

# The angle $\gamma$ of the Cabibbo-Kobayashi-Maskawa ansatz: a journey towards precision at LHCb

Martin Tat, on behalf of the LHCb collaboration

University of Oxford

CERN LHC seminar

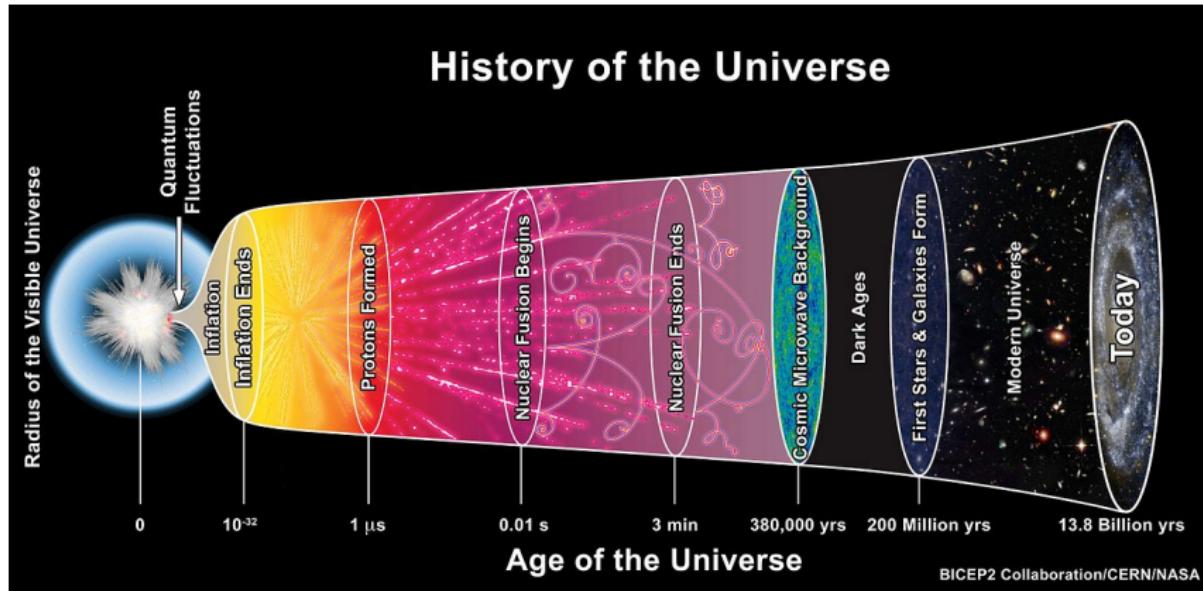
25th July 2023



## Introduction to $CP$ violation

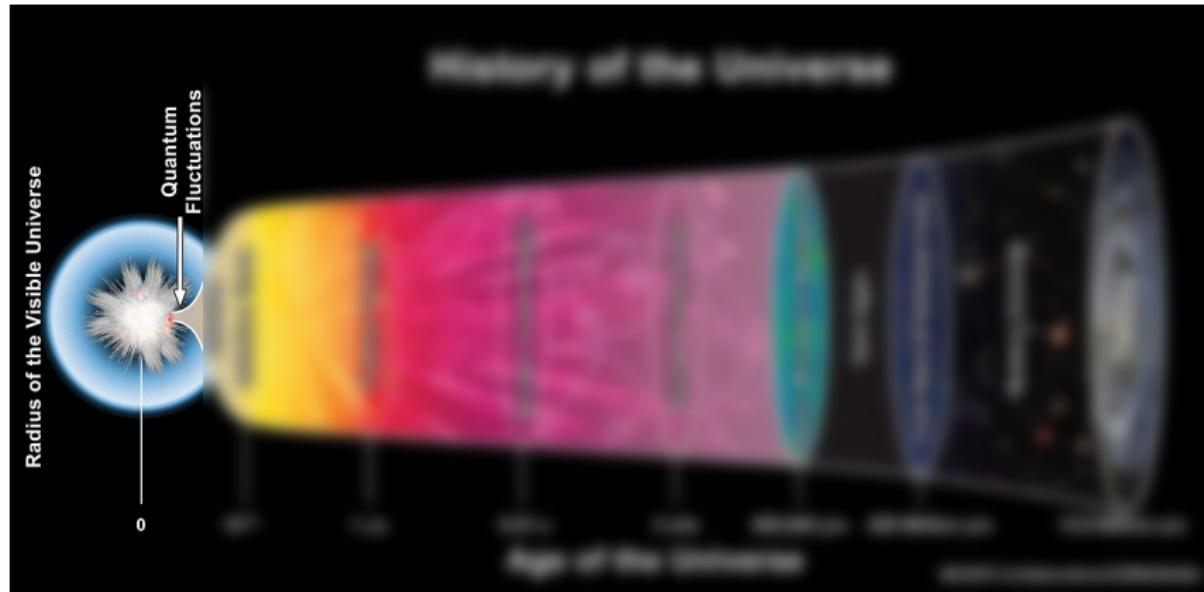
What is  $\gamma$  and why measure it?

# Big Bang and matter-antimatter asymmetry



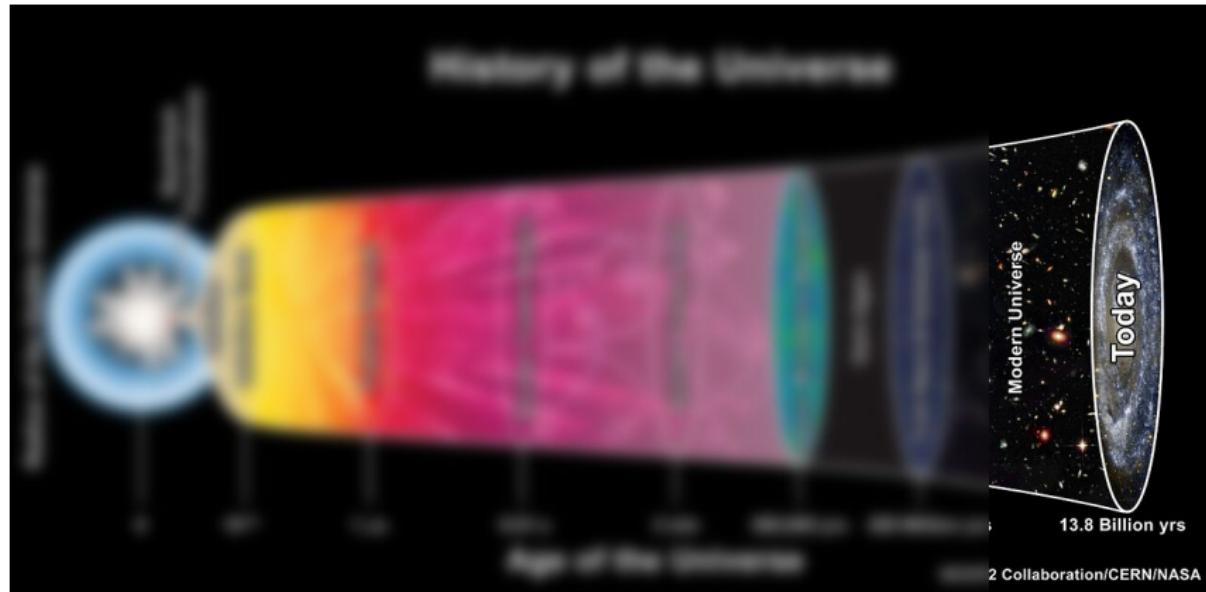
Where is the antimatter in the universe?

# Big Bang and matter-antimatter asymmetry



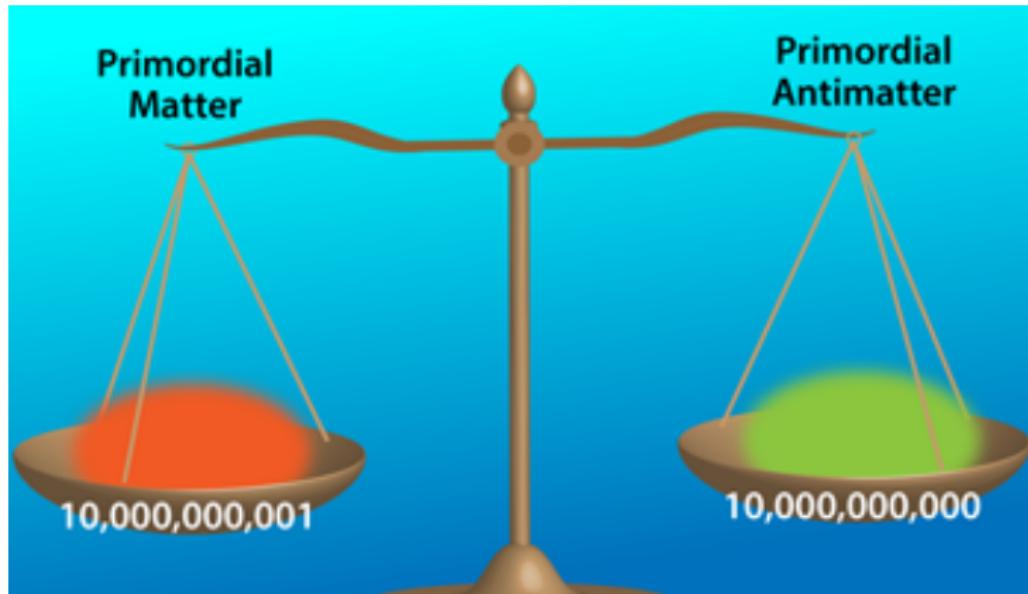
Initially equal amounts of matter and antimatter...

# Big Bang and matter-antimatter asymmetry



... but today we only see matter!

# Big Bang and matter-antimatter asymmetry



APS/Alan Stonebraker

The difference is very small...

# Big Bang and matter-antimatter asymmetry



Quantum Diaries: "Why B physics? Why not A Physics?"

... but the effects we observe today are obviously huge!  
How can we explain this?

# *CP* violation

## The Nobel Prize in Physics 1980



Photo from the Nobel Foundation archive.  
James Watson Cronin  
Prize share: 1/2

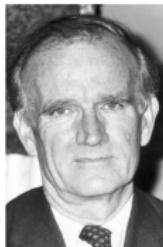


Photo from the Nobel Foundation archive.  
Val Logsdon Fitch  
Prize share: 1/2

The Nobel Prize in Physics 1980 was awarded jointly to James Watson Cronin and Val Logsdon Fitch "for the discovery of violations of fundamental symmetry principles in the decay of neutral K-mesons"

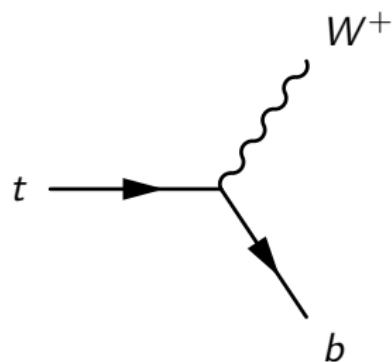
- *CP* violation discovery in 1964
- Phys. Rev. Lett. **13**, 138
- Observed  $K_L^0 \rightarrow \pi^+ \pi^-$
- Since, *CP* violation has also been observed in the  $B$ ,  $B_s$  and  $D$  systems

Can Standard Model CPV explain the matter-antimatter asymmetry?  
Or, could it be physics beyond the SM?

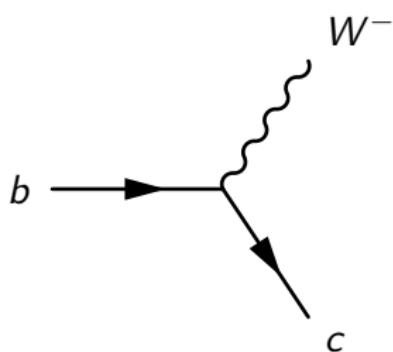
# The CKM matrix and the Unitary Triangle

In SM, the charged current  $W^\pm$  interactions couple (left-handed) up- and down-type quarks, given by

$$\frac{-g}{\sqrt{2}} \begin{bmatrix} \bar{u}_L & \bar{c}_L & \bar{t}_L \end{bmatrix} \gamma^\mu W_\mu V_{CKM} \begin{bmatrix} d_L \\ s_L \\ b_L \end{bmatrix} + \text{h.c.}$$



(a)  $t \rightarrow b W^+$



(b)  $b \rightarrow c W^-$

# The CKM matrix and the Unitary Triangle

The Cabibbo-Kobayashi-Maskawa matrix  $V_{\text{CKM}}$  has a single complex phase that is responsible for all CPV in SM

$$\begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} = \begin{bmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix} + \mathcal{O}(\lambda^4)$$

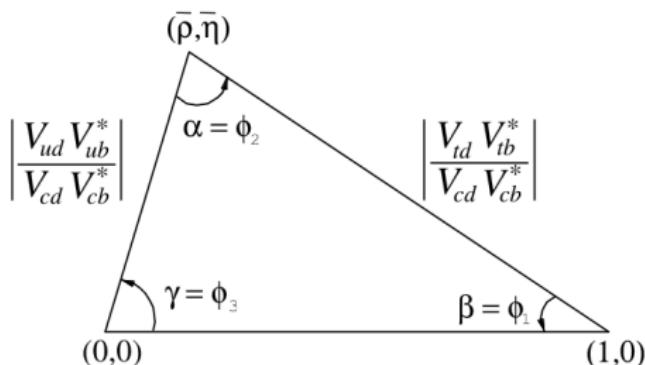
$V_{\text{CKM}}$  must be unitary, which gives us 9 constraints, such as:

$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

# The CKM matrix and the Unitary Triangle

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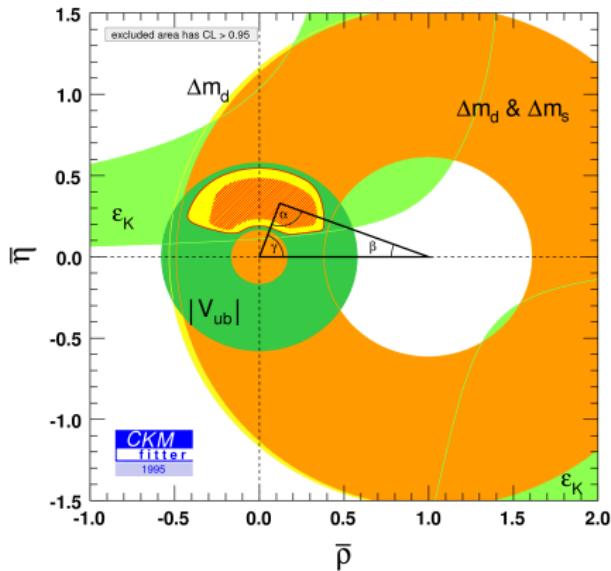
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R. L. Workman *et al.* (Particle Data Group), Prog. Theor. Exp. Phys. 2022, 083C01 (2022)

# The CKM matrix and the Unitary Triangle

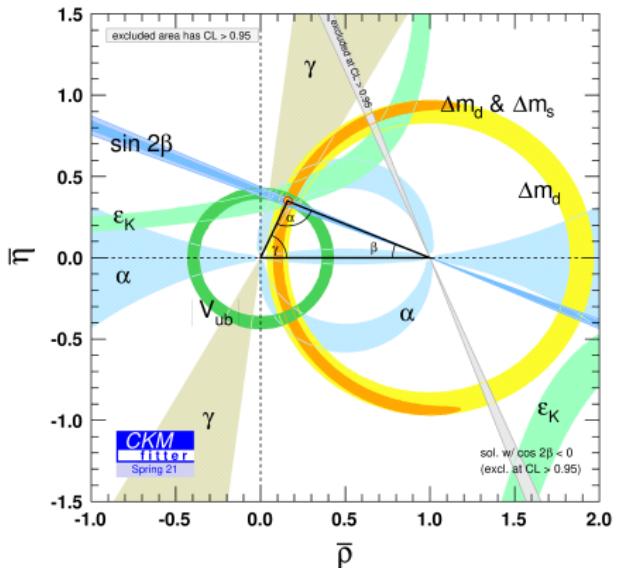
Initially, the Unitary Triangle was poorly constrained...



CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), updated results and plots available at:  
<http://ckmfitter.in2p3.fr>

# The CKM matrix and the Unitary Triangle

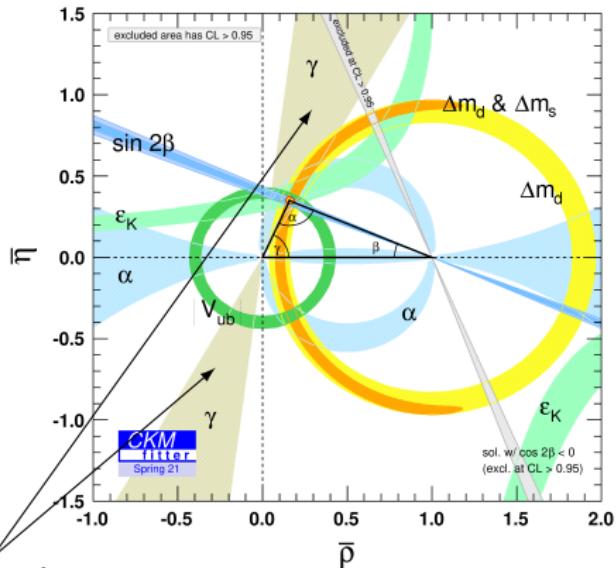
... but a tremendous effort has been made the last 30 years!



CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), updated results and plots available at:  
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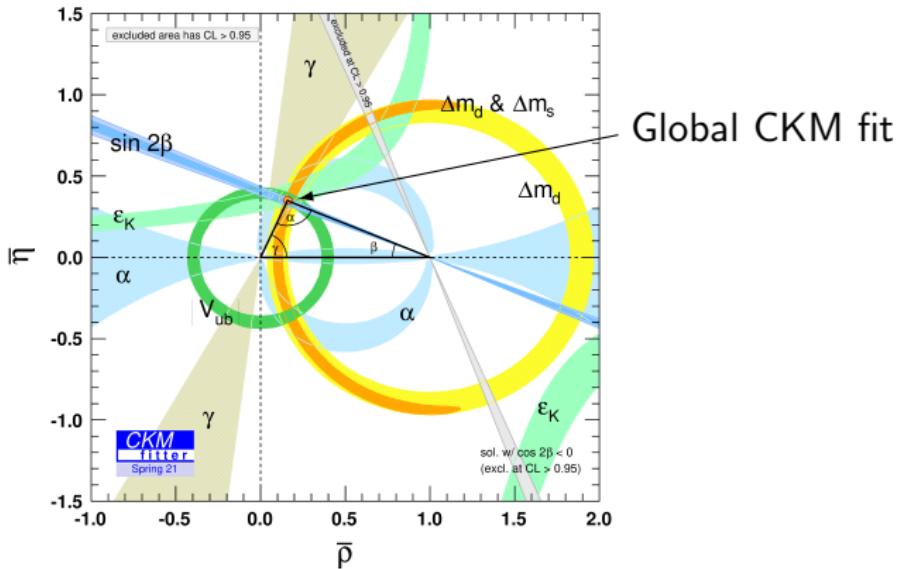


Direct  $\gamma$  measurements

CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), updated results and plots available at:  
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# The CKM matrix and the Unitary Triangle

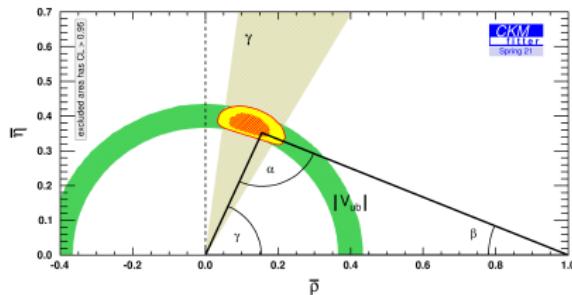
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# The CKM matrix and the Unitary Triangle

- Why is  $\gamma = \arg\left(-\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*}\right)$  important?
  - ① Negligible theoretical uncertainties: Ideal SM benchmark
  - ② Only CKM angle accessible in tree level decays
  - ③ Compare with global CKM fit: Is the Unitary Triangle a triangle?

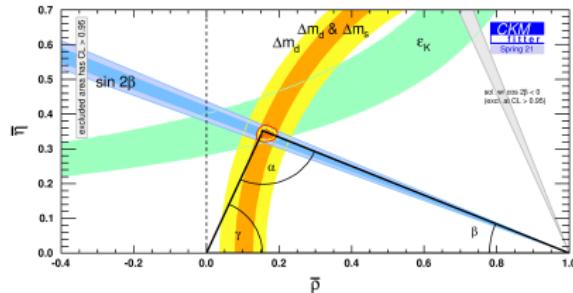


Tree level:  $\gamma = (72.1^{+5.4}_{-5.7})^\circ$

CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), updated results and plots available at:  
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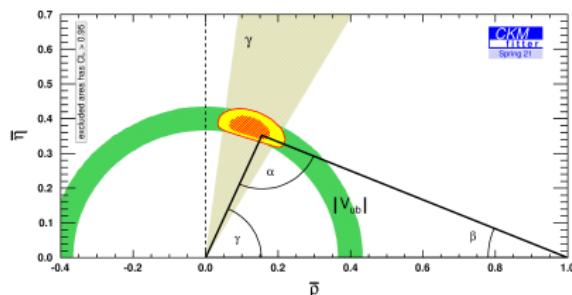


Loop level:  $\gamma = (65.5^{+1.1}_{-2.7})^\circ$

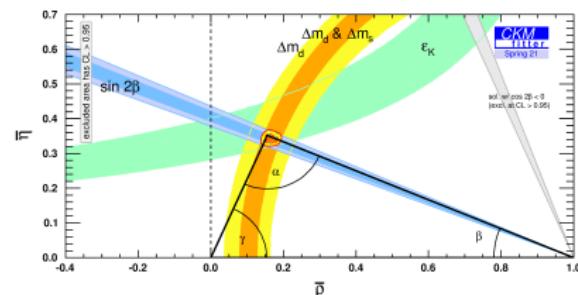
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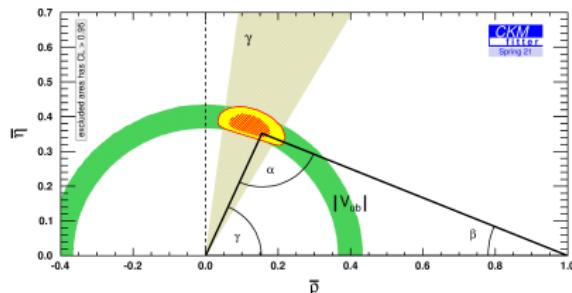


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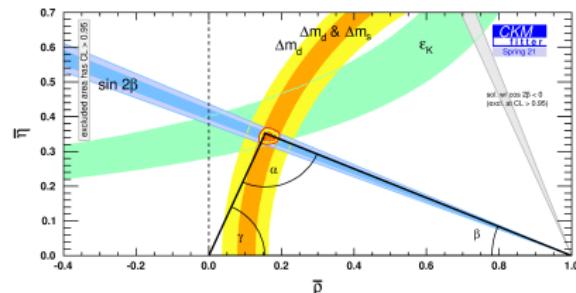
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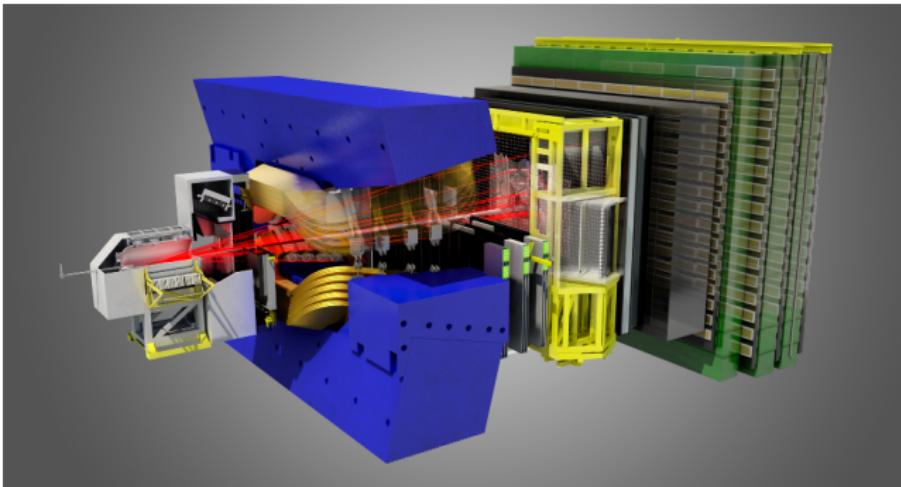
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CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41, 1-131 (2005), updated results and plots available at:  
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Need to reduce uncertainties significantly in direct measurements!

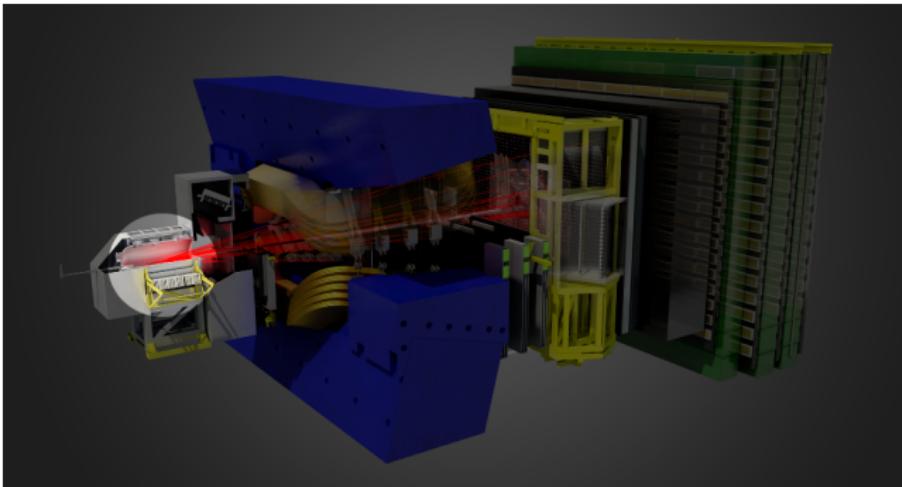
# The LHCb detector

# The LHCb detector



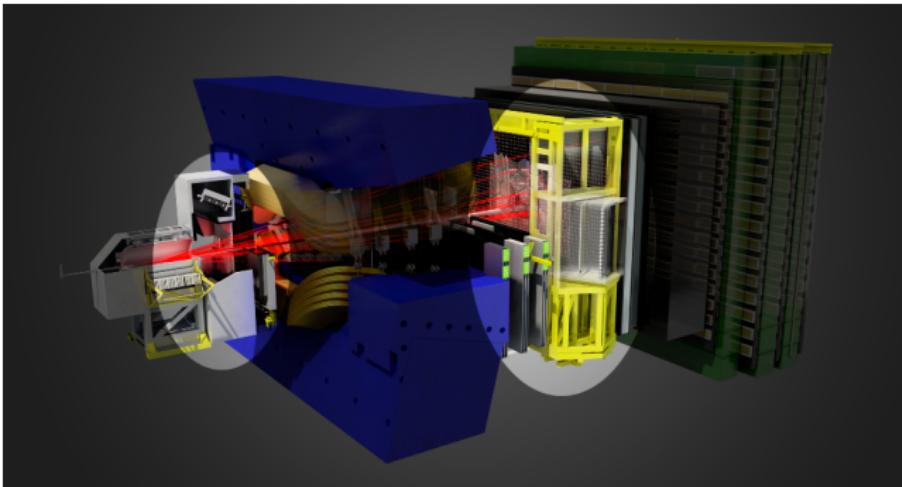
LHCb: A beauty experiment with a lot of charm

# The LHCb detector



VELO: Vertex locator to reconstruct  $B$  and  $D$  vertices

# The LHCb detector



RICH: Identify  $B$  and  $D$  daughter particles

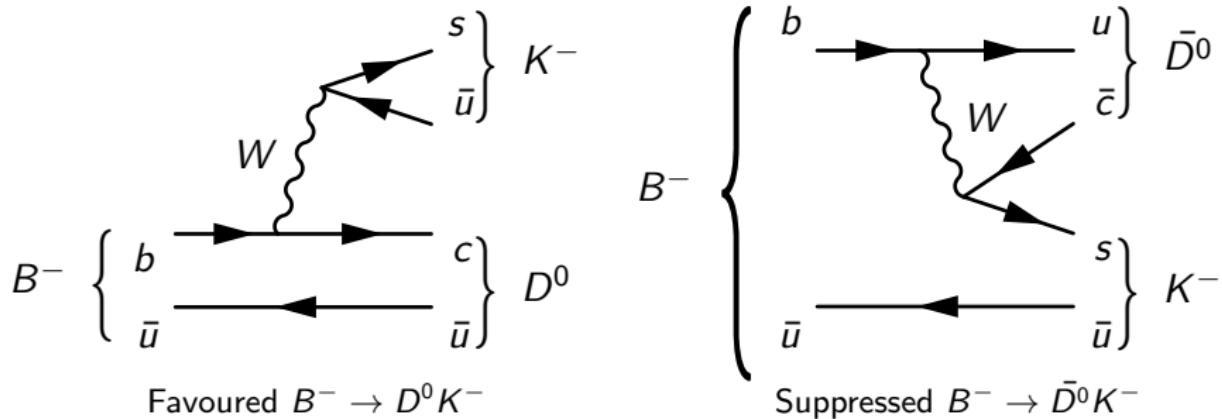
How to measure  $\gamma$ ?

# How to measure $\gamma$ ?

It's all about interferences!

# Sensitivity through interference

Measure  $\gamma$  through interference effects in  $B^\pm \rightarrow DK^\pm$



- Superposition of  $D^0$  and  $\bar{D}^0$

- $b \rightarrow u \bar{c} s$  and  $b \rightarrow c \bar{u} s$  interference  $\rightarrow$  Sensitivity to  $\gamma$

$$\mathcal{A}(B^-) = \mathcal{A}_B \left( \mathcal{A}_{D^0} + r_B e^{i(\delta_B - \gamma)} \mathcal{A}_{\bar{D}^0} \right)$$

$$\mathcal{A}(B^+) = \mathcal{A}_B \left( \mathcal{A}_{\bar{D}^0} + r_B e^{i(\delta_B + \gamma)} \mathcal{A}_{D^0} \right)$$

- The magnitude of interference effects governed by  $r_B \approx 0.1$

# What $D$ final states?

What  $D$  final states should we consider?

No single method is sufficient to determine  $\gamma$  precisely!

- ① GLW method: CP eigenstates

- $D \rightarrow K^+K^-, \pi^+\pi^-, \dots$

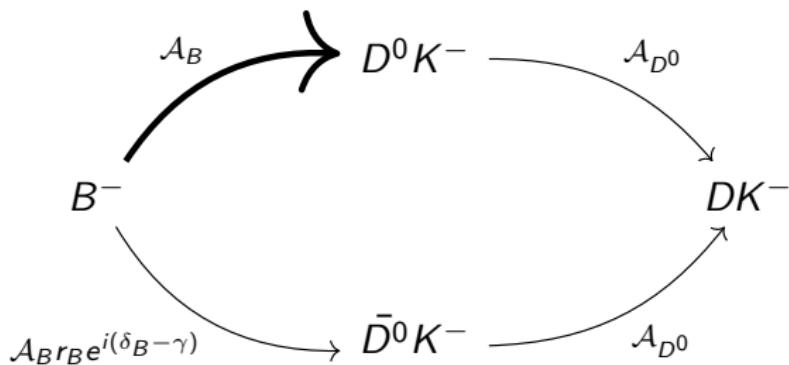
②

③

## $D$ decays to a $CP$ eigenstate

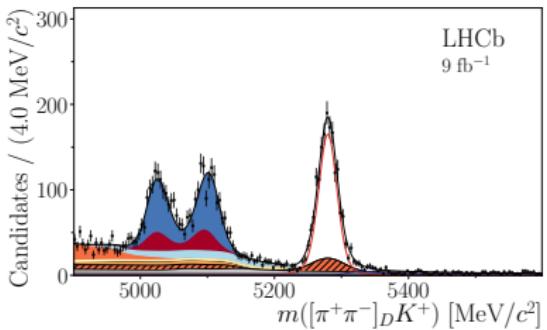
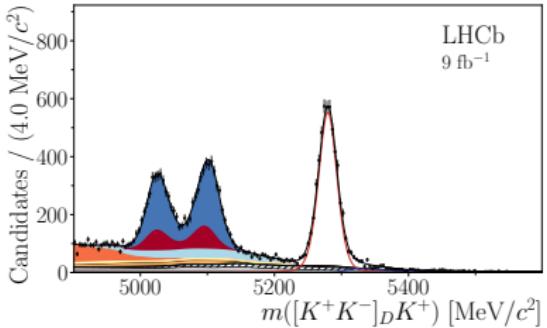
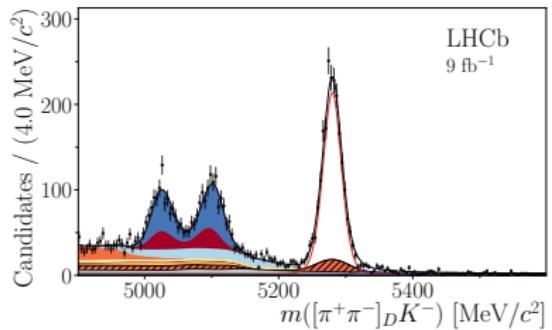
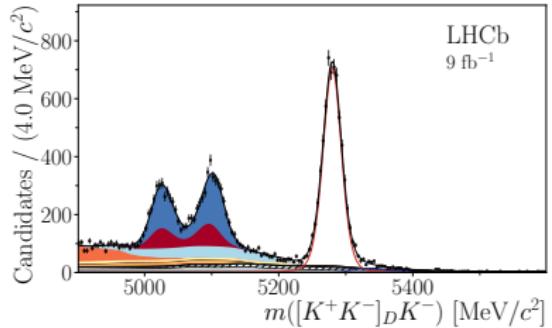
Naively, we expect the size of CPV effects to be around  $r_B \approx 10\%$

For  $CP$  eigenstates,  $\mathcal{A}_{D^0} = \mathcal{A}_{\bar{D}^0}$



$$|\mathcal{A}(B^-)|^2 \propto 1 + r_B^2 + 2r_B \cos(\delta_B - \gamma)$$

# $D$ decays to a $CP$ eigenstate



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In  $B^\pm \rightarrow [h^+ h^-]_D K^\pm$ , we see significant CPV effects

# What $D$ final states?

What  $D$  final states should we consider?

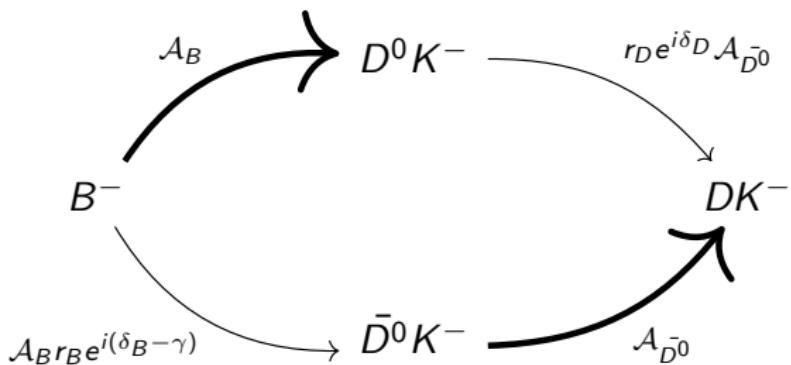
No single method is sufficient to determine  $\gamma$  precisely!

- ① GLW method: CP eigenstates
  - $D \rightarrow K^+K^-, \pi^+\pi^-, \dots$
- ② ADS method: Doubly-Cabbibo Suppressed decays
  - $D \rightarrow K^-\pi^+, K^-\pi^+\pi^-\pi^+, \dots$
- ③

# Doubly Suppressed Cabibbo $D$ decays

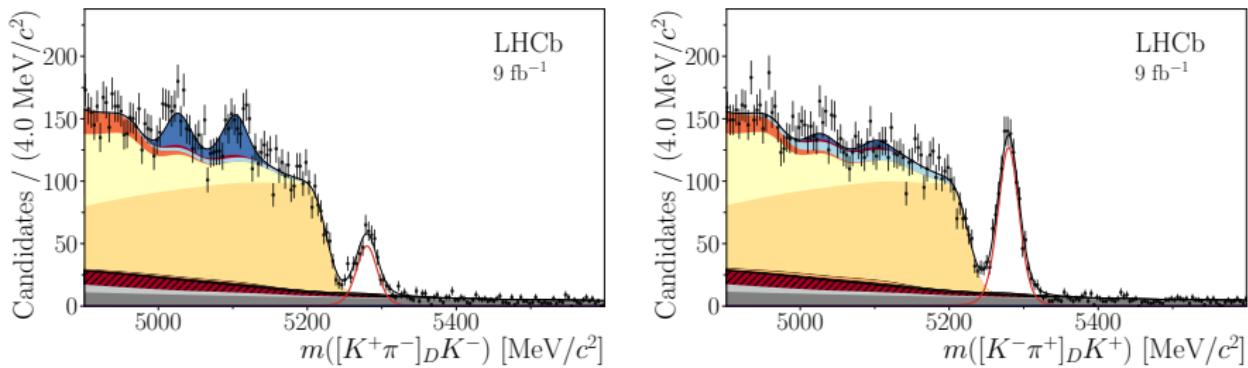
Interference effects can be greatly enhanced

Use a Doubly Suppressed Cabibbo decay:  $\mathcal{A}_{D^0} = r_D e^{i\delta_D} \mathcal{A}_{\bar{D}^0}$



$$|\mathcal{A}(B^-)|^2 \propto r_D^2 + r_B^2 + 2r_B r_D \cos(\delta_B - \gamma + \delta_D)$$

# Doubly Suppressed Cabibbo $D$ decays



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$B^\pm \rightarrow [K^\mp \pi^\pm]_D K^\pm$  has lower statistics, but a spectacular asymmetry!

Additionally, the partially reconstructed background has an equal but opposite asymmetry

# What $D$ final states?

What  $D$  final states should we consider?

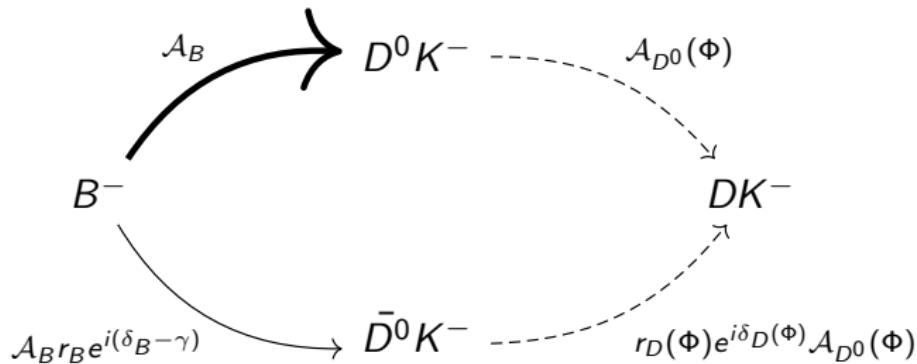
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  - $D \rightarrow K^+K^-, \pi^+\pi^-, \dots$
- ② ADS method: Doubly-Cabbibo Suppressed decays
  - $D \rightarrow K^-\pi^+, K^-\pi^+\pi^-\pi^+, \dots$
- ③ BPGGSZ method: Multi-body final states
  - $D \rightarrow K_S^0\pi^+\pi^-, K_S^0K^+K^-, \dots$

# Multi-body $D$ decays

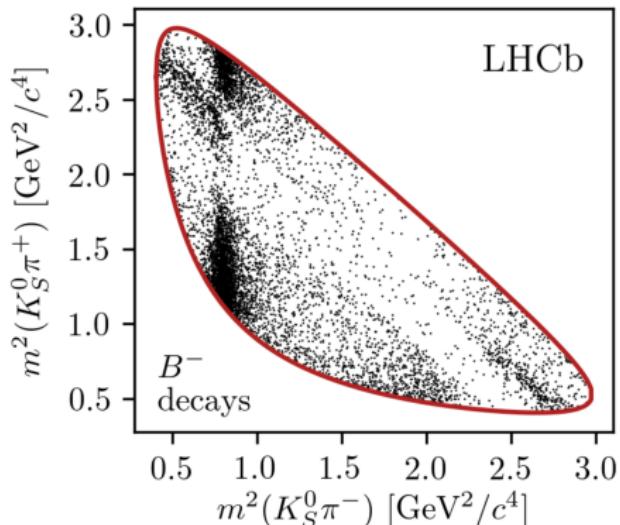
The three-body  $D$ -decay phase space is two-dimensional  $\implies$  Dalitz plot

$B^\pm$  decay rate depends on the  $D^0$  and  $\bar{D}^0$  strong-phase difference

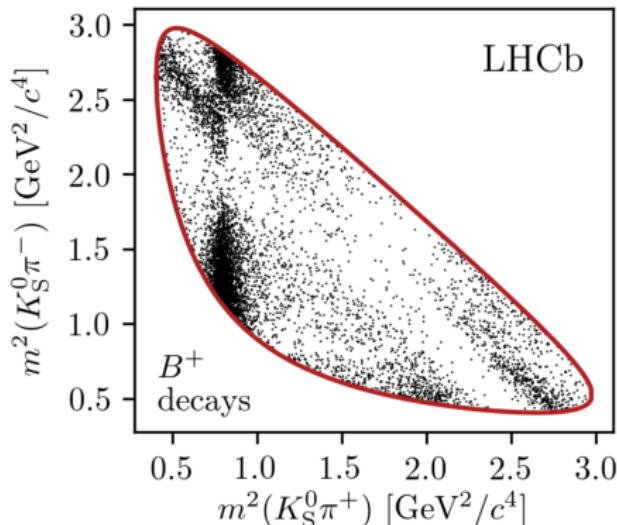


$$|\mathcal{A}(B^-)|^2 \propto 1 + r_B^2 r_{\bar{D}}^2(\Phi) + 2r_B r_{\bar{D}}(\Phi) \cos(\delta_B - \gamma + \delta_D(\Phi))$$

# Multi-body $D$ decays



$$B^- \rightarrow [K_S^0 \pi^+ \pi^-]_D K^-$$

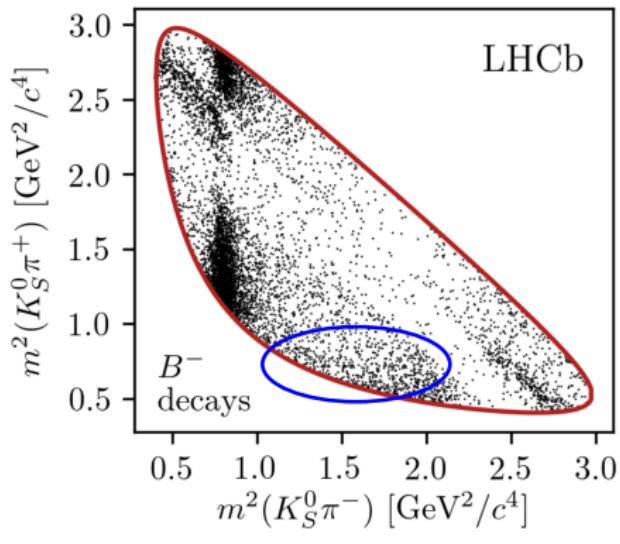


$$B^+ \rightarrow [K_S^0 \pi^+ \pi^-]_D K^+$$

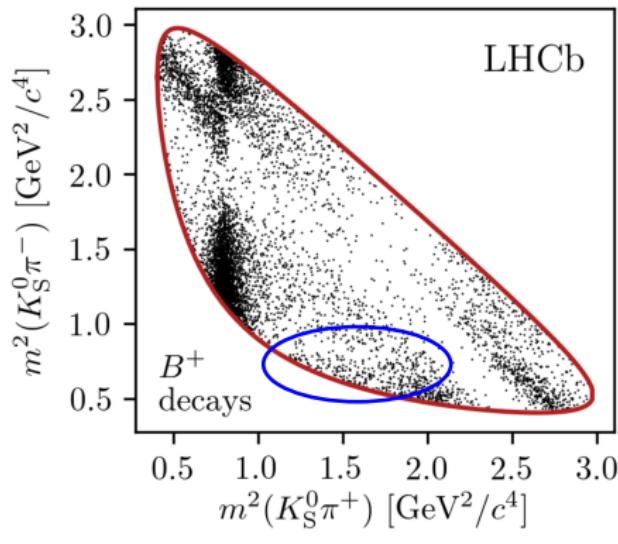
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Can you find the asymmetries?

# Multi-body $D$ decays



$$B^- \rightarrow [K_S^0 \pi^+ \pi^-]_D K^-$$

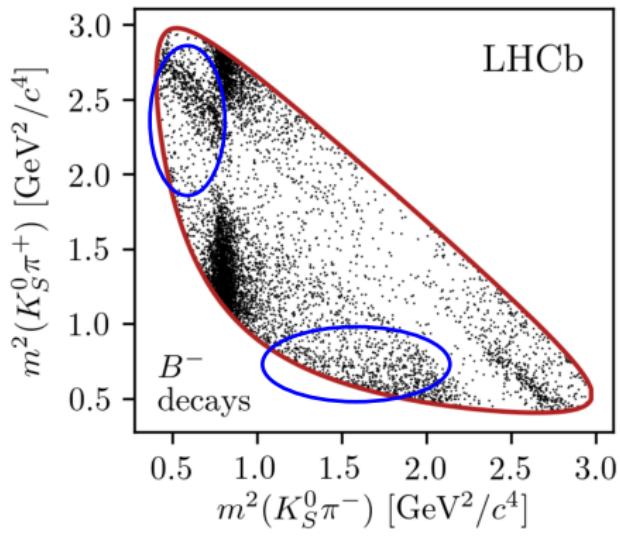


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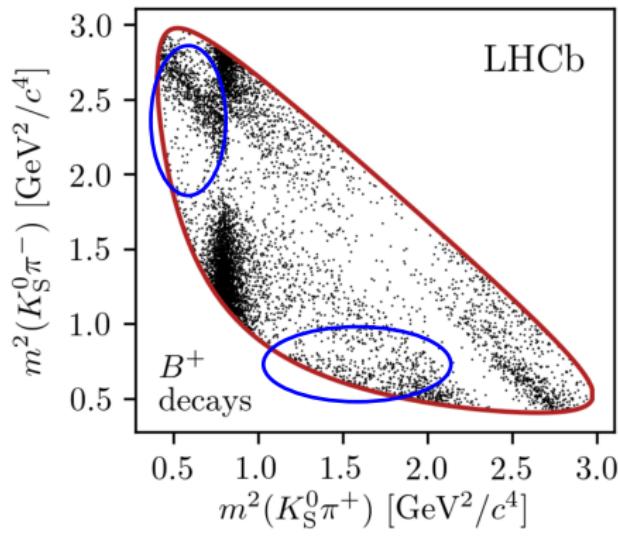
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# Multi-body $D$ decays



$$B^- \rightarrow [K_S^0\pi^+\pi^-]_D K^-$$



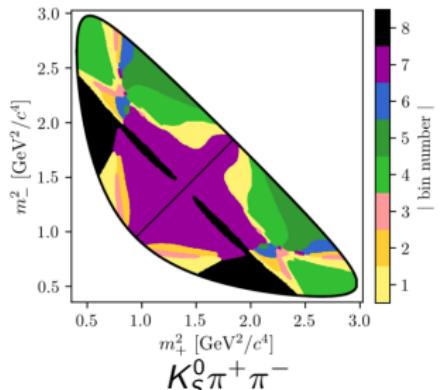
$$B^+ \rightarrow [K_S^0\pi^+\pi^-]_D K^+$$

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Can you find the asymmetries?

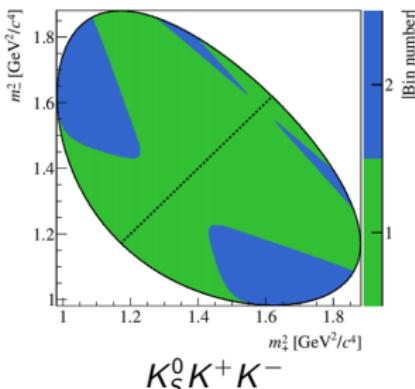
# Multi-body $D$ decays

- Interpretation of  $\gamma$  from the multi-body charm decays require external inputs of the charm strong-phase differences
- The three-body decays  $D \rightarrow K_S^0 h^+ h^-$  have been studied extensively, using an optimised phase-space binning:
  - CLEO [Phys. Rev. D82 \(2010\) 112006](#)
  - BESIII [Phys. Rev. D101 \(2020\) 112002](#)
- With charm inputs from CLEO and BESIII, the measurement of  $\gamma$  becomes model independent



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CERN LHC seminar



25th July 2023

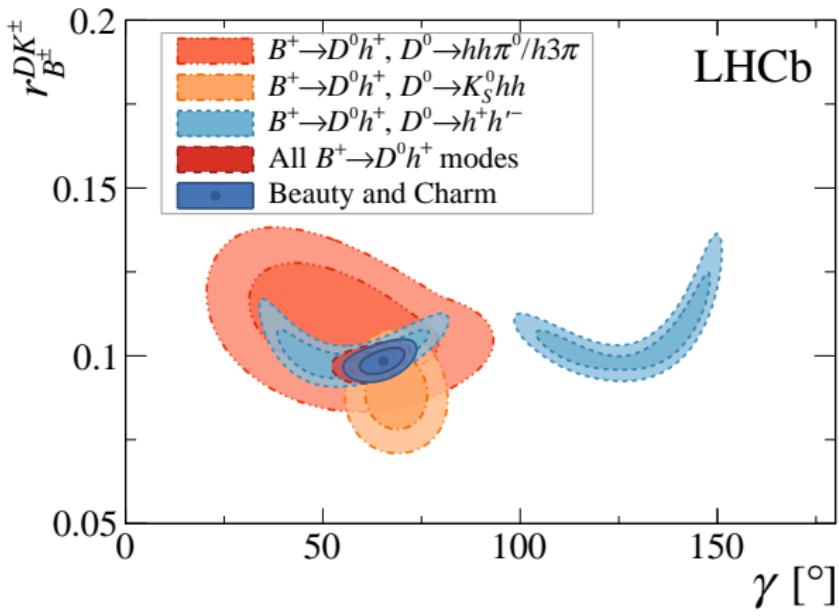
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# The LHCb $\gamma$ and charm combination

## The LHCb $\gamma$ and charm combination

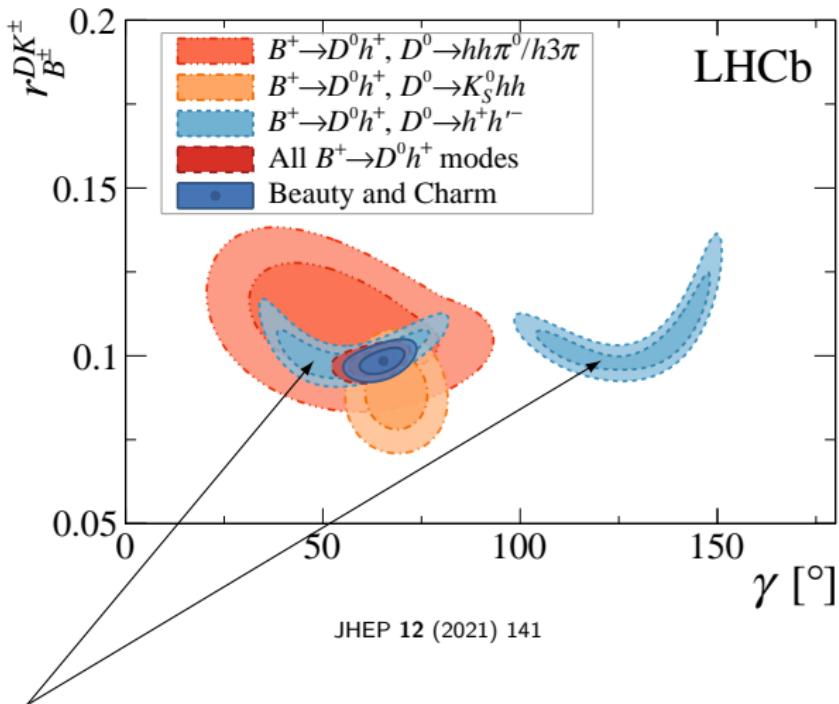
Lots of beauty in charm!

# The LHCb $\gamma$ and charm combination



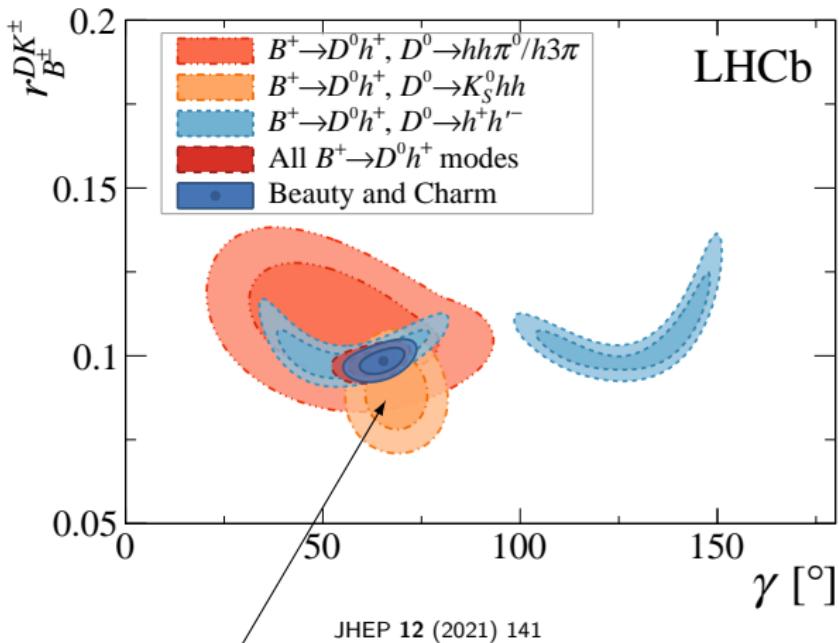
Currently,  $\gamma$  measurements are dominated by  $B^\pm \rightarrow Dh^\pm$

# The LHCb $\gamma$ and charm combination



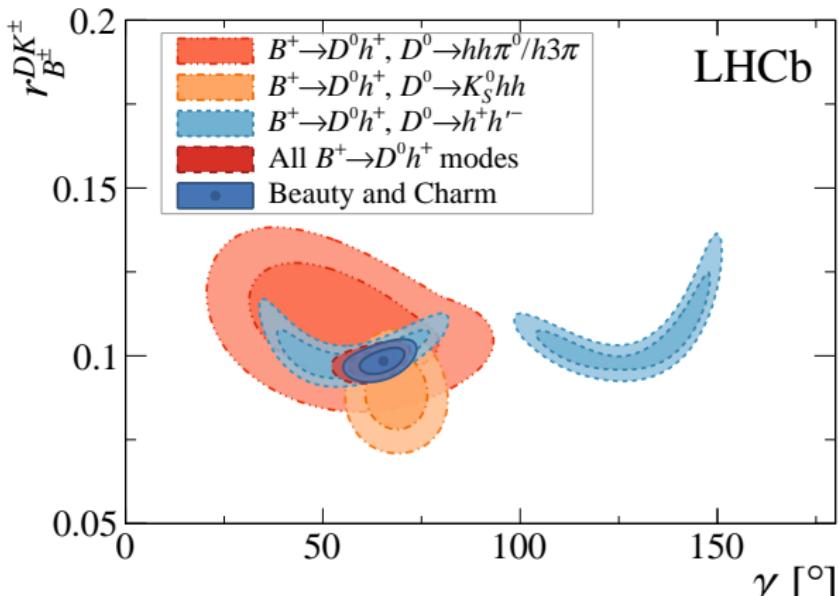
GLW+ADS: Narrow bands, but many degenerate solutions

# The LHCb $\gamma$ and charm combination



BPGGSZ: Wider, but unique solution

# The LHCb $\gamma$ and charm combination



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A combination of direct  $\gamma$  measurements is necessary!

# The LHCb $\gamma$ and charm combination

Other  $B$  decays are also interesting for  $\gamma$  measurements:

- ①  $B^\pm \rightarrow DK^\pm$
- ②  $B^\pm \rightarrow D^{*0} K^\pm$
- ③  $B^0 \rightarrow DK^{*0}$
- ④  $B_s^0 \rightarrow D_s^- K^+$
- And many more...

# The LHCb $\gamma$ and charm combination

Other  $B$  decays are also interesting for  $\gamma$  measurements:

- ①  $B^\pm \rightarrow DK^\pm \leftarrow$  Golden mode
- ②  $B^\pm \rightarrow D^{*0} K^\pm$
- ③  $B^0 \rightarrow DK^{*0}$
- ④  $B_s^0 \rightarrow D_s^- K^+$ 
  - And many more...

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- ①  $B^\pm \rightarrow DK^\pm \leftarrow$  Golden mode
- ②  $B^\pm \rightarrow D^{*0} K^\pm \leftarrow$  New results!
- ③  $B^0 \rightarrow DK^{*0} \leftarrow$  New results!
- ④  $B_s^0 \rightarrow D_s^- K^+$ 
  - And many more...

# The LHCb $\gamma$ and charm combination

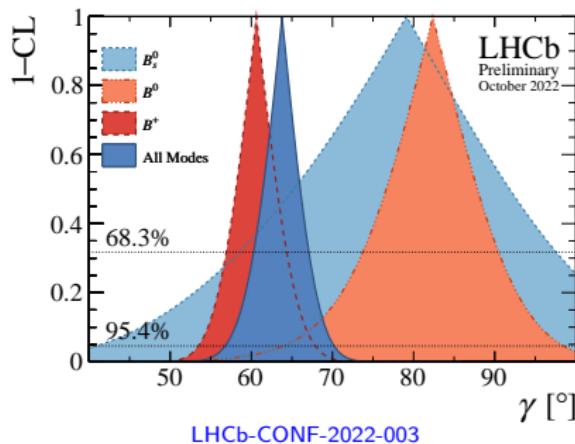
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- ③  $B^0 \rightarrow DK^{*0} \leftarrow$  New results!
- ④  $B_s^0 \rightarrow D_s^- K^+ \leftarrow$  Not covered today
  - And many more...

# The LHCb $\gamma$ and charm combination

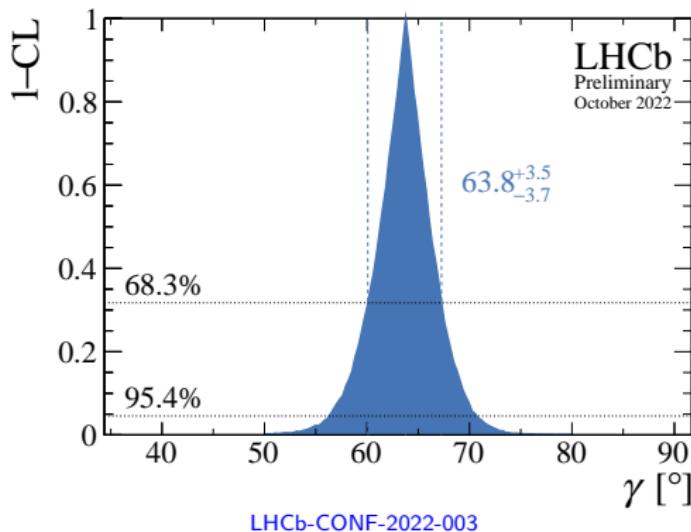
Other  $B$  decays are also interesting for  $\gamma$  measurements:

- ①  $B^\pm \rightarrow DK^\pm$
- ②  $B^\pm \rightarrow D^{*0} K^\pm$
- ③  $B^0 \rightarrow DK^{*0}$
- ④  $B_s^0 \rightarrow D_s^- K^+$
- And many more...



# The LHCb $\gamma$ and charm combination

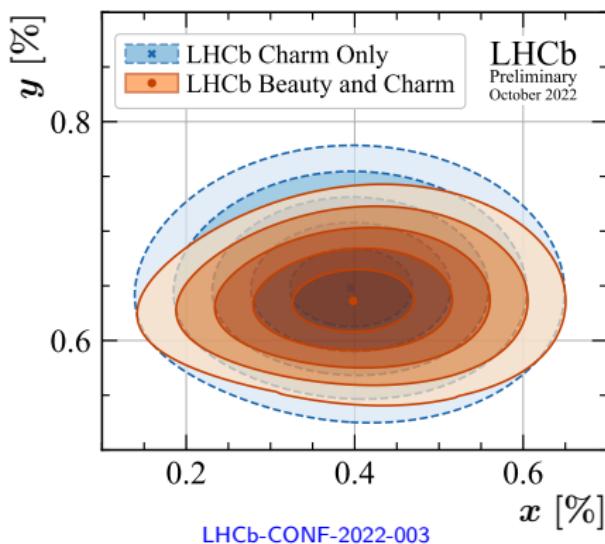
Our most precise knowledge of  $\gamma$  comes from the combination of  $\gamma$  and charm mixing parameters



This is the most precise determination of  $\gamma$  by a single experiment!  
Charged  $B^\pm$  modes dominate, but all measurements are consistent

# The LHCb $\gamma$ and charm combination

Our most precise knowledge of  $\gamma$  comes from the combination of  $\gamma$  and charm mixing parameters



Mixing effects to  $\gamma$  measurements are approaching the statistical sensitivity  
Knowledge of  $y$  is significantly improved through correlations with  $\delta_D^{K\pi}$

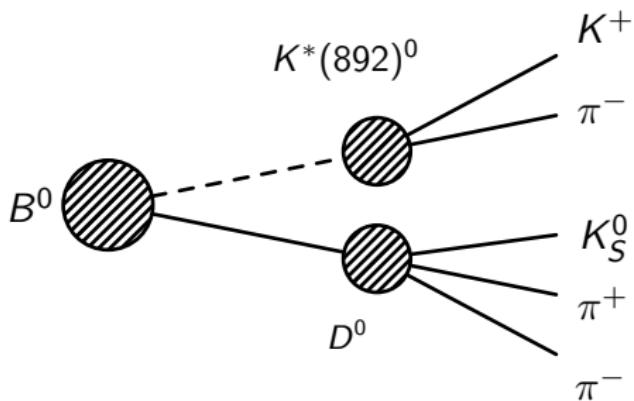
# Neutral $B$ decays

More interference with less statistics

## Neutral $B$ decays

Neutral  $B$  decays are analysed with an identical strategy:

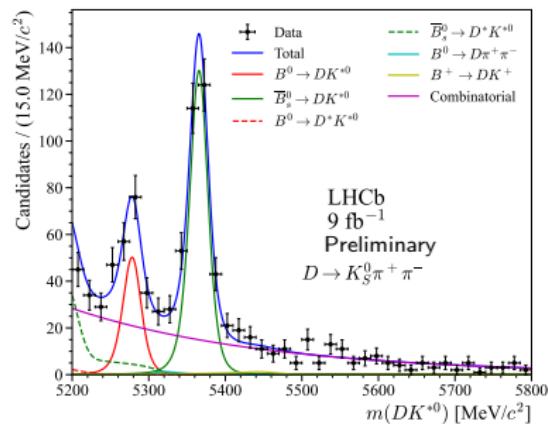
LHCb-PAPER-2023-009 (in preparation) New results!



$$B^0 \rightarrow (K_S^0 h^+ h^-)_D (K^+ \pi^-)_{K^*}$$

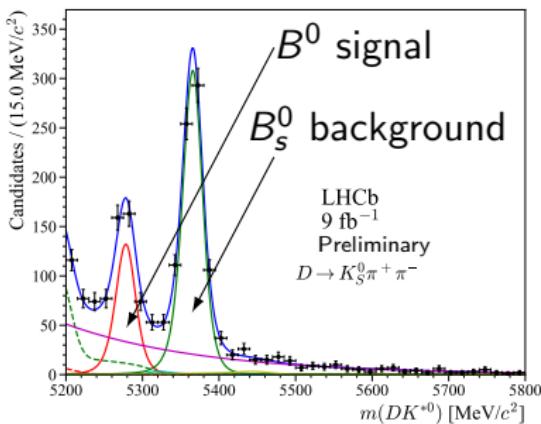
This results supersedes that of JHEP **08** (2016) 137

# Neutral $B$ decays



LL

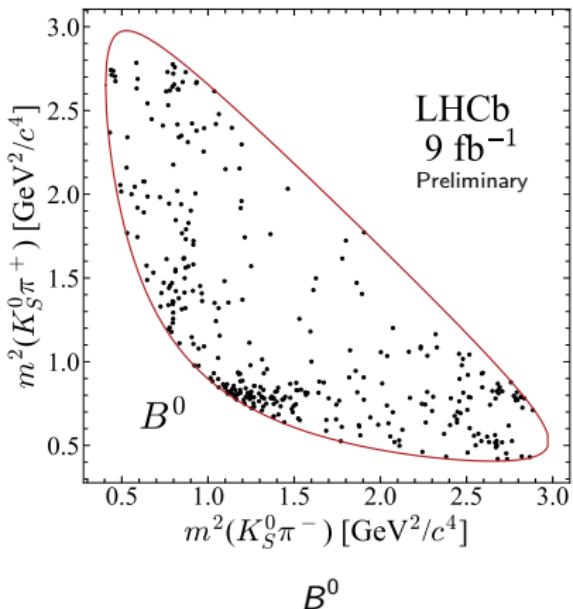
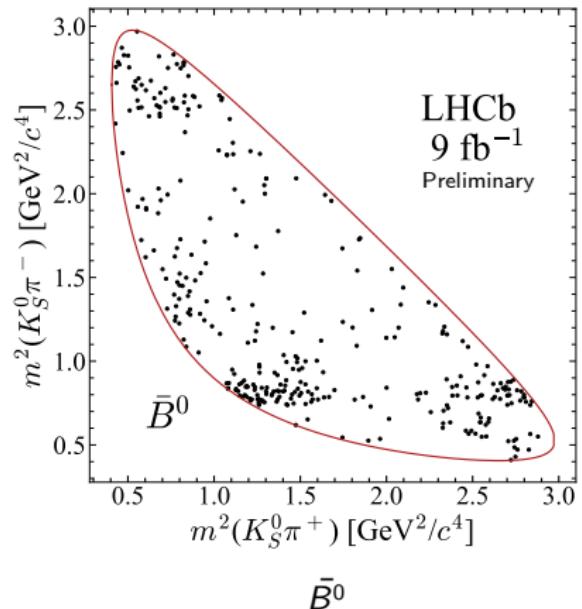
LHCb-PAPER-2023-009



DD

- Two separate selections of  $K_S^0$ :
  - LL (long tracks):  $K_S^0$  decays in the VELO
  - DD (downstream tracks):  $K_S^0$  decays downstream of the VELO
- $B^0 \rightarrow DK^{*0}$  candidates with  $D \rightarrow K_S^0 \pi^+ \pi^-$  ( $D \rightarrow K_S^0 K^+ K^-$ ):
  - LL:  $102 \pm 17$  ( $12 \pm 6$ )
  - DD:  $288 \pm 25$  ( $32 \pm 8$ )

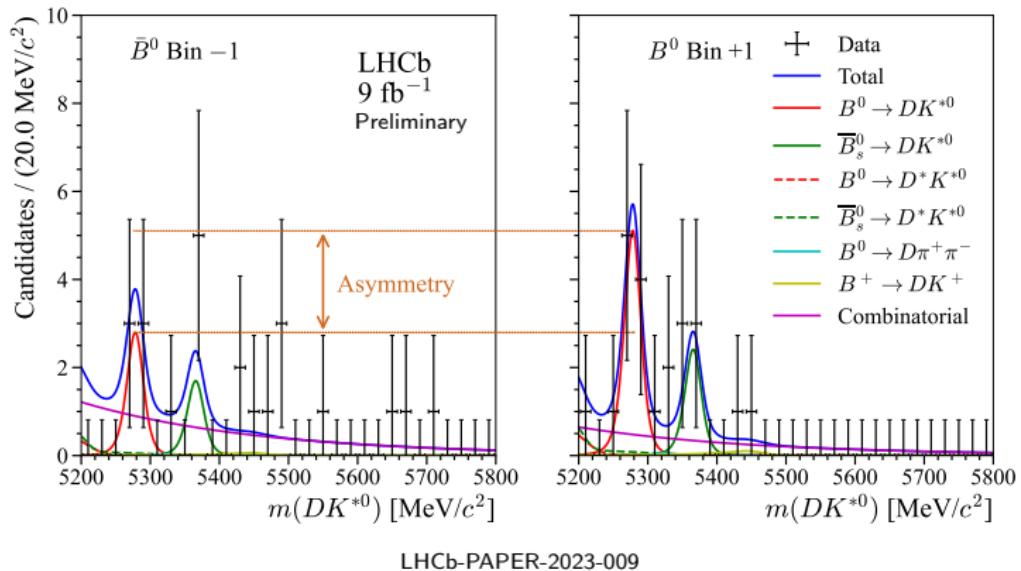
# Neutral $B$ decays



$$B^0 \rightarrow [K_S^0\pi^+\pi^-]_D K^{*0} \text{ Dalitz plots}$$

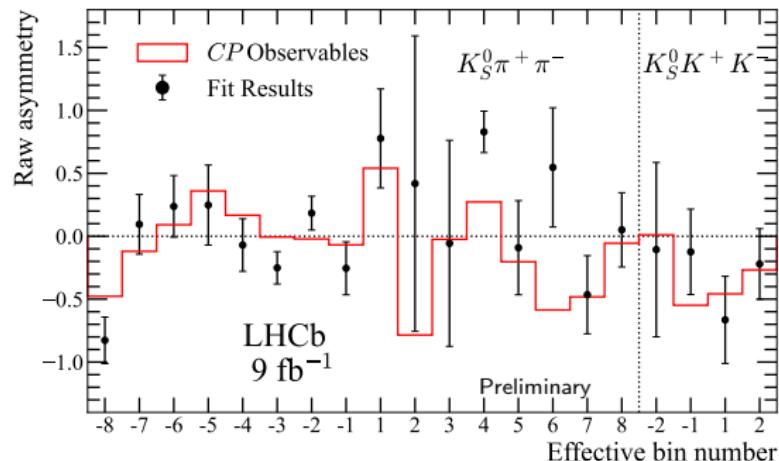
Can you find the asymmetries?

# Neutral $B$ decays



- Non-zero bin asymmetries are seen:
  - Large asymmetries between  $B^0$  ( $\bar{B}^0$ ) bin pairs
  - Very small CPV is expected in  $B_s^0$  decays, and these are not looked for

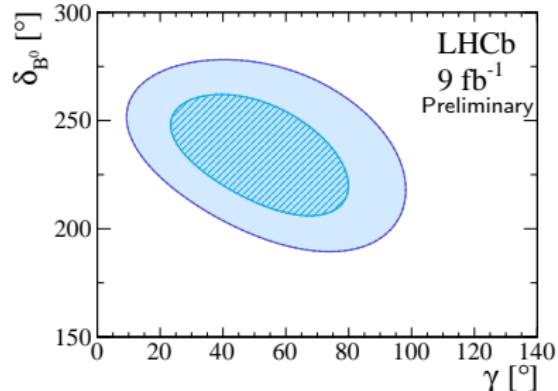
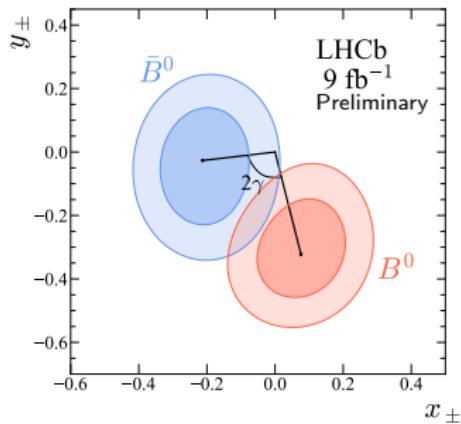
# Neutral $B$ decays



LHCb-PAPER-2023-009

- Non-zero bin asymmetries are seen:
  - Large asymmetries between  $B^0$  ( $\bar{B}^0$ ) bin pairs
  - Very small CPV is expected in  $B_s^0$  decays, and these are not looked for
- Asymmetries differ in size and magnitude across bins of phase space

# Neutral $B$ decays



LHCb-PAPER-2023-009

- Measured  $CP$ -violating observables:

$$x_{\pm} \equiv r_{B^0} \cos(\delta_{B^0} \pm \gamma) \text{ and } y_{\pm} \equiv r_{B^0} \sin(\delta_{B^0} \pm \gamma)$$

- Measured value of  $\gamma$  is consistent with world average:

- $\gamma = (49 \pm 20)^\circ$
- $\delta_{B^0} = (236 \pm 19)^\circ$
- $r_{B^0} = 0.27 \pm 0.07 \leftarrow 3 \text{ times larger than the } B^{\pm} \text{ modes!}$

B decays to excited  $D^*$  final states

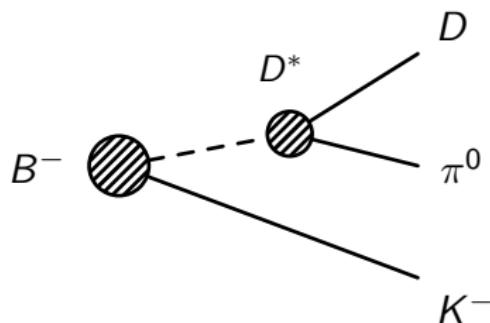
# B decays to excited $D^*$ final states

A measurements with neutral particles

## $B$ decays to excited $D^*$ final states

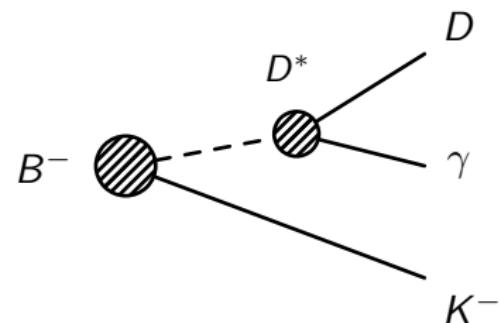
$B^- \rightarrow D^* K^-$  decays are also a powerful probe of CPV:

LHCb-PAPER-2023-012 (in preparation) New results!



$$B^- \rightarrow [D\pi^0]_{D^*} K^-$$

$$\mathcal{A}(D^0) + r_B e^{i(\delta_B - \gamma)} \mathcal{A}(\bar{D}^0)$$

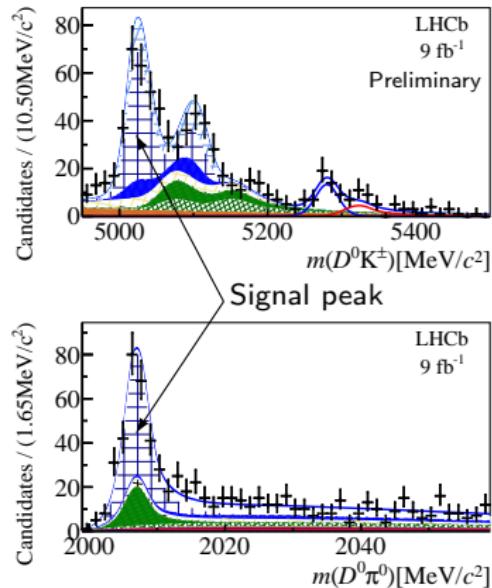


$$B^- \rightarrow [D\gamma]_{D^*} K^-$$

$$\mathcal{A}(D^0) - r_B e^{i(\delta_B - \gamma)} \mathcal{A}(\bar{D}^0)$$

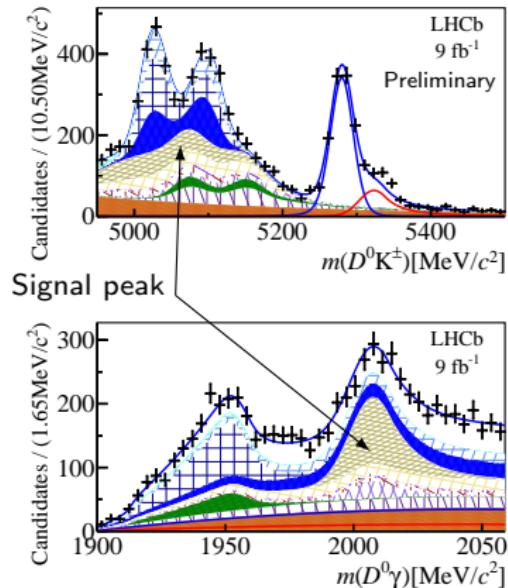
The relative sign swap, due to the phase difference of  $\pi$  between  $\pi^0$  and  $\gamma$ , results in opposite  $CP$  asymmetries between  $D^* \rightarrow D\pi^0$  and  $D^* \rightarrow D\gamma$

# $B$ decays to excited $D^*$ final states



$$D^* \rightarrow D\pi^0$$

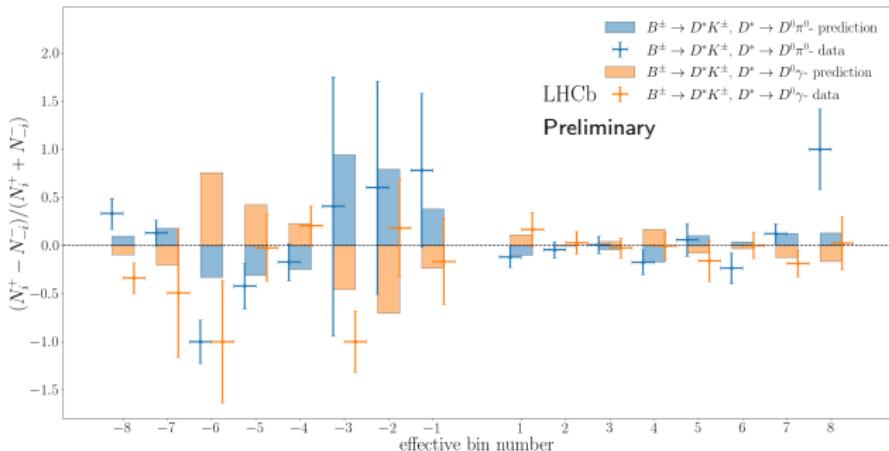
LHCb-PAPER-2023-012



(a)  $D^* \rightarrow D\gamma$

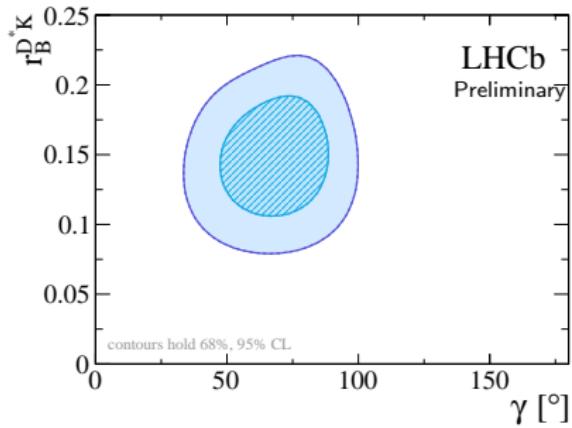
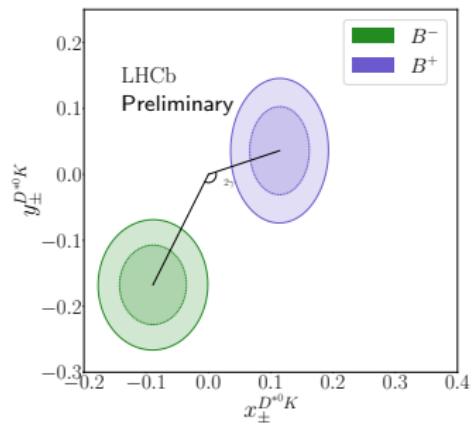
A 2D fit is necessary to separate signal from background

# $B$ decays to excited $D^*$ final states



- Good agreement between individual bin asymmetries and the combined  $CP$  fit
- Bin asymmetries between  $D^* \rightarrow D\pi^0$  and  $D^* \rightarrow D\gamma$  are generally opposite in sign

# $B$ decays to excited $D^*$ final states



LHCb-PAPER-2023-012

These results provide strong constraints on  $\gamma$ :

- $\gamma = (69 \pm 14)^\circ$
- $\delta_B^{D^*K} = (311 \pm 15)^\circ$
- $r_B^{D^*K} = 0.15 \pm 0.03$

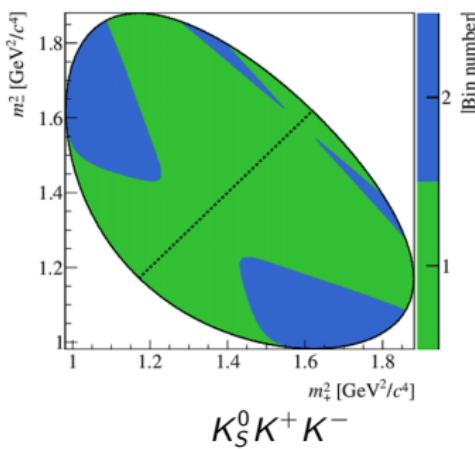
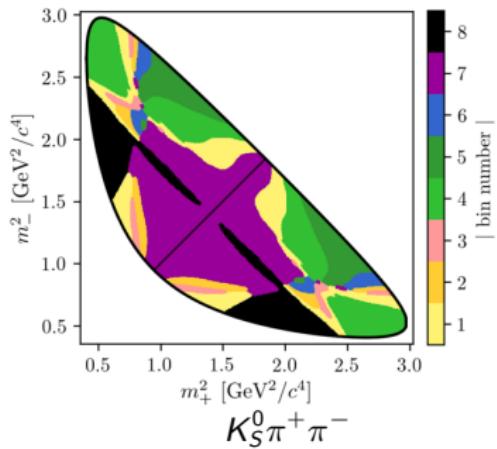
Binned four-body decay:  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

Binned four-body decay:  
 $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

A journey through five dimensions

# Binned four-body decay: $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

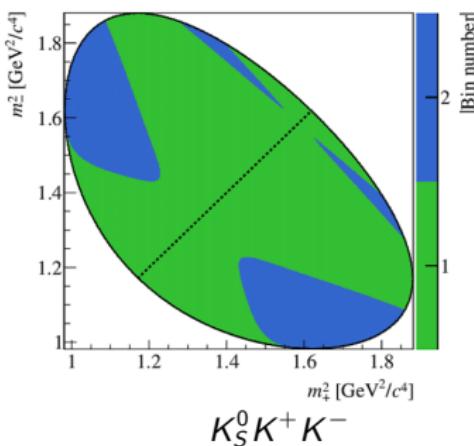
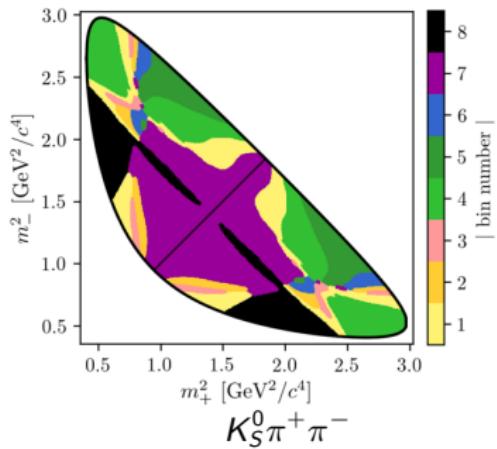
Back to  $D^0 \rightarrow K_S^0 h^+ h^-$  binning schemes, visualised on a Dalitz plot:



- Bin boundaries are optimised for sensitivity to  $\gamma$  by CLEO
- We would like to do this for  $D^0 \rightarrow K^+ K^- \pi^+ \pi^- \dots$

# Binned four-body decay: $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

Back to  $D^0 \rightarrow K_S^0 h^+ h^-$  binning schemes, visualised on a Dalitz plot:



- Bin boundaries are optimised for sensitivity to  $\gamma$  by CLEO
- We would like to do this for  $D^0 \rightarrow K^+ K^- \pi^+ \pi^- \dots$
- ... but the four-body phase space is five-dimensional!

Binned four-body decay:  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

But how do we navigate through a 5D space?



Use an amplitude model! JHEP **02** (2019) 126

Ultimately, the charm strong-phase differences will be measured directly at BESIII, and the  $\gamma$  measurement will be model independent

Binned four-body decay:  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

A binning scheme must satisfy the following:

- Minimal dilution of strong phases when integrating over bins
- Enhance interference between  $B^\pm \rightarrow D^0 K^\pm$  and  $B^\pm \rightarrow \bar{D}^0 K^\pm$

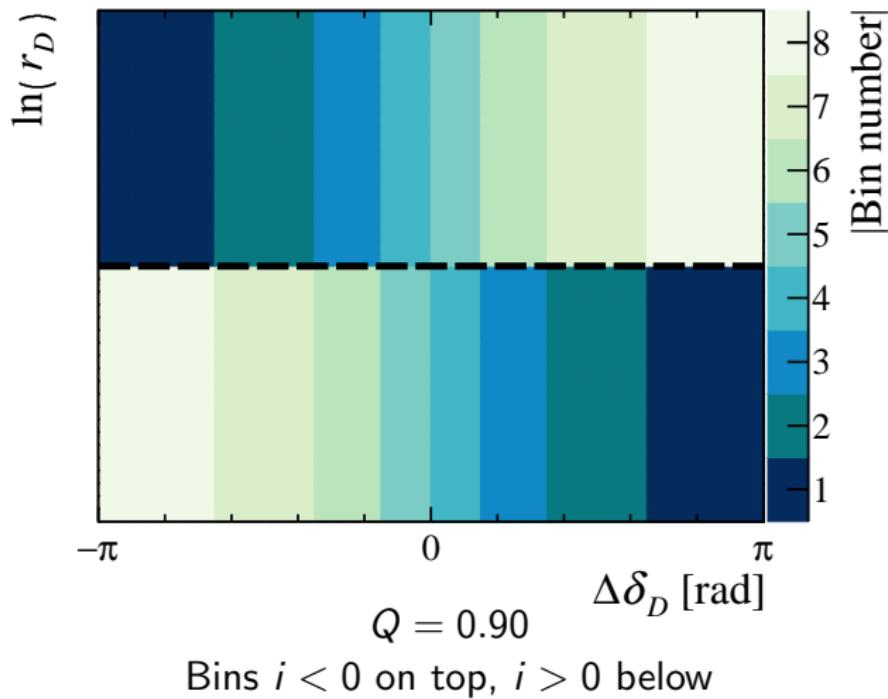
How to bin a 5-dimensional phase space?

- ① For each  $B^\pm$  candidate, use the amplitude model to calculate

$$\frac{\mathcal{A}(D^0)}{\mathcal{A}(\bar{D}^0)} = r_D e^{i\delta_D}$$

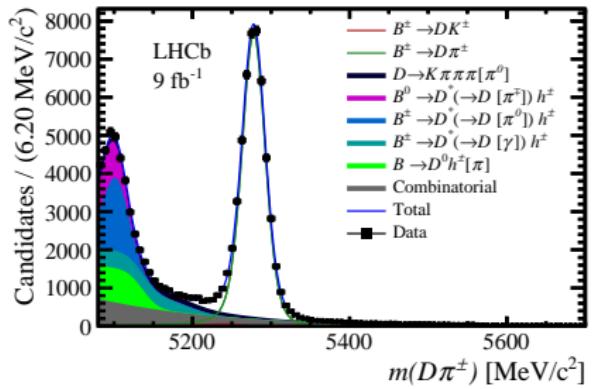
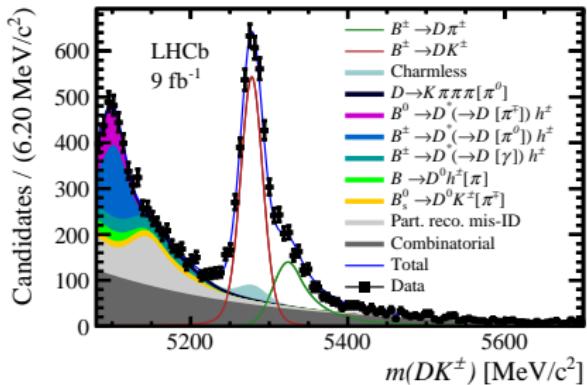
- ② Split  $\delta_D$  into uniformly spaced bins
- ③ Use the symmetry line  $r_D = 1$  to separate bin  $+i$  from  $-i$
- ④ Optimise the binning scheme by adjusting the bin boundaries in  $\delta_D$

# Binning scheme



# Phase-space binned $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D K^\pm$

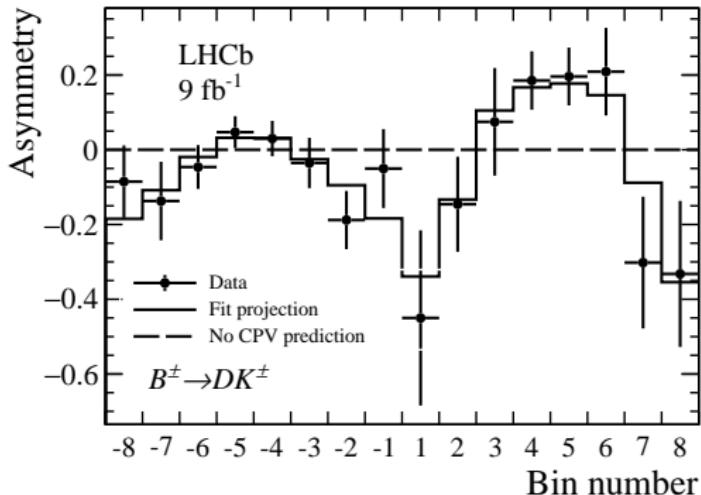
Fully charged final state  $\implies$  Highly suitable for LHCb



Eur. Phys. J. C 83, 547 (2023)

- $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^\pm$  signal yield:
  - $B^\pm \rightarrow DK^\pm$ :  $3026 \pm 38$
  - $B^\pm \rightarrow D\pi^\pm$ :  $44349 \pm 218$

# Phase-space binned $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D K^\pm$



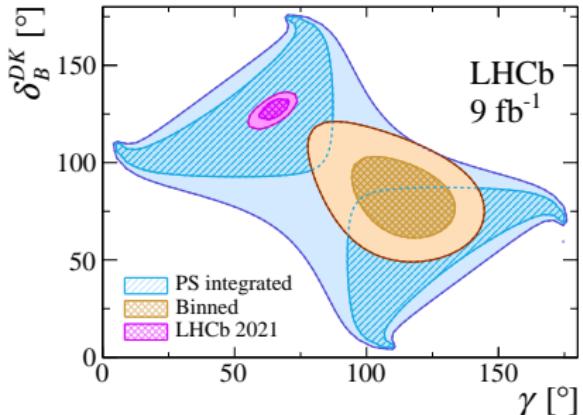
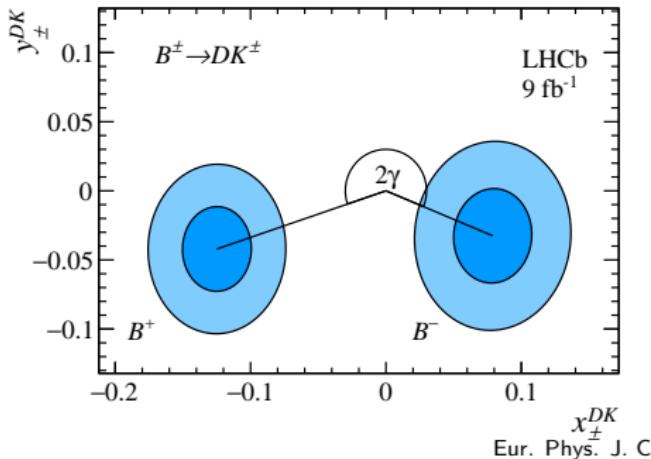
Eur. Phys. J. C 83, 547 (2023)

- Clear bin asymmetries are seen, and the non-trivial distribution is driven by the change in strong-phase differences across phase space
- While the interpretation of  $\gamma$  require charm inputs, the observed bin asymmetries are model independent

# Phase-space binned $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D K^\pm$

From the phase-space binned asymmetries, we obtain:

- $\gamma = (116^{+12}_{-14})^\circ$
- $\delta_B^{DK} = (81^{+12}_{-14})^\circ$
- $r_B^{DK} = 0.110^{+0.020}_{-0.020}$



These results are model dependent, and will be updated once BESIII  
strong-phase inputs are available

# The angle $\gamma$ of the Cabibbo-Kobayashi-Maskawa ansatz

Almost at the end of this seminar, but not the end of the journey!

# Future prospects

## Future prospects:

- The measurements presented today will make valuable improvements to future  $\gamma$  combinations
- Several interesting Run 1+2 results are in the pipeline:
  - ① Update of  $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^\pm$  with charm inputs from BESIII
  - ② GLW and ADS results from  $B^0 \rightarrow D K^{*0}$  could shed light on the (slight) tension between  $B^0$  and  $B^\pm$
  - ③ Time-dependent measurements, such as  $B_s^0 \rightarrow D_s^- K^+$  with Run 2, will be interesting to compare with results from  $B^\pm/B^0$
- LHCb, during Run 3 and 4 anticipates to collect five times more data!

# Summary and future prospects

## In summary:

- ① Long journey towards a precise determination of  $\gamma$
- ② Two recent results of  $B^\pm \rightarrow D^* h^\pm$  and  $B^0 \rightarrow DK^{*0}$  with  $D \rightarrow K_S^0 h^+ h^-$ , using external inputs from BESIII
  - Unique synergy between beauty and charm factories
- ③ A binned measurement with the channel  $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D K^\pm$  has been performed for the first time
  - Need external inputs for charm strong-phases from BESIII
- ④ LHCb is on track to reach a  $1^\circ$  precision after Run 3 and 4!

# Summary and future prospects

In summary:

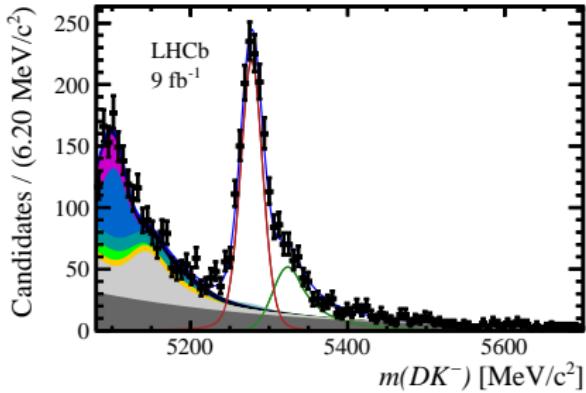
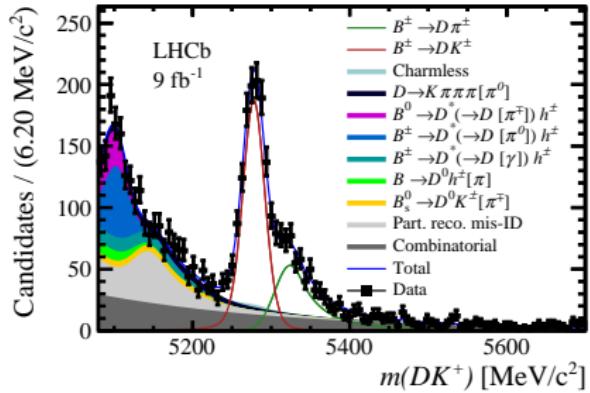
- ① Long journey towards a precise determination of  $\gamma$
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  - Need external inputs for charm strong-phases from BESIII
- ④ LHCb is on track to reach a  $1^\circ$  precision after Run 3 and 4!

Thanks for your attention!

# Backup slides

# Phase-space integrated $B^\pm \rightarrow [K^+ K^- \pi^+ \pi^-]_D K^\pm$

Additionally, one can measure the phase-space integrated asymmetries and measure additional  $CP$ -violating observables



Eur. Phys. J. C 83, 547 (2023)

More  $B^-$  candidates because  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  is predominantly  $CP$ -even (Phys. Rev. D 107 032009)