# BESIII Charm Meeting Phase-space binned analysis of strong-phase parameters in $D^0 \to K^+ K^- \pi^+ \pi^-$

Martin Tat Guy Wilkinson Sneha Malde

University of Oxford

23rd May 2023





#### Outline

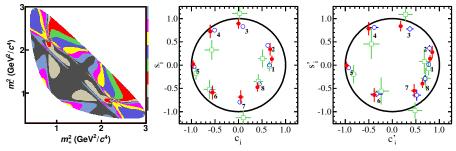
- Introduction
- 2 Binning scheme
- 3 Formalism and measurement strategy
- 4 Selection
- 5 Fit of double tag yields
- 6 Likelihood fit
- Input-output checks
- 8 Fit results
- 9 Bonus measurements
- Summary and conclusion

#### Introduction

- Aim of analysis: Measure the strong-phase difference between  $D^0$  and  $\bar{D^0} \to K^+ K^- \pi^+ \pi^-$  using the 8 fb $^{-1}$   $\psi(3770)$  dataset
- Important input to:
  - **1** Measurement of the CKM angle  $\gamma$
  - 2 Studies of charm mixing and CPV
  - Strong-phase measurements of other decay modes
- Phase-space integrated measurement: Phys. Rev. D 107 032009
- $\bullet$  Strong phases are diluted when integrated over the 5D phase space  $\implies$  Perform analysis in bins
- ullet Binned model-dependent analysis of  $\gamma$  has already been performed at LHCb: arXiv:2301.10328

## Introduction - Previous strong-phase analyses

- ullet Golden mode for strong-phase analysis:  $D^0 o K^0_{S,L} \pi^+ \pi^-$
- Measurements of  $c_i$  and  $s_i$  performed in  $2 \times 8$  bins:
  - CLEO: Phys. Rev. D 82 112006
  - BESIII: Phys. Rev. D 101 112002

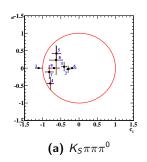


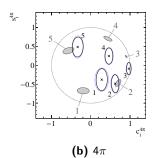
(a) Equal  $\delta$  binning scheme

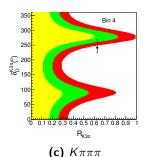
**(b)**  $c_i$  and  $s_i$  for (left)  $K_S\pi\pi$  and (right)  $K_L\pi\pi$ 

#### Introduction - Previous strong-phase analyses

- Four-body decays have a 5D phase space, but their binning scheme may be defined analogously
- Many different strategies for 5D binning schemes:
  - $D \to K_5^0 \pi^+ \pi^- \pi^0$ : JHEP **2018** 82 (2018)
  - $D \to \pi^{+}\pi^{-}\pi^{+}\pi^{-}$ : JHEP **2018** 144 (2018)
  - $D \to K^- \pi^+ \pi^- \pi^+$ : JHEP **2021** 164 (2021)

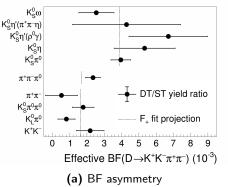


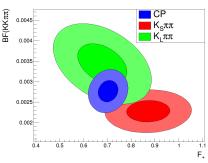




#### Introduction - Previous strong-phase analyses

- The phase-space integrated strong phase of  $D \to K^+K^-\pi^+\pi^-$  was previously measured: Phys. Rev. D **107** 032009
- The asymmetry in the branching fraction measured using CP even and CP odd tags is sensitive to the CP-even fraction  $F_+$ 
  - $c_i = 2F_+ 1$  is the amplitude-average cosine of the strong phase





**(b)**  $F_+$  combination

## Binning scheme

#### Strong phases are determined from yield asymmetries

- When integrating over a phase-space region, variations in strong phases dilute the asymmetries
- We require a binning scheme to minimise the dilution and enhance asymmetry effects

## How to bin a 5-dimensional phase space?

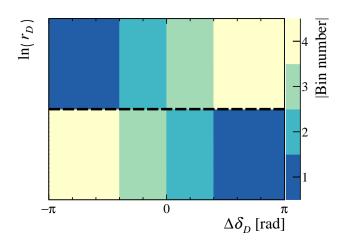
- Generate C++ code for LHCb amplitude model using AmpGen<sup>1</sup>
- For each D event, calculate

$$\frac{A(D^0)}{A(\bar{D^0})} = r_D e^{i\delta_D}$$

- Bin along  $\delta_D$  and  $r_D$ 
  - ullet Bin boundaries in  $\delta_D$  are moved to maximise sensitivity to  $\gamma$

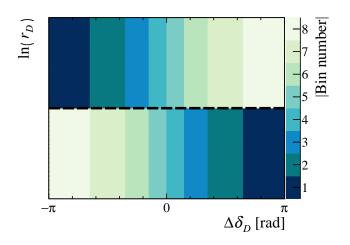
<sup>&</sup>lt;sup>1</sup>AmpGen by Tim Evans

## Binning scheme



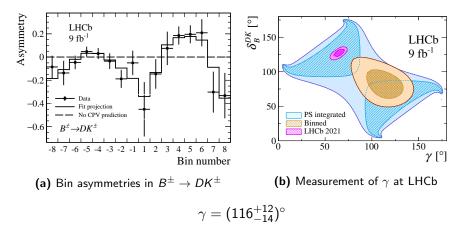
Bins i < 0 on top, i > 0 below Binning scheme used in this analysis

## Binning scheme



 $\mbox{Bins } i < 0 \mbox{ on top, } i > 0 \mbox{ below}$  Possible binning scheme for future analysis with  $20 \mbox{ fb}^{-1}$  dataset

#### Model-dependent LHCb measurement



This result is currently model independent, and needs external inputs from BESIII to become a proper model-independent measurement!

## $Dar{D}$ pair from $\psi(3770)$ is prepared in a $\mathcal{C}=-1$ state

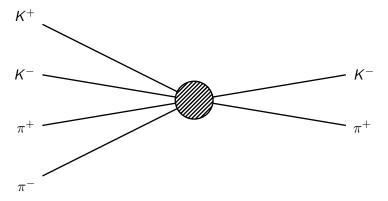
- D mesons are "quantum correlated"
- The decay of  $D \to K^+ K^- \pi^+ \pi^-$  is correlated by the CP content of the tag mode

$$|D\bar{D}\rangle = \frac{1}{\sqrt{2}} \left( |D^0\rangle |\bar{D^0}\rangle - |\bar{D^0}\rangle |D^0\rangle \right)$$

• Equivalently, in terms of *CP* eigenstates:

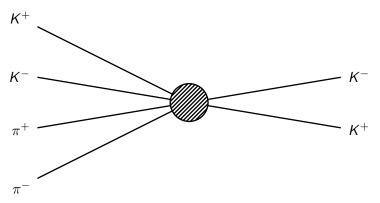
$$|Dar{D}
angle = rac{1}{\sqrt{2}}ig(|D_{-}
angle|D_{+}
angle - |D_{+}
angle|D_{-}
angleig)$$

- Tag mode can be a flavour tag
  - $K\pi$ ,  $K\pi\pi^0$ ,  $K\pi\pi\pi$ ,  $Ke\nu$



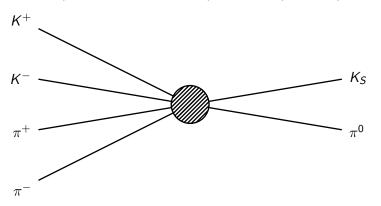
Use flavour tags to measure fraction of  $D^0 \to KK\pi\pi$  decays in bin i:  $\frac{N_i^{\rm DT}}{M^{\rm NST}} = \mathcal{B} \times K_i$ 

- Tag mode can be a CP even tag
  - KK (fully and part. reco KK $\pi\pi$ ),  $\pi\pi$ ,  $\pi\pi\pi^0$ , K<sub>S</sub> $\pi^0\pi^0$ , K<sub>L</sub> $\pi^0$



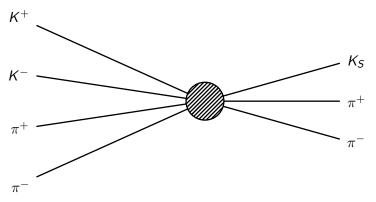
 $D o K^+ K^-$ , which is CP even, forces  $D o K^+ K^- \pi^+ \pi^-$  to be CP odd:  $rac{N_i^{\mathrm{DT}}}{N^{\mathrm{ST}}} = \mathcal{B} imes (K_i + K_{-i} - 2\sqrt{K_i K_{-i}} c_i)$ 

- Tag mode can be a CP odd tag
  - $K_S\pi^0$  (fully and part.reco  $KK\pi\pi$ ),  $K_S\eta$ ,  $K_S\eta'(\pi\pi\eta,\pi\pi\gamma)$ ,  $K_S\pi\pi\pi^0$



 $D o K_S^0 \pi^0$ , which is *CP* odd, forces  $D o K^+ K^- \pi^+ \pi^-$  to be *CP* even:  $\frac{N_i^{\rm DT}}{N^{\rm ST}} = \mathcal{B} imes (K_i + K_{-i} + 2 \sqrt{K_i K_{-i}} c_i)$ 

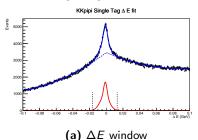
- Tag mode can be a self-conjugate multi-body tag
  - $K_S\pi\pi$  (fully and part.reco  $KK\pi\pi$ ),  $K_L\pi\pi$

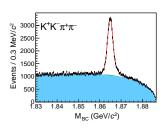


 $D \to \mathcal{K}_{S}^{0} \pi^{+} \pi^{-} \text{ has different strong phases in different bins of phase space:} \\ \frac{\mathcal{N}_{j}^{\text{DT}}}{\mathcal{N}^{\text{ST}}} = \mathcal{B} \times \left( \mathcal{K}_{i} \mathcal{K}_{-j}' + \mathcal{K}_{-i} \mathcal{K}_{j}' - 2 \sqrt{\mathcal{K}_{i} \mathcal{K}_{-i} \mathcal{K}_{j}' \mathcal{K}_{-j}'} (c_{i} c_{j}' + s_{i} s_{j}') \right)$ 

#### Selection

- Double-tag analysis: Select  $D \to KK\pi\pi$  events tagged with flavour, *CP* and multi-body tags
- Selection more or less identical to previous double-tag analyses
- $D \to KK\pi\pi$  selection:
  - 4 good charged tracks
  - $3\sigma$  window around signal peak in  $\Delta E$
  - Asymmetric  $K_SKK$  veto for  $m(\pi^+\pi^-) \in [477, 507]$  MeV
  - Flight significance cut at 2

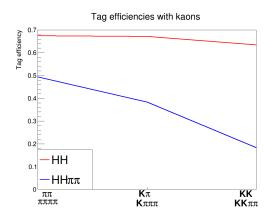




**(b)** Single tag yield:  $29227 \pm 268$ 

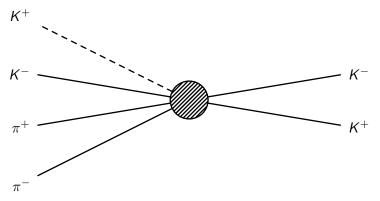
## Partially reconstructed $D^0 o K^+ K^- \pi^+ \pi^-$

- The reconstruction efficiency of  $D o KK\pi\pi$  is less than 20%
- For comparison, the efficiency of  $D \to \pi\pi\pi\pi$  is around 50%!
- Poor kaon tracking efficiency at low momentum



## Partially reconstructed $D^0 o K^+K^-\pi^+\pi^-$

- Solution: Only reconstruct 3 of the charged D daughters
  - Presence of missing kaon is inferred from the missing momentum
  - Yields are similar to fully reconstructed sample
  - ullet Large, but non-peaking background from  $D o K\pi\pi\pi\pi^0$

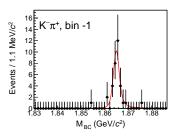


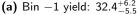
Use this technique with the  $K^+K^-$ ,  $K^0_S\pi^0$  and  $K^0_S\pi^+\pi^-$  tags

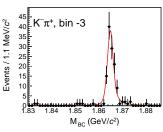
## Double tag fits

- ullet Fit strategy: Only fit signal side  $M_{
  m BC}$  because of low statistics
- Fit model:
  - Signal: PDF from signal MC, convolved with a Gaussian
  - Flat background: Argus PDF
  - Peaking background: Shape and efficiency from MC, correct for quantum correlation
  - Simple sideband subtraction for correct signal but wrong tag event
- Fit all bins simultaneously
  - Shape is floated and shared across all bins
  - Yield of signal and combinatorial background is floated in each bin

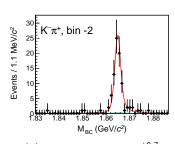
## Double tag fit of $KK\pi\pi$ vs $K\pi$



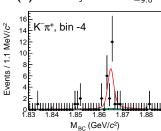




(c) Bin -3 yield:  $120.3^{+11.6}_{-10.9}$ 

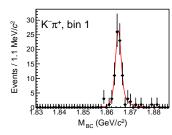


**(b)** Bin -2 yield:  $82.7^{+9.7}_{-9.0}$ 

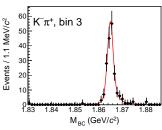


(d) Bin -4 yield:  $22.4^{+5.2}_{-4.6}$ 

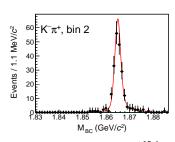
## Double tag fit of $KK\pi\pi$ vs $K\pi$



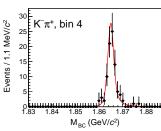




(c) Bin 3 yield:  $181.0^{+14.0}_{-13.3}$ 

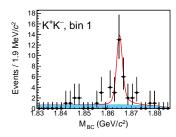


**(b)** Bin 2 yield: 211.2<sup>+15.4</sup><sub>-14.8</sub>

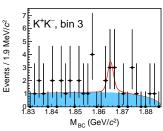


(d) Bin 4 yield:  $88.6^{+9.7}_{-9.0}$ 

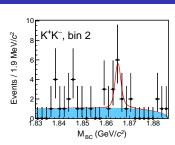
#### Double tag fit of $KK\pi\pi$ vs KK



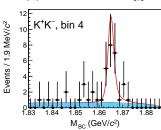
(a) Bin 1 yield: 25.3<sup>+6.2</sup><sub>-5.5</sub>



(c) Bin 3 yield:  $4.5^{+3.3}_{-2.6}$ 

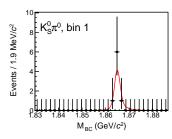


**(b)** Bin 2 yield:  $8.8^{+4.0}_{-3.3}$ 

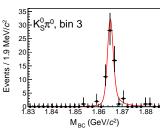


(d) Bin 4 yield:  $21.1_{-4.8}^{+5.5}$ 

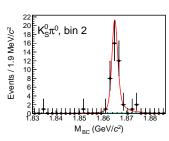
## Double tag fit of $KK\pi\pi$ vs $K_S\pi^0$



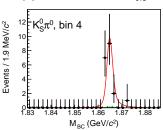
(a) Bin 1 yield:  $7.9^{+3.1}_{-2.5}$ 



(c) Bin 3 yield:  $61.1^{+8.3}_{-7.8}$ 

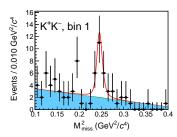


**(b)** Bin 2 yield:  $40.4^{+6.8}_{-6.3}$ 

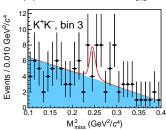


**(d)** Bin 4 yield: 18.3<sup>+4.5</sup><sub>-3.9</sub>

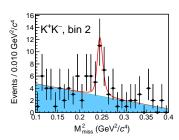
#### Double tag fit of partially reconstructed $KK\pi\pi$ vs KK



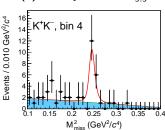




(c) Bin 3 yield:  $7.5^{+6.6}_{-5.7}$ 

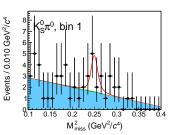


**(b)** Bin 2 yield:  $17.7^{+6.1}_{-5.3}$ 

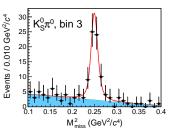


**(d)** Bin 4 yield: 17.3<sup>+5.1</sup><sub>-4.4</sub>

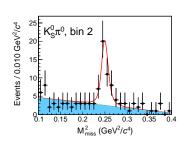
## Double tag fit of partially reconstructed $KK\pi\pi$ vs $K_S\pi^0$



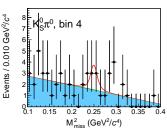
(a) Bin 1 yield:  $7.2^{+4.3}_{-3.5}$ 



(c) Bin 3 yield:  $64.4^{+9.6}_{-9.0}$ 



**(b)** Bin 2 yield: 39.3<sup>+8.3</sup><sub>-7.3</sub>



(d) Bin 4 yield:  $4.8^{+3.8}_{-3.0}$ 

## How to put this all together to determine $c_i$ and $s_i$ ?

• In the past, the strategy has been to measure normalised  $K_i$  first, using flavour tags:

$$rac{N_i^{ ext{DT}}}{N^{ ext{ST}}} \propto K_i, ~~ \sum_i K_i = 1$$

② However, the  $K_i$  must be corrected for efficiencies and bin migrations, in addition to Doubly Cabbibo Suppressed decays:

$$f_i = \left(1 + r_D^2 \frac{\kappa_{-i}}{\kappa_i} - 2\sqrt{\frac{\kappa_{-i}}{\kappa_i}} \left(c_i \cos(\delta_D) + s_i \sin(\delta_D)\right)\right)^{-1}$$

Finally, with the K<sub>i</sub> fixed, c<sub>i</sub> and s<sub>i</sub> are determined from CP and multi-body tags:

$$\frac{N_i^{\rm DT}}{N^{\rm ST}} \propto K_i + K_{-i} \mp 2\sqrt{K_i K_{-i}} c_i \\ \frac{N_{ij}^{\rm DT}}{N^{\rm ST}} \propto K_i K_{-j}' + K_{-i} K_j' - 2\sqrt{K_i K_{-i} K_j' K_{-j}'} (c_i c_j' + s_i s_j')$$

## We propose a simpler, but more powerful strategy:

- **1** Treat all  $K_i$ ,  $c_i$  and  $s_i$  as free parameters
- ② Include the  $D o KK\pi\pi$  branching fraction  ${\cal B}$  as a free parameter
- Fit flavour, CP and multi-body tags simultaneously

#### Master equations:

$$\begin{split} \hat{N}_{i}^{\mathrm{DT}} = & N^{\mathrm{ST}} \mathcal{B} \epsilon_{ij} \Big( K_{-j} + r_{D}^{2} K_{j} - 2 r_{D} \sqrt{K_{j} K_{-j}} \big( c_{j} \cos(\delta_{D}) + s_{j} \sin(\delta_{D}) \big) \Big) \\ \hat{N}_{i}^{\mathrm{DT}} = & N^{\mathrm{ST}} \mathcal{B} \epsilon_{ij} \big( K_{j} + K_{-j} \mp 2 \sqrt{K_{j} K_{-j}} c_{j} \big) \\ \hat{N}_{ij}^{\mathrm{DT}} = & N^{\mathrm{ST}} \mathcal{B} \epsilon_{ijkl} \big( K_{k} K_{-l}' + K_{-k} K_{l}' - 2 \sqrt{K_{k} K_{-k} K_{l}' K_{-l}'} (c_{k} c_{l}' + s_{k} s_{l}') \big) \end{split}$$

 $\epsilon_{ij}$  are combined efficiency and bin migration matrices

The master equations use the free parameters  $\mathcal{B}$ ,  $K_i$ ,  $c_i$  and  $s_i$  to make a prediction  $\hat{N}^{\mathrm{DT}}$  to the measured double tag yields  $N^{\mathrm{DT}}$ 

#### Master equations

$$\hat{N}_{i}^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ij} \left( K_{-j} + r_{D}^{2} K_{j} - 2r_{D} \sqrt{K_{j} K_{-j}} \left( c_{j} \cos(\delta_{D}) + s_{j} \sin(\delta_{D}) \right) \right) 
\hat{N}_{i}^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ij} \left( K_{j} + K_{-j} \mp 2 \sqrt{K_{j} K_{-j}} c_{j} \right) 
\hat{N}_{ij}^{\text{DT}} = N^{\text{ST}} \mathcal{B} \epsilon_{ijkl} \left( K_{k} K_{-l}' + K_{-k} K_{l}' - 2 \sqrt{K_{k} K_{-k} K_{l}' K_{-l}'} \left( c_{k} c_{l}' + s_{k} s_{l}' \right) \right)$$

Ordinarily, we would construct a Gaussian (log)likelihood function  $\implies$ Obtain  $\mathcal{B}$ ,  $K_i$ ,  $c_i$  and  $s_i$  by minimising the following function:

$$-\ln(\mathcal{L}) = \frac{1}{2} \sum_{\mathrm{Tag}} \sum_{jj} (V^{-1})_{ij} (N_i^{\mathrm{DT}} - \hat{N}_i^{\mathrm{DT}}) (N_j^{\mathrm{DT}} - \hat{N}_j^{\mathrm{DT}})$$

$$V_{ij} = \rho_{ij}\sigma_i\sigma_i$$

 $<sup>^{1}\</sup>rho$  are correlation coefficients

#### Master equations

$$\begin{split} \hat{N}_{i}^{\mathrm{DT}} &= N^{\mathrm{ST}} \mathcal{B} \epsilon_{ij} \Big( K_{-j} + r_{D}^{2} K_{j} - 2 r_{D} \sqrt{K_{j} K_{-j}} \big( c_{j} \cos(\delta_{D}) + s_{j} \sin(\delta_{D}) \big) \Big) \\ \hat{N}_{i}^{\mathrm{DT}} &= N^{\mathrm{ST}} \mathcal{B} \epsilon_{ij} \big( K_{j} + K_{-j} \mp 2 \sqrt{K_{j} K_{-j}} c_{j} \big) \\ \hat{N}_{ij}^{\mathrm{DT}} &= N^{\mathrm{ST}} \mathcal{B} \epsilon_{ijkl} \big( K_{k} K_{-l}' + K_{-k} K_{l}' - 2 \sqrt{K_{k} K_{-k} K_{l}' K_{-l}'} (c_{k} c_{l}' + s_{k} s_{l}') \big) \end{split}$$

Our DT yields are very small, so their uncertainties are asymmetric  $\implies$ Approximate covariance matrix from the asymmetric uncertainties<sup>2</sup>:

$$\begin{split} -\ln(\mathcal{L}) = & \frac{1}{2} \sum_{\mathrm{Tag}} \sum_{ij} (V^{-1})_{ij} (N_i^{\mathrm{DT}} - \hat{N}_i^{\mathrm{DT}}) (N_j^{\mathrm{DT}} - \hat{N}_j^{\mathrm{DT}}) \\ V_{ij} = & \rho_{ij} \sigma_i \sigma_j, \quad \sigma = \sqrt{\sigma_- \sigma_+ - (\sigma_+ - \sigma_-) (N^{\mathrm{DT}} - \hat{N}^{\mathrm{DT}})} \end{split}$$

<sup>&</sup>lt;sup>2</sup>arXiv:physics/0406120

$$\begin{split} -\ln(\mathcal{L}) = & \frac{1}{2} \sum_{\mathrm{Tag}} \sum_{ij} (V^{-1})_{ij} (N_i^{\mathrm{DT}} - \hat{N}_i^{\mathrm{DT}}) (N_j^{\mathrm{DT}} - \hat{N}_j^{\mathrm{DT}}) \\ V_{ij} = & \rho_{ij} \sigma_i \sigma_j, \quad \sigma = \sqrt{\sigma_- \sigma_+ - (\sigma_+ - \sigma_-) (N^{\mathrm{DT}} - \hat{N}^{\mathrm{DT}})} \end{split}$$

- The above likelihood has good coverage for flavour and CP tags...
- ... but not for multi-body decays
  - Bins with  $\sigma_- \approx 0$  make the fit unstable
  - Fit convergence was found to be less than 60%
- In multi-body decays, use the full unbinned likelihood directly
  - Fit convergence improves to over 95%
  - Much slower, but much more accurate

#### Parameterisation of $K_{ii}$

#### One last technical detail...

- 8  $K_i$  parameters, but  $\sum_i K_i = 1$ , so only 7 are independent
- We could set  $K_4 = 1 \sum_{i \neq 4} K_i$ , but such a parameterisation is unstable due to strong correlations
- Use a recursive fraction parameterisation (JHEP 2021 169 (2021))
- Recursive fractions R<sub>i</sub> are defined as:

$$R_{i} \equiv \begin{cases} K_{i}, & i = -4 \\ K_{i} / \sum_{j \geq i} K_{j}, & -4 < i < +4 \\ 1, & i = +4. \end{cases}$$

#### Parameterisation of $K_i$

Visualisation of this parameterisation:

#### 1. Definition of $a_1'$

$$\underbrace{a_1}_{\equiv a_1'} + \underbrace{a_2 + a_3 + \dots}_{\equiv 1 - a_1'} = 1$$

#### 2. Eliminate $a_1$

$$a_2 + a_3 + a_4 + \dots = 1 - a_1'$$

#### 3. Definition of $a_2'$

$$\underbrace{\frac{a_2}{1 - a_1'}}_{\equiv a_2'} + \underbrace{\frac{a_3 + a_4 + \dots}{1 - a_1'}}_{\equiv 1 - a_2'} = 1$$

Repeat this procedure until the last coefficient, which is 1

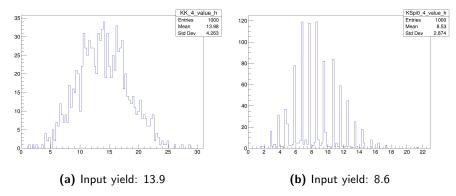
Input-output checks are performed by generating toy datasets

- Check fit convergence
- ② Check error coverage
- Orrect any biases in fitted parameters

Unfortunately, in this analysis we cannot simply generate Poisson-distributed DT yields:

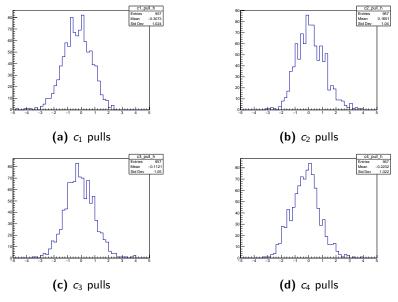
- Asymmetric uncertainties
- 2 Large background-to-signal
- Multi-body tags with low yields require a full unbinned likelihood

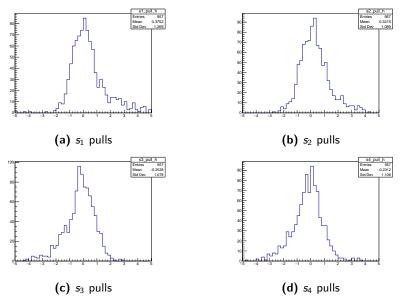
Solution: Generate toy datasets for each double tag fit



**Figure 12:** Fitted yields in toy datasets for the (left) KK and (right)  $K_S\pi^0$  tags

- No biases are observed
- Oistributions are asymmetric support uncertainties are asymmetric
- ullet KK uncertainty is non-Poisson because of large backgrounds from  $qar{q}$
- $K_S\pi^0$  has small backgrounds, so observed yields Poisson distributed



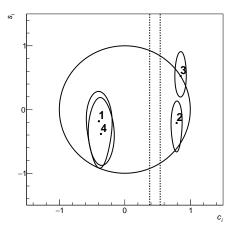


#### What do the toy fits tell us?

- **1** Both  $c_i$  and  $s_i$  pulls has asymmetric tails
- 2 Effect is small in  $c_i$  and a small bias correction will be sufficient
- $\odot$   $s_i$  has a small overcoverage, but a large tail, so special care is required

#### Fit results

$$c_1 = -0.40 \pm 0.13$$
  
 $s_1 = -0.18 \pm 0.32$   
 $c_2 = 0.79 \pm 0.06$   
 $s_2 = -0.21 \pm 0.24$   
 $c_3 = 0.85 \pm 0.06$   
 $s_3 = 0.53 \pm 0.24$   
 $c_4 = -0.36 \pm 0.13$   
 $s_4 = -0.38 \pm 0.31$ 



**Figure 15:** Contours of  $c_i$  vs  $s_i$ , corresponding to  $\Delta \log(\mathcal{L}) = 2.30$ , or 68% confidence level. Dashed lines show  $F_+$  measurement.

#### Fit results

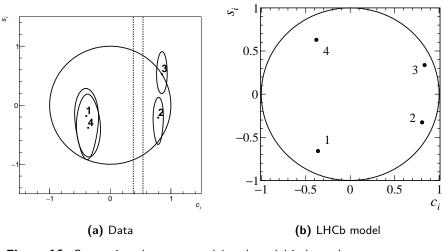


Figure 16: Comparison between model and model-independent measurements

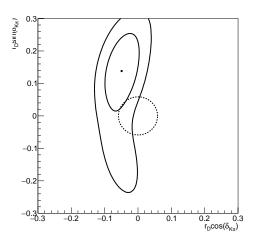
#### $\delta_{K\pi}$ measurement

- $\bullet$   $K_i$ , which are constrained by flavour tags, are free parameters
- Corrections for DCS decays, which depend on the strong phases  $\delta_D$ , are part of the fit
- ullet We could instead treat  $\delta_D$  as a free parameter, and make a simultaneous measurement
- In this analysis we can measure  $\delta_{K\pi}$  with negligible reduction in sensitivity to  $c_i$  and  $s_i$

$$\hat{N}_{i}^{\mathrm{DT}} = N^{\mathrm{ST}} \mathcal{B} \epsilon_{ij} \Big( K_{-j} + r_{D}^{2} K_{j} - 2 r_{D} \sqrt{K_{j} K_{-j}} \big( c_{j} \cos(\delta_{D}) + s_{j} \sin(\delta_{D}) \big) \Big)$$

• Free parameters:  $r_D \cos(\delta_{K\pi})$  and  $r_D \sin(\delta_{K\pi})$ 

#### $\delta_{K\pi}$ measurement



**Figure 17:** Contours of  $r_D \cos(\delta_{K\pi})$  vs  $r_D \sin(\delta_{K\pi})$ , corresponding to  $\Delta \log(\mathcal{L}) = 2.30$  and 6.18, or 68% and 95% confidence level. Dashed line indicates measured value of  $r_D$ .

#### Branching fraction measurement

In addition, we have another nuisance parameter: The  $D^0 \to K^+ K^- \pi^+ \pi^-$  branching fraction

Current PDG value: 
$$\mathcal{B} = (2.47 \pm 0.11) \times 10^{-3}$$
 Fit result:

 $\mathcal{B} = (2.765 \pm 0.046) \times 10^{-3}$ 

Much better precision than the current world average!

Do we have the necessary estimates of tracking and PID systematics to also publish this measurement?

## Summary

- A phase-space binned measurement of the  $D^0 o K^+K^-\pi^+\pi^-$  strong-phases has been performed
  - Analysis uses a 2 × 4 binning scheme
  - ullet Measurement will improve significantly with full  $20\,\mathrm{fb^{-1}}$  dataset
  - New strategy: Treat  $c_i$ ,  $s_i$  and  $K_i$  on an equal footing
  - Additionally,  $\delta_D^{K\pi}$  can also be determined
  - The  $D^0 \to K^+ K^- \pi^+ \pi^-$  branching fraction more than 2 times more precise than the current PDG value
- Next steps:
  - Perform systematics studies
  - Write up MEMO

# Thank you!