# Model independent measurement of the CKM angle $\gamma$ with $B^{\pm} \rightarrow [K^{+}K^{-}\pi^{+}\pi^{-}]_{D}K^{\pm}$ at LHCb and BESIII

Martin Tat

University of Oxford

IOP Joint APP and HEPP Conference

3rd-5th April 2023





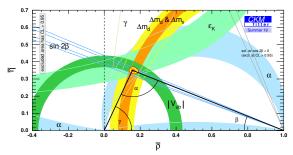


#### Introduction to $\gamma$ and CP violation

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#### Introduction to $\gamma$ and *CP* violation

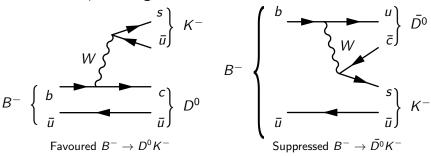
- ullet CPV in SM is described by the Unitary Triangle, with angles lpha, eta,  $\gamma$
- The angle  $\gamma = \arg \Bigl( \frac{V_{ud} \, V_{ub}^*}{V_{cc} V_{cb}^*} \Bigr)$  is very important:
  - Negligible theoretical uncertainties: Ideal SM benchmark
  - Accessible at tree level: Indirectly probe New Physics that enter loops
  - **3** Compare with  $\alpha$ ,  $\beta$  measurements: Is the Unitary Triangle a triangle?



CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005)

#### Sensitivity through interference

Measure  $\gamma$  through interference effects in  $B^{\pm} \to DK^{\pm}$ 



- ullet Superposition of  $D^0$  and  $ar{D^0}$
- $b \rightarrow u\bar{c}s$  and  $b \rightarrow c\bar{u}s$  interference  $\rightarrow$  Sensitivity to  $\gamma$

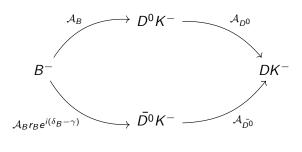
$$\mathcal{A}(B^{-}) = \mathcal{A}_{B} \Big( \mathcal{A}_{D^{0}} + r_{B} e^{i(\delta_{B} - \gamma)} \mathcal{A}_{\bar{D^{0}}} \Big)$$
 $\mathcal{A}(B^{+}) = \mathcal{A}_{B} \Big( \mathcal{A}_{\bar{D^{0}}} + r_{B} e^{i(\delta_{B} + \gamma)} \mathcal{A}_{D^{0}} \Big)$ 

ullet The magnitude of interference effects governed by  $r_B pprox 0.1$ 

#### Sensitivity through interference

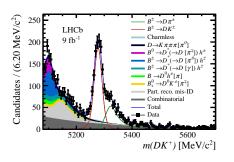
Phase space integrated analysis: Compare yields of  $B^+$  and  $B^-$ 

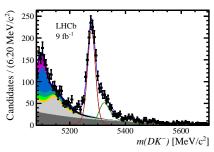
Interference depends on  $\gamma$ , diluted by  $\kappa$  when integrated over phase space



$$|\mathcal{A}(B^-)|^2 \propto 1 + r_B^2 + 2r_B\kappa\cos(\delta_B - \gamma)$$

## First look at $B^{\pm} \rightarrow [K^+K^-\pi^+\pi^-]_D K^{\pm}$





arXiv:2301.10328

Measuring the  $B^{\pm} \to DK^{\pm}$  asymmetry  $\mathcal{A}$  provide useful constraints on  $\gamma$ , but with some caveats:

- **①** Interference effects are diluted by a factor  $\kappa = 0.46 \pm 0.08$ 
  - Phys. Rev. D 107, 032009
- Pour-fold degeneracy:
  - $(\gamma, \delta_B) \rightarrow (\delta_B, \gamma)$
  - $(\gamma, \delta_B) \rightarrow (\pi \gamma, \pi \delta_B)$

## Phase space binned analysis of $B^{\pm} \rightarrow [K^+K^-\pi^+\pi^-]_D K^{\pm}$

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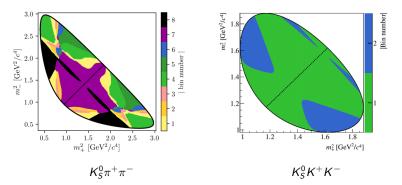
## Phase space binned analysis of $B^{\pm} \rightarrow [K^+K^-\pi^+\pi^-]_D K^{\pm}$

Our aim: First binned model independent measurement of  $\gamma$  in  $B^{\pm} \rightarrow [K^+K^-\pi^+\pi^-]_D K^{\pm}$ 

- Model dependent measurement first proposed in 2007
  - Phys. Lett. B647 (2007) 400 by J. Rademacker and G. Wilkinson
  - FOCUS amplitude model: 14° precision with 1000 candidates
- State of the art amplitude analysis by LHCb:
  - JHEP **02** (2019) 126
  - Exploits the huge dataset of charm decays collected by LHCb
  - Useful for understanding the 5-dimensional phase space
- Use LHCb model to identify regions with large asymmetries
  - Place regions with similar asymmetries together in bins

## Phase space binned analysis of $B^{\pm} \rightarrow [K^+K^-\pi^+\pi^-]_D K^{\pm}$

• Analogous to the decays  $D^0 \to K_S^0 \pi^+ \pi^-$  and  $K_S^0 K^+ K^-$ , where the binning scheme may be visualised on a Dalitz plot



- The strong phase  $\delta_D$  varies across phase space, leading to different asymmetries in each bin
- The amplitude-averaged cosine and sine of  $\delta_D$ ,  $c_i$  and  $s_i$ , are measured directly in each bin at charm factories

#### Binning scheme

## Binning scheme

#### Binning scheme

#### A binning scheme must satisfy the following:

- Minimal dilution of strong phases when integrating over bins
- Enhance interference between  $B^\pm o D^0 K^\pm$  and  $B^\pm o ar{D^0} K^\pm$

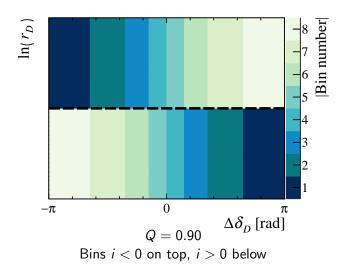
#### How to bin a 5-dimensional phase space?

• For each  $B^{\pm}$  candidate, use the amplitude model to calculate

$$\frac{A(D^0)}{A(\bar{D^0})} = r_D e^{i\delta_D}$$

- ② Split  $\delta_D$  into uniformly spaced bins
- **3** Use the symmetry line  $r_D = 1$  to separate bin +i from -i
- lacktriangle Optimise the binning scheme by adjusting the bin boundaries in  $\delta_D$

#### Binning scheme

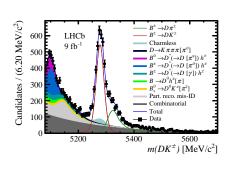


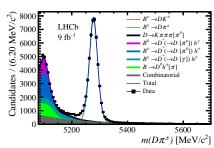
#### Mass fits, *CP* fit and $\gamma$

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#### Mass fits, CP fit and $\gamma$

# Global invariant mass fit: Fit everything to understand mass shapes and background yields





#### Signal yield:

 $B^{\pm} \to DK^{\pm}: 3026 \pm 38$ 

 $B^{\pm} \to D\pi^{\pm}$ : 44 349  $\pm$  218

#### Interpretation of $\gamma$

CP fit: Simultaneously fit bin yields and determine physics parameters, using model predictions of  $c_i$  and  $s_i$ :

$$\begin{split} \gamma &= (116^{+12}_{-14})^{\circ}, \\ \delta^{DK}_{B} &= (81^{+14}_{-13})^{\circ}, \\ r^{DK}_{B} &= 0.110^{+0.020}_{-0.020}, \\ \delta^{D\pi}_{B} &= (298^{+62}_{-118})^{\circ}, \\ r^{D\pi}_{B} &= 0.0041^{+0.0054}_{-0.0041}, \end{split}$$

However, the latest  $\gamma$  and charm combination result is:

$$\gamma = (63.8^{+3.5}_{-3.7})^{\circ}$$

What went wrong?!

Do we trust the model predicted  $c_i$  and  $s_i$ , or their uncertainties?

#### Interpretation of $\gamma$

# There are several reasons why amplitude models <u>cannot</u> be trusted

- Amplitude models are just models, which may not reflect reality
- ② In fact, the model is fitted to data that knows nothing about  $\delta_D$
- It is impossible to assign an objective error to a model!

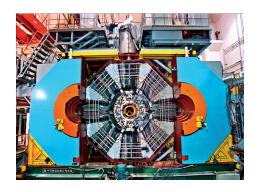
We wish to do a model independent measurement Let's go and measure  $c_i$  and  $s_i$  at BESIII!

## Strong phase analysis of $D^0 \to K^+K^-\pi^+\pi^-$ at BESIII

Strong phase analysis of 
$$D^0 o K^+K^-\pi^+\pi^-$$
 at BESIII

## Strong phase analysis of $D^0 \to K^+K^-\pi^+\pi^-$ at BESIII

- BESIII: Beijing Spectrometer III, a detector at the Beijing Electron-Positron Collider II, located at IHEP
- ullet  $e^+e^-$  collider at the  $\psi(3770) o D^0 ar{D^0}$  threshold
  - 2010-2011:  $3 \, \text{fb}^{-1}$
  - 2022: 5 fb<sup>-1</sup>
  - Expect  $20 \, \text{fb}^{-1}$  in total by end of 2024

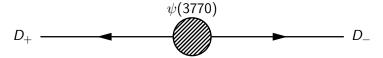


## Strong phase analysis of $D^0 o K^+ K^- \pi^+ \pi^-$ at BESIII

- Double-tag analysis: Reconstruct signal  $(KK\pi\pi)$  and tag mode
- $D^0\bar{D^0}$  pair is quantum correlated



- Equivalently, we can consider  $D_+D_-$ 
  - $D_{\pm}=rac{1}{\sqrt{2}}(D^0\pm ar{D^0})$  are CP eigenstates

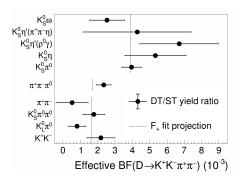


The DD pair is quantum correlated, spooky action at a distance!

## Strong phase analysis of $D^0 \to K^+K^-\pi^+\pi^-$ at BESIII

Quantum correlation: The *CP* content of the tag can modify the effective branching fraction:

$$rac{ extstyle N^{ ext{DT}}}{ extstyle N^{ ext{ST}}} = \mathcal{B}(D^0 o extstyle extstyle extstyle extstyle K extstyle \pi \pi) ig(1 \pm c_1ig)$$



Phys. Rev. D 107, 032009

 $c_1$  is the cosine of the strong phase, averaged over the whole phase space

## Strong phase analysis of $D^0 \to K^+K^-\pi^+\pi^-$ at BESIII

Our next task is to change the phase space inclusive analysis,

$$\begin{split} \frac{\textit{N}^{\rm DT}}{\textit{N}^{\rm ST}} = & \mathcal{B}(\textit{D}^0 \to \textit{KK}\pi\pi) \quad \text{(flavour tag)} \\ \frac{\textit{N}^{\rm DT}}{\textit{N}^{\rm ST}} = & \mathcal{B}(\textit{D}^0 \to \textit{KK}\pi\pi) \big(1 \pm c_1\big) \quad \text{(CP tag)} \\ \text{into a binned phase space analysis:} \end{split}$$

$$\begin{split} \frac{N_i^{\rm DT}}{N^{\rm ST}} = & \mathcal{B}(D^0 \to KK\pi\pi) F_i \quad \text{(flavour tag)} \\ \frac{N_i^{\rm DT}}{N^{\rm ST}} = & \mathcal{B}(D^0 \to KK\pi\pi) (F_i + \bar{F}_i \pm 2\sqrt{F_i \bar{F}_i} c_i) \quad \text{(CP tag)} \end{split}$$

- ②  $c_i$ : Determine from asymmetry of CP even and odd tags
- $\circ$   $s_i$ : Analogous to  $c_i$ , but requires binning of tag mode

## Strong phase analysis of $D^0 o K^+ K^- \pi^+ \pi^-$ at BESIII

Our next task is to change the phase space inclusive analysis,

$$\frac{N^{\mathrm{DT}}}{N^{\mathrm{ST}}} = \mathcal{B}(D^0 \to KK\pi\pi) \quad \text{(f)} \quad \text{(CP tag)}$$

$$\frac{N^{\mathrm{DT}}}{N^{\mathrm{ST}}} = \mathcal{B}(D^0 \to KK\pi) \quad \text{(ce analysis:}$$

$$\frac{N^{\mathrm{DT}}_i}{N^{\mathrm{ST}}} = \mathcal{B}(D^0 \to KV) \quad \text{(avour tag)}$$

$$\frac{N^{\mathrm{DT}}_i}{N^{\mathrm{ST}}} = \mathcal{B}(D^0 \to KV) \quad \text{(CP tag)}$$

- $F_i$ : Measure using 1. our tags
- ②  $c_i$ : Determine from asymmetry of CP even and odd tags

## Summary and conclusion

- ① I have presented the first model-independent measurement of  $\gamma$  using  $B^\pm \to [K^+K^-\pi^+\pi^-]_D K^\pm$
- ② The optimised binning scheme, developed with an amplitude model, successfully identified regions with large, local *CP* asymmetries
- ullet However, amplitude model predictions of  $\delta_D$  should not be trusted
- LHCb analysis has  $3\sigma$  tension with world average



Making binning scheme with amplitude model

Predicting strong phases with amplitude model

 External inputs from charm factories, such as BESIII, are crucial to eliminate biases due to modelling

#### Summary and conclusion

Thanks for your attention!

#### Backup slides

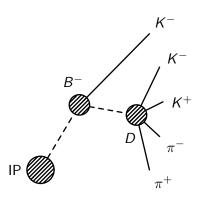
## Backup slides

## Event selection

### Decay topology

#### Look for:

- 5 charged tracks
- Oisplaced B vertex
- 1 bachelor track with good PID information
- Displaced D vertex with invariant mass within
   MeV of the D<sup>0</sup> mass



#### Offline selection has 3 stages

#### Initial cuts:

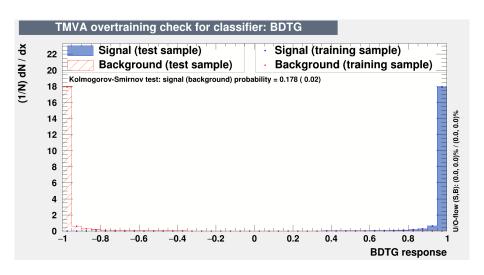
- Invariant D and B mass cuts
- Momentum and RICH requirements

#### Boosted Decision Tree (BDT)

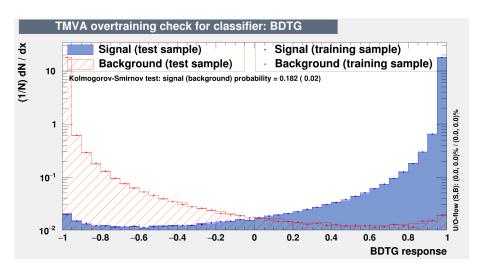
- Signal sample: Simulation samples
- Background sample: Upper B mass sideband
- 28 variables describing kinematics, impact parameters, vertex quality

#### Final selection

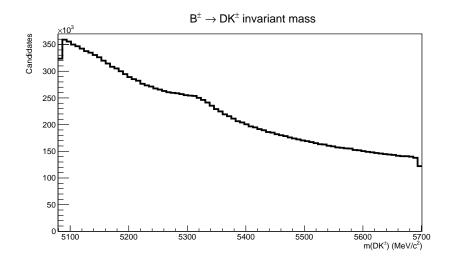
- D Flight distance
- Particle Identification of bachelor



BDT is highly efficient at rejecting combinatorial background

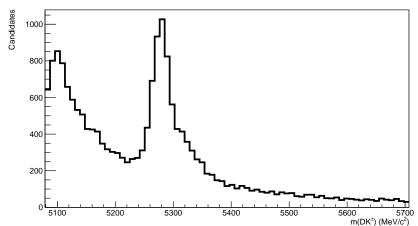


Very important, combinatorial background is large in multi-body decays



The invariant B mass, after online selection, show no visible signal...





... but the BDT does a great job cleaning this up!