

ARC: A RICH detector proposal for FCC-ee

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Analysis work summary

- For my PhD thesis I'm working on two related projects:
 - Measurement of $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ strong phases c_i and s_i at BESIII
 - Binned GGSZ analysis of γ in $B^\pm \rightarrow D h^\pm$, $D \rightarrow K^+ K^- \pi^+ \pi^-$ at LHCb
- Over the summer I have worked on two papers:

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH (CERN)



CERN-EP-2022-222
LHCb-PAPER-2022-037
October 19, 2022

Measurement of the CP -even fraction of $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ (based: October 21, 2022)

A determination of the CP -even fraction F_0 in the decay $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ is presented. Using 2.05×10^{10} $K^+ K^- \pi^+ \pi^-$ events recorded by the LHCb experiment in the BESIII detector, one observes F_0 reconstructed in the signal mode and the other in a CP eigenstate or the decay $D \rightarrow K^+ K^- \pi^+ \pi^-$.
Another measurement of F_0 is obtained by fitting the ratio of the CP -even fraction to the CP -odd fraction F_1 to 0.021, where the first uncertainty is statistical and the second is systematic. This is the first model-independent measurement of F_0 in $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ decays.

I. INTRODUCTION

The Standard Model description of CP violation may be tested by measuring the length and the sign of the Unitary Triangle of the CKM matrix [1, 2]. One of these ways is to measure the CP -even fraction F_0 in the $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ decay. This is a model-independent process, with negligible theoretical uncertainties [3]. The CP -even fraction F_0 is defined as the fraction of D^0 mesons which decay into an even number of K^+ mesons. Direct measurements of F_0 can be compared with indirect measurements, which requires to measure the length of long lived D^0 .

The tag tag is automatically present in $D^0 \rightarrow DK^-$ decays, where D is a superposition of the flavor eigenstates D_s^0 and D^0 . An important class of D -meson decays for this analysis is the $D \rightarrow K^+ K^-$ decay. This decay can also be used along with mixed CP content to measure F_0 , provided that this content is known [4].

This paper presents the first model-independent measurement of the CP -even fraction F_0 in the decay $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ using 2.05 $\times 10^{10}$ CP -mixed $K^+ K^- \pi^+ \pi^-$ events recorded by the BESIII experiment [5].

The paper is organized as follows. In Sec. II we present the global analysis of the $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ decay. The global fit is composed of three parts: the global fit of the CP -even fraction, the global fit of the CP -odd fraction and the global fit of the CP -odd fraction. The global fit of the CP -even fraction is performed with data collected by CLEO [6] and BESIII [5]. The global fit of the CP -odd fraction is input to future analyses of $D^0 \rightarrow D^0 \pi^0$ oscillations using this dataset.

II. MEASUREMENT STRATEGY

The strong decay of $D^0 \rightarrow D$ conserves the CP quantum number. The D meson can be in an CP -even wave function. This quantum correlation allows for a direct access to the CP -even fraction F_0 in the decay $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ through a double-tag (DT) analysis. The method uses single-tag (ST) analysis, which consists of the reconstruction of a signal mode in a CP eigenstate with its requirements on the decay of the other meson, and DT analysis, which consists of the reconstruction of a tag mode and the other in the signal mode.

In Table I lists all the tag modes used for this analysis. The analysis can be split into three categories: CP tag, ST

= $A_{K^+ K^-}^{CP}$ and $A_{\pi^+ \pi^-}^{CP}$. The tag tag are modes in which the D^0 meson decays to a CP eigenstate. The modes $D \rightarrow K^+ K^-$ and $D \rightarrow \pi^+ \pi^-$ are CP -odd eigenstates. A tag tag can proceed through both CP -even and CP -odd eigenstates. The tag tag is $D \rightarrow K^+ K^-$ because it is the CP -eigenstate of the tag tag. The tag tag is $D \rightarrow \pi^+ \pi^-$ because it is the CP -eigenstate of the tag tag. It is very close to unity. The modes $D \rightarrow A_{K^+ K^-}^{CP}$, $D \rightarrow A_{\pi^+ \pi^-}^{CP}$ and $D \rightarrow \bar{D}^0 \pi^0$ are CP -odd eigenstates. These modes are included because their yields are low and the inclusion of these tag modes will significantly improve the precision of the measurement.

Table I. Tag modes used in this analysis.

CP tag	CP tag	CP tag	CP tag	CP tag	CP tag	CP tag	CP tag
$D \rightarrow K^+ K^-$	$D \rightarrow \pi^+ \pi^-$	$D \rightarrow A_{K^+ K^-}^{CP}$	$D \rightarrow A_{\pi^+ \pi^-}^{CP}$	$D \rightarrow \bar{D}^0 \pi^0$	$D \rightarrow K^+ K^-$	$D \rightarrow \pi^+ \pi^-$	$D \rightarrow A_{K^+ K^-}^{CP}$
$D \rightarrow K^+ K^-$	$D \rightarrow \pi^+ \pi^-$	$D \rightarrow A_{K^+ K^-}^{CP}$	$D \rightarrow A_{\pi^+ \pi^-}^{CP}$	$D \rightarrow \bar{D}^0 \pi^0$	$D \rightarrow K^+ K^-$	$D \rightarrow \pi^+ \pi^-$	$D \rightarrow A_{\pi^+ \pi^-}^{CP}$

The predicted ST yield of a tag mode $D \rightarrow f$ with CP -even fraction F_f^0 is given by

$$N^{ST}(f) = 2N_{DD}R_f\epsilon_{tag}F_fN(f)(1 - (2F_f^0 - 1)) \quad (1)$$

where N_{DD} is the total number of DD pairs, R_f is the branching fraction, ϵ_{tag} is the reconstruction efficiency of the tag tag, $N(f)$ is the total number of f decays, F_f^0 is the CP -even fraction of the f decay, ϵ_{tag} is the tag tag reconstruction efficiency and ϵ_{tag} is the tag tag reconstruction efficiency. In Eq. (1) and subsequent equations, CP -odd fractions F_f^1 are neglected. For pure CP -even tag tag, $F_f^1 = 1 - F_f^0$.

Events where one tag meson is reconstructed in the signal mode and the other in the tag tag mode are reconstructed as a pair of mixed CP tag mode. If there is no reconstructed DT yield

$$N^{DT}(K K \bar{K} \bar{K} f) = N_{DD}R_f\epsilon_{tag}K\epsilon_{tag}(K K \bar{K} \bar{K} f) \quad (2)$$

where $R_f\epsilon_{tag}K\epsilon_{tag}$ is the branching fraction of $D \rightarrow K^+ K^- \pi^+ \pi^-$, ϵ_{tag} is the tag tag reconstruction efficiency of the DT event, and F_f^0 denotes the CP -even fraction of $D \rightarrow K^+ K^- \pi^+ \pi^-$. Equations (1) and (2) can be combined to obtain

$$\frac{N^{ST}(f)}{N^{DT}(K K \bar{K} \bar{K} f)} = \frac{\epsilon_{tag}(f)}{\epsilon_{tag}(K K \bar{K} \bar{K} f)} \quad (3)$$

Equation (3) indicates that the ratio of the DT to ST yields, after efficiency corrections, is sensitive to

A study of CP violation in the decays $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^\pm$ ($h = K, \pi$) and $B^\pm \rightarrow [\pi^+ \pi^- \pi^+ \pi^-]_D h^\pm$

LHCb collaboration

Abstract

The first study of CP violation in the decay mode $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^\pm$ ($h = K, \pi$) is presented, exploiting a data sample of proton-proton collisions collected by the LHCb experiment that corresponds to an integrated luminosity of 9 fb^{-1} . The analysis is performed in bins of phase space, which are optimised for sensitivity to local CP violation. The results are compared to the predictions of the Standard Model and the Unitarity Triangle is determined. The analysis requires external information on decay-decay parameters, which are currently taken from an amplitude model, but can be updated in future when direct measurements become available. Measurements are also performed of phase-space integrated observables for $B^0 \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^0$ and $B^0 \rightarrow [\pi^+ \pi^- \pi^+ \pi^-]_D h^0$.

To be submitted to Eur. Phys. J. C

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Authors are listed at the end of this paper.

Analysis work summary

- BESIII paper: Measurement of the CP-even fraction of $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$
 - Has gone through all BESIII reviews, except for Spokesperson approval
 - We were unfortunately assigned the spokesperson with a backlog of 5 papers, expect 2-3 months delay...

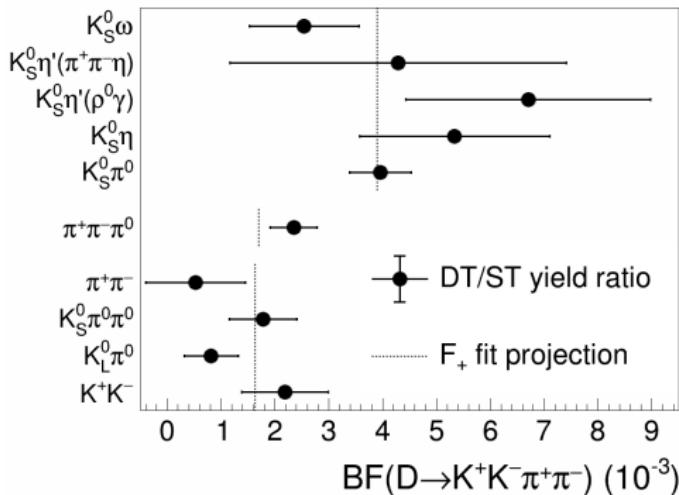


Figure 2: CP-even and CP-odd branching fractions of $D \rightarrow K^+ K^- \pi^+ \pi^-$

Analysis work summary

- LHCb paper: A study of CP violation in the decays $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^\pm$ ($h = K, \pi$) and $B^\pm \rightarrow [\pi^+ \pi^- \pi^+ \pi^-]_D h^\pm$
 - Binned model dependent GGSZ analysis of $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^\pm$
 - Phase space integrated GLW analysis of $B^\pm \rightarrow [h^+ h^- \pi^+ \pi^-]_D h^\pm$
 - Paper has just been through 1st collaboration wide circulation, currently waiting for RC and EB to review the new draft

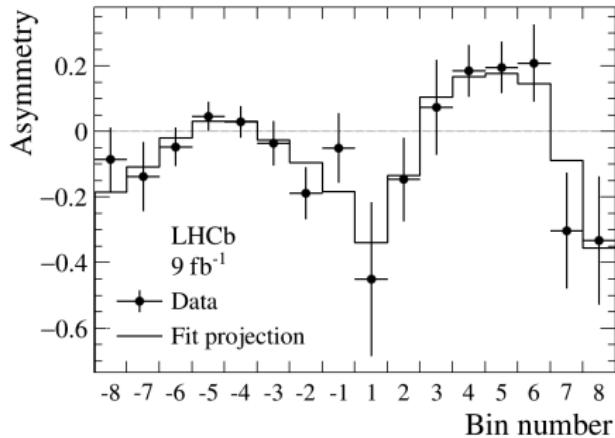


Figure 3: $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D K^\pm$ bin asymmetries

ARC: A solution for PID at FCC-ee Array of RICH Cells

A side project I've worked on while my BESIII and LHCb papers go through review

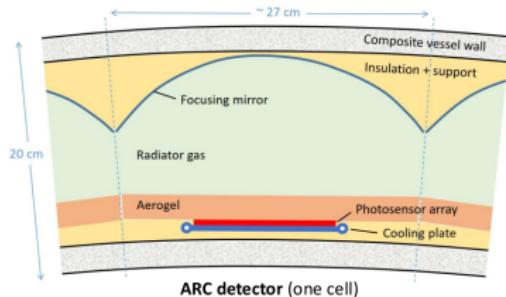
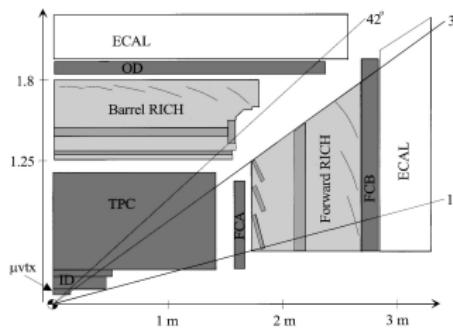


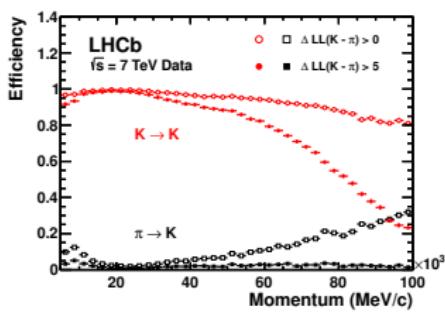
Figure 4: ARC: Array of RICH Cells

RICH detectors

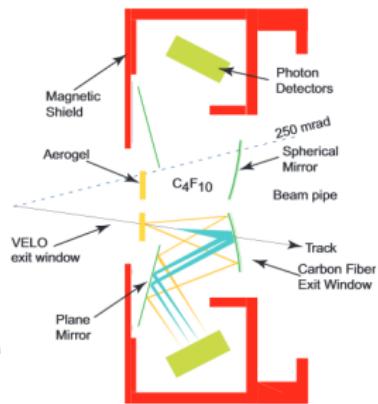
- Excellent hadron PID is crucial in flavour physics
- RICH detectors are very powerful for particle ID at high momentum
- At LHCb, π - K separation is excellent up to 100 GeV
- A 4π collider RICH layout was previously used at DELPHI and SLD
 - Challenging because of the space required



(a) DELPHI RICH layout



(b) LHCb RICH performance



(c) LHCb RICH layout

Motivation for RICH at FCC-ee

- FCC-ee will collect 5×10^{12} Z boson decays in 4 years
 - Allows for a world-leading flavour physics programme
 - Combined with excellent PID capabilities, FCC-ee will reach an unprecedented precision
- Good PID performance is also required for Higgs, WW and $t\bar{t}$ physics
 - In particular, kaon ID is crucial for $H \rightarrow s\bar{s}$

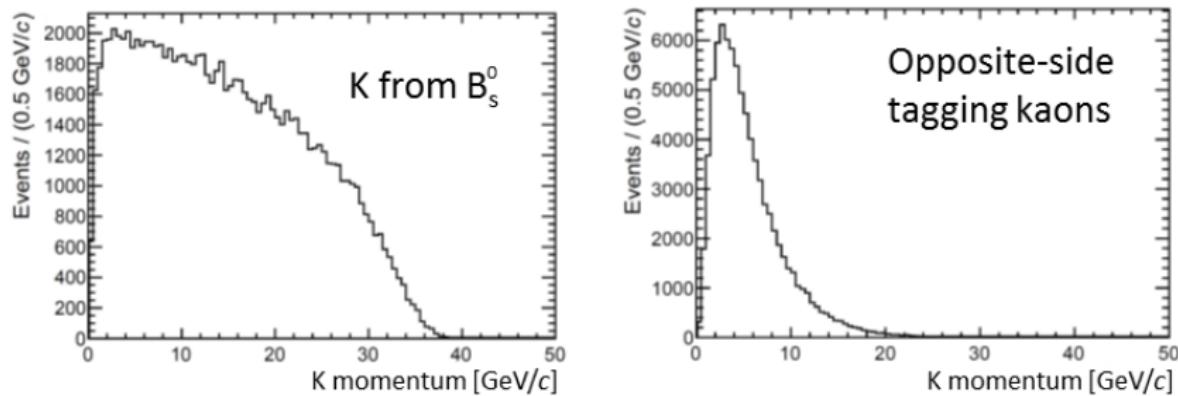


Figure 6: $B_s^0 \rightarrow D_s^\pm K^\mp$

B physics requires pion-kaon separation from low momentum up to 40 GeV

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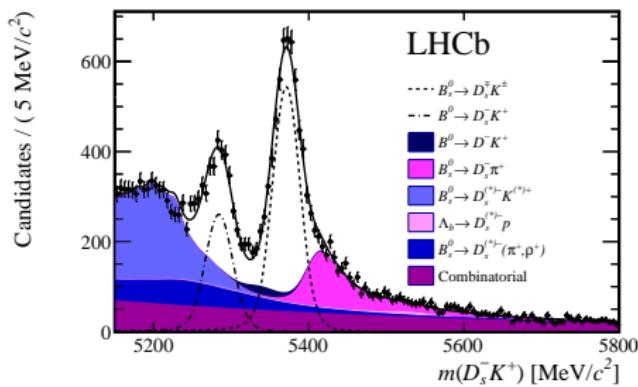


Figure 7: $B_s^0 \rightarrow D_s^{\pm} K^{\mp}$

The $B_s^0 \rightarrow D_s^{\pm} \pi^{\mp}$ background would be 10 times larger without PID capabilities!

Array of RICH Cells

- **Array of RICH Cells (ARC):** A novel RICH detector concept
 - First presented by R. Forty at [FCC week 2021](#)
 - Compact, low-mass solution for particle ID for FCC-ee
 - Concept inspired by the compound eyes of an insect
- Adapted to fit into the [CLD experiment](#) concept, taking 10% from the tracker volume
 - Radial depth of 20 cm, radius of 2.1 m and a length of 4.4 m
 - Aim to keep material budget below $0.1X_0$
- Aerogel and gas radiators with a spherical mirror
 - Aerogel also acts as thermal insulation between gas and detector

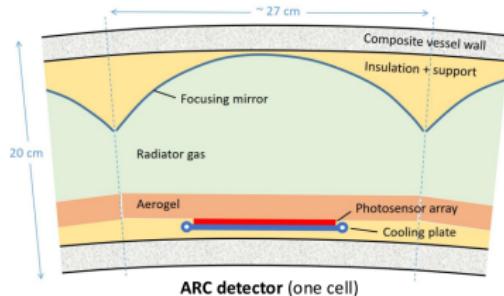


Figure 8: ARC has a cellular structure, similar to an insect's compound eyes

Array of RICH Cells

- All cells are the same size, organised on a hexagonal grid
 - Barrel (endcap) has 945 (384) cells in total, where 18 (21) are unique
 - Hexagonal shape avoids the corners, where performance is worse

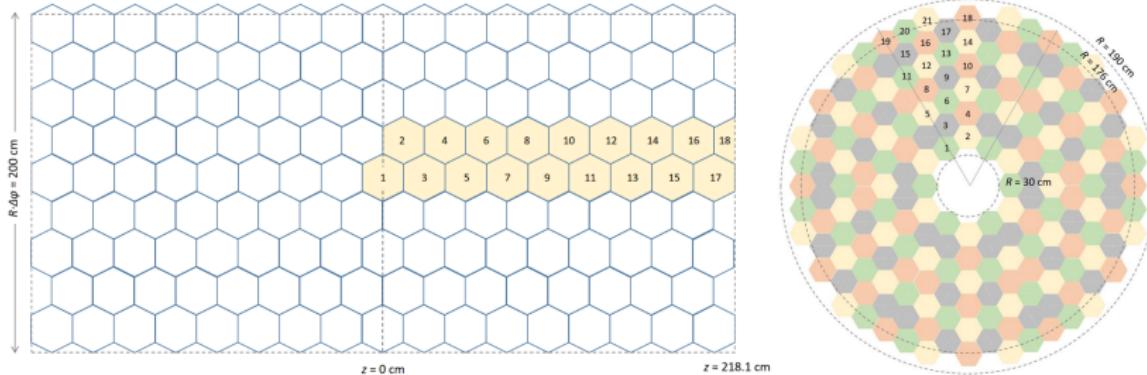


Figure 9: Barrel (left) and endcap (right) cells

ARC radiators

- C_4F_{10} :
 - Baseline assumption, well known from LHCb RICH1
 - $n = 1.0014 \Rightarrow \theta_c = 53 \text{ mrad}$, suitable for high momentum particles
 - C_4F_{10} is a greenhouse gas, plan to replace with suitable Novec gas, such as $C_5F_{10}O$
- Aerogel:
 - Well known as a RICH radiator, e.g. from ARICH at Belle II
 - $n = 1.01-1.10 \Rightarrow \theta_c = 141-430 \text{ mrad}$, suitable at low momentum
 - Very low thermal conductivity
 - Suitable to separate gas from detector, which must be cooled
 - Cherenkov photons come for “free” and are focused by the same mirror
 - Drawback: Some loss of photons from scattering

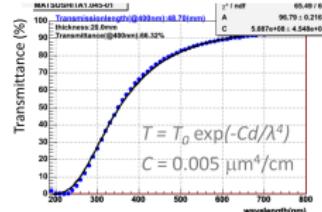


Figure 10: Belle aerogel tiles (left) and aerogel transmission function (right).

Photon hits

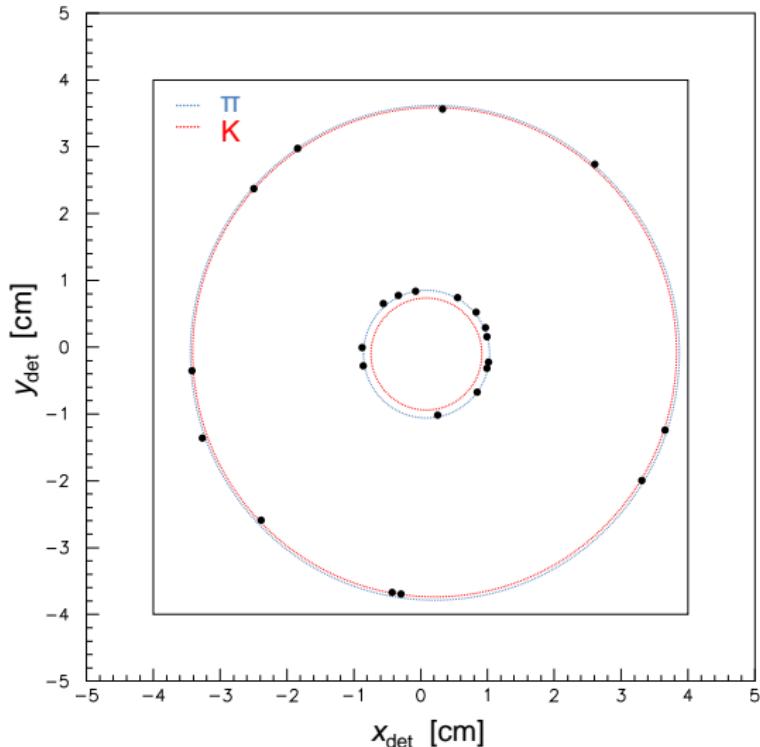


Figure 11: Photon hits on photodetector

Event display

Display of a simulated $B_s \rightarrow D_s K$ event in ARC

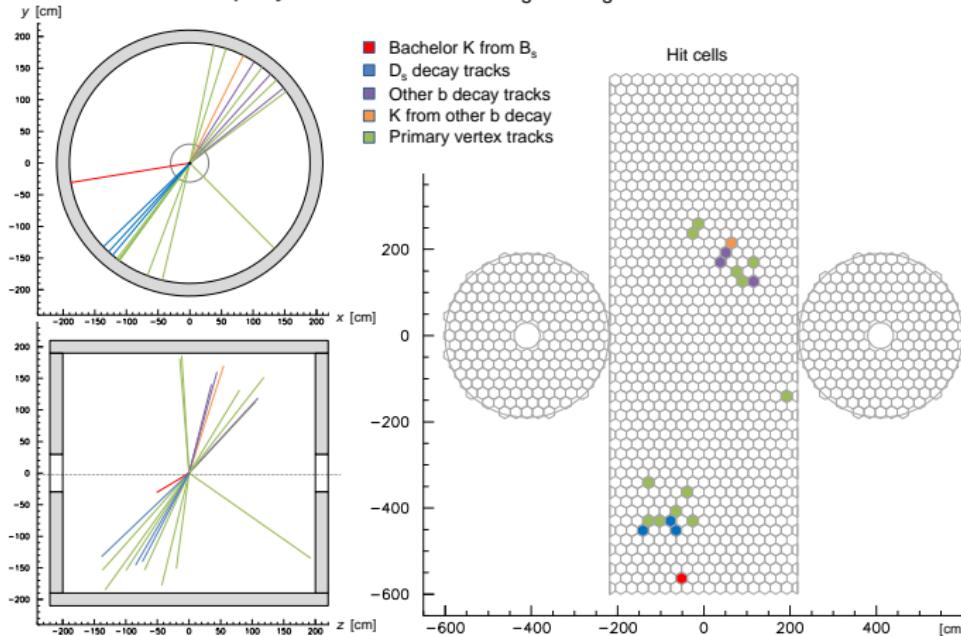


Figure 12: $B_s \rightarrow D_s K$ (no magnetic field yet)

Optimisation of ARC layout

- The following procedure is used to evaluate the ARC performance:
 - ① Generate straight particle track from IP and trace it through ARC
 - ② Generate Cherenkov photons from gas radiator
 - ③ Track photons through the optics and to detector
 - ④ Reconstruct Cherenkov angles and calculate standard deviation
- Three sources of uncertainty are considered:
 - ① Emission point uncertainty: Emission point is assumed to be the mid-point of the track inside the gaseous radiator
 - ② Chromatic dispersion uncertainty: Spread in Cherenkov angle due to wavelength dependence on refractive index
 - ③ Pixel size: Will be chosen so that it does not limit the performance

Minimise the Cherenkov angle uncertainty:

$$\Delta\theta = \frac{1}{\sqrt{N}} \times \frac{1}{1-N} \times \sum_{i=0}^{N-1} (\theta - \bar{\theta})^2$$

Examples of photon tracking through optimised layout

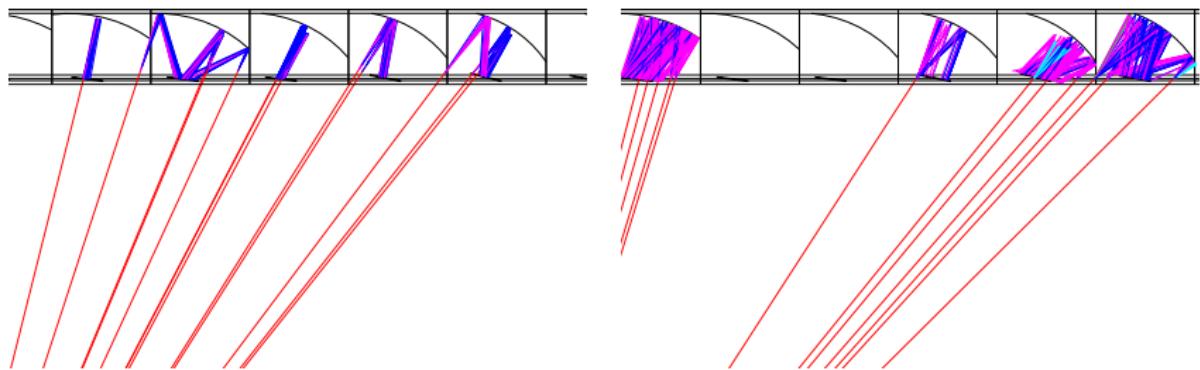


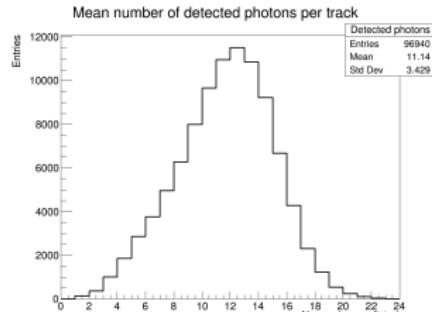
Figure 13: Tracking of photons from gas radiator (left) and aerogel radiator (right) through the ARC optics

- Parameters that are optimised:
 - Mirror curvature
 - Mirror vertical and horizontal position
 - Detector horizontal position and tilt

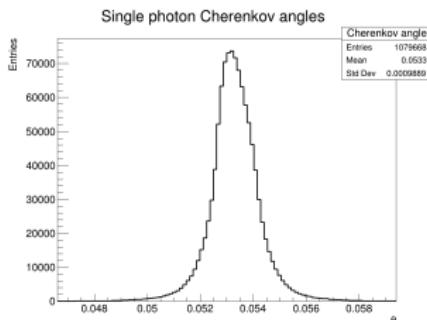
Technical details about minimisation

- Total Cherenkov angle uncertainty $\sigma(\vec{x})$ is not easy to calculate analytically
 - \vec{x} are the 5 parameters we want to optimise
- Finite number of photons $\implies \sigma(\vec{x})$ is not differentiable
 - Cannot be minimised using conventional methods (Minuit, etc)
- I have experimented with a new type of minimisation algorithms:
Stochastic optimisation
 - Differential evolution
 - Start with a population of possible solutions, form new solutions by combining (mutating) existing solutions
 - Advantage: Doesn't require initial guess, robust against functions that are not continuous, noisy, change over time, etc
 - Disadvantage: No way to tell if optimal solution has been found, so it requires many iterations

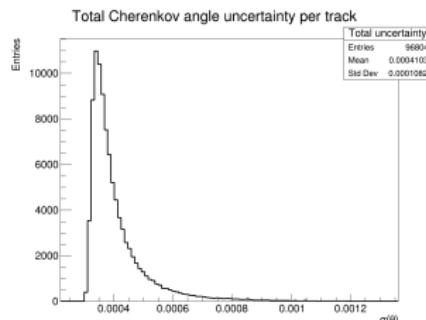
Cherenkov angle uncertainty for gas radiator



(a) Mean number of photons detected



(b) Single photon uncertainty:
1.0 mrad



(c) Total uncertainty:
0.4 mrad

Figure 14: Gas radiator performance averaged over all barrel cells

Cherenkov angle uncertainty for aerogel radiator

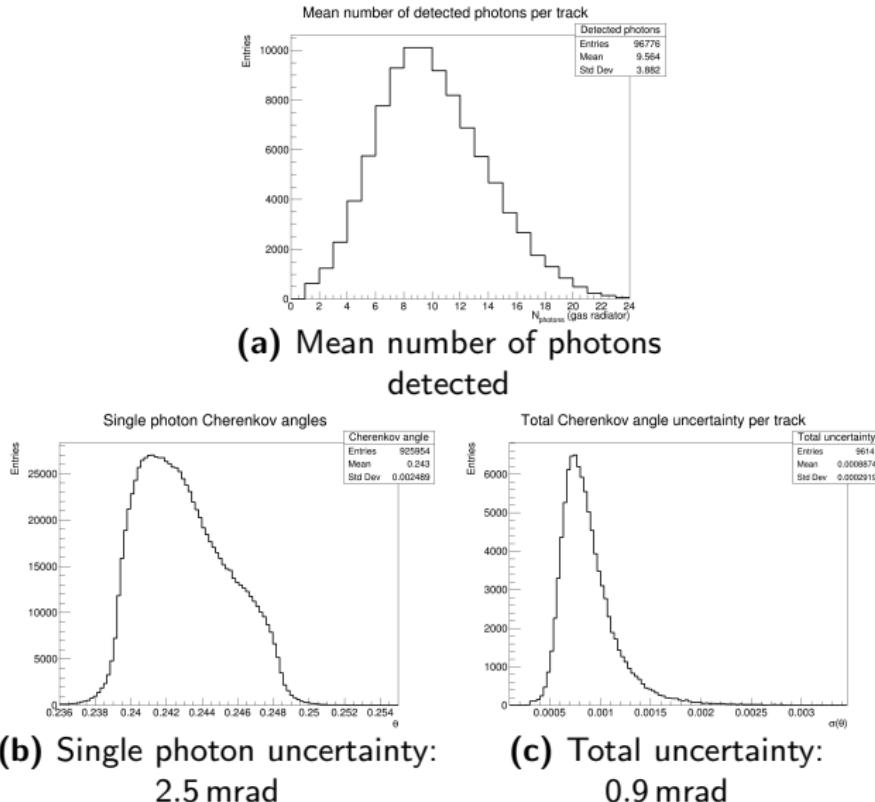


Figure 15: Aerogel radiator performance averaged over all barrel cells

Performance of optimised ARC

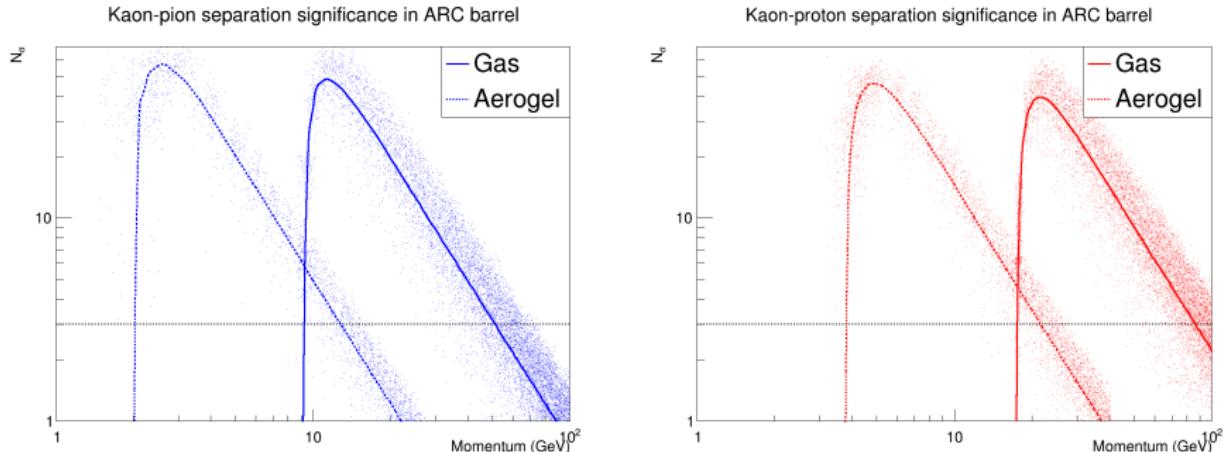


Figure 16: Separation significance per track for π - K (left) and p - K (right)

- Gas (aerogel) provides over 3σ pion-kaon separation in the range 10-50 GeV (2-10 GeV)
 - Effect of magnetic field not yet included in these studies
- Combined, the aerogel and gas ensure excellent PID performance over the whole range of interest to flavour physics

Summary and next steps

- ARC is a low mass and compact cellular PID detector designed to occupy minimum space (20 cm in the radial dimension) in a 4π detector at an e^+e^- collider such as FCC-ee
- We have developed an optimised layout that should achieve a 3σ kaon-pion separation in the range 2-50 GeV
 - Our studies focus mainly on flavour physics at the Z-pole
- ARC will allow us to fully exploit the full range of flavour physics potential at future e^+e^- colliders
 - Will enhance the capabilities in Higgs, WW and top physics
- Next steps will include completing the optimisation, including magnetic field effects, and R&D on photodetectors

Thank you for listening!