

# BESIII Physics & Software Meeting

## Phase-space binned analysis of strong-phase parameters in $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

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# Outline

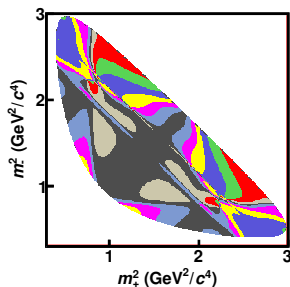
- 1 Introduction
- 2 Binning scheme
- 3 Formalism and measurement strategy
- 4 Selection
- 5 Fit of double tag yields
- 6 Likelihood fit
- 7 Input-output checks
- 8 Fit results
- 9 Bonus measurements
- 10 Systematic uncertainties
- 11 Summary and conclusion

# Introduction

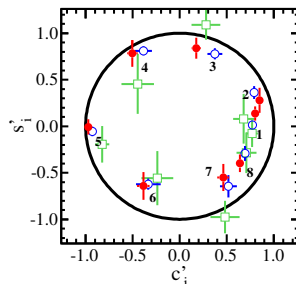
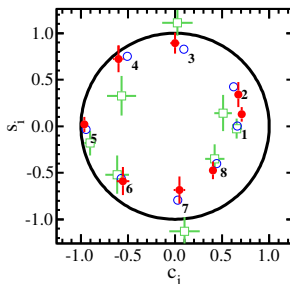
- Aim of analysis: Measure the strong-phase difference between  $D^0$  and  $\bar{D}^0 \rightarrow K^+ K^- \pi^+ \pi^-$  using the  $8 \text{ fb}^{-1} \psi(3770)$  dataset
- Important input to:
  - ① Measurement of the CKM angle  $\gamma$
  - ② Studies of charm mixing and CPV
  - ③ Strong-phase measurements of other decay modes
- Phase-space integrated measurement: [Phys. Rev. D \*\*107\*\* 032009](#)
- Strong phases are diluted when integrated over the 5D phase space  
 $\implies$  Perform analysis in bins
- Binned model-dependent analysis of  $\gamma$  has already been performed at LHCb: [Eur. Phys. J. C \*\*83\*\*, 547 \(2023\)](#)
  - BESIII measurement will allow for a model-independent  $\gamma$  update

# Introduction - Previous strong-phase analyses

- Golden mode for strong-phase analysis:  $D^0 \rightarrow K_{S,L}^0 \pi^+ \pi^-$
- Measurements of  $c_i$  and  $s_i$  performed in  $2 \times 8$  bins:
  - CLEO: [Phys. Rev. D \*\*82\*\* 112006](#)
  - BESIII: [Phys. Rev. D \*\*101\*\* 112002](#)



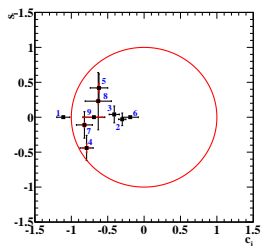
(a) Equal  $\delta$  binning scheme



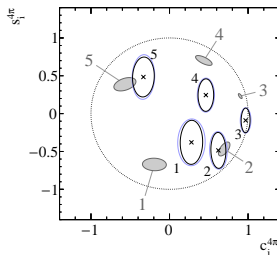
(b)  $c_i$  and  $s_i$  for (left)  $K_S \pi \pi$  and (right)  $K_L \pi \pi$

# Introduction - Previous strong-phase analyses

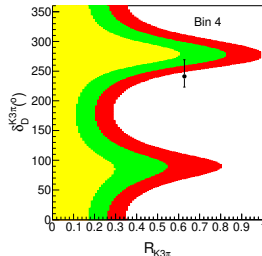
- Four-body decays have a 5D phase space, but their binning scheme may be defined analogously
- Many different strategies for 5D binning schemes:
  - $D \rightarrow K_S^0 \pi^+ \pi^- \pi^0$ : [JHEP 2018 82 \(2018\)](#) (CLEO-c data)
  - $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ : [JHEP 2018 144 \(2018\)](#) (CLEO-c data)
  - $D \rightarrow K^- \pi^+ \pi^- \pi^+$ : [JHEP 2021 164 \(2021\)](#) (BESIII collaboration)



(a)  $K_S \pi \pi \pi^0$



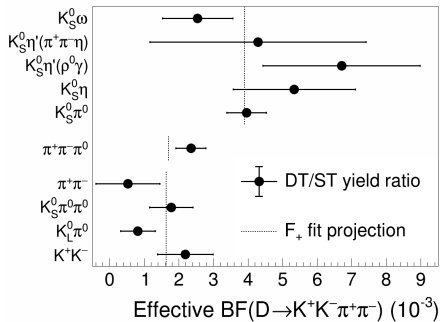
(b)  $4\pi$



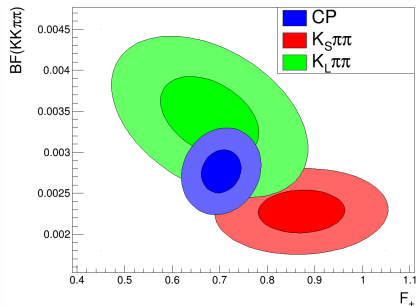
(c)  $K \pi \pi \pi$

# Introduction - Previous strong-phase analyses

- The phase-space integrated strong phase of  $D \rightarrow K^+ K^- \pi^+ \pi^-$  was previously measured ([Phys. Rev. D \*\*107\*\* 032009](#)):  $F_+ = 0.73 \pm 0.04$
- The asymmetry in the branching fraction measured using CP even and CP odd tags is sensitive to the CP-even fraction  $F_+$ 
  - $c_i = 2F_+ - 1$  is the amplitude-average cosine of the strong phase



(a) BF asymmetry



(b)  $F_+$  combination

## Strong phases are determined from yield asymmetries

- When integrating over a phase-space region, variations in strong phases dilute the asymmetries
- We require a binning scheme to minimise the dilution and enhance asymmetry effects

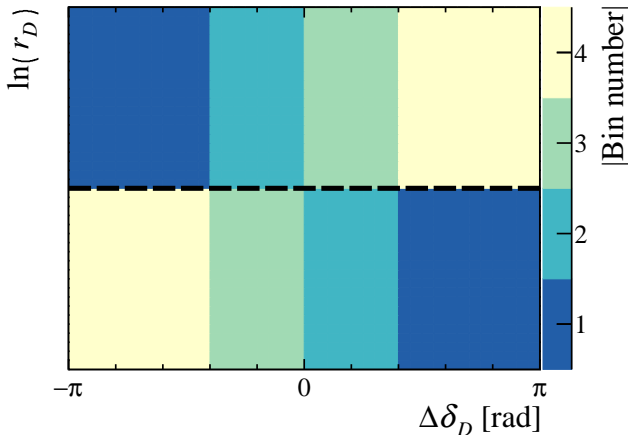
## How to bin a 5-dimensional phase space?

- For each  $D$  event, use LHCb model to calculate

$$\frac{\mathcal{A}(D^0)}{\mathcal{A}(\bar{D}^0)} = r_D e^{i\delta_D}$$

- Bin along  $\delta_D$  and  $r_D$ 
  - Bin boundaries in  $\delta_D$  are moved to maximise sensitivity to  $\gamma$
- The  $KK\pi\pi$  binning scheme has been fixed by LHCb analysis of  $\gamma$  in [Eur. Phys. J. C \*\*83\*\*, 547 \(2023\)](#)

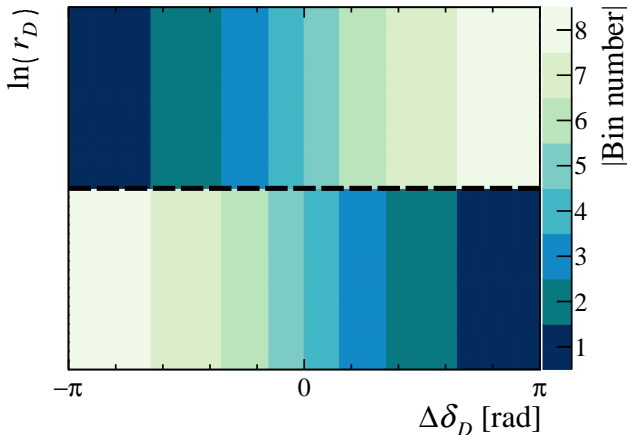
# $KK\pi\pi$ binning scheme



Bins  $i < 0$  on top,  $i > 0$  below  
Binning scheme used in this analysis



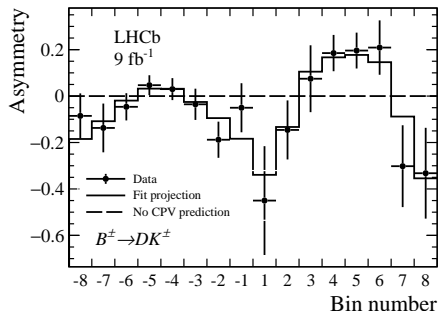
# $KK\pi\pi$ binning scheme



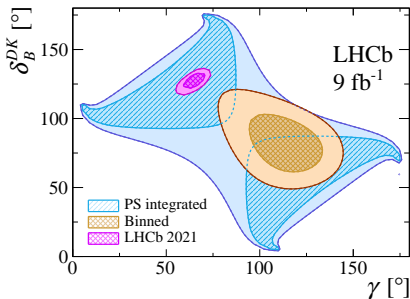
Bins  $i < 0$  on top,  $i > 0$  below

Possible binning scheme for future analysis with  $20 \text{ fb}^{-1}$  dataset

# Model-dependent LHCb measurement



(a) Bin asymmetries in  $B^\pm \rightarrow DK^\pm$



(b) Measurement of  $\gamma$  at LHCb

$$\gamma = (116_{-14}^{+12})^\circ$$

This result is currently model independent, and needs external inputs from BESIII to become a proper model-independent measurement!

$D\bar{D}$  pair from  $\psi(3770)$  is prepared in a  $\mathcal{C} = -1$  state

- $D$  mesons are “quantum correlated”
- The decay of  $D \rightarrow K^+ K^- \pi^+ \pi^-$  is correlated by the  $CP$  content of the tag mode

$$|D\bar{D}\rangle = \frac{1}{\sqrt{2}}(|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle)$$

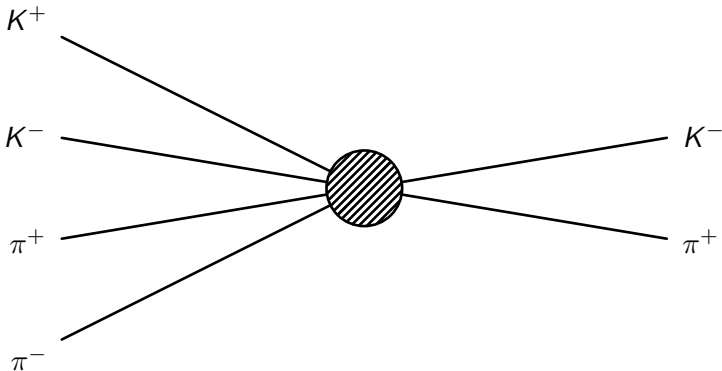
- Equivalently, in terms of  $CP$  eigenstates:

$$|D\bar{D}\rangle = \frac{1}{\sqrt{2}}(|D_-\rangle|D_+\rangle - |D_+\rangle|D_-\rangle)$$

# Formalism and measurement strategy

- Tag mode can be a flavour tag

- $K\pi$ ,  $K\pi\pi^0$ ,  $K\pi\pi\pi$ ,  $Ke\nu$

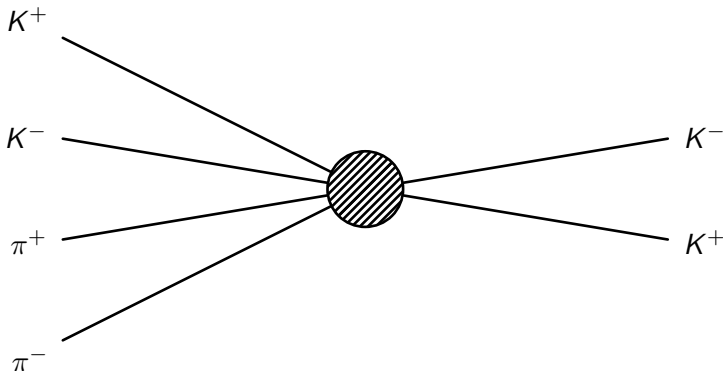


Use flavour tags to measure fraction of  $D^0 \rightarrow KK\pi\pi$  decays in bin  $i$ :

$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B} \times K_i$$

# Formalism and measurement strategy

- Tag mode can be a CP even tag
  - $KK$  (fully and part. reco  $KK\pi\pi$ ),  $\pi\pi$ ,  $\pi\pi\pi^0$ ,  $K_S\pi^0\pi^0$ ,  $K_L\pi^0$

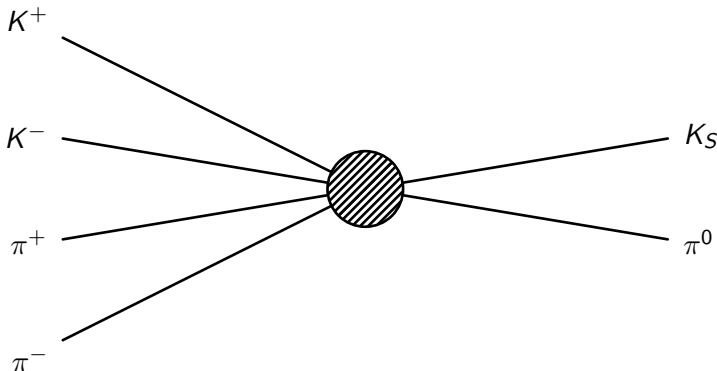


$D \rightarrow K^+K^-$ , which is  $CP$  even, forces  $D \rightarrow K^+K^-\pi^+\pi^-$  to be  $CP$  odd:

$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B} \times (K_i + K_{-i} - 2\sqrt{K_i K_{-i}} c_i)$$

# Formalism and measurement strategy

- Tag mode can be a CP odd tag
  - $K_S\pi^0$  (fully and part.reco  $KK\pi\pi$ ),  $K_S\eta$ ,  $K_S\eta'(\pi\pi\eta, \pi\pi\gamma)$ ,  $K_S\pi\pi\pi^0$

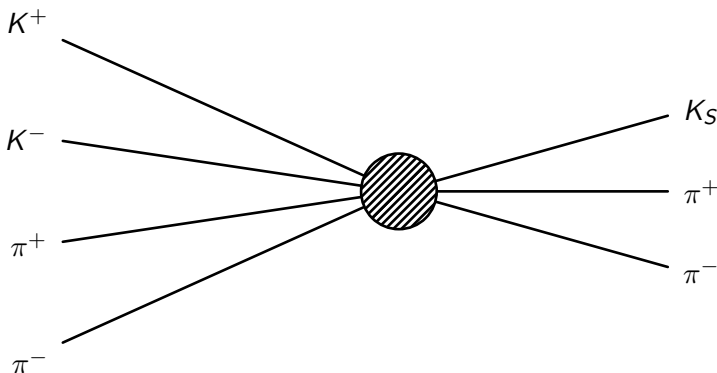


$D \rightarrow K_S^0\pi^0$ , which is  $CP$  odd, forces  $D \rightarrow K^+K^-\pi^+\pi^-$  to be  $CP$  even:

$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B} \times (K_i + K_{-i} + 2\sqrt{K_i K_{-i}} c_i)$$

# Formalism and measurement strategy

- Tag mode can be a self-conjugate multi-body tag
  - $K_S \pi \pi$  (fully and part.reco  $KK\pi\pi$ ),  $K_L \pi \pi$

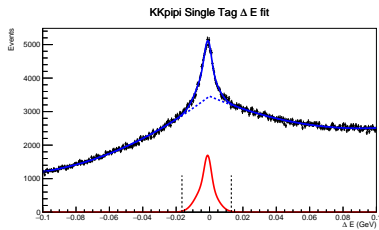


$D \rightarrow K_S^0 \pi^+ \pi^-$  has different strong phases in different bins of phase space:

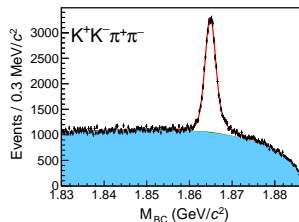
$$\frac{N_{ij}^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B} \times (K_i K'_{-j} + K_{-i} K'_j - 2\sqrt{K_i K_{-i} K'_j K'_{-j}}(c_i c'_j + s_i s'_j))$$

# Selection

- Double-tag analysis: Select  $D \rightarrow KK\pi\pi$  events tagged with flavour,  $CP$  and multi-body tags
- Selection more or less identical to previous double-tag analyses
- $D \rightarrow KK\pi\pi$  selection:
  - 4 good charged tracks
  - $3\sigma$  window around signal peak in  $\Delta E$
  - Asymmetric  $K_S KK$  veto for  $m(\pi^+\pi^-) \in [477, 507]$  MeV
  - Flight significance cut at 2



(a)  $\Delta E$  window

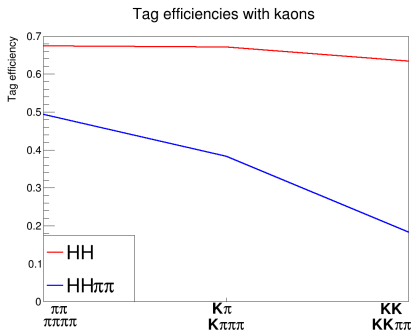


(b) Single tag yield:  $29\,227 \pm 268$



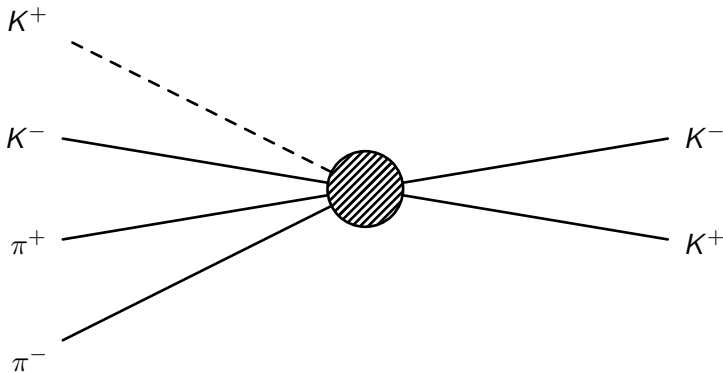
# Partially reconstructed $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

- The reconstruction efficiency of  $D \rightarrow KK\pi\pi$  is less than 20%
- For comparison, the efficiency of  $D \rightarrow \pi\pi\pi\pi$  is around 50%!
- Poor kaon tracking efficiency at low momentum
  - (Can be recovered in reconstruction, e.g. by requiring fewer hits on tracks? This was done at CLEO-c)



# Partially reconstructed $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

- Solution: Only reconstruct 3 of the charged  $D$  daughters
  - Presence of missing kaon is inferred from the missing momentum
  - Yields are similar to fully reconstructed sample
  - Large, but non-peaking background from  $D \rightarrow K \pi \pi \pi \pi^0$

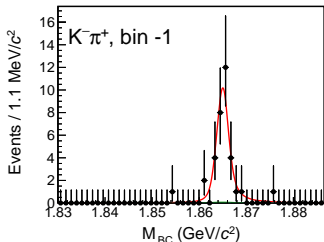


Use this technique with the  $K^+ K^-$ ,  $K_S^0 \pi^0$  and  $K_S^0 \pi^+ \pi^-$  tags

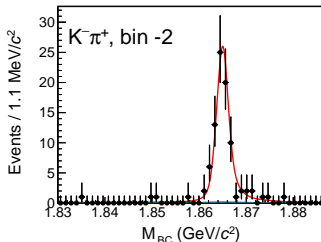
# Double tag fits

- Fit strategy: Only fit signal side  $M_{BC}$  because of low statistics
- Fit model:
  - Signal: PDF from signal MC, convolved with a Gaussian
  - Flat background: Argus PDF
  - Peaking background: Shape and efficiency from MC, correct for quantum correlation
  - Simple sideband subtraction for correct signal but wrong tag event
- Fit all bins simultaneously
  - Shape is floated and shared across all bins
  - Yield of signal and combinatorial background is floated in each bin

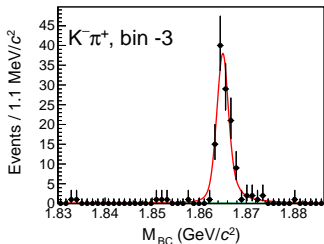
# Double tag fit of $KK\pi\pi$ vs $K\pi$



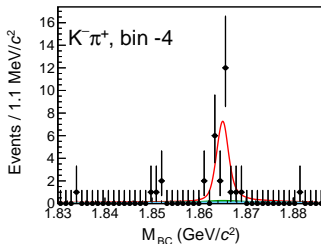
(a) Bin -1 yield:  $32.4^{+6.2}_{-5.5}$



(b) Bin -2 yield:  $82.7^{+9.7}_{-9.0}$

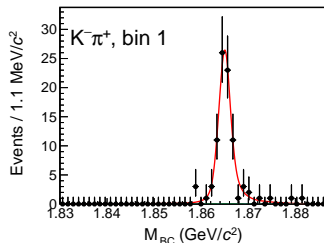


(c) Bin -3 yield:  $120.3^{+11.6}_{-10.9}$

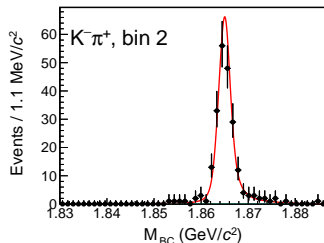


(d) Bin -4 yield:  $22.4^{+5.2}_{-4.6}$

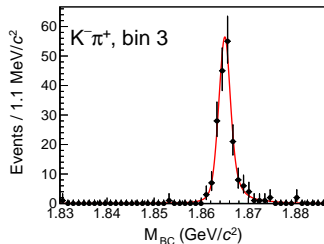
# Double tag fit of $KK\pi\pi$ vs $K\pi$



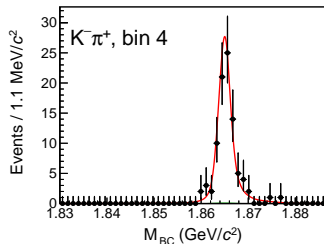
(a) Bin 1 yield:  $84.5^{+9.8}_{-9.1}$



(b) Bin 2 yield:  $211.2^{+15.4}_{-14.8}$

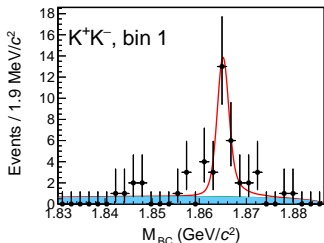


(c) Bin 3 yield:  $181.0^{+14.0}_{-13.3}$

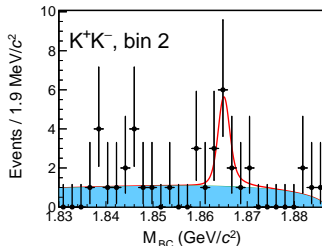


(d) Bin 4 yield:  $88.6^{+9.7}_{-9.0}$

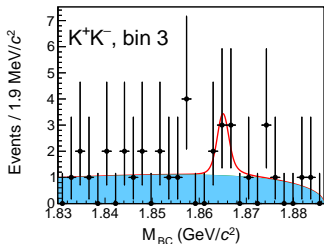
# Double tag fit of $KK\pi\pi$ vs $KK$



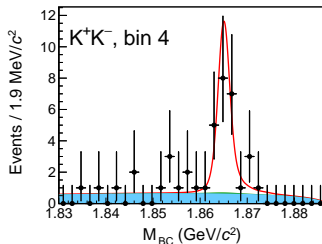
(a) Bin 1 yield:  $25.3^{+6.2}_{-5.5}$



(b) Bin 2 yield:  $8.8^{+4.0}_{-3.3}$

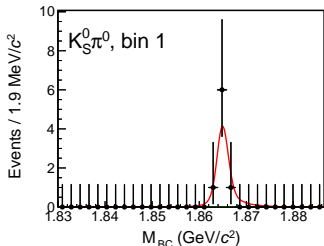


(c) Bin 3 yield:  $4.5^{+3.3}_{-2.6}$

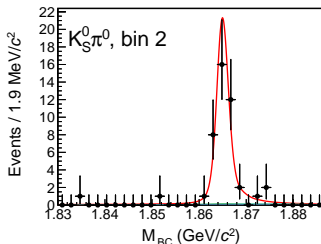


(d) Bin 4 yield:  $21.1^{+5.5}_{-4.8}$

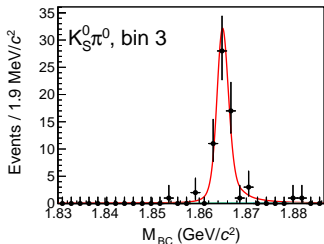
# Double tag fit of $KK\pi\pi$ vs $K_S\pi^0$



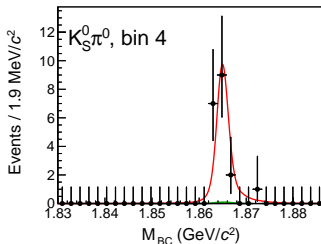
(a) Bin 1 yield:  $7.9^{+3.1}_{-2.5}$



(b) Bin 2 yield:  $40.4^{+6.8}_{-6.3}$

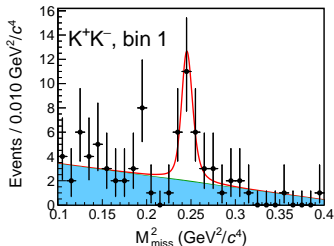


(c) Bin 3 yield:  $61.1^{+8.3}_{-7.8}$

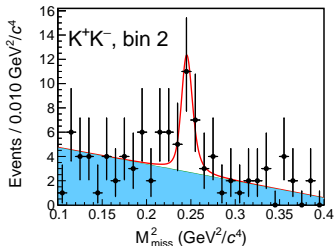


(d) Bin 4 yield:  $18.3^{+4.5}_{-3.9}$

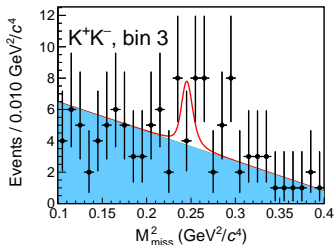
# Double tag fit of partially reconstructed $KK\pi\pi$ vs $KK$



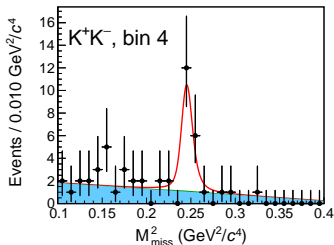
(a) Bin 1 yield:  $19.6^{+5.8}_{-5.1}$



(b) Bin 2 yield:  $17.7^{+6.1}_{-5.3}$



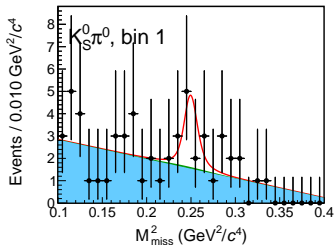
(c) Bin 3 yield:  $7.5^{+6.6}_{-5.7}$



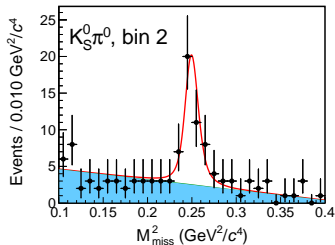
(d) Bin 4 yield:  $17.3^{+5.1}_{-4.4}$



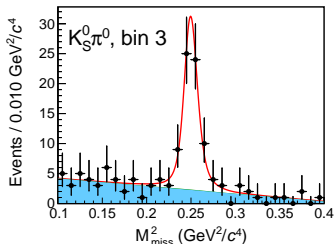
# Double tag fit of partially reconstructed $KK\pi\pi$ vs $K_S\pi^0$



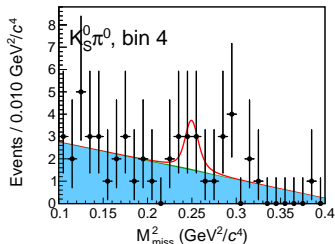
(a) Bin 1 yield:  $7.2^{+4.3}_{-3.5}$



(b) Bin 2 yield:  $39.3^{+8.3}_{-7.3}$



(c) Bin 3 yield:  $64.4^{+9.6}_{-9.0}$



(d) Bin 4 yield:  $4.8^{+3.8}_{-3.0}$

## How to put this all together to determine $c_i$ and $s_i$ ?

- 1 In the past, the strategy has been to measure normalised  $K_i$  first, using flavour tags:

$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} \propto K_i, \quad \sum_i K_i = 1$$

- 2 However, the  $K_i$  must be corrected for efficiencies and bin migrations, in addition to Doubly Cabibbo Suppressed decays:

$$f_i = \left( 1 + r_D^2 \frac{K_{-i}}{K_i} - 2\sqrt{\frac{K_{-i}}{K_i}} (c_i \cos(\delta_D) + s_i \sin(\delta_D)) \right)^{-1}$$

- 3 Finally, with the  $K_i$  fixed,  $c_i$  and  $s_i$  are determined from CP and multi-body tags:

$$\begin{aligned} \frac{N_i^{\text{DT}}}{N^{\text{ST}}} &\propto K_i + K_{-i} \mp 2\sqrt{K_i K_{-i}} c_i \\ \frac{N_{ij}^{\text{DT}}}{N^{\text{ST}}} &\propto K_i K'_{-j} + K_{-i} K'_j - 2\sqrt{K_i K_{-i} K'_j K'_{-j}} (c_i c'_j + s_i s'_j) \end{aligned}$$

We propose a simpler, but more powerful strategy:

- 1 Treat all  $K_i$ ,  $c_i$  and  $s_i$  as free parameters
- 2 Include the  $D \rightarrow KK\pi\pi$  branching fraction  $\mathcal{B}$  as a free parameter
- 3 Fit flavour,  $CP$  and multi-body tags simultaneously

Master equations:

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ij} \left( K_{-j} + r_D^2 K_j - 2r_D \sqrt{K_j K_{-j}} (c_j \cos(\delta_D) + s_j \sin(\delta_D)) \right)$$

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ij} (K_j + K_{-j} \mp 2\sqrt{K_j K_{-j}} c_j)$$

$$\hat{N}_{ij}^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ijkl} (K_k K'_{-l} + K_{-k} K'_l - 2\sqrt{K_k K_{-k} K'_l K'_{-l}} (c_k c'_l + s_k s'_l))$$

$\epsilon_{ij}$  are combined efficiency and bin migration matrices

The master equations use the free parameters  $\mathcal{B}$ ,  $K_i$ ,  $c_i$  and  $s_i$  to make a prediction  $\hat{N}^{\text{DT}}$  to the measured double tag yields  $N^{\text{DT}}$

## Master equations

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B}_{\epsilon ij} \left( K_{-j} + r_D^2 K_j - 2r_D \sqrt{K_j K_{-j}} (c_j \cos(\delta_D) + s_j \sin(\delta_D)) \right)$$

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B}_{\epsilon ij} (K_j + K_{-j} \mp 2\sqrt{K_j K_{-j}} c_j)$$

$$\hat{N}_{ij}^{\text{DT}} = N^{\text{ST}} \mathcal{B}_{\epsilon ijkl} (K_k K'_{-l} + K_{-k} K'_l - 2\sqrt{K_k K_{-k} K'_l K'_{-l}} (c_k c'_l + s_k s'_l))$$

Ordinarily, we would construct a Gaussian (log)likelihood function  $\Rightarrow$   
Obtain  $\mathcal{B}$ ,  $K_i$ ,  $c_i$  and  $s_i$  by minimising the following function:

$$-\ln(\mathcal{L}) = \frac{1}{2} \sum_{\text{Tag}} \sum_{ij} (V^{-1})_{ij} (N_i^{\text{DT}} - \hat{N}_i^{\text{DT}}) (N_j^{\text{DT}} - \hat{N}_j^{\text{DT}})$$

$$V_{ij} = \rho_{ij} \sigma_i \sigma_j$$

---

<sup>0</sup> $\rho$  are correlation coefficients

## Master equations

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B}_{\epsilon ij} \left( K_{-j} + r_D^2 K_j - 2r_D \sqrt{K_j K_{-j}} (c_j \cos(\delta_D) + s_j \sin(\delta_D)) \right)$$

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B}_{\epsilon ij} (K_j + K_{-j} \mp 2\sqrt{K_j K_{-j}} c_j)$$

$$\hat{N}_{ij}^{\text{DT}} = N^{\text{ST}} \mathcal{B}_{\epsilon ijkl} (K_k K'_{-l} + K_{-k} K'_l - 2\sqrt{K_k K_{-k} K'_l K'_{-l}} (c_k c'_l + s_k s'_l))$$

Our DT yields are very small, so their uncertainties are asymmetric  $\implies$   
 Approximate covariance matrix from the asymmetric uncertainties<sup>1</sup>:

$$-\ln(\mathcal{L}) = \frac{1}{2} \sum_{\text{Tag}} \sum_{ij} (V^{-1})_{ij} (N_i^{\text{DT}} - \hat{N}_i^{\text{DT}}) (N_j^{\text{DT}} - \hat{N}_j^{\text{DT}})$$

$$V_{ij} = \rho_{ij} \sigma_i \sigma_j, \quad \sigma = \sqrt{\sigma_- \sigma_+ - (\sigma_+ - \sigma_-)(N^{\text{DT}} - \hat{N}^{\text{DT}})}$$

<sup>1</sup>[arXiv:physics/0406120](https://arxiv.org/abs/physics/0406120)

$$-\ln(\mathcal{L}) = \frac{1}{2} \sum_{\text{Tag}} \sum_{ij} (V^{-1})_{ij} (N_i^{\text{DT}} - \hat{N}_i^{\text{DT}})(N_j^{\text{DT}} - \hat{N}_j^{\text{DT}})$$

$$V_{ij} = \rho_{ij} \sigma_i \sigma_j, \quad \sigma = \sqrt{\sigma_- \sigma_+ - (\sigma_+ - \sigma_-)(N^{\text{DT}} - \hat{N}^{\text{DT}})}$$

- The above likelihood has good coverage for flavour and  $CP$  tags...
- ... but not for multi-body decays
  - Bins with  $\sigma_- \approx 0$  make the fit unstable
  - Fit convergence was found to be less than 60%
- In multi-body decays, use the full unbinned likelihood directly
  - Fit convergence improves to over 95%
  - Much slower, but much more accurate

## One last technical detail...

- 8  $K_i$  parameters, but  $\sum_i K_i = 1$ , so only 7 are independent
- We could set  $K_4 = 1 - \sum_{i \neq 4} K_i$ , but such a parameterisation is unstable due to strong correlations
- Use a recursive fraction parameterisation ([JHEP 2021 169 \(2021\)](#))
- Recursive fractions  $R_i$  are defined as:

$$R_i \equiv \begin{cases} K_i, & i = -4 \\ K_i / \sum_{j \geq i} K_j, & -4 < i < +4 \\ 1, & i = +4. \end{cases}$$

# Parameterisation of $K_i$

Visualisation of this parameterisation:

## 1. Definition of $a'_1$

$$\underbrace{a_1}_{\equiv a'_1} + \underbrace{a_2 + a_3 + \dots}_{\equiv 1 - a'_1} = 1$$

## 2. Eliminate $a_1$

$$a_2 + a_3 + a_4 + \dots = 1 - a'_1$$

## 3. Definition of $a'_2$

$$\underbrace{\frac{a_2}{1 - a'_1}}_{\equiv a'_2} + \underbrace{\frac{a_3 + a_4 + \dots}{1 - a'_1}}_{\equiv 1 - a'_2} = 1$$

Repeat this procedure until the last coefficient, which is 1



# Input-output checks

Input-output checks are performed by generating toy datasets

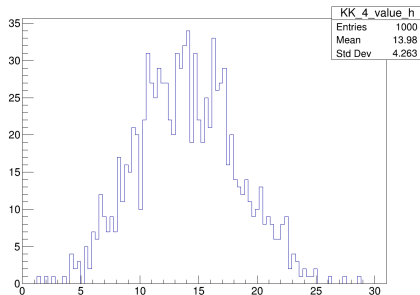
- 1 Check fit convergence
- 2 Check error coverage
- 3 Correct any biases in fitted parameters

Unfortunately, in this analysis we cannot simply generate Poisson-distributed DT yields:

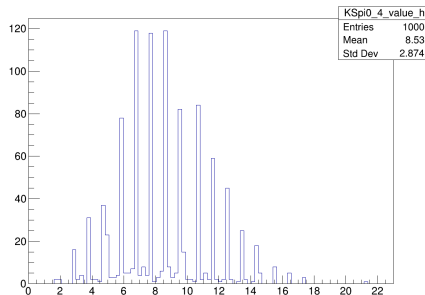
- 1 Asymmetric uncertainties
- 2 Large background-to-signal
- 3 Multi-body tags with low yields require a full unbinned likelihood

Solution: Generate toy datasets for each double tag fit

# Input-output checks



(a) Input yield: 13.9

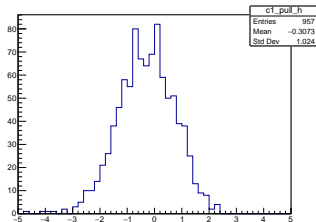


(b) Input yield: 8.6

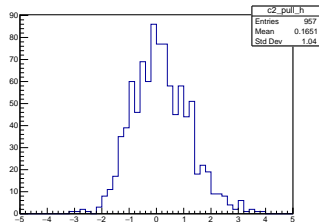
**Figure 12:** Fitted yields in toy datasets for the (left)  $KK$  and (right)  $K_S\pi^0$  tags

- 1 No biases are observed
- 2 Distributions are asymmetric  $\implies$  uncertainties are asymmetric
- 3  $KK$  uncertainty is non-Poisson because of large backgrounds from  $q\bar{q}$
- 4  $K_S\pi^0$  has small backgrounds, so observed yields Poisson distributed

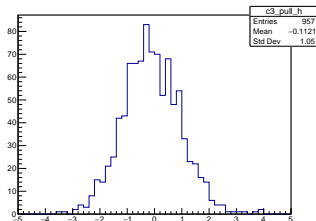
# Input-output checks



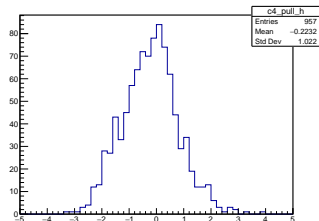
(a)  $c_1$  pulls



(b)  $c_2$  pulls

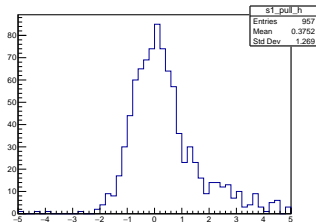


(c)  $c_3$  pulls

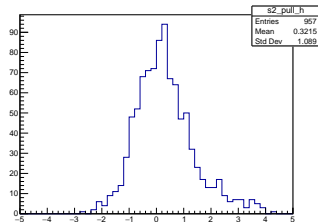


(d)  $c_4$  pulls

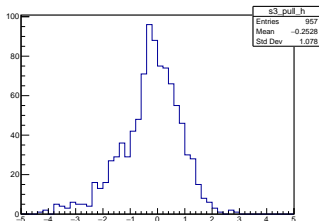
# Input-output checks



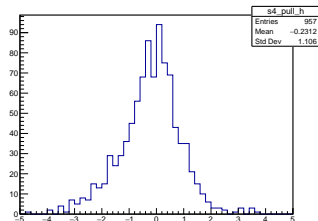
(a)  $s_1$  pulls



(b)  $s_2$  pulls



(c)  $s_3$  pulls



(d)  $s_4$  pulls

## What do the toy fits tell us?

- ① Small bias in the  $c_i$  pull distribution
  - A small bias correction is sufficient
- ②  $s_i$  has a small overcoverage, but a large tail
  - Use the median pull to correct bias in central value
  - Assign asymmetric uncertainty to correct the pull width to unity
- ③ Small bias corrections are also applied to  $\mathcal{B}$  and  $R_i$

# Fit results

Bias-corrected results:

$$c_1 = -0.37 \pm 0.13 \pm 0.01$$

$$s_1 = -0.21^{+0.45}_{-0.28} \pm 0.04$$

$$c_2 = 0.79 \pm 0.06 \pm 0.01$$

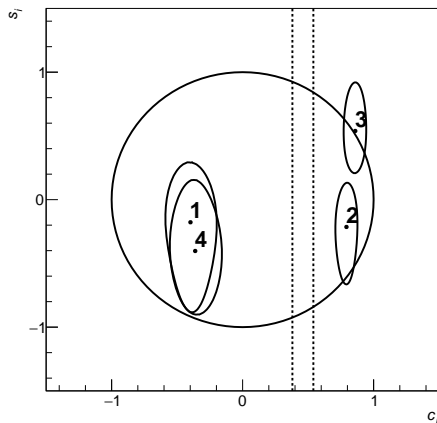
$$s_2 = -0.27^{+0.37}_{-0.18} \pm 0.03$$

$$c_3 = 0.87 \pm 0.06 \pm 0.01$$

$$s_3 = 0.58^{+0.17}_{-0.45} \pm 0.06$$

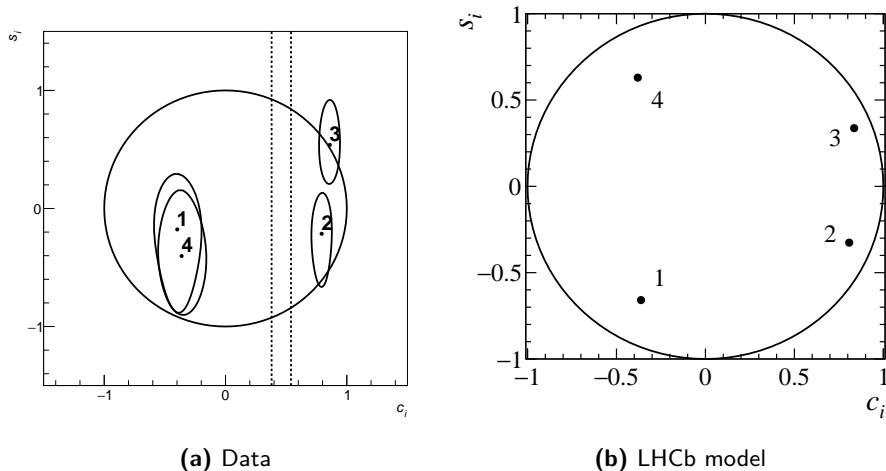
$$c_4 = -0.33 \pm 0.13 \pm 0.01$$

$$s_4 = -0.48^{+0.48}_{-0.26} \pm 0.05$$



**Figure 15:** Contours of  $c_i$  vs  $s_i$ , corresponding to  $\Delta \log(\mathcal{L}) = 2.30$ , or 68% confidence level. Dashed lines show  $F_+$  measurement.

# Fit results



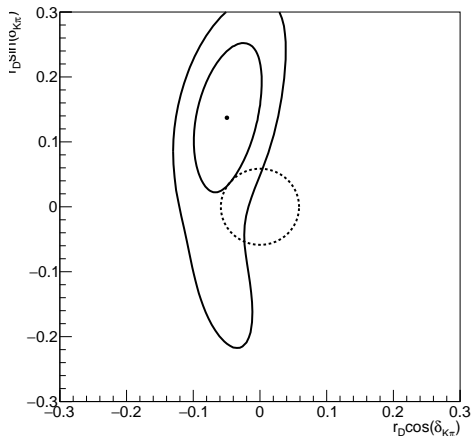
**Figure 16:** Comparison between model and model-independent measurements

- $K_i$ , which are constrained by flavour tags, are free parameters
- Corrections for DCS decays, which depend on the strong phases  $\delta_D$ , are part of the fit
- We could instead treat  $\delta_D$  as a free parameter, and make a simultaneous measurement
- In this analysis we can measure  $\delta_{K\pi}$  with negligible reduction in sensitivity to  $c_i$  and  $s_i$

$$\hat{N}_i^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ij} \left( K_{-j} + r_D^2 K_j - 2r_D \sqrt{K_j K_{-j}} (c_j \cos(\delta_D) + s_j \sin(\delta_D)) \right)$$

- Free parameters:  $r_D \cos(\delta_{K\pi})$  and  $r_D \sin(\delta_{K\pi})$





**Figure 17:** Contours of  $r_D \cos(\delta_{K\pi})$  vs  $r_D \sin(\delta_{K\pi})$ , corresponding to  $\Delta \log(\mathcal{L}) = 2.30$  and  $6.18$ , or 68% and 95% confidence level. Dashed line indicates measured value of  $r_D$ .

In addition, we have another nuisance parameter: The  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  branching fraction

Current PDG value:

$$\mathcal{B} = (2.47 \pm 0.11) \times 10^{-3}$$

Fit result:

$$\mathcal{B} = (2.76 \pm 0.05 \pm 0.03) \times 10^{-3}$$

Much better precision than the current world average

Systematics due to tracking and PID efficiencies have been included, so that we can publish this result as well!

Results completely dominated by statistical uncertainties:

	BF	$c_1$	$c_2$	$c_3$	$c_4$	$s_1$	$s_2$	$s_3$	$s_4$
ST yield	0.2	1.4	0.5	0.5	1.1	1.1	0.7	1.1	0.4
$K_L^0 \pi^0$ ST yield	0.1	4.3	4.8	2.8	3.7	0.3	0.3	0.4	0.6
$K^- e^+ \nu_e$ ST yield	0.7	0.3	0.1	0.4	0.6	3.3	1.1	3.2	0.9
External strong phases	0.2	5.7	3.7	5.0	3.1	18.4	25.4	36.4	24.5
Finite MC size	0.6	5.2	2.9	2.1	4.8	36.8	18.6	42.6	44.9
Single and double tag fit	0.3	3.1	4.6	3.0	5.2	4.8	4.2	2.5	4.3
$K_S^0$ veto	0.0	0.2	4.3	1.9	5.8	0.4	6.2	1.7	2.8
Tracking and PID efficiency	3.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
Total systematic	3.1	9.4	9.2	7.1	10.5	41.6	32.4	56.3	51.4
Statistical	4.6	130.2	55.5	57.2	131.0	319.6	239.9	239.7	312.4

# Summary

- A phase-space binned measurement of the  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  strong-phases has been performed
  - Analysis uses a  $2 \times 4$  binning scheme
  - Measurement will improve significantly with full  $20 \text{ fb}^{-1}$  dataset
  - New strategy: Treat  $c_i$ ,  $s_i$  and  $K_i$  on an equal footing
  - Additionally,  $\delta_D^{K\pi}$  can also be determined
  - The  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  branching fraction is almost 2 times more precise than the current PDG value
- Analysis is ready for review

Thank you!