DYNAMIC AFM



-1- Amplitude Modulation Mode (or <u>tapping</u> mode) AM-AFM

THIS IS WHAT WE HAVE STUDIED IN THE QBIO COURSE

- Constant Amplitude (feedback loop)
- Intermittent mechanical contact between tip and sample
- Tip position (z) described by non-linear harmonic oscillator → no analytical solutions, but numerical, often multiple, solutions exist

$$m\ddot{z} + \delta\dot{z} + kz = F_{interaction}(z) + F_0(\omega t)$$

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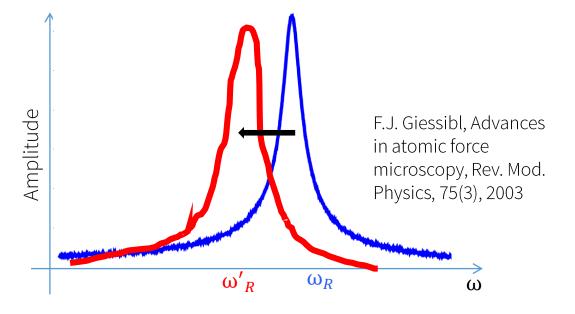
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-2- Frequency Modulation Mode – FM-AFM WE DID NOT MENTION THIS ONE IN THE QBIO COURSE

T. R. Albrecht, P. Gritter, F. Horne, D. Rugar, "Frequency modulation detection using high-q cantilevers for enhanced force microscope sensitivity", J. Appl. Phys, vol. 69, pp. 668-673, 1991

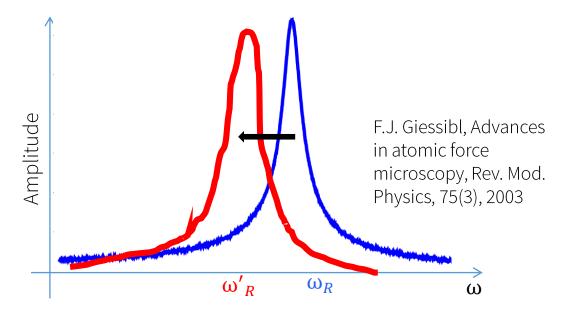


If amplitude of tip oscillation: 0.1nm – 1nm (ten times lower than in AM-AFM), Because of $F_{interaction}(z)$, the resonance frequeny ω_R shifts to a new frequency ω'_R



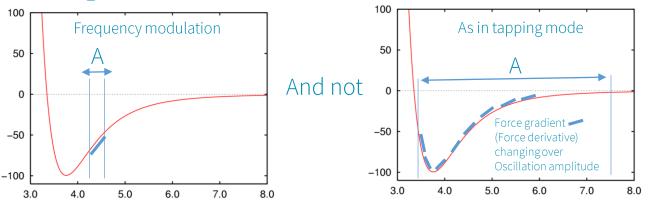


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 $m\ddot{z} + \delta\dot{z} + kz = F_{interaction}(z) + F_{0(\omega t)}$ can be rewritten if two conditions are fulfilled

1. If oscillation amplitude A explores a constant Force gradient over the oscillation, meaning



2. If cantilever spring constant k is lower than the Force gradient $k < dF_{interaction}/dz$



If the two conditions are fulfilled,

$$m\ddot{z} + \delta\dot{z} + kz = F_{interaction}(z) + F_0(\omega t)$$

We can use Taylor expansion for
$$F_{interaction} = F_{int}$$
 around its oscillation center Z0 $m\ddot{z} + \delta \dot{z} + kz = F_{int}(z - z0) + \left(\frac{dF_{int}(z - z0)}{dz}\right)(z - z0) + \dots + F_0(\omega t)$

Where $\left(\frac{dF_{int}(z-z_0)}{dz}\right)$ is the force gradient. Now it is the equation of a linear harmonic oscillator that can be solved analytically as we have done in absence of interaction force, leading to an harmonic solution equal to the one we obtained, except for the new resonance. We can indeed define a new resonance frequency ω'_R that replaces ω_R

$$\omega'_{R} \approx \sqrt{\frac{k - \frac{dF_{int}}{dz}}{m}}$$

$$A = |\mathbf{z}| = \frac{\frac{F_{0}}{m}}{\sqrt{(\omega_{R}^{2} - \omega^{2})^{2} + \omega^{2}\delta^{2}}}$$
Phase = $\arctan(\varphi) = \frac{\delta\omega}{\omega_{R}^{2} - \omega^{2}}$



Defining the frequency shift as df = ω'_R - ω_R

$$df = \sqrt{\frac{k - \frac{dF_{int}}{dz}}{m}} - \sqrt{\frac{k}{m}}$$

Using k < d
$$F_{int}$$
/dz, we can write $\sqrt{\frac{k - \frac{dF_{int}}{dz}}{m}}$ as $\approx (1 - \frac{\frac{dF_{int}}{dz}}{2k})\sqrt{\frac{k}{m}}$

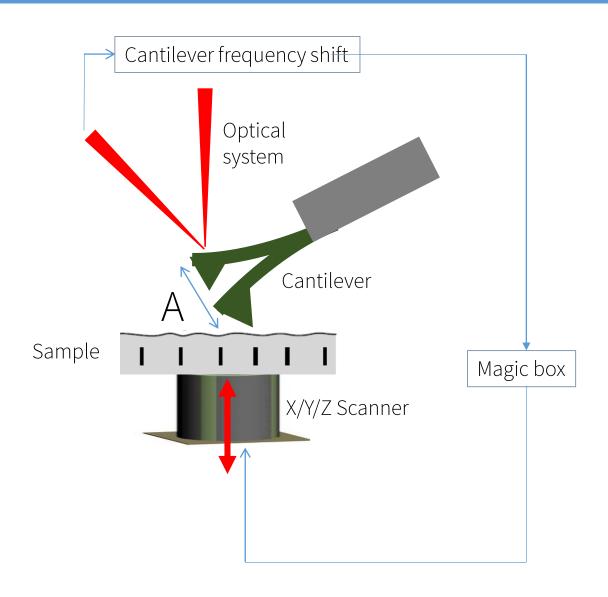
leading to df =
$$(1 - \frac{\frac{dF_{int}}{dz}}{2k})\sqrt{\frac{k}{m}} - \sqrt{\frac{k}{m}} = -\frac{\frac{dF_{int}}{dz}}{2k}\sqrt{\frac{k}{m}} = -\frac{\frac{dF_{int}}{dz}}{2k}\omega_R$$

Which means that in frequency modulation mode,

the frequency shift is proportional to the tip-sample interaction force gradient $\frac{dF_{int}}{dz}$

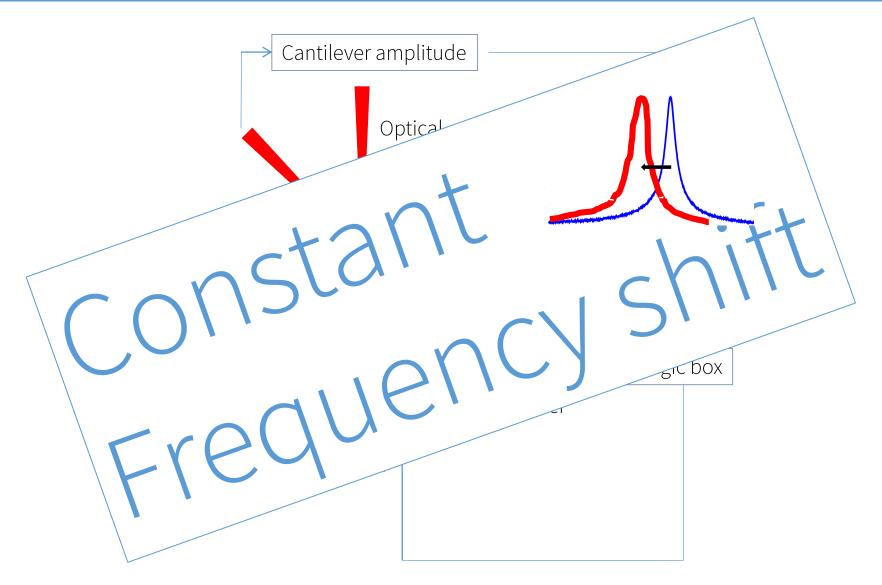
AFM operational scheme





AFM operational scheme





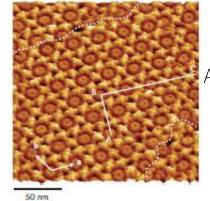
Frequency modulation mode



Mainly used for AFM in vacuum High Q-factor is important

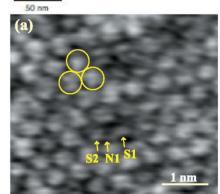
Bio (in liquid) application on FLAT samples

Hirofumi Yamada



Antibody crystal S. Ido et al. Nature Materials, 13, 264-270, 2014

Takeshi Fukuma



DPPC bilayer T. Fukuma et al. Physical Review Letters, 98(106101), 2007