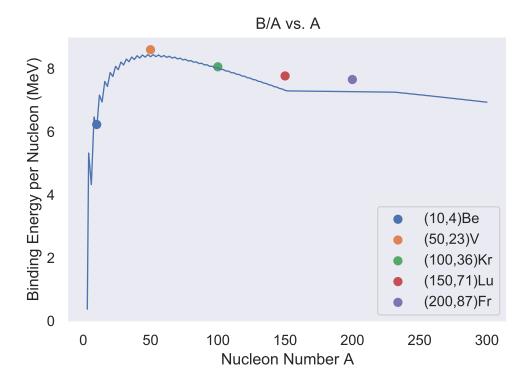
## Week 1

1. Q — **Problem 6: The semi-emperical mass formula**. Make a script to calculate the binding energy B in MeV for a given mass number A, proton number Z (and N neutrons). Compare B/A to figure 5.1.

A — Plot



2. Q — **Problem 1: Isospin properties**. In a two-nucleon state with quantum numbrs T and  $T_3$  for the total isospin and its third component, show that the expectation value of the vector product of  $t_1$  and  $t_2$  is [2T(T+1)-3]/4.

$$T = t_1 + t_2 \tag{1}$$

Taking the square to get an expression for the vector product

$$\left\langle T^{2}\right\rangle =\left\langle t_{1}^{2}\right\rangle +\left\langle t_{2}^{2}\right\rangle +2\left\langle t_{1}\cdot t_{2}\right\rangle \tag{2}$$

And using the relations

$$\langle S^2 \rangle = S(S+1), \quad \langle S_i^2 \rangle = \frac{1}{2} \left( \frac{1}{2} + 1 \right)$$
 (3)

This leads to

$$\langle t_1 \cdot t_2 \rangle = \frac{T(T+1) - \frac{1}{2} \left(\frac{1}{2} + 1\right) - \frac{1}{2} \left(\frac{1}{2} + 1\right)}{2}$$
$$= \frac{2T(T+1) - 3}{4}$$

Which is the final expression.

3. Q — Problem 5: The Deuteron and other s-waves in nuclei. Consider only a central potential between the proton and the neutron. Use U(r) and write down the Schrödinger equation in term of the radial deuteon wavefunction u(r)/r. Solve this either analytically or numerically and plot the wavefunction

A — Considering the potential between the proton and the neutron given by

$$U(r) = \begin{cases} -U_0, & r \le R \\ 0 & r > R \end{cases}$$

The radial equation is given by

$$-\frac{\hbar^2}{2m}\frac{d^2u(r)}{dr^2} + \left[U + \frac{\hbar^2l(l+1)}{2mr^2}\right]u(r) = Eu(r)$$
 (4)

This is identical to the one-dimensional Schrodinger equation with an effective potential, where the centrifugal term pushes the particle outward. To solve this analytically I rewrite the equation and consider the boundary conditions.

$$\frac{d^2 u(r)}{dr^2} = \left[ -k^2 + \frac{l(l+1)}{r^2} \right] u(r), \quad \text{where} \quad k = \frac{\sqrt{2m(E+U_0)}}{\hbar}$$
 (5)

For l = 0

$$\frac{\mathrm{d}^2 u(r)}{\mathrm{d}r^2} = -ku(r) \Rightarrow u(r) = A\sin(kr) + B\cos(kr) \tag{6}$$

Since R(r) = u(r)/r and  $\cos(kr)/r$  blows up as  $r \to 0 \Rightarrow B = 0$  and the solution is

$$u(r) = A\sin(kr), \quad r \le R \tag{7}$$

For r > R, l = 0 and U(r) = 0

$$\frac{\mathrm{d}^2 u(r)}{\mathrm{d}r^2} = -\frac{2m}{\hbar^2} = \kappa^2 u(r) \Rightarrow u(r) = Ce^{\kappa r} + De^{-\kappa r}, \quad \kappa = \frac{\sqrt{-2mE}}{\hbar}$$
 (8)

Since  $Ce^{\kappa r}$  blows up as  $r \to \infty$ . Now I can use the fact that the solutions must match at r = R. This must be true for both u(r) and u(r)'. This leads to two equations for r = R

$$A\sin(kR) = De^{-\kappa R} \tag{9}$$

$$Ak\cos(kR) = -D\kappa e^{-\kappa R} \tag{10}$$

This leads to

$$-\cot(kR) = \frac{\kappa}{k} \tag{11}$$

Plugging in an appropiate value for R = 2.127 fm yields

$$U_0 = \frac{\hbar^2 \pi^2}{2mR^2} - E$$
$$= \frac{\hbar^2 \pi^2}{2mR^2} + B$$
$$= 24.82 \,\text{MeV}$$

