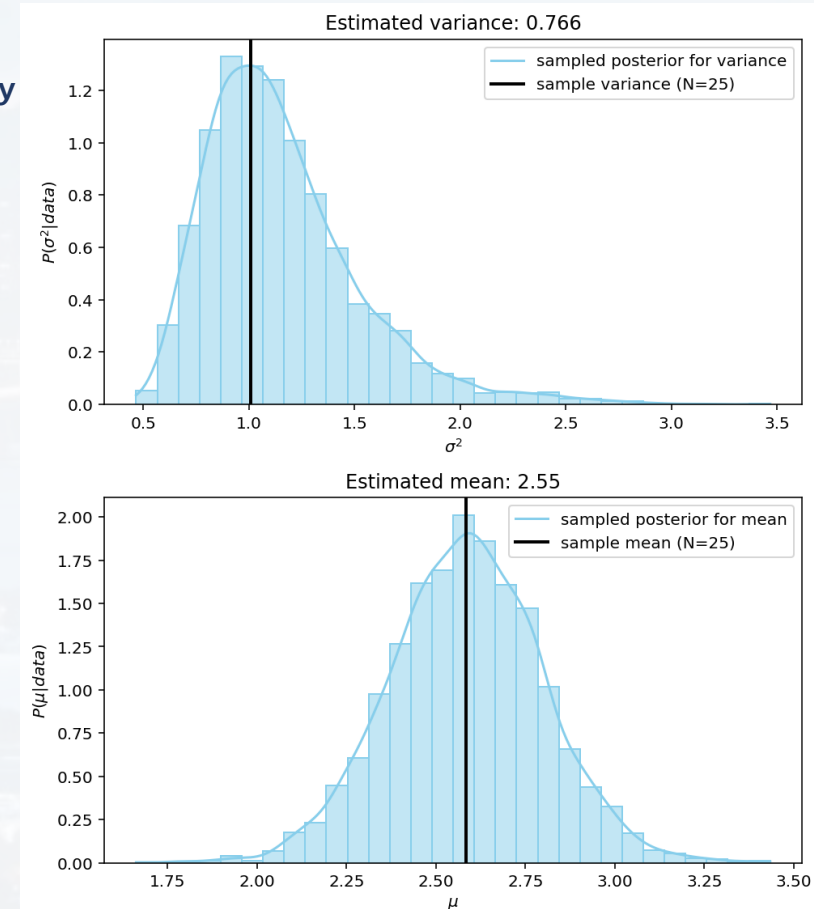
$q_i(Z_i|\{Z_{j \neq i}\})$  sampling for all  $i$  in a particular order

$D$	: data set of size $K$
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**Variational Bayes**  
`Var_Bayes_Example.py`



**Gibbs Sampling**  
`Gibbs_NormGamma_Example.py`







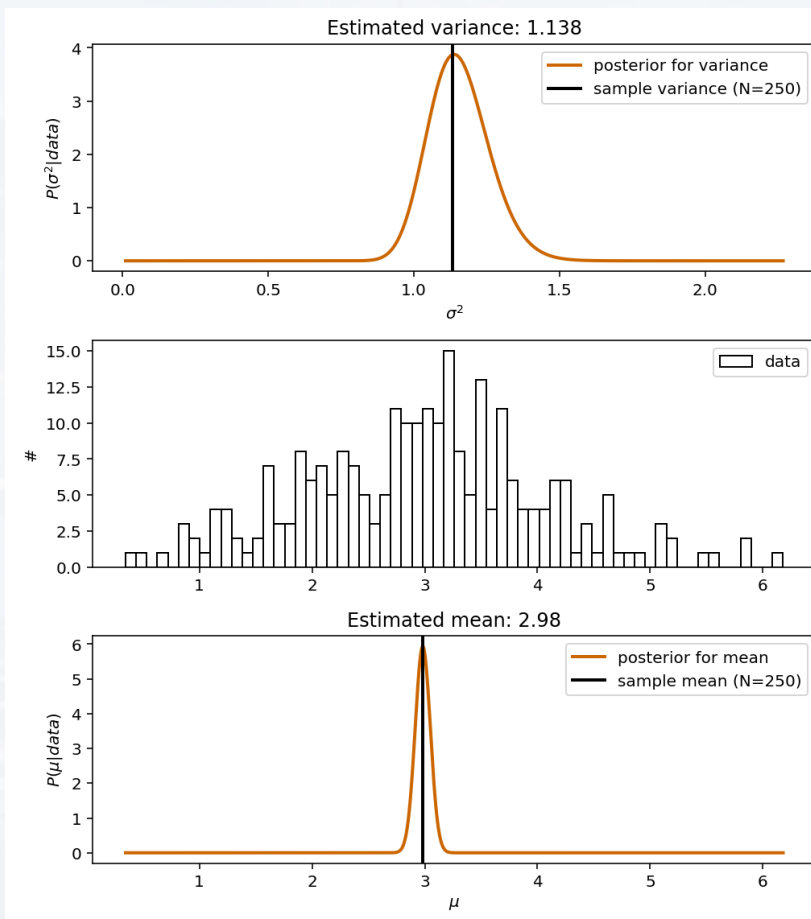
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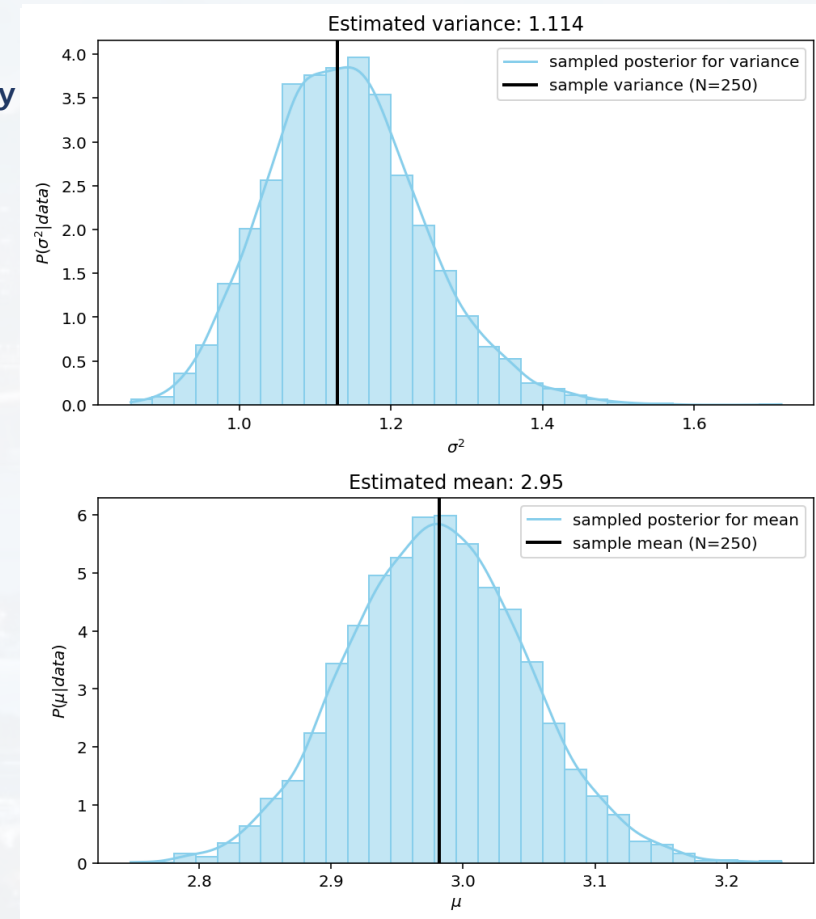
$q_i(Z_i|\{Z_{j \neq i}\})$  sampling for all  $i$  in a particular order

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**Variational Bayes**  
`Var_Bayes_Example.py`



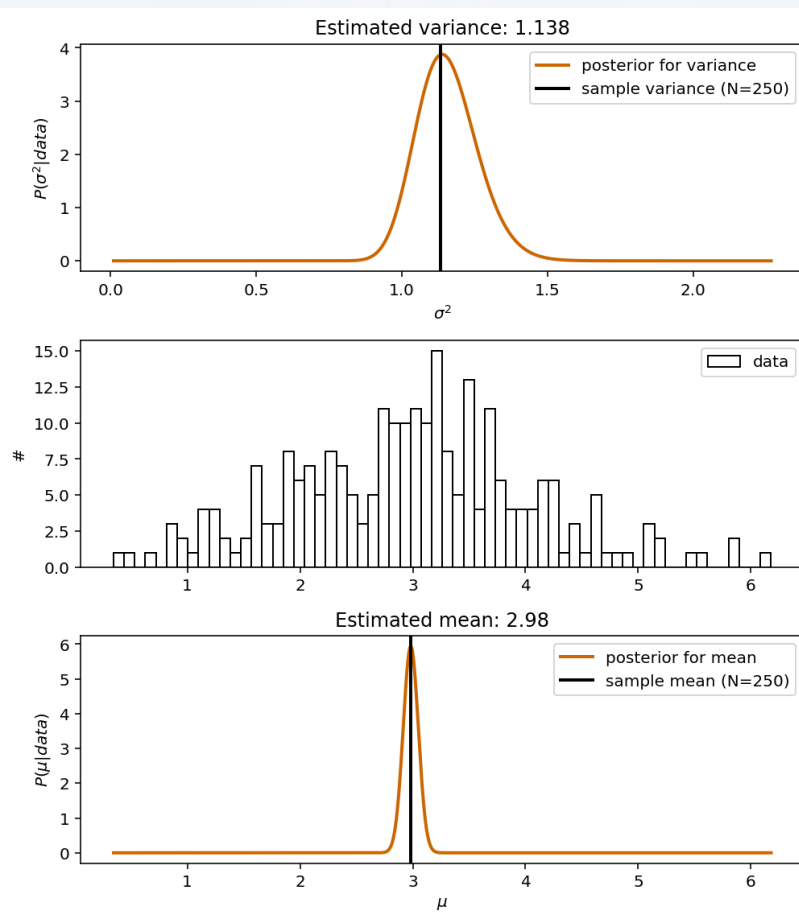
**Gibbs Sampling**  
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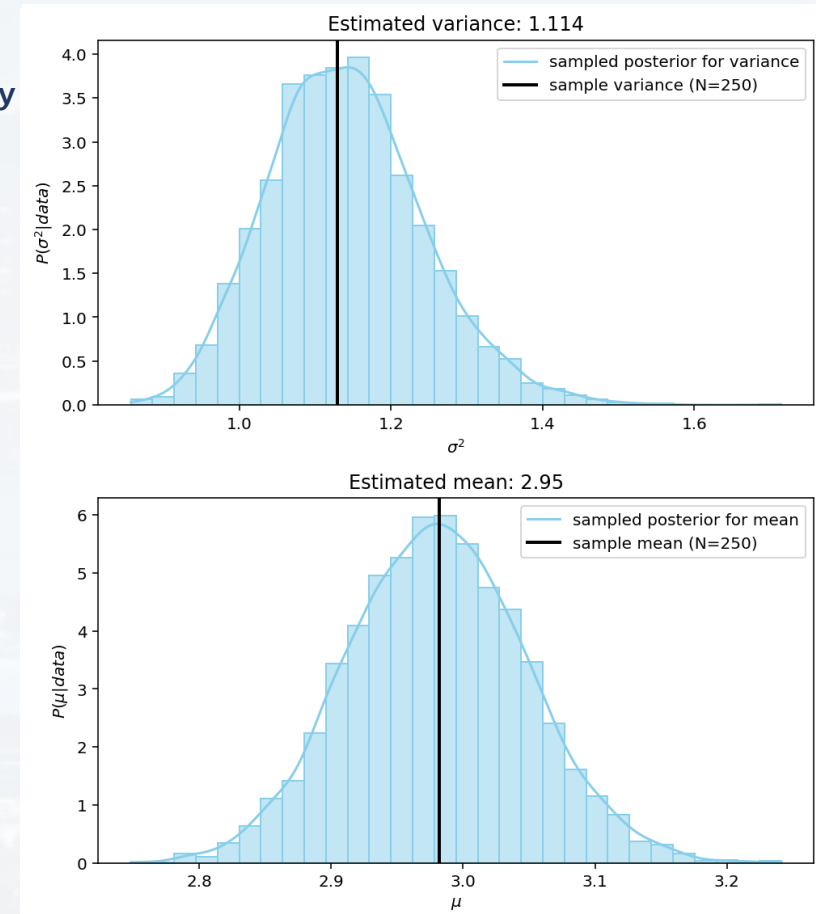


- note:**
- Gibbs sampling does **not guarantee** to find a **global solution!**
  - circular dependencies: **has to be ergodic**, otherwise interruption if  $q_i(Z_i|\{Z_{j \neq i}\}) = 0$  for any  $i$

### Variational Bayes Var\_Bayes\_Example.py



### Gibbs Sampling Gibbs\_NormGamma\_Example.py





## CHECKING IN



## Outline

Basic Idea & Finding Pi

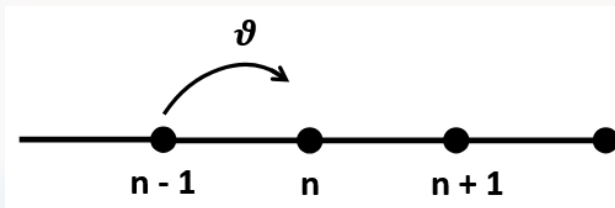
Mapping Distributions & Gibbs Sampling

**Gillespie Algorithm**

Metropolis (- Hastings) Algorithm

Bootstrapping





$$\tau = -\frac{1}{\vartheta} \ln[\rho]$$

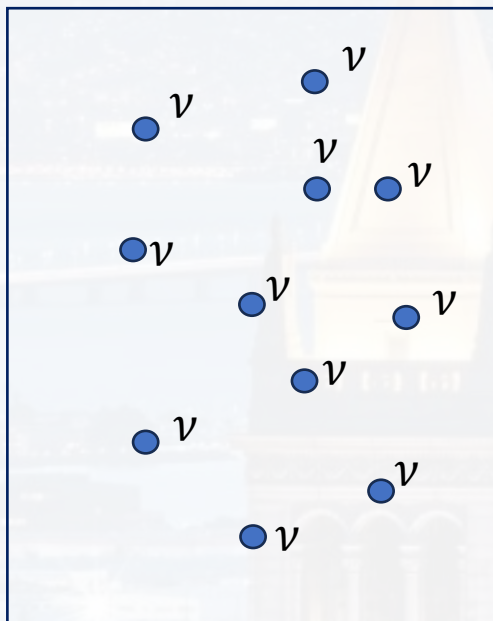
$$w(\tau) = \vartheta e^{-\vartheta \tau}$$

$n$ :	different states
$\vartheta$ :	hopping rate ( <i>probability/time</i> )
$\tau$ :	waiting time
$w(\tau)$ :	waiting time distribution

mean waiting time:  $t^* = \frac{1}{\vartheta}$

$A \xrightarrow{k} \emptyset$  stochastic scenario: number  $n$  of particles  $A$

$t$



for  $t = 0$  **many** atoms  
 $\rightarrow \tau$  is small

$\Delta t$

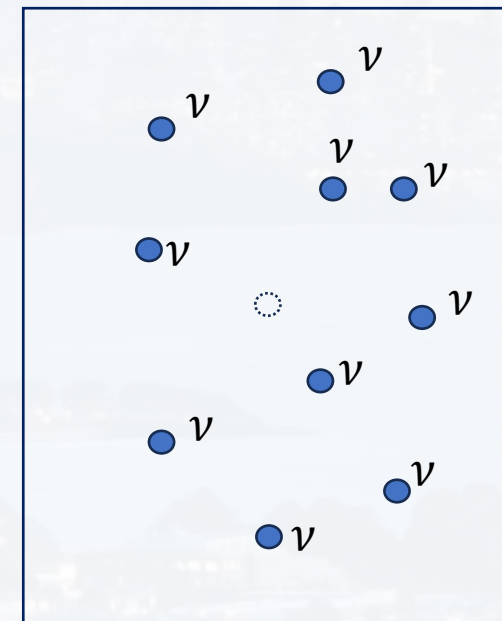
**each** atom has the probability  $\nu$  to decay per time

logical **or**  $\rightarrow$  **adding** the probabilities

$$\nu \rightarrow \nu n(t)$$

$$\Delta t = -\frac{1}{\nu n(t)} \ln[P(0|t)]$$

$t + \Delta t$





**Gillespie:**

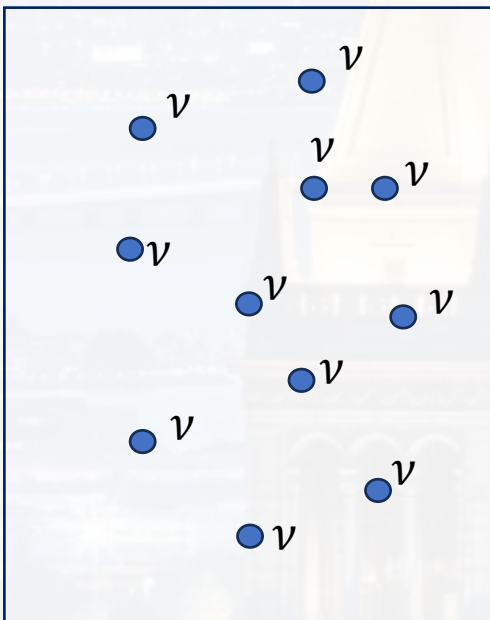
1) draw a **random number**  $\rho$  from a **uniform distribution** in the interval **(0, 1)**

2) calculate the time  $\Delta t$  that elapses until the next decay

$$\Delta t = -\frac{1}{\nu n(t)} \ln \rho$$

3) set  $t \rightarrow t + \Delta t$  and  $n(t + \Delta t) = n(t) - 1$

4) repeat



$$\Delta t = -\frac{1}{\nu n(t)} \ln[\rho]$$



**Gillespie:**

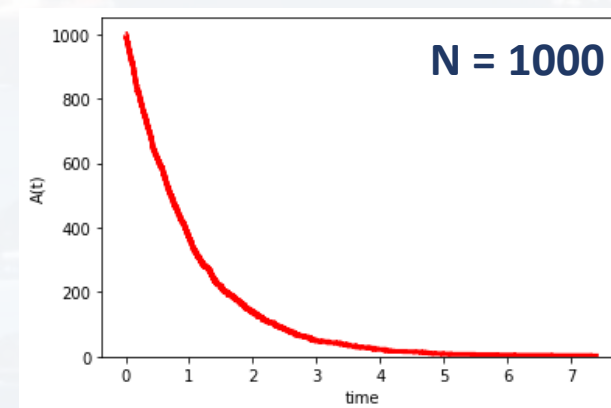
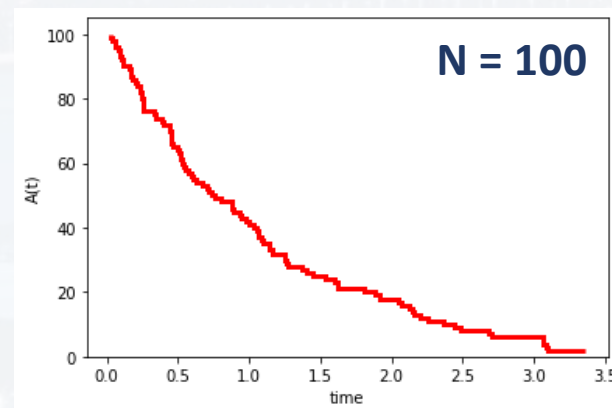
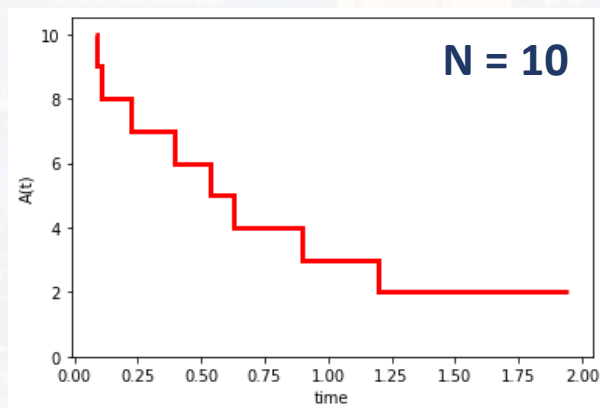
1) draw a **random number**  $\rho$  from a **uniform distribution** in the interval **(0, 1)**

2) calculate the time  $\Delta t$  that elapses until the next decay

$$\Delta t = -\frac{1}{v n(t)} \ln \rho$$

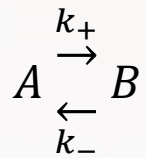
3) set  $t \rightarrow t + \Delta t$  and  $n(t + \Delta t) = n(t) - 1$

4) repeat



see Decay.py





$n$ : number of particles of A  
 $m$ : number of particles of B

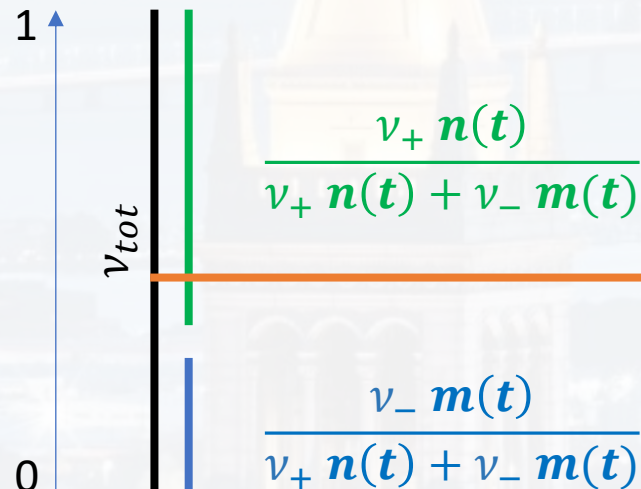
$n$ : different states  
 $\vartheta$ : hopping rate (*probability/time*)  
 $\tau$ : waiting time  
 $w(\tau)$ : waiting time distribution

$$v(A) \rightarrow v_+ n(t) \quad v(B) \rightarrow v_- m(t) \quad v_{tot} = v(A) + v(B) = v_+ n(t) + v_- m(t)$$

$$\Delta t = - \frac{1}{v_+ n(t) + v_- m(t)} \ln[\rho]$$

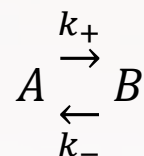
time that elapses until **a** reaction to occurs

next: deciding **which** reaction should occur



depending into which fraction  
 this random number falls  
 → this reaction occurs

generating a random number from a  
**uniform distribution** in the interval **(0, 1)**



n: number of particles of A  
m: number of particles of B

n: different states  
ϑ: hopping rate (*probability/time*)  
τ: waiting time  
w(τ): waiting time distribution

**Gillespie:**

1) draw a **random number**  $\rho_1$  from a **uniform distribution** in the interval **(0, 1)**

2) calculate the time  $\Delta t$  that elapses until the next reaction

$$\Delta t = - \frac{1}{\nu_+ n(t) + \nu_- m(t)} \ln \rho_1$$

3) draw a **second random number**  $\rho_2$  from a **uniform distribution** in the interval **(0, 1)**

4) decide which reaction occurs:

$$\text{if } \rho_2 < \frac{\nu_+ n(t)}{\nu_+ n(t) + \nu_- m(t)} :$$

reaction  $A \rightarrow B$  is more likely

$$n(t + \Delta t) = n(t) - 1$$

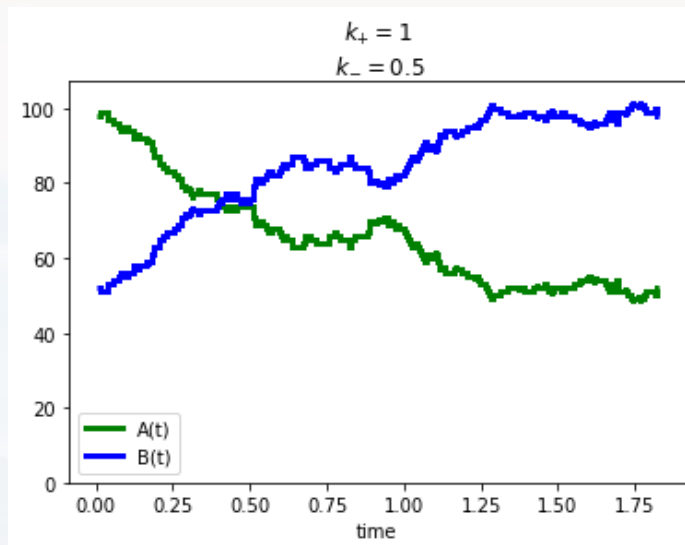
$$m(t + \Delta t) = m(t) + 1$$

else:

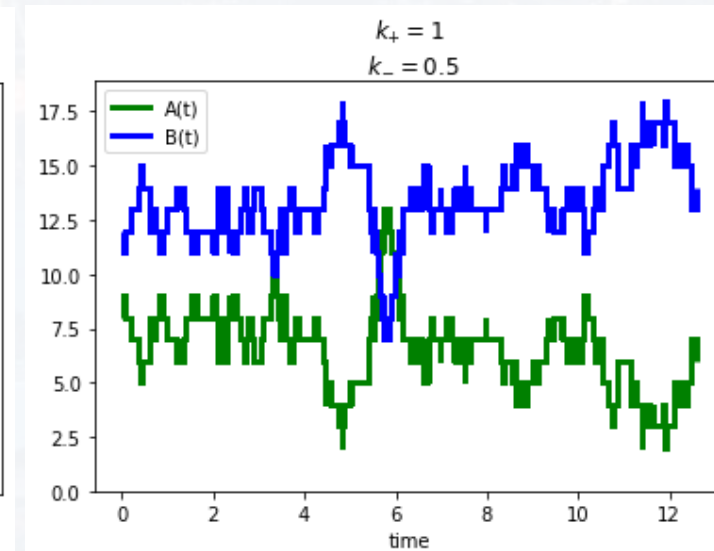
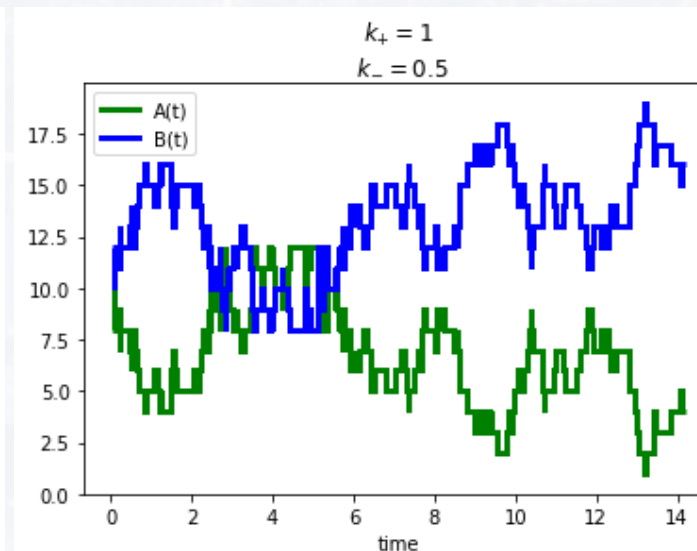
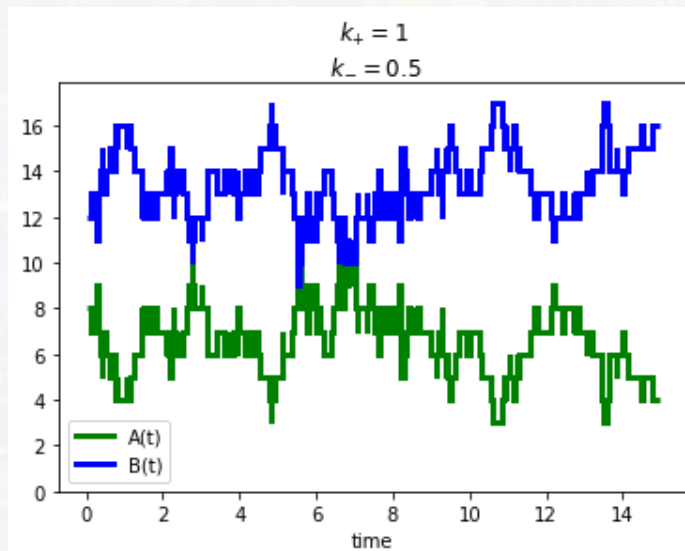
reaction  $B \rightarrow A$  is more likely

$$n(t + \Delta t) = n(t) + 1$$

$$m(t + \Delta t) = m(t) - 1$$



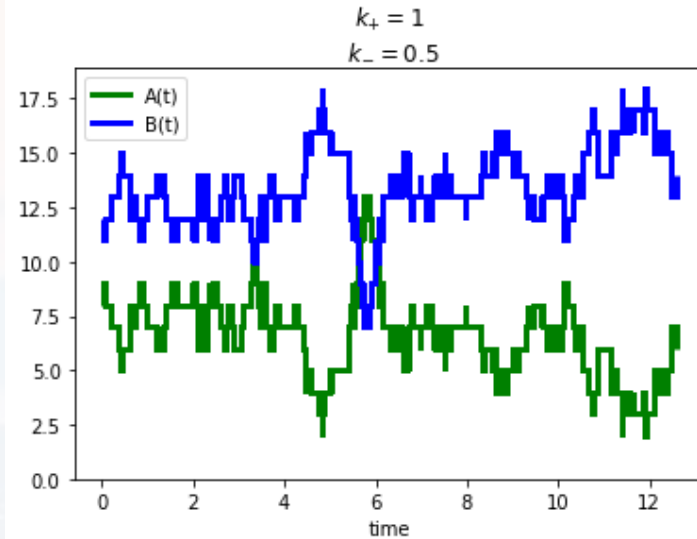
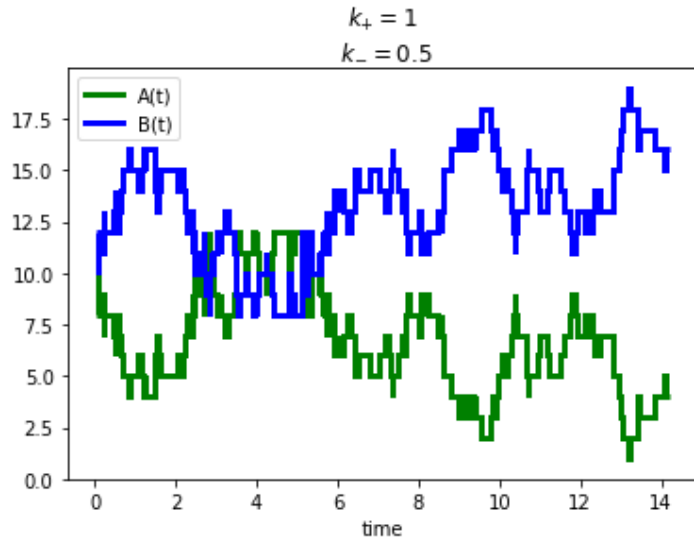
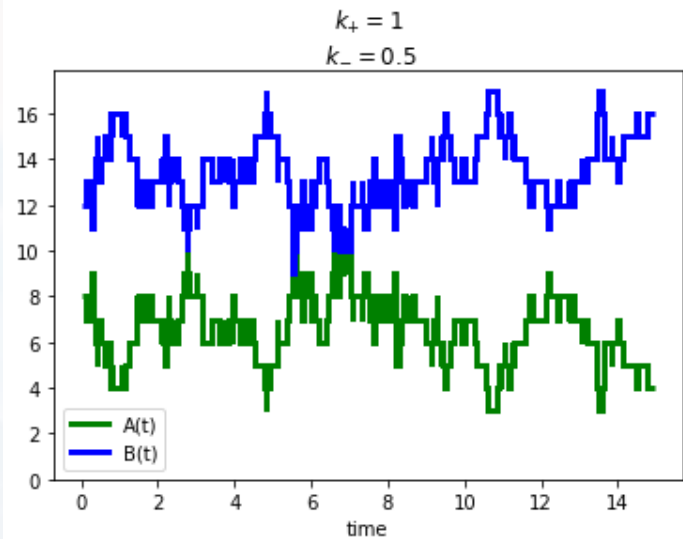
$N = 10$



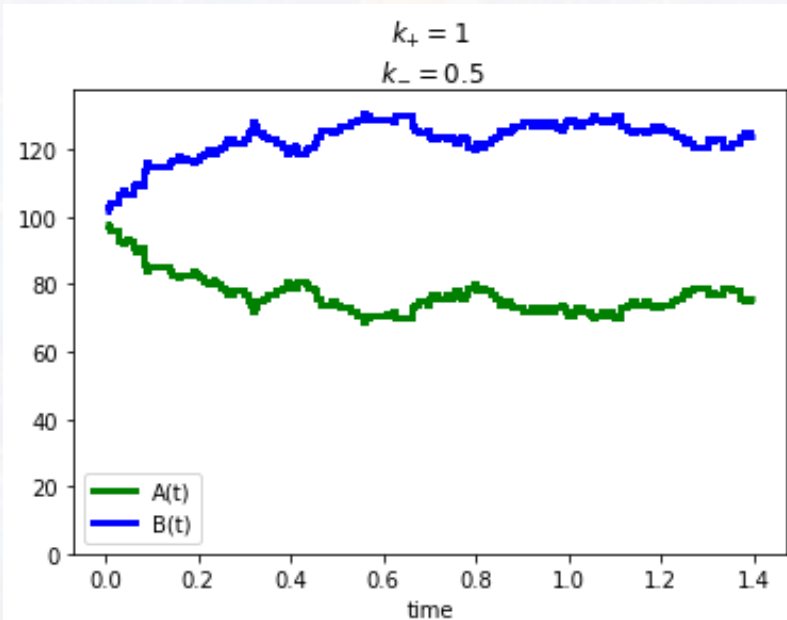




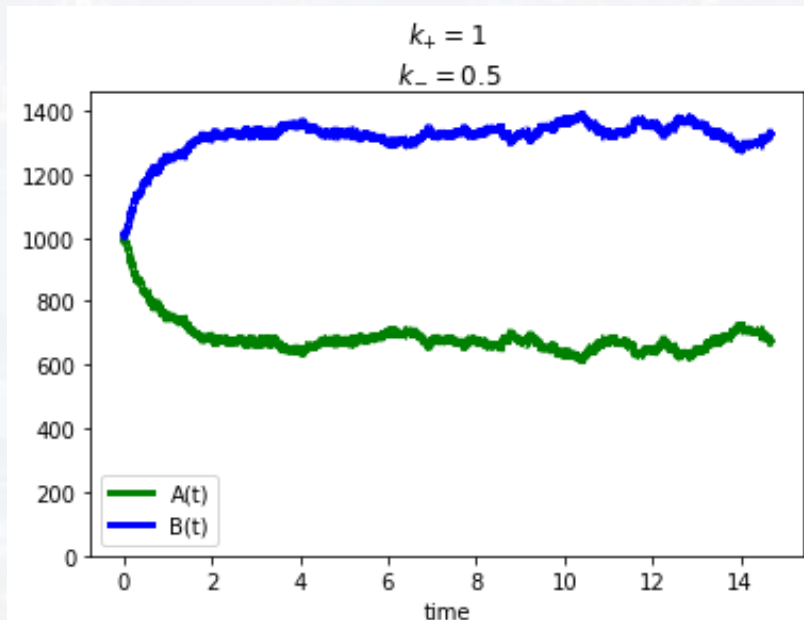
$N = 10$

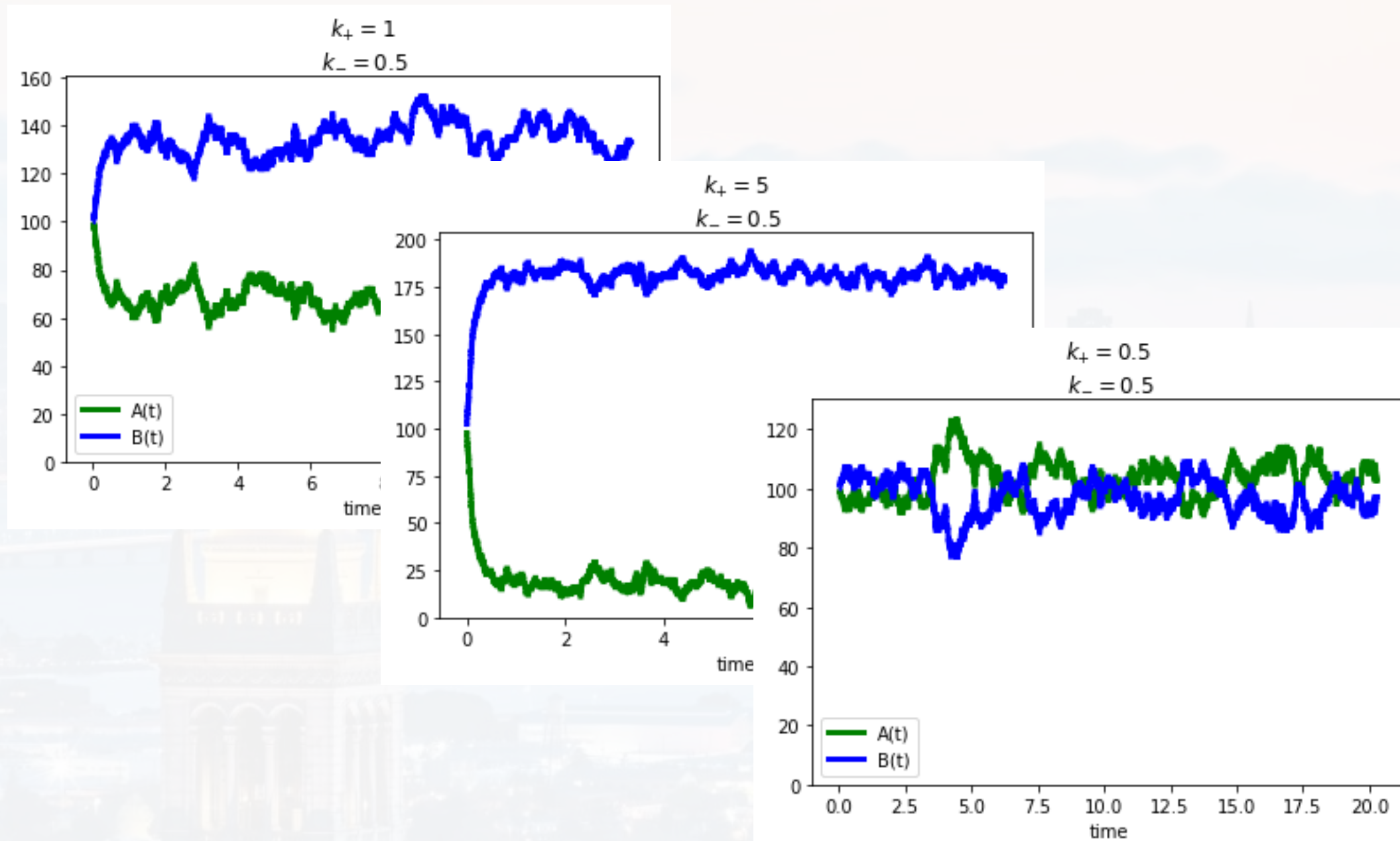


$N = 100$



$N = 1000$





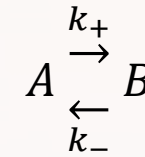
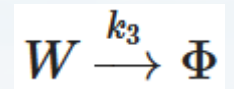
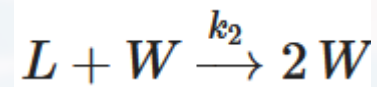
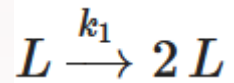


We also know how to solve the Predator-Prey model now!

L: sheep (lambs)

W: wolves

E: "empty"



$n$ :

$\vartheta$ :

$dt$ :

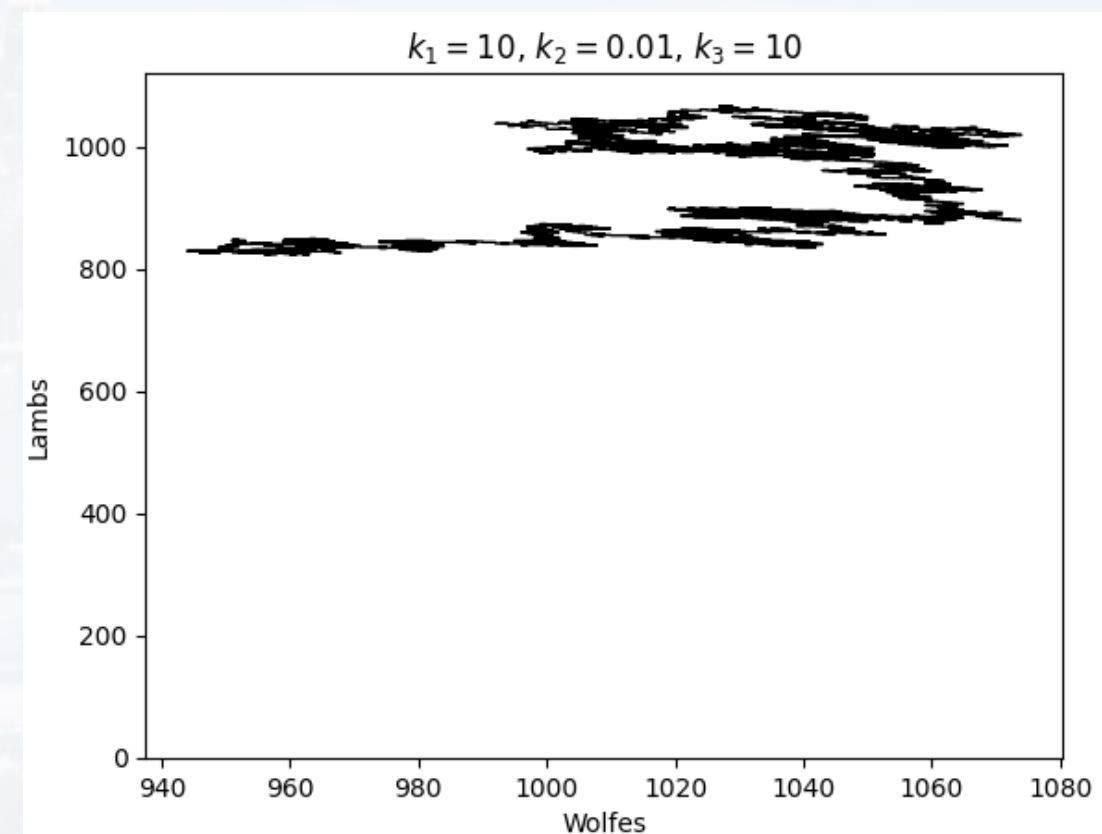
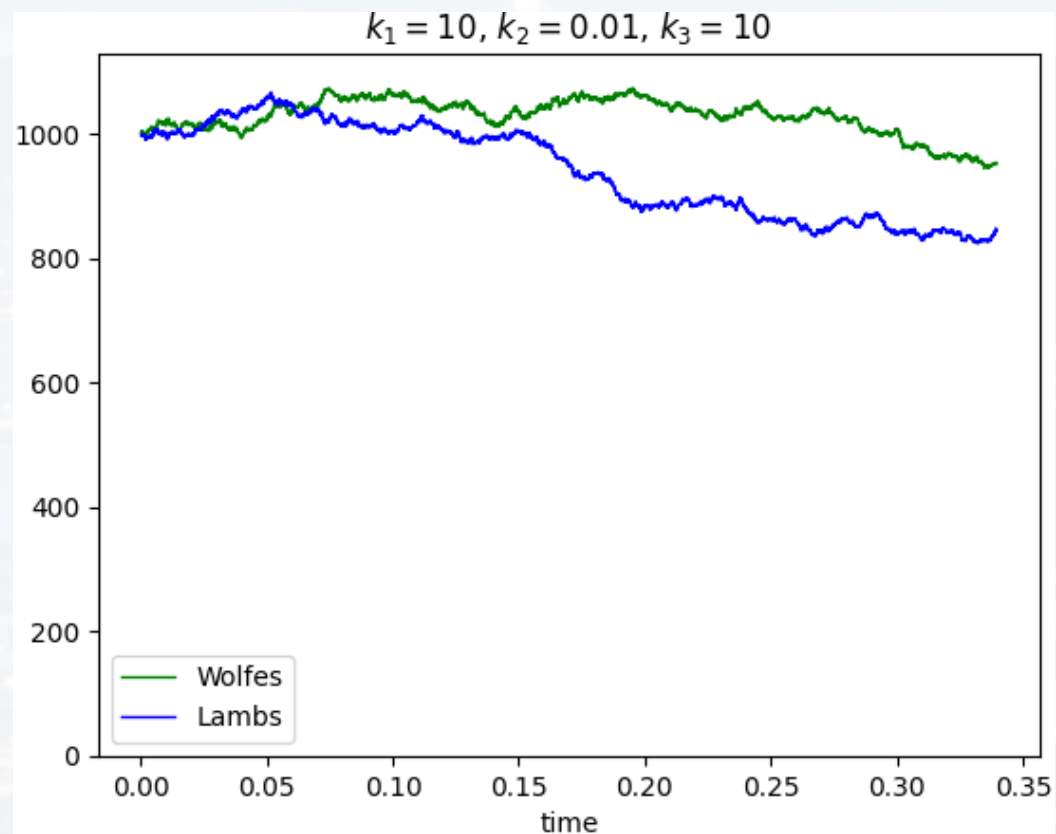
$\tau = -\ln[P(0, t)] / \nu$ :

different states

hopping rate

time increment

waiting time





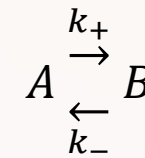
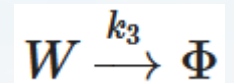
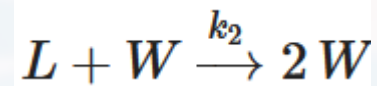
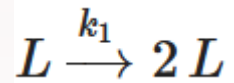


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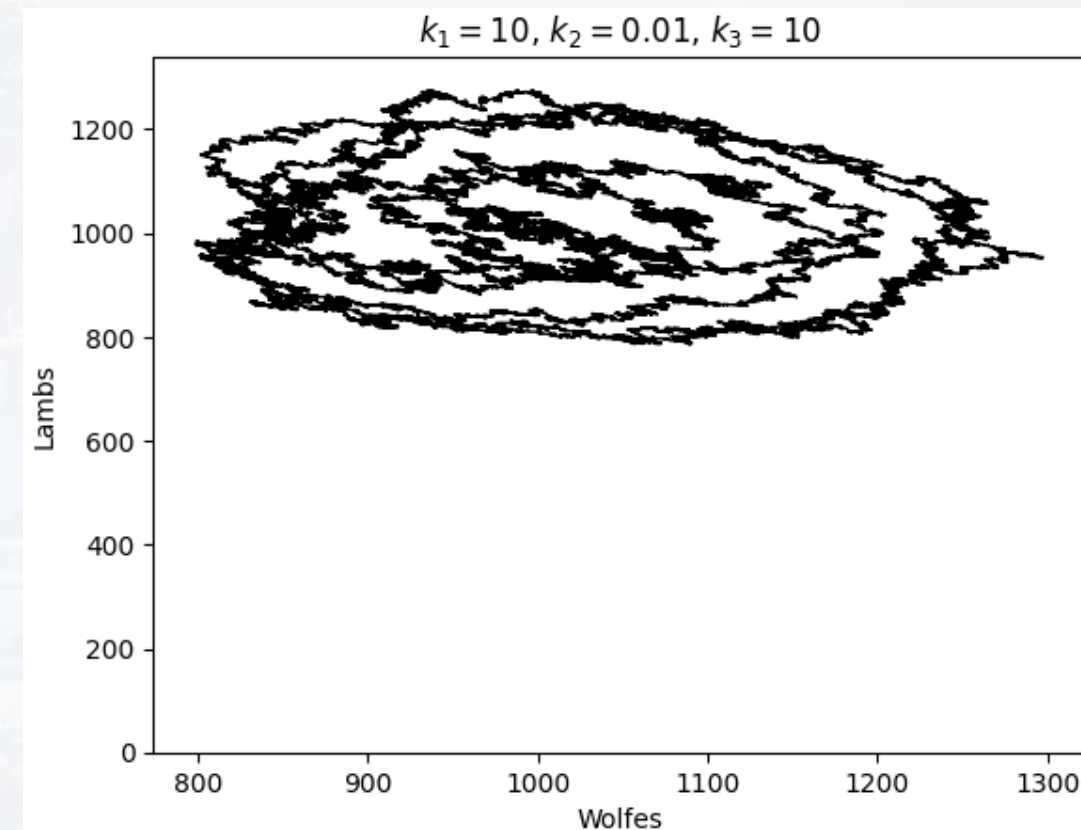
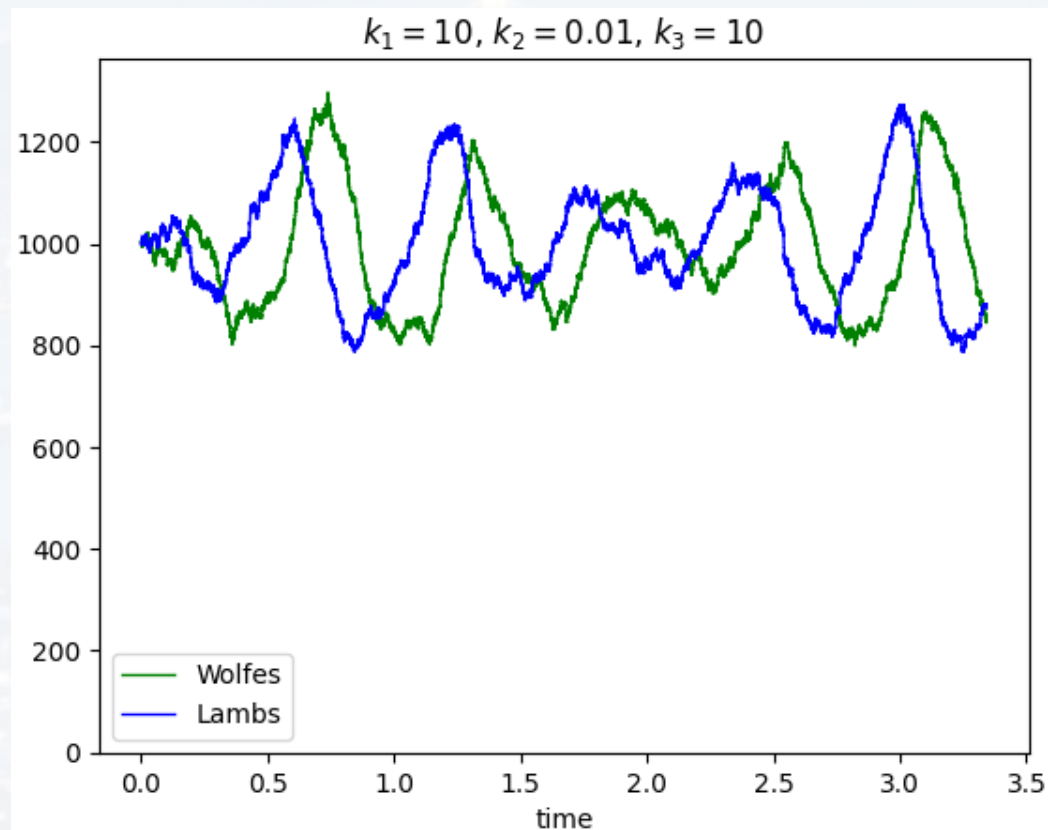
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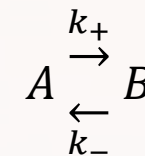
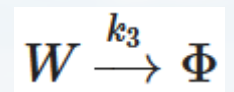
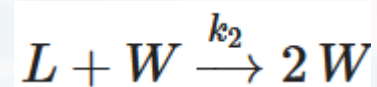
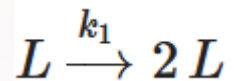


We also know how to solve the Predator-Prey model now!

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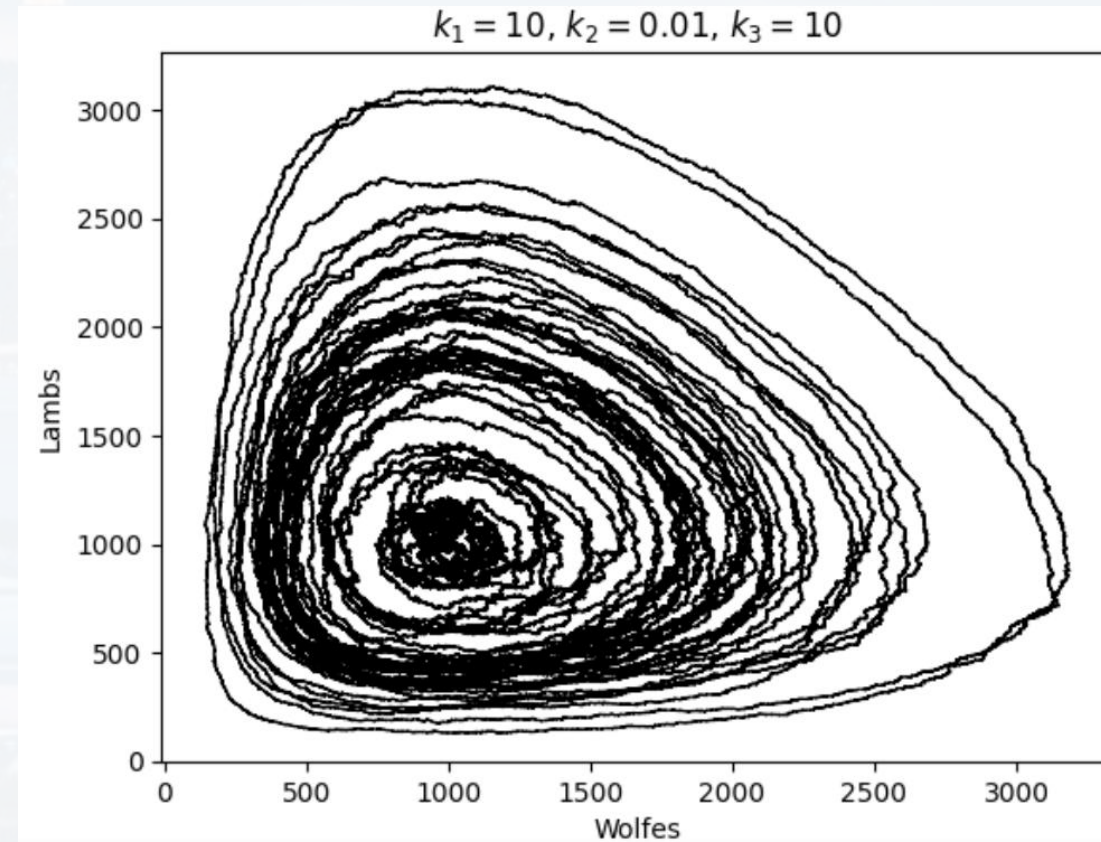
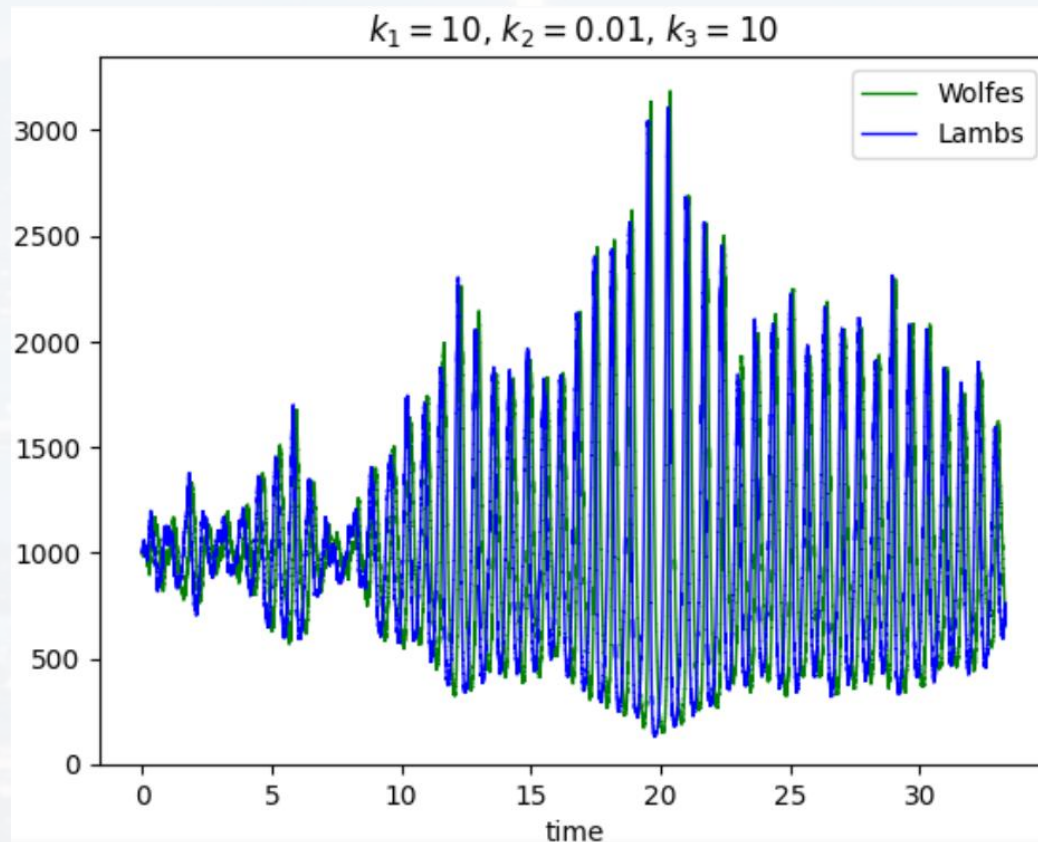
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## CHECKING IN



## Outline

Basic Idea & Finding Pi

Mapping Distributions & Gibbs Sampling

Gillespie Algorithm

**Metropolis (- Hastings) Algorithm**

Bootstrapping





**problem:** sometimes we need to draw from a probability distribution (= **target**) that is **difficult to sample** from directly

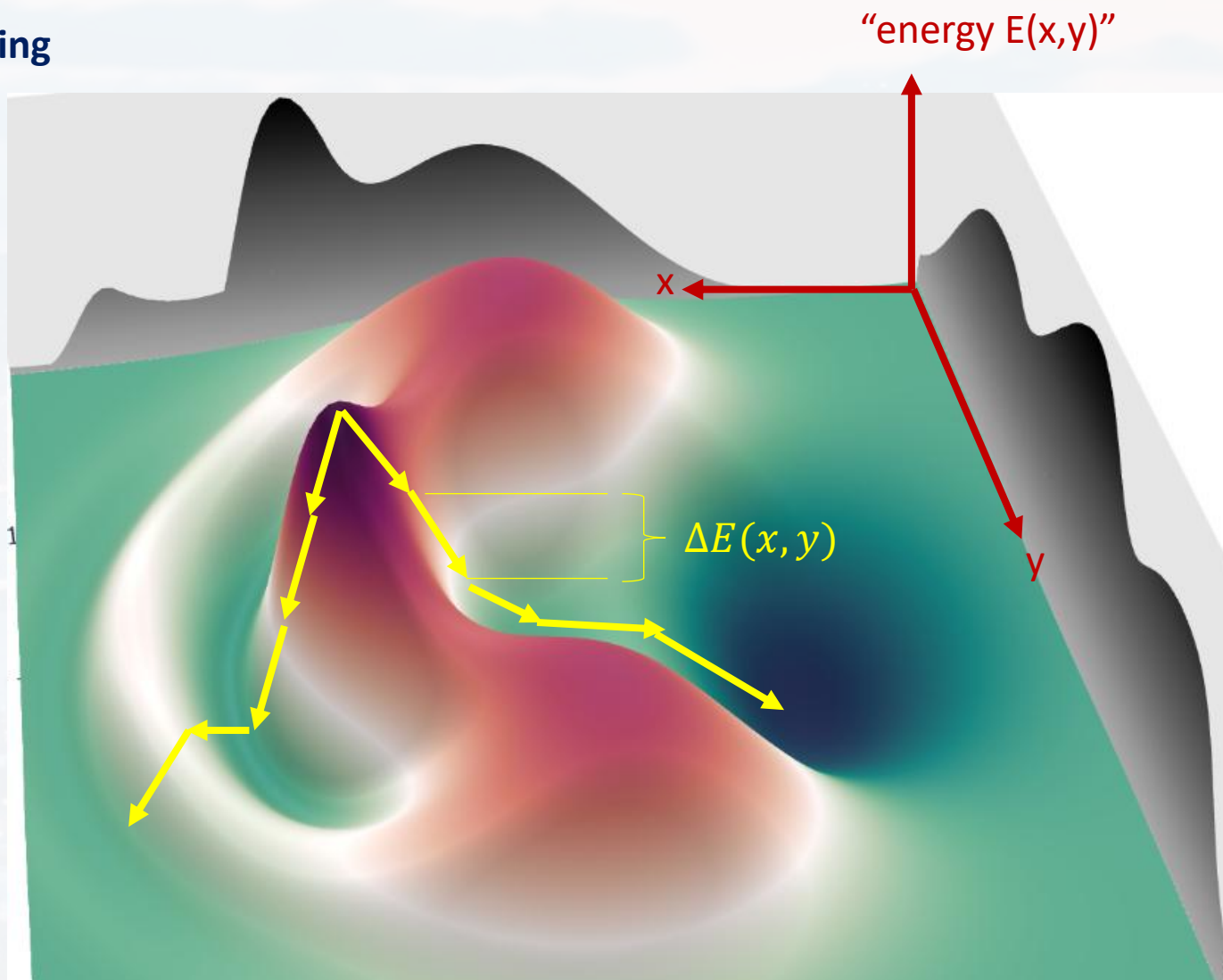
**idea:** drawing from **a proposal distribution**  
**accept/reject the proposal** based on an acceptance probability

**application:** in Physics: often **energy-based models**  
sampling from  $p_i(x|\vartheta) \sim \exp\{-E(x, \vartheta)\}$   
essentially samples the **partition function**  $\mathcal{Z}$



Any algorithm needs a “goal” aka **objective function** that has to be **optimized** (finding an **extreme**)

**example:**  
**simulated annealing**



If  $\Delta E(x, y)$  is **negative**:  
→ **always move**  
(a ball always rolls down the hill)

If  $\Delta E(x, y)$  is **positive**:  
→ calculate the **probability to move**  
→ leaves some chance to escape local minimum

$T$ : temperature

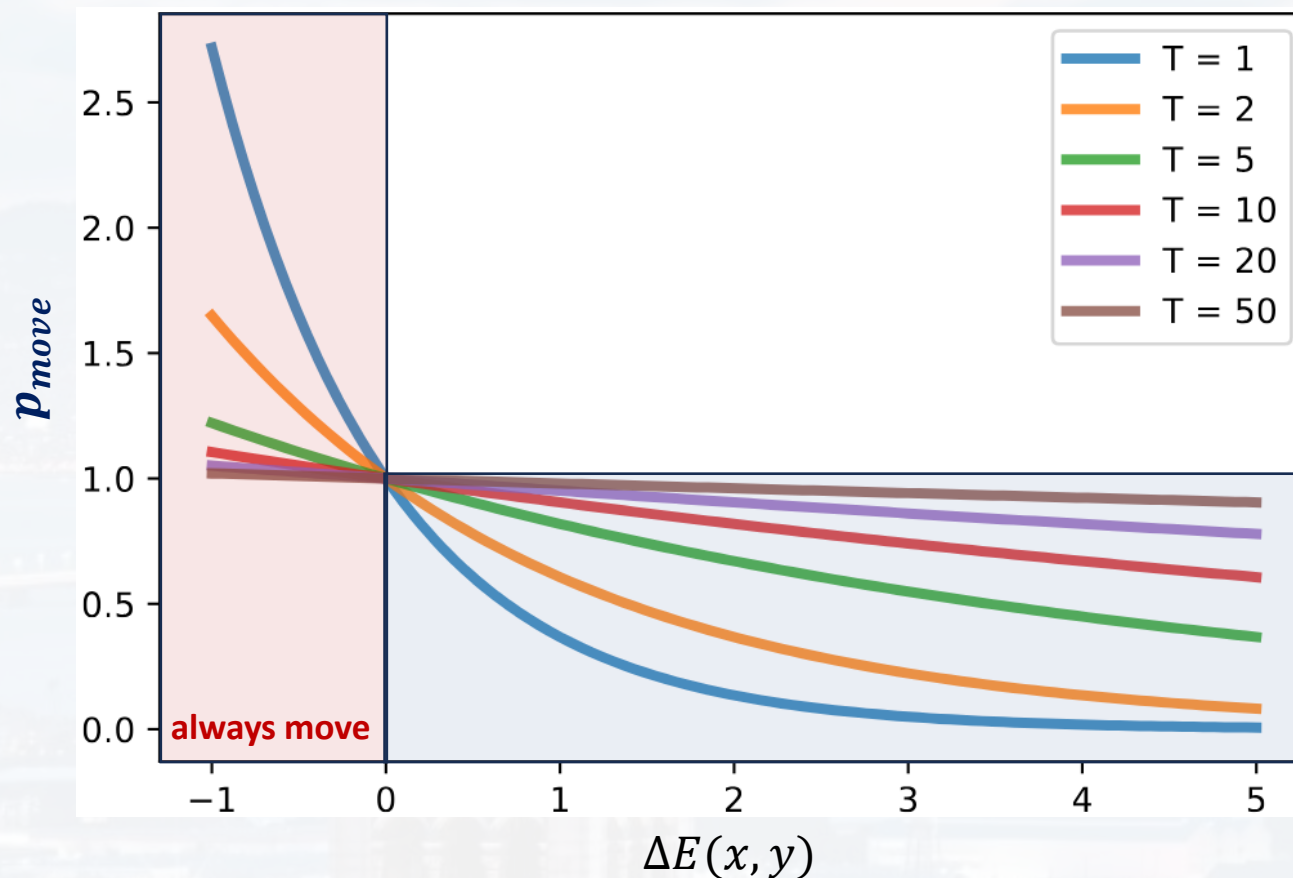
Boltzmann factor

$$p_{move} \sim \exp \left[ -\frac{\Delta E(x, y)}{T} \right]$$



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If  $\Delta E(x, y)$  is **negative**:  
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If  $\Delta E(x, y)$  is **positive**:  
→ calculate the **probability to move**  
→ leaves some chance to escape local minimum

$T$ : temperature

Boltzmann factor

$$p_{move} \sim \exp \left[ -\frac{\Delta E(x, y)}{T} \right]$$

slowly reducing  $T$  → making larger jumps ( $\Delta E(x, y)$ ) less likely over time





Any algorithm needs a “goal” aka **objective function** that has to be **optimized** (finding an **extreme**)

**Metropolis:**

1) suggest a random move  $\Delta\vec{r}$

2) calculate  $\Delta E = E(\vec{r}) - E(\vec{r} + \Delta\vec{r})$

3) move or not:

a) move if  $\Delta E < 0$

b) if  $\Delta E > 0$

- draw a **random number**  $\rho$  from a **uniform distribution** in the interval  $(0, 1)$

- move if  $\rho < \exp\left[-\frac{\Delta E}{T}\right]$

4) reduce  $T$  and repeat



Any algorithm needs a “goal” aka **objective function** that has to be *optimized* (finding an **extreme**)

**Metropolis:** 1) suggest a random move  $\Delta\vec{r}$

2) calculate  $\Delta E = E(\vec{r}) - E(\vec{r} + \Delta\vec{r})$

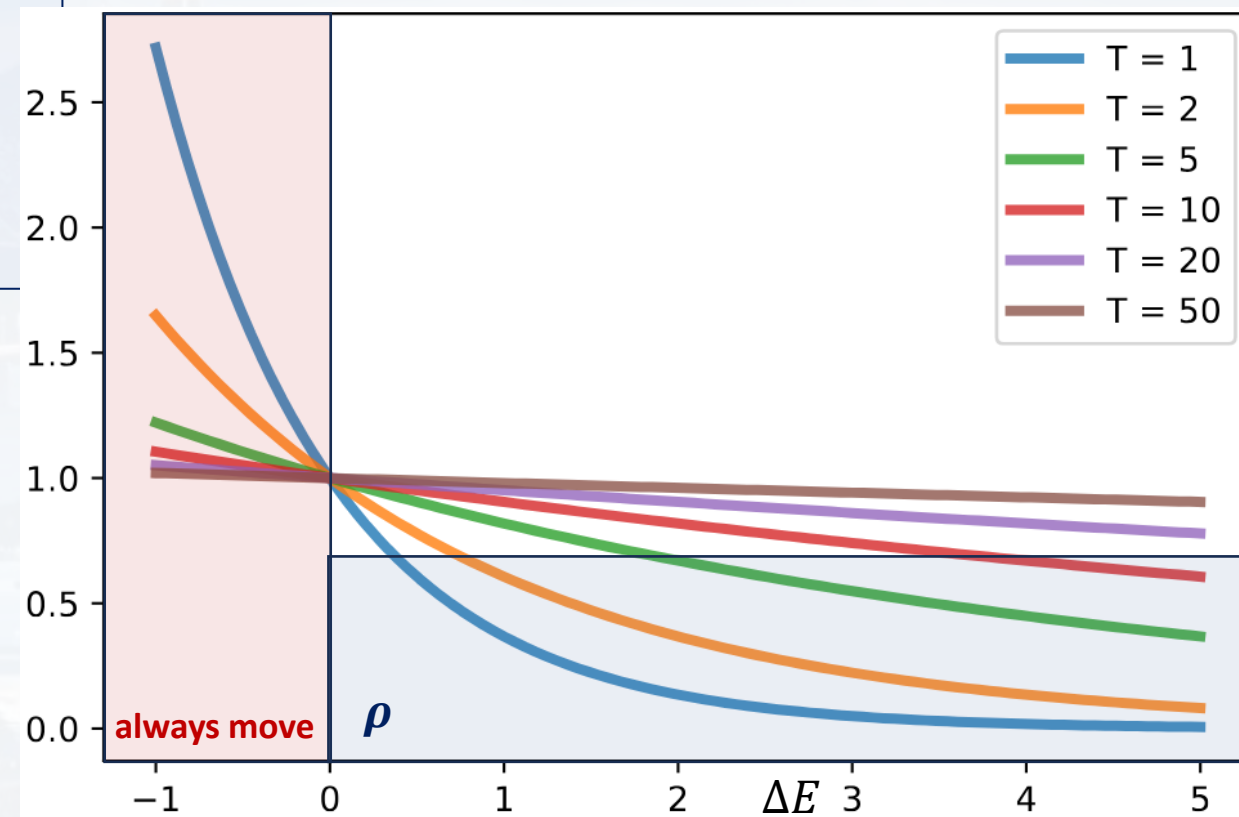
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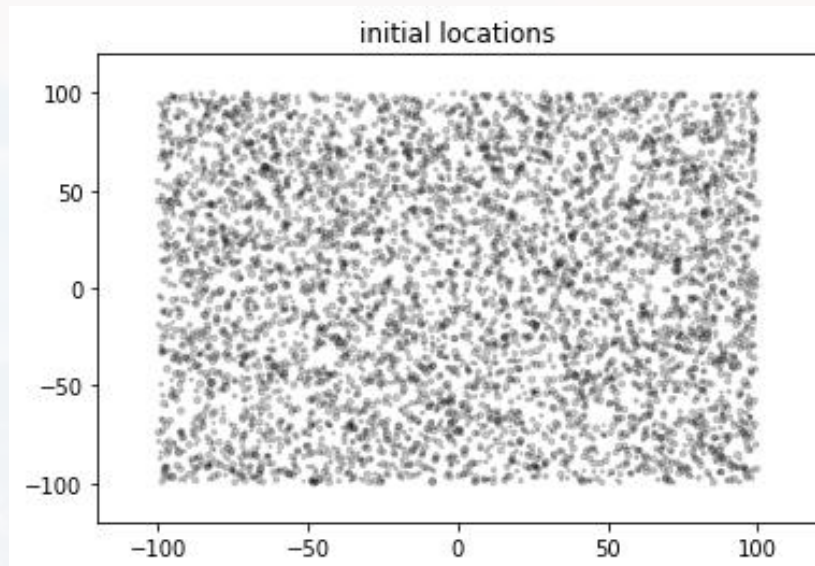
- draw a **random number**  $\rho$  from a  
**uniform distribution** in the interval  $(0, 1)$   
- move if  $\rho < \exp\left[-\frac{\Delta E}{T}\right]$

4) reduce  $T$  and repeat





example: alternative to finite differences

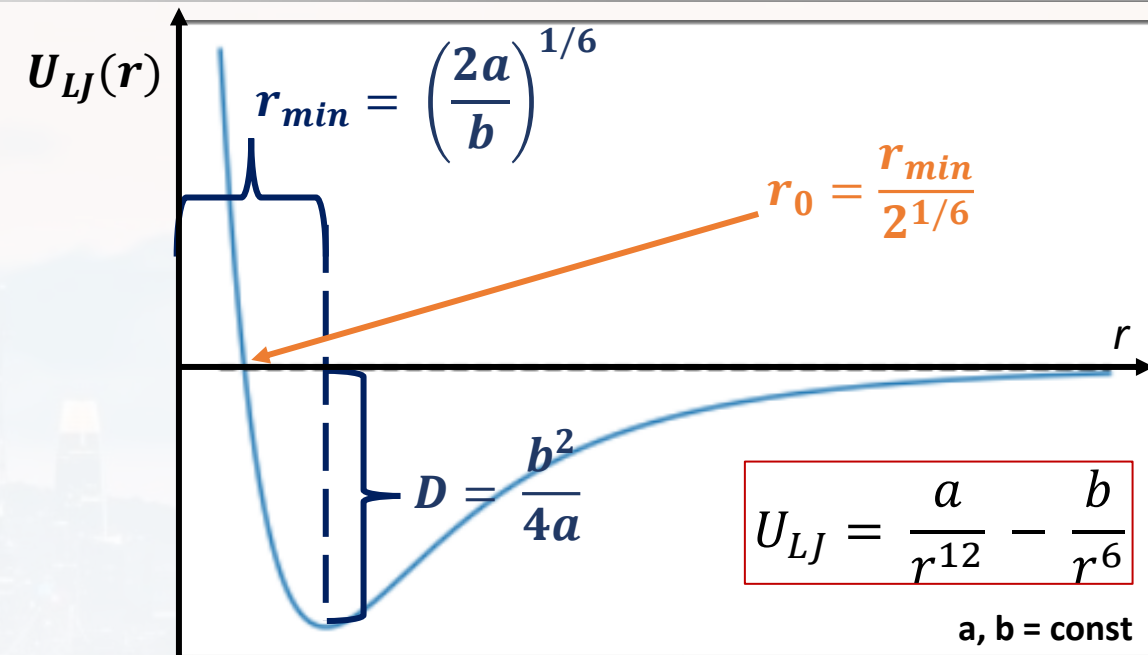


simulating many particles

Naïve solution: solving **Newton's equation of motion**

$$x_{t+\Delta t} = x_t + v(x)_t \cdot \Delta t + \frac{1}{2} a(x)_t \Delta t^2$$

$$y_{t+\Delta t} = y_t + v(y)_t \cdot \Delta t + \frac{1}{2} a(y)_t \Delta t^2$$



total force/potential  
that acts on the particle

$$a(x) = \frac{F(x)_{tot}}{m} = \frac{1}{m} \frac{\partial U_{tot}(x, y)_{LJ}}{\partial x}$$

$$a(y) = \frac{F(y)_{tot}}{m} = \frac{1}{m} \frac{\partial U_{tot}(x, y)_{LJ}}{\partial y}$$



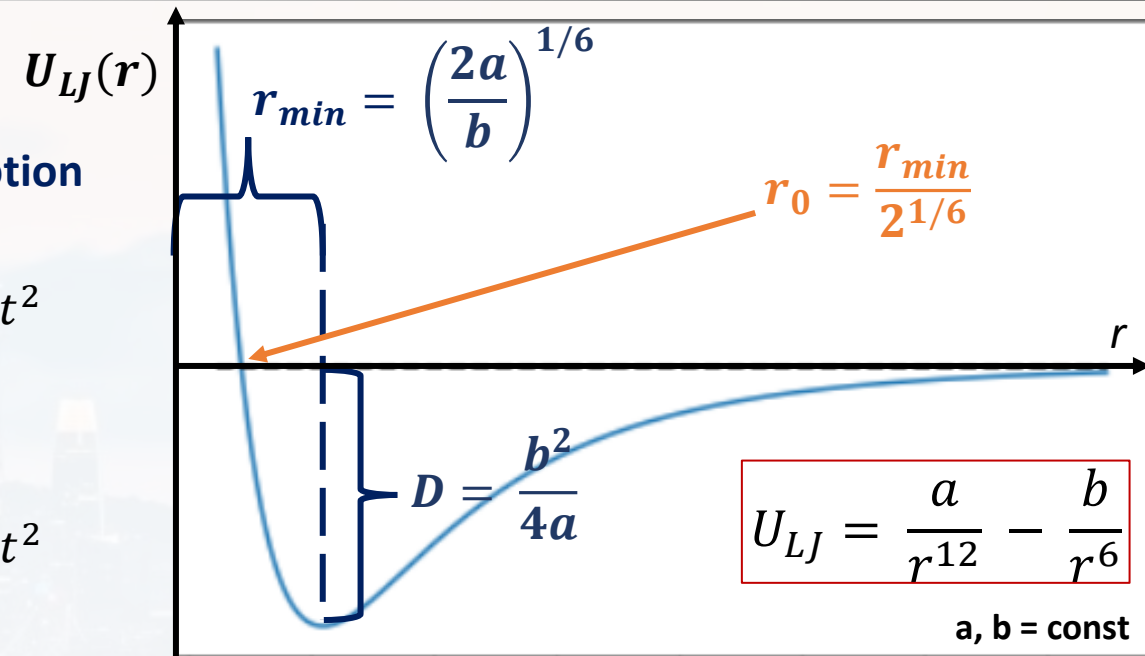


example: alternative to finite differences

Naïve solution: solving **Newton's equation of motion**

$$x_{t+\Delta t} = x_t + v(x)_t \cdot \Delta t + \frac{1}{2m} \frac{\partial U_{tot}(x, y)_{LJ}}{\partial x} \bigg|_t \Delta t^2$$

$$y_{t+\Delta t} = y_t + v(y)_t \cdot \Delta t + \frac{1}{2m} \frac{\partial U_{tot}(x, y)_{LJ}}{\partial y} \bigg|_t \Delta t^2$$



We pick a specific value for  $\Delta t$  and update locations, velocities and acceleration

for particles with  $r \approx r_0$ :

- $\frac{\partial U(x,y)_{LJ}}{\partial y}$  or  $\frac{\partial U(x,y)_{LJ}}{\partial x}$  explode
- particles get kicked out
- wouldn't have gotten so close in the first place  $\rightarrow \Delta t$  **too large**

for particles with  $r \gg r_0$ :

- $\frac{\partial U(x,y)_{LJ}}{\partial y}$  or  $\frac{\partial U(x,y)_{LJ}}{\partial x} \approx 0$
- nothing happens, very inefficient  $\rightarrow \Delta t$  **too small**



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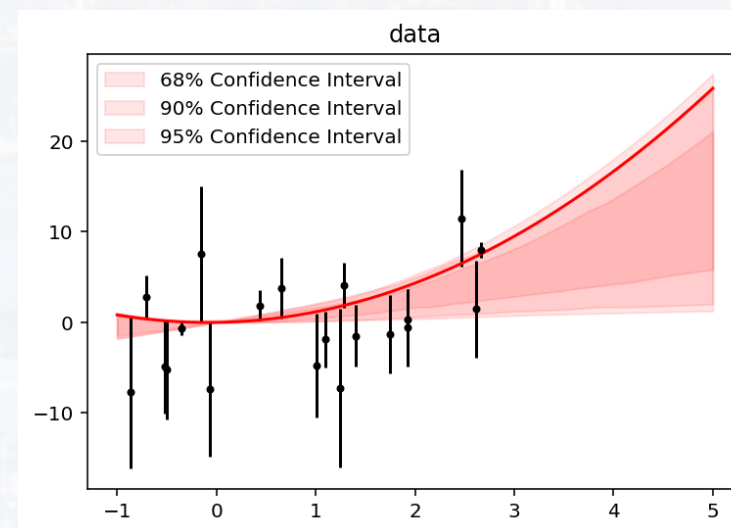
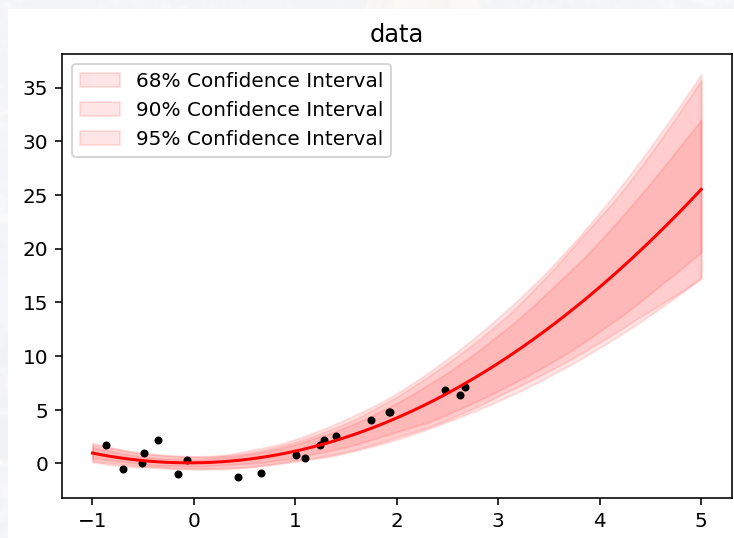
**Bootstrapping**



**problem:** sometimes datapoints  $x_i$  don't have error bars

**idea:** → sampling from  $x_i$  in order to estimate **confidence intervals**  
 $x_i | model \sim ?$

if  $x_i$  have error bars  $\sigma_i$ , usually sampling from normal distribution  
 $x_i | \sigma \sim \mathcal{N}(\bar{x}_i, \sigma_i^2)$







Physics 77/88

```
from scipy.optimize import curve_fit
```

```
def fun_to_fit(x, a, b, c):  
    return a*x**2 + b*x + c
```

```
ValsBest, Cov = curve_fit(fun_to_fit, x, y)
```



Physics 77/88

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```
def fun_to_fit(x, a, b, c):  
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```

```
ValsBest, Cov = curve_fit(fun_to_fit, x, y)
```

bootstrapping:

if **no errors of  $y_i$  known**

→ assuming that fitted parameters follow a **normal distribution**, i.e.  $a = 1.37 \pm 0.39$  where  $\mu_a = 1.37$  and  $\sigma_a = 0.39$  and so on...

→ **varying** the parameters **within their errors** using `np.random.normal( $\mu_a$ ,  $\sigma_a$ , N)` **N** times

→ for each **N**, **generating a curve fit**

→ from set of N curve fits → **calculating percentiles** for confidence band/ interval



Physics 77/88

```
from scipy.optimize import curve_fit
```

```
def fun_to_fit(x, a, b, c):  
    return a*x**2 + b*x + c
```

```
ValsBest, Cov = curve_fit(fun_to_fit, x, y)
```

bootstrapping:

if errors of  $y_i$  known

→ assuming that errors of  $y_i$  follow a **normal distribution**, i.e.  $y_i(\text{boot}) = y_i \pm \sigma_i$

→ **varying** all  $y_i$  **within their errors** using `np.random.normal(y_i,  $\sigma_i$ , N)`  
N times

→ for each N, **generating a curve fit**

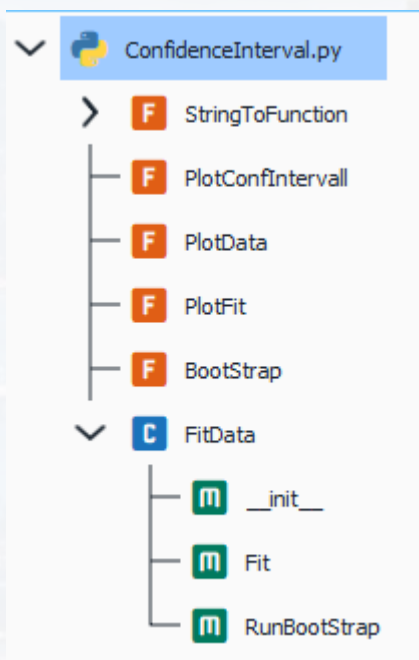
→ from set of N curve fits → **calculating percentiles** for confidence band/ interval





## bootstrapping

## ConfidenceInterval.py



### USAGE:

```

...generating a test sample:
....
....x.....=np.linspace(-1,3,20)
....err.....=np.random.normal(0,1,(len(x),))#1sigma errorbars
....y.....=x**2+.err
....errorbars=abs(err)
....
....
....1).plotting data
.....
....F1=FitData(x,y)
....F2=FitData(x,y,errorbars)
....F3=FitData(x,y,errorbars,time='[s]',pressure='[MPa]')
....
....
....2).fitting data (returns best values of fitted params, 1sigma confidence and
.....reduced chi2 if errorbars given, MSE else)
.....
....res1=F1.Fit()
....res2=F2.Fit()
....res2=F3.Fit()
....
....res12=F1.Fit("a*x**2",[1],(-0.5,10))
....
....
....3).Bootstrapping (either varying within errorbars or within conf of fitted
.....params)
.....
....F1.RunBootStrap()
....F2.RunBootStrap()
....F3.RunBootStrap()
....
....F1.RunBootStrap(100,[90,95],np.linspace(-1,5,200))
....F3.RunBootStrap(100,[90,95],np.linspace(-1,5,200))
....

```



Thank you very much for your attention!

## CHECKING IN

