

# $D \rightarrow K^+ K^- \pi^+ \pi^-$ strong phase analysis at BESIII

Martin Tat

Oxford LHCb

27th February 2023



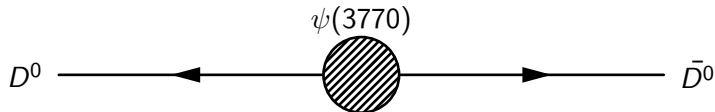
## Recap of BESIII analysis

## Analysis of $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

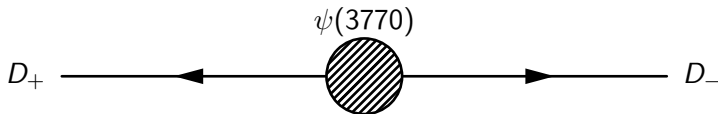
- Study  $D^0$ - $\bar{D}^0$  strong phase difference in bins of the 5D phase space
- Measurement of amplitude averaged strong phases  $c_i$  and  $s_i$
- $c_i$  and  $s_i$  are important inputs to the  $\gamma$  measurement at LHCb
  - LHCb result:  $\gamma = (116^{+12}_{-14})^\circ$  with model dependent inputs
  - $\gamma$  may change when updated with model independent  $c_i$  and  $s_i$
- Measurement technique unique to charm factories: Study decays of quantum correlated  $D\bar{D}$  pairs using a double tag method

# Recap of BESIII analysis

- $\psi(3770) \rightarrow D^0 \bar{D}^0$  decay conserves  $\mathcal{C} = -1$



- But since they are quantum correlated, we must consider their CP eigenstates  $D_{\pm} = (|D^0\rangle \pm |\bar{D}^0\rangle)/\sqrt{2}$
- Total wavefunction is  $|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle = |D_+\rangle|D_-\rangle + |D_-\rangle|D_+\rangle$

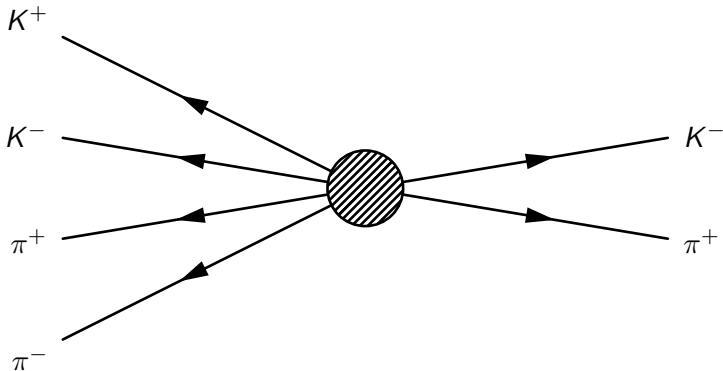


The two  $D$  mesons do not communicate, but the  $D \rightarrow KK\pi\pi$  decay is perfectly correlated with the tagged  $D$

# Recap of BESIII analysis

- Tag mode can be a flavour tag

- $K\pi$ ,  $K\pi\pi^0$ ,  $K\pi\pi\pi$ ,  $Ke\nu$

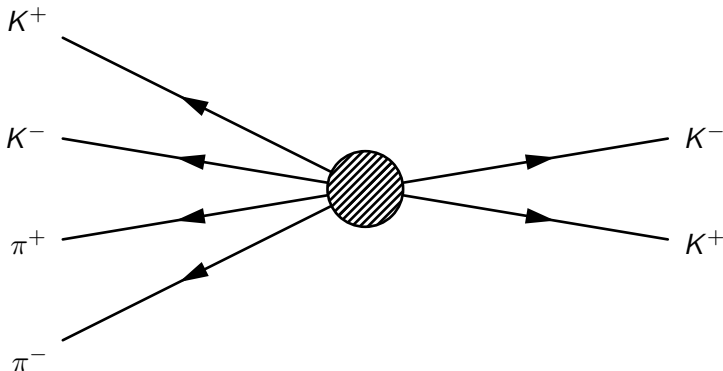


Use flavour tags to measure fraction of  $D^0 \rightarrow KK\pi\pi$  decays in bin  $i$ :

$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B}(KK\pi\pi) \times K_i$$

# Recap of BESIII analysis

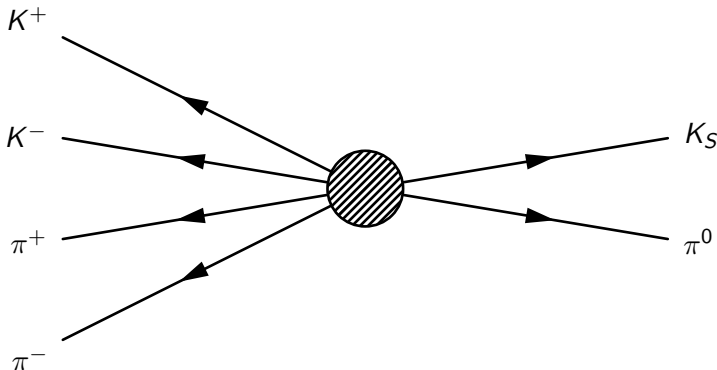
- Tag mode can be a CP even tag
  - $KK, \pi\pi, \pi\pi\pi^0, K_S^0\pi^0, K_L\pi^0, K_L\omega$



$D \rightarrow K^+K^-$ , which is  $CP$  even, forces  $D \rightarrow K^+K^-\pi^+\pi^-$  to be  $CP$  odd

# Recap of BESIII analysis

- Tag mode can be a CP odd tag
  - $K_S\pi^0$ ,  $K_S\omega$ ,  $K_S\eta$ ,  $K_S\eta'$ ,  $K_L\pi^0\pi^0$



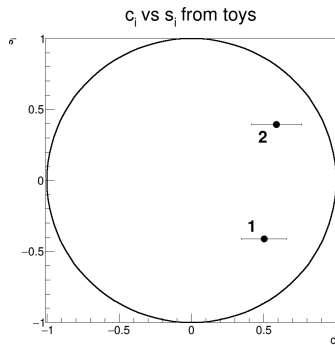
$D \rightarrow K_S^0\pi^0$ , which is  $CP$  odd, forces  $D \rightarrow K^+K^-\pi^+\pi^-$  to be  $CP$  even

# Recap of BESIII analysis

The asymmetry between CP even and CP odd double tag yields is sensitive to  $c_i$ :

$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B}(KK\pi\pi) \times (K_i + \bar{K}_i \pm 2\sqrt{K_i\bar{K}_i}c_i)$$

- ①  $K_i$ : Fraction of  $D^0$  decays in bin  $i$
- ②  $\bar{K}_i$ : Fraction of  $\bar{D}^0$  decays in bin  $i$
- ③  $2\sqrt{K_i\bar{K}_i}c_i$ : Interference term, which depend on the cosine of the phase difference between  $D^0$  and  $\bar{D}^0$  decays

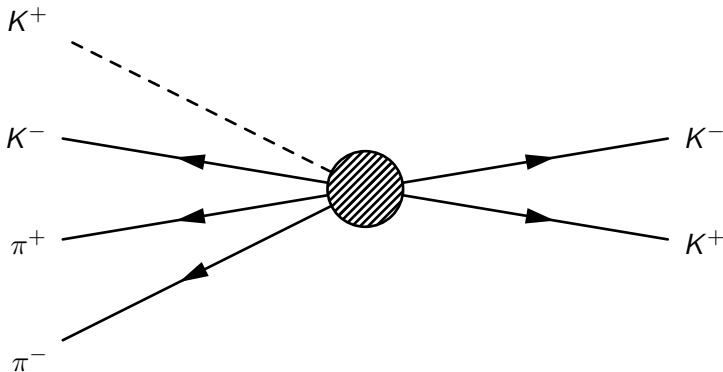




## Double tag yields of partially reconstructed events

# Double tag yields of partially reconstructed events

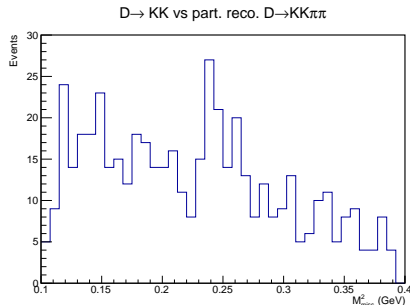
- Reconstruction efficiency of  $D \rightarrow KK\pi\pi$  is very low (18%), compared to other four body modes such as  $D \rightarrow \pi\pi\pi\pi$  (49%)
  - Soft kaons get stuck in magnetic field...
- Try out partially reconstructed  $D \rightarrow KK\pi\pi$  technique
  - No overlap with the fully reconstructed sample



In MC, the yields were found to be similar to fully reconstructed DTs

# Double tag yields of partially reconstructed events

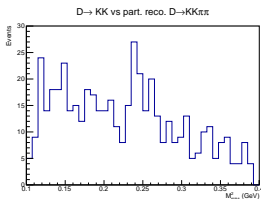
First look at partially reconstructed  $D \rightarrow KK\pi\pi$  vs  $D \rightarrow KK$ :



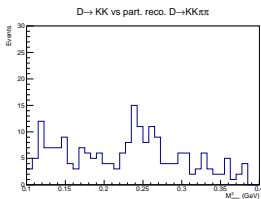
- Large background from  $D \rightarrow K\pi\pi\pi\pi^0$ ,  $\pi\pi^0$  not reconstructed
- Background is non-peaking since two particles are missing
- No reliable MC because no model for this decay exists!

# Double tag yields of partially reconstructed events

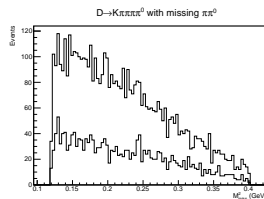
Method for removing  $\pi^0$  backgrounds: Veto events with  $\pi^0$  candidates  
Signal efficiency is reduced by 30%, but 65% of  $K\pi\pi\pi\pi^0$  is rejected



(a) Data, without veto



(b) Data, with  $\pi^0$  veto



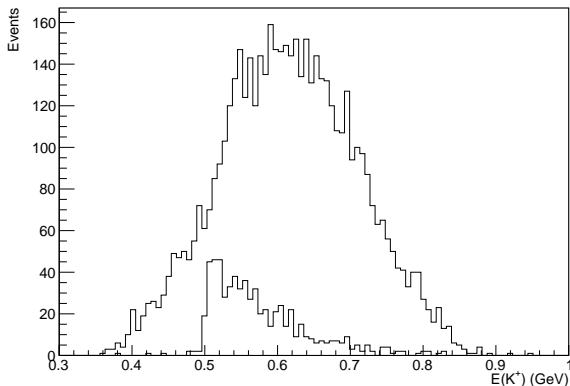
(c) MC, with/without veto

**Figure 1:** Study of  $\pi^0$  veto

# Double tag yields of partially reconstructed events

Using the method of trial and error, I also discovered that a cut on the kaon energy is highly efficient at removing the  $K\pi\pi\pi\pi^0$  background:

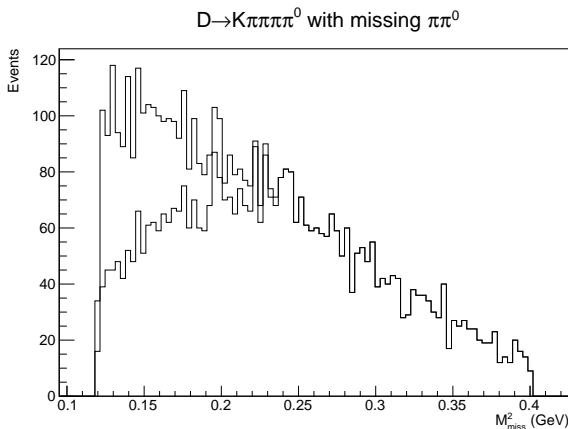
$D \rightarrow K\pi\pi\pi^0$  with missing  $\pi\pi^0$



**Figure 2:** The energy spectrum of the kaon from the background  $D \rightarrow K\pi\pi\pi\pi^0$  and the signal  $D \rightarrow KK\pi\pi$

# Double tag yields of partially reconstructed events

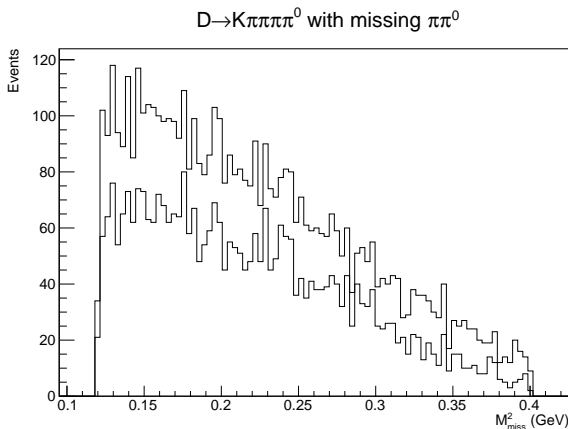
Unfortunately a cut of  $E > m_K$  affects the mass shape...



**Figure 3:** Mass shape before and after a  $E > m_K$  cut

# Double tag yields of partially reconstructed events

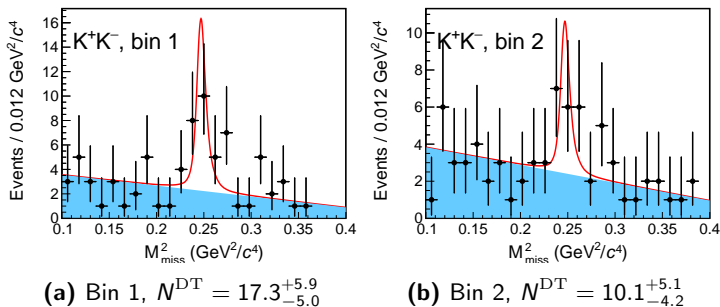
... but luckily a  $E < 0.7$  GeV cut does not change the mass shape!



**Figure 4:** Mass shape before and after a  $E < 0.7$  GeV cut

# Double tag yields of partially reconstructed events

Finally, I can run the fit and the yields were found to be similar to the fully reconstructed samples

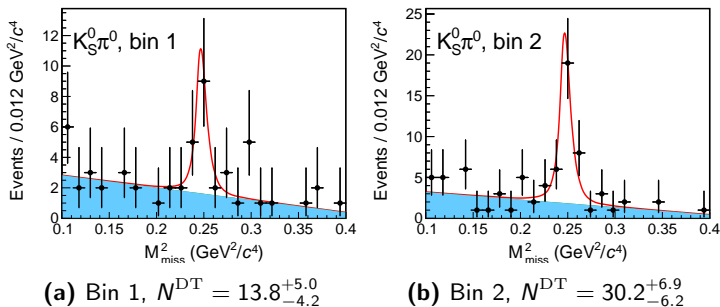


**Figure 5:** Double tag fit of partially reconstructed  $KK\pi\pi$  vs  $KK$



# Double tag yields of partially reconstructed events

Finally, I can run the fit and the yields were found to be similar to the fully reconstructed samples



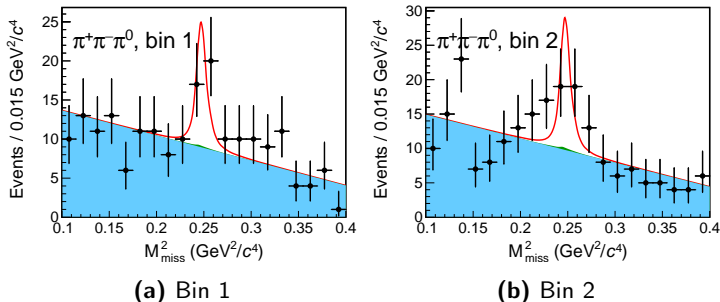
**Figure 6:** Double tag fit of partially reconstructed  $KK\pi\pi$  vs  $K_S^0\pi^0$

# Double tag yields of partially reconstructed events

Partially reconstructed  $KK\pi\pi$  vs  $\pi\pi\pi^0$  mode not understood yet...

Probably background from light hadrons, such as

$$e^+e^- \rightarrow K_1\bar{K}^*\rho \rightarrow KK\pi\pi\pi\pi^0$$



**Figure 7:** Double tag fit of partially reconstructed  $KK\pi\pi$  vs  $\pi\pi\pi^0$

$c_i$  and  $s_i$  fit

## Equations for $c_i$ and $s_i$

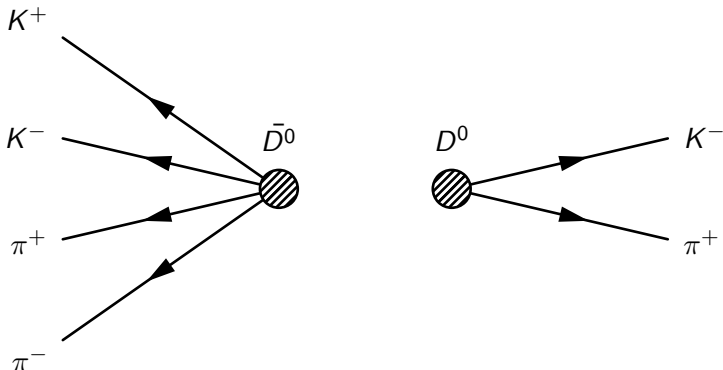
$$\frac{N_i^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B}(KK\pi\pi) \times (K_i + \bar{K}_i \pm 2\sqrt{K_i\bar{K}_i}c_i)$$
$$\frac{N_{ij}^{\text{DT}}}{N^{\text{ST}}} = \mathcal{B}(KK\pi\pi) \times (K_i\bar{K}_j' + \bar{K}_iK_j' - 2\sqrt{K_i\bar{K}_iK_j'\bar{K}_j'}(c_ic_j' + s_is_j'))$$

Fit setup:

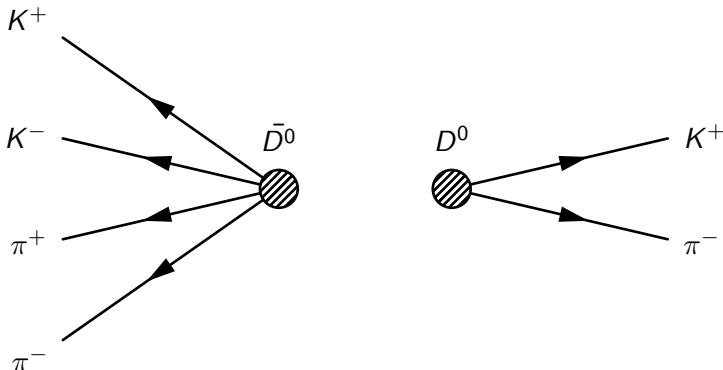
- 1  $K_i$  and  $N^{\text{ST}}$  are fixed, apply systematic or Gaussian constrain
- 2  $\mathcal{B}(KK\pi\pi)$ ,  $c_i$  and  $s_i$  are floating
- 3 Construct likelihood from DT yields and their (asymmetric) uncertainties, using the equations above
- 4 Correct  $K_i$  for DCS decays directly in the fit

- What are DCS corrections?

- Naively, we expect a  $K^-\pi^+$  to originate from  $D^0$  (CF mode)
- Therefore the other decay must be  $\bar{D}^0 \rightarrow KK\pi\pi$



- What are DCS corrections?
  - But what if  $D^0 \rightarrow K^+\pi^-$  (DCS mode)?
  - This leads to a bias in  $K_i$  because the flavour of  $D \rightarrow KK\pi\pi$  is wrong



- Since  $|\mathcal{A}(CF)| = r_D |\mathcal{A}(DCS)|$ , where  $r_D \approx 5\%$ , DCS decays are suppressed by  $r_D^2 = \mathcal{O}(10^{-3})$
- However, since  $D\bar{D}$  are quantum correlated, there is an interference!
- Interference depends on the CF/DCS ratio  $r_D$ , coherence factor  $R$  and strong phase  $\delta_D$
- It also depends on  $c_i$  and  $s_i$ , which we are trying to determine...

The true  $K_i$  has contributions from CF and DCS decays, and their interference:

$$K_i^{\text{true}} = K_i^{\text{obs}} + r_D^2 \bar{K}_i^{\text{obs}} - 2r_D R \sqrt{K_i^{\text{obs}} \bar{K}_i^{\text{obs}}} (c_i \cos(\delta_D) - s_i \sin(\delta_D))$$

## DCS correction factors:

$$f_i = \frac{1}{1 + r_D^2 \frac{\bar{K}_i^{\text{obs}}}{K_i^{\text{obs}}} - 2r_D R \sqrt{\frac{\bar{K}_i^{\text{obs}}}{K_i^{\text{obs}}}} (c_i \cos(\delta_D) - s_i \sin(\delta_D))}$$

- In previous analyses of  $D \rightarrow K_S \pi \pi$  and  $K_S K K$ ,  $f_i$  were calculated using model predictions of  $c_i$  and  $s_i$
- Unfortunately, we don't know if the model prediction of  $D \rightarrow K K \pi \pi$  strong phases is accurate
- In the  $D \rightarrow K_S \pi \pi \pi^0$  analysis, only the  $D \rightarrow K e \nu$  flavour tag was used, because no amplitude model exists (yet)

Solution: Add DCS corrections directly into the fit  
 Calculate  $K_i^{\text{true}}$ , using 3 iterations, with the fitted  $c_i/s_i$



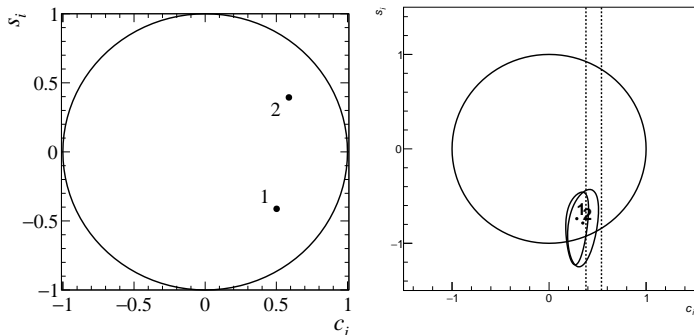
With this strategy, the fit will determine the DCS corrections for us  
Here, using  $D \rightarrow K\pi$  as an example:

Bin	$f_i$	$f_{-i}$
1	0.95	0.98
2	0.94	0.97

Bonus measurement: Instead of using the  $K\pi$  tag to measure  $K_i$ , we can actually float  $\delta_D^{K\pi}$  directly in the fit, and hence simultaneously measure the strong phase of the  $D \rightarrow K\pi$  decay!

## $c_i$ and $s_i$ fit

Perform fit of  $c_i$  and  $s_i$ , using the fully reconstructed  $KK$ ,  $\pi\pi\pi^0$ ,  $K_S\pi^0$  and  $K_S\pi\pi$  tags, as well as partially reconstructed  $KK\pi\pi$  vs  $KK$  and  $K_S\pi^0$



**Figure 8:** Comparison of fitted values of  $c_i$  and  $s_i$  (left), and the model predictions (right)

$c_i$  agrees well with  $F_+$  measurement (as it should), but  $s_i$  clearly shows a tension with the model!

# Summary and next steps

- BESIII measurement of  $c_i$  and  $s_i$  is progressing well
- A partially reconstructed  $D \rightarrow KK\pi\pi$  method has been tested, but there were some challenges with large  $D \rightarrow K\pi\pi\pi\pi^0$  backgrounds
- The preliminary fit of  $c_i$  and  $s_i$  shows promising results
  - A method for direct DCS decay corrections is working well
  - Results of  $c_i$  agree with the  $F_+$  measurement
  - $s_i$  shows tensions with the LHCb model
- Next steps:
  - 1 Finish calculation of peaking backgrounds in each bin
  - 2 Reprocess all data and generate new MC once new data is available
  - 3 Add the rest of the tags
  - 4 Charm WG review

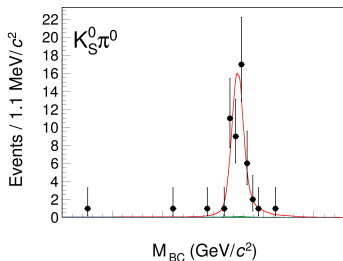
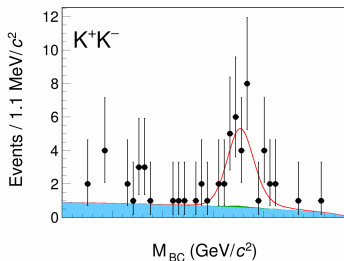
Thank you for listening!

Backup: Double tag yields of fully  
reconstructed events

# Backup: Double tag yields of fully reconstructed events

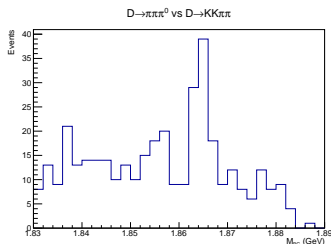
Strategy for measuring DT yields of fully reconstructed events:

- Fit the beam constrained mass  $M_{\text{BC}} = \sqrt{E_{\text{beam}}^2 - \sum_i |\vec{p}_i|^2}$
- Signal shape: MC-derived shape, convolved with Gaussian
- Flat background: Argus function
- Peaking background: Fixed from MC

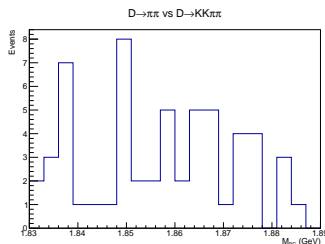


# Backup: Double tag yields of fully reconstructed events

Problem: Fully pionic tag modes have very large backgrounds!



(a)  $D \rightarrow \pi\pi\pi^0$  tag



(b)  $D \rightarrow \pi\pi$  tag

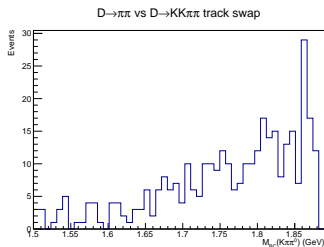
**Figure 9:**  $D \rightarrow KK\pi\pi$  with pionic tags

Flat background mainly from light hadrons, such as

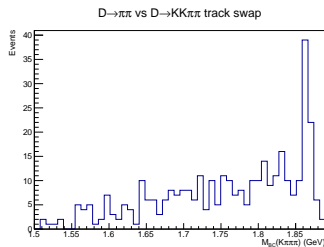
$$e^+e^- \rightarrow K^*\bar{K}^*\rho \rightarrow KK\pi\pi\pi\pi$$

# Backup: Double tag yields of fully reconstructed events

However, in the  $\pi\pi\pi^0$  tag there is track swap background we can veto  
 $K\pi\pi^0$  vs  $K\pi\pi\pi \rightarrow \pi\pi\pi^0$  vs  $KK\pi\pi$   
Consider  $M_{BC}$  with a track swap:



(a)  $D \rightarrow K\pi\pi^0$



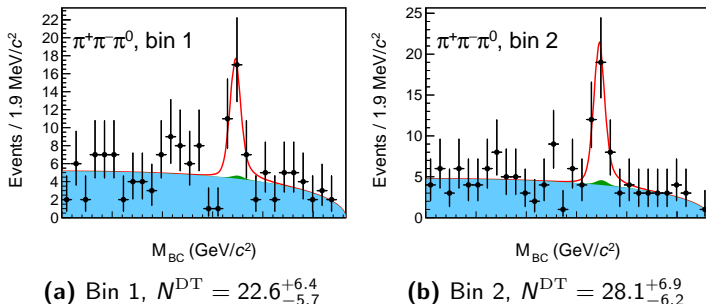
(b)  $D \rightarrow K\pi\pi\pi$

**Figure 10:** Track swapped  $M_{BC}$

A veto in the region  $[1.86, 1.87]\text{GeV}$  removes 20% of the flat background

# Backup: Double tag yields of fully reconstructed events

Double tag yield uncertainties improved, but still lots of background



**Figure 11:** Double tag fit of  $\pi\pi\pi^0$  tag